Spacecraft Dynamics and Control

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Lecture 18: Feedback Control of Attitude Dynamics

Attitude Dynamics

In this Lecture we will cover:

An Review of Feedback control for Attitude Dynamics

- Transfer Functions
- PID Control
- Root Locus

Problem: 3-axis Stabilization

- Detumble (if $\vec{\omega} \not\cong 0$)
- Attitude Tracking (assuming $\vec{\omega} \cong 0$)

Linearizing Euler's Equations

Recall:

$$\begin{bmatrix} L \\ M \\ N \end{bmatrix} = \begin{bmatrix} I_{xx}\dot{\omega}_x + \omega_y\omega_z(I_{zz} - I_{yy}) \\ I_{yy}\dot{\omega}_y + \omega_x\omega_z(I_{xx} - I_{zz}) \\ I_{zz}\dot{\omega}_z + \omega_x\omega_y(I_{yy} - I_{xx}) \end{bmatrix}$$

Non-axisymmetric Case $I_x \neq I_y \neq I_z$.

We first assume the spacecraft has been detumbled, so we have Small Spin Assumption: $\omega_x = \omega_y = \omega_z \cong 0$.

Nominal motion is

$$\omega_0(t) = \begin{bmatrix} \omega_{x,0}(t) \\ \omega_{y,0}(t) \\ \omega_{z,0}(t) \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \\ 0 \end{bmatrix}$$

In this case, the Linearized dynamics become:

$$\begin{bmatrix} L \\ M \\ N \end{bmatrix} = \begin{bmatrix} I_x \dot{\omega}_x \\ I_y \dot{\omega}_y \\ I_z \dot{\omega}_z \end{bmatrix} \qquad \text{or} \qquad \begin{bmatrix} \dot{\omega}_x \\ \dot{\omega}_y \\ \dot{\omega}_z \end{bmatrix} = \begin{bmatrix} \frac{M}{I_z} \\ \frac{M}{I_y} \\ \frac{N}{I_z} \end{bmatrix}$$

The Dynamics are all uncoupled.

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Kinematics and Euler Angles

Problem: We don't measure rotation rates. We measure rotation angles.

• Now we need to choose an inertial coordinate system.

The Euler Angles define the transformation from the body-fixed to inertial coordinates

$$\begin{bmatrix} p \\ q \\ r \end{bmatrix} = \begin{bmatrix} \omega_x \\ \omega_y \\ \omega_z \end{bmatrix} = \vec{\omega}_{/B} = R_1(\phi)R_2(\theta)R_3(\psi)\vec{\omega}_{/I} = R(\phi)R(\theta)R(\psi) \begin{bmatrix} \dot{\phi} \\ \dot{\theta} \\ \dot{\psi} \end{bmatrix}$$
$$= \begin{bmatrix} \dot{\phi} - \dot{\psi}\sin\theta \\ \dot{\theta}\cos\phi + \dot{\psi}\cos\theta\sin\phi \\ \dot{\psi}\cos\theta\cos\psi - \dot{\theta}\sin\phi \end{bmatrix}$$

Of course, we often have $\vec{\omega}_{/B}$ and are trying to find $\vec{\omega}_{/I}$. In this case, the rotation matrices can be inverted to obtain

$$\begin{bmatrix} \dot{\phi} \\ \dot{\theta} \\ \dot{\psi} \end{bmatrix} = \begin{bmatrix} p + (q\sin\phi r\cos\phi)\tan\theta \\ q\cos\phi - r\sin\phi \\ (q\sin\phi + r\cos\phi)\sec\theta \end{bmatrix}$$

Notice the Singularity at $\theta=\pm90^{\circ}$ (can be avoided with quaternions). Equations are also different for 2-1-3 and 1-2-3 rotation sequences.

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Kinematics and Euler Angles

$$\begin{bmatrix} \dot{\phi} \\ \dot{\theta} \\ \dot{\psi} \end{bmatrix} = \begin{bmatrix} p + (q\sin\phi r\cos\phi)\tan\theta \\ q\cos\phi - r\sin\phi \\ (q\sin\phi + r\cos\phi)\sec\theta \end{bmatrix}$$

where

$$\begin{bmatrix} \dot{p} \\ \dot{q} \\ \dot{w} \end{bmatrix} = \begin{bmatrix} \dot{\omega}_x \\ \dot{\omega}_y \\ \dot{\omega}_z \end{bmatrix} = \begin{bmatrix} \frac{L}{I_x} \\ \frac{M}{I_y} \\ \frac{N}{I_z} \end{bmatrix}$$

Which is a set of 6 nonlinear coupled differential equations.

