

DOING PHYSICS WITH PYTHON

COMPUTATIONAL OPTICS

RAYLEIGH-SOMMERFELD 1

DIFFRACTION INTEGRAL:

GAUSSIAN BEAM PROPAGATION

FROM A CIRCULAR APERTURE

Ian Cooper

Please email me any corrections, comments, suggestions or additions: **matlabvisualphysics@gmail.com**

DOWNLOAD DIRECTORIES FOR PYTHON CODE

[Google drive](#)

[GitHub](#)

[emRSGb01.py](#) Irradiance: planes

[emRSGb01Z.py](#) Irradiance: optical axis

Reference: [Gaussian Beams](#)

INTRODUCTION

The **Rayleigh-Sommerfeld diffraction integral of the first kind** is used to calculate the intensity of a **Gaussian beam** diffracted by a circular aperture.

The geometry of the aperture and observation spaces is shown in figure 1 and figure 2 shows an outline of how the RS1 diffraction integral is computed in Python.

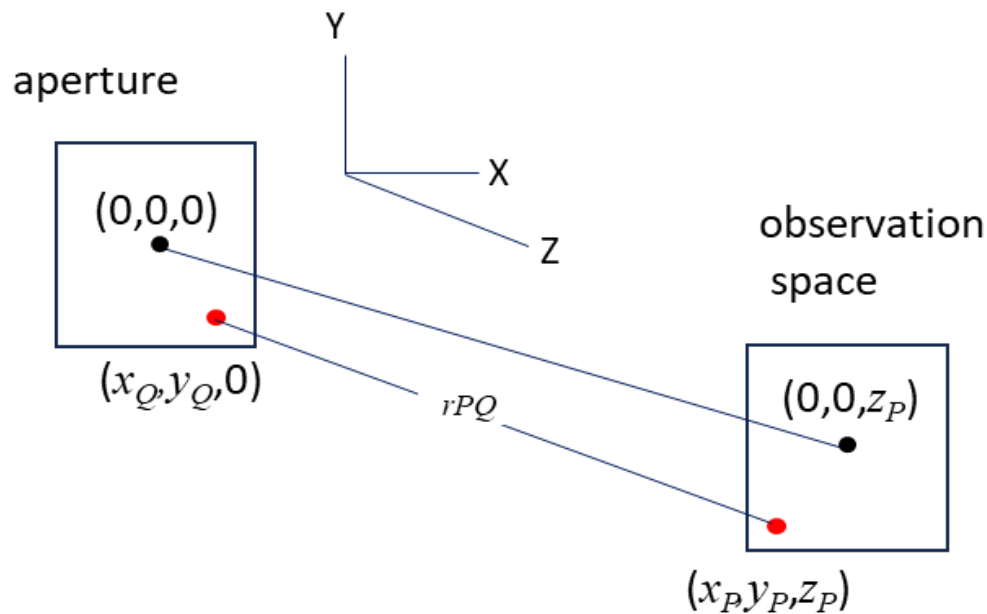




Fig. 1. Geometry of the aperture and observation spaces.

Rayleigh-Sommerfeld diffraction integral


$$E_P = \frac{1}{2\pi} \iint_{S_A} E_Q \frac{e^{jk r_{PQ}}}{r_{PQ}^3} z_P (jk r_{PQ} - 1) dS$$


 numerical integration:
 [2D] Simpson's 1/3 rule


$$E_P(x_P, y_P, z_P) = z_P \sum_{m=1}^{n_Q} \sum_{n=1}^{n_Q} \left(\left(\frac{e^{jk r_{PQmn}}}{r_{PQmn}^3} \right) (jk r_{PQmn} - 1) (E_{Qmn} S_{mn}) \right)$$




nPxnP matrix **EP**



nQxnQ
matrix **MP**



nQxnQ
matrix **EQ**



nQxnQ
matrix **S**

Fig. 2. Matrices used in computed the diffraction integral.

The electric field E_Q within the circular aperture of radius a has a Gaussian [2D] profile.

$$E_Q(x_Q, y_Q) = \exp\left(-\frac{x_Q^2 + y_Q^2}{2s^2}\right)$$

where s^2 is the variance of the Gaussian [2D] profile. Figure 3 shows a [3D] view of the Gaussian beam. In the far-field, all XY planes have a Gaussian profile for the irradiance. As the the z distance increases from the centre of the aperture, the beam spreads and the peak irradiance decreases.

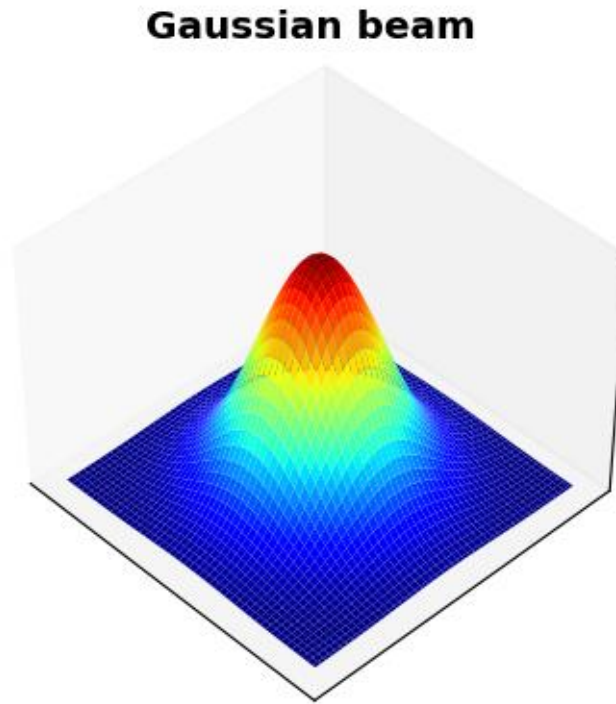


Fig. 3. [3D] view for the irradiance in an XY plane for the Gaussian beam. **emRSGB01.py**