- We have already linearized the second set.
- We should also linearize the first set.
 - Assume $\theta \cong 0$, $\phi \cong 0$, and $\psi \cong 0$.

$$\begin{bmatrix} \dot{\phi} \\ \dot{\theta} \\ \dot{\psi} \end{bmatrix} = \begin{bmatrix} p \\ q \\ r \end{bmatrix}$$

Decoupling the Equations

Combining these two sets of equations, we get

$$\begin{bmatrix} \ddot{\phi} = \frac{L}{I_x} \\ \ddot{\theta} = \frac{M}{I_y} \\ \ddot{\psi} = \frac{N}{I_z} \end{bmatrix}$$

If L, M, and N are decoupled, we can decouple the equations

• Orthogonal reaction wheel for each body-fixed axis.

Lets design a controller for roll

$$\ddot{\phi} = \frac{L}{I_x}$$

Consider Proportional Feedback

$$\frac{L(t)}{I_x} = -K(\phi(t) - \phi_0)$$

Then the closed-loop poles are at $s=\pm i\sqrt{K}$

Neutrally Stable

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PD Control

2nd-order system

Lets look at the effect of PD control on a 2nd-order system:

$$\hat{G}(s) = \frac{1}{s^2 + bs + c}$$

Controller: $\hat{K}(s) = -K [1 + T_D s]$ Closed Loop Transfer Function:

$$\frac{\hat{K}(s)\hat{G}(s)}{1+\hat{K}(s)\hat{G}(s)} = \frac{K[1+T_Ds]}{s^2+bs+c+K[1+T_Ds]}$$
$$= \frac{K[1+T_Ds]}{s^2+(b+KT_D)s+(c+K)}$$

The poles of the system are freely assignable for a 2nd order system.

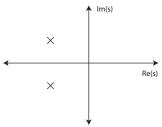
• T_D and K allow us to construct any denominator we desire.

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PD Control

2nd-order system

Suppose we want poles at $s = p_1, p_2$.



• We want the closed loop of the form:

$$\frac{1}{(s-p_1)(s-p_2)} = \frac{1}{(s^2 - (p_1 + p_2)s + p_1p_2)}$$

Thus we want

•
$$c + K = p_1 p_2$$

•
$$b + KT_D = -(p_1 + p_2)$$

$$\begin{array}{ll} \bullet \ c + K = p_1 p_2 & \text{which means } K = p_1 p_2 - c. \\ \bullet \ b + K T_D = -(p_1 + p_2) & \text{which means } T_D = -\frac{p_1 + p_2 + b}{K} = -\frac{p_1 + p_2 + b}{p_1 p_2 - c} \end{array}$$

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PD feedback gives Total Control over a 2nd-order system.

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Pole Locations

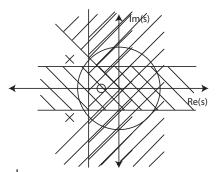
Multiple Constraints

Factor in Constraints on the Step Response:

$$\sigma < -\frac{4.6}{T_{s,desired}}$$

$$\omega_d < \frac{\pi}{\ln(M_{p,desired})} \sigma$$

$$\omega_n > \frac{1.8}{T_{r,desired}}$$



Any pole locations not prohibited are allowed.

- σ is real part of s
- ω_d is imaginary part of s
- ω_n is magnitude of s

- T_s is setlling time
- M_p is percent overshoot
- T_r is rise time

Steady-State Error

Definition 1.

Steady-State Error (e_{ss}) for a stable system G is the final difference between input and output.

$$e_{ss} = \lim_{t \to \infty} u(t) - y(t)$$

Theorem 2 (Final Value Theorem).

$$\lim_{t\to\infty}y(t)=\lim_{s\to 0}s\hat{y}(s)$$

For Step response: $\hat{u} = \frac{1}{s}$. So if G(s) is the transfer function of G,

$$e_{ss} = \lim_{s \to 0} \frac{1}{1 + G(s)K(s)}$$

Now, for roll control, $G(s) = \frac{1}{s^2}$,

$$e_{ss} = \lim_{s \to 0} \frac{1}{1 + G(s)K(s)} = \lim_{s \to 0} \frac{s^2}{s^2 + KT_D s + K} = \frac{0}{K} = 0$$

So $e_{ss} = 0!$

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Steady-State Error

Ramp Response - Important for tracking (Since spacecraft is moving). In this case the input is

$$\hat{u}(s) = \frac{1}{s^2}$$

The stead-state error can be found as:

$$e_{ss} = \lim_{s \to 0} \left(\frac{1}{s^2} - \frac{G(s)K(s)}{1 + G(s)K(s)} \frac{1}{s^2} \right) s = \lim_{s \to 0} \frac{1}{(1 + G(s)K(s)) s}$$

Now, for roll control, $G(s) = \frac{1}{s^2}$,

$$e_{ss} = \lim_{s \to 0} \frac{1}{s(1 + G(s)K(s))} = \lim_{s \to 0} \frac{s}{s + KT_D s + K} = \frac{0}{K} = 0$$

So still, $e_{ss} = 0!$

Conclusion: Integral Control is not necessary in space!!

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Lyapunov Stability for Detumbling Spacecraft

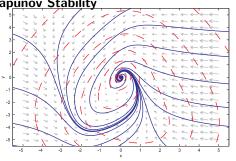
Controller Design for Nonlinear Dynamics

A VERY Brief Introduction to Lyapunov Stability

Consider a Nonlinear ODE

$$\dot{x}(t) = f(x(t))$$

with $x(0) \in \mathbb{R}^n$.



Theorem 3 (Lyapunov Stability).

Suppose there exists a continuous V and $\alpha, \beta, \gamma > 0$ where

$$V(x) > 0$$
 and

$$\dot{V}(x(t)) = \nabla V(x(t))^T f(x(t)) < 0$$

for all $x \in \mathbb{R}^n$. Then $\dot{x} = f(x)$ is Globally Stable.

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-Lyapunov Stability for Detumbling Spacecraft



The literature on Lyapunov functions is vast

- BY FAR the most commonly used tool for control of nonlinear systems.
- Represents the potential energy of a particular state.
 - Potential Energy as measured by the size (square integral) of the resulting trajectory.

Lyapunov Stability for Detumbling Spacecraft

The Dynamics of a tumbling spacecraft with input torque u(t):

ullet Uses the matrix form of cross-product $(\omega_{ imes})$ and inertia tensor (I)

$$I\dot{\omega}(t) = -\omega_{\times}(t)I\omega(t) + u(t)$$

Now we Propose a Lyapunov Function:

$$V(\omega) = \frac{1}{2}\omega^T I\omega = \frac{1}{2} ||I^{\frac{1}{2}}\omega||^2 > 0$$

Take the time-derivative of this Lyapuonv Function:

• Uses the matrix form of cross-product (ω_{\times})

$$\dot{V}(t) = \omega(t)^T I \dot{\omega}(t) = \omega(t)^T (-\omega_{\times}(t)\omega(t) + u(t))$$
$$= -\omega(t)^T (\omega(t) \times \omega(t)) + \omega(t)^T u(t))$$
$$= \omega(t)^T u(t)$$

Choose Controller

$$u(t) = -P\omega(t)$$
 where $P > 0$

Since all the eigenvalues of P are positive, $\omega^T P \omega > 0$ and hence

$$\dot{V}(t) = -\omega(t)P\omega(t) < 0$$

Which proves global stability!

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Lecture 18

Lyapunov Stability for Detumbling Spacecraft

Lyapunov Stability for Detumbling Spacecraft

The Dynamics of a tumbling spacecraft with input torque u(t):

• Uses the matrix form of rosus-poduct (ω_s) and inertia tensor (I) $L^2(t) = -\omega_s(t)L_0(t) + u(t)$ Now we Proson a Lyapunov Eurocise.

 $V(\omega) = \frac{1}{2}\omega^T L \omega = \frac{1}{2}\|I^2\omega\|^2 > 0$ Take the time-derivative of this Lyapuonv Function:
• Uses the matrix form of cross-product (ω_n) $\dot{V}(t) = \omega_t(t)^T I\dot{\omega}(t) = \omega(t)^T (-\omega_n(t)\omega(t) + n(t))$ $= -\omega_t(t)^T I\dot{\omega}(t) \times \omega(t) + \omega_t(t)^T u(t)$

Choose Controller

 $u(t) = -P\omega(t) \qquad \text{where} \quad P>0$ Since all the eigenvalues of P are positive, $\omega^TP\omega>0$ and hence $V(t) = -\omega(t)P\omega(t) < 0$ Which proves global stability!

- ullet I is a positive matrix, so we can take its square root.
- Note this does not control to any particular orientation
- Assumes spacecraft capable of large torques
- Typically we use nonlinear control for detumble and piecewise-linearized control for attitude tracking.

Attitude Dynamics

In this Lecture we have covered: Kinematics Coupled with Dynamics

Linearized to uncoupled version

Feedback Control

- Proportion-Differential Feedback
- Steady-State Error.