For our Gaussian beam

$w(z)$ **beam spot** and w_0 is the **beam waist**

$w(z)$ is the radius of the beam at position z at which the irradiance is $1/e^2$ of its axial value. At position z along the axis, the beam spot $w(z)$ is given by

$$(1) \quad w(z) = w_0 \sqrt{1 + \left(\frac{\lambda z}{\pi w_0^2} \right)^2} \quad w(0) = w_0$$

Also, the beam spot is calculated numerically in the Python Code **emRSGB01.py**

SIMULATIONS

In running the Python code **emRSGB01.py** a summary of the input and output parameters are displayed in the Console Window.

$nQ = 199$ $nP = 237$

aperture radius = 0.400 mm

Gaussian beam: width $s = 0.200$ mm

Wavelength $wL = 632.8$ nm

Observation space: $zP = 1.500$ mm

Gaussian beam waist $w0/a = 0.707$

Gaussian beam spot (theoretical) $wT/a = 2.763$

Relative max irradiance (normalized to 1 at $zP = 1.0$ m) = 62.550

Execution time 165 s

The results of each simulation are displayed a set of Figure Windows. The irradiance **IXY** is normalized such that the peak value of the irradiance for **zP** = 1.00m is set to 1.

The Aperture Space

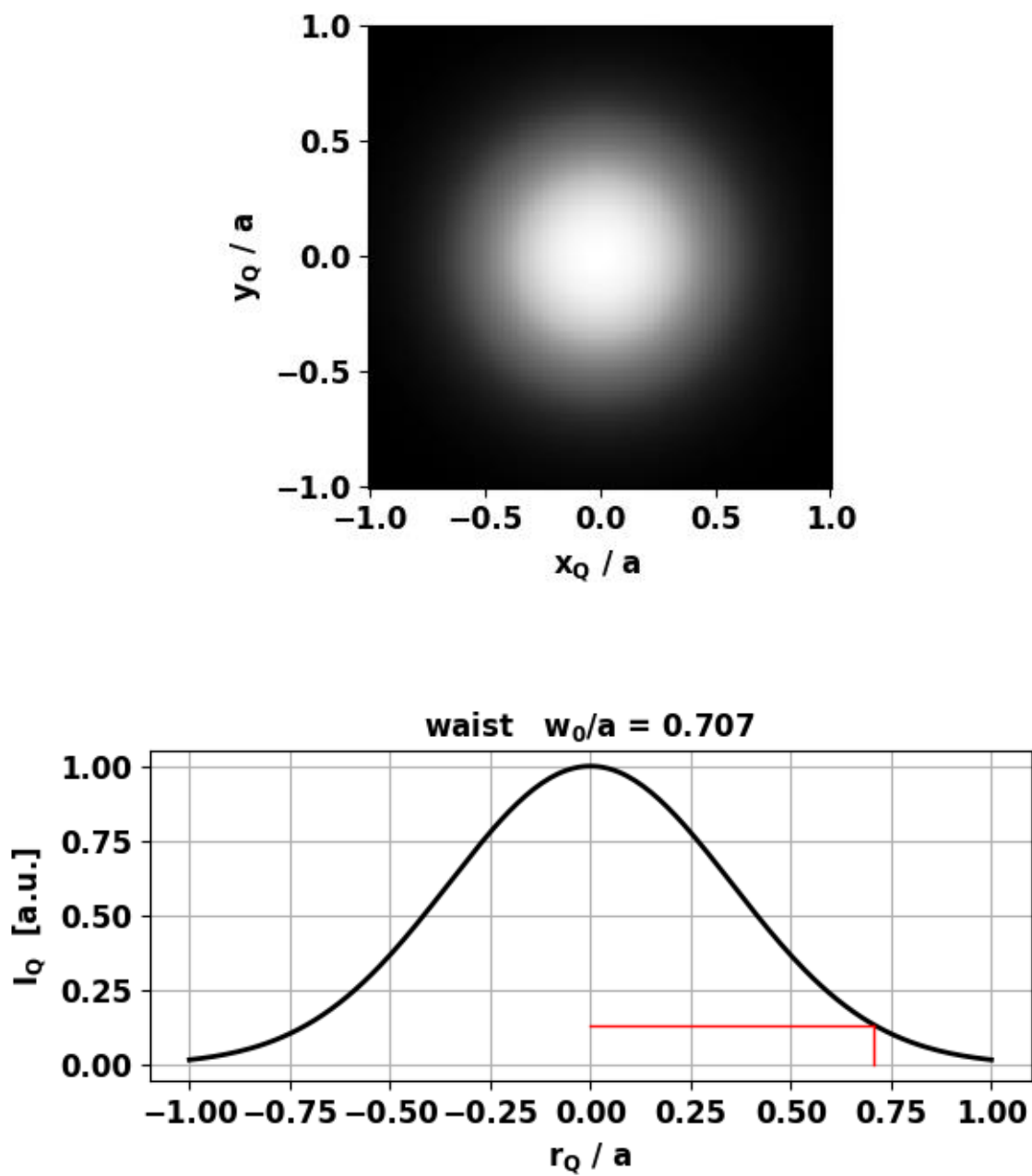


Fig. 4. The irradiance distribution for the circular aperture space of radius a . [emRSGB01.py](#)

Propagation space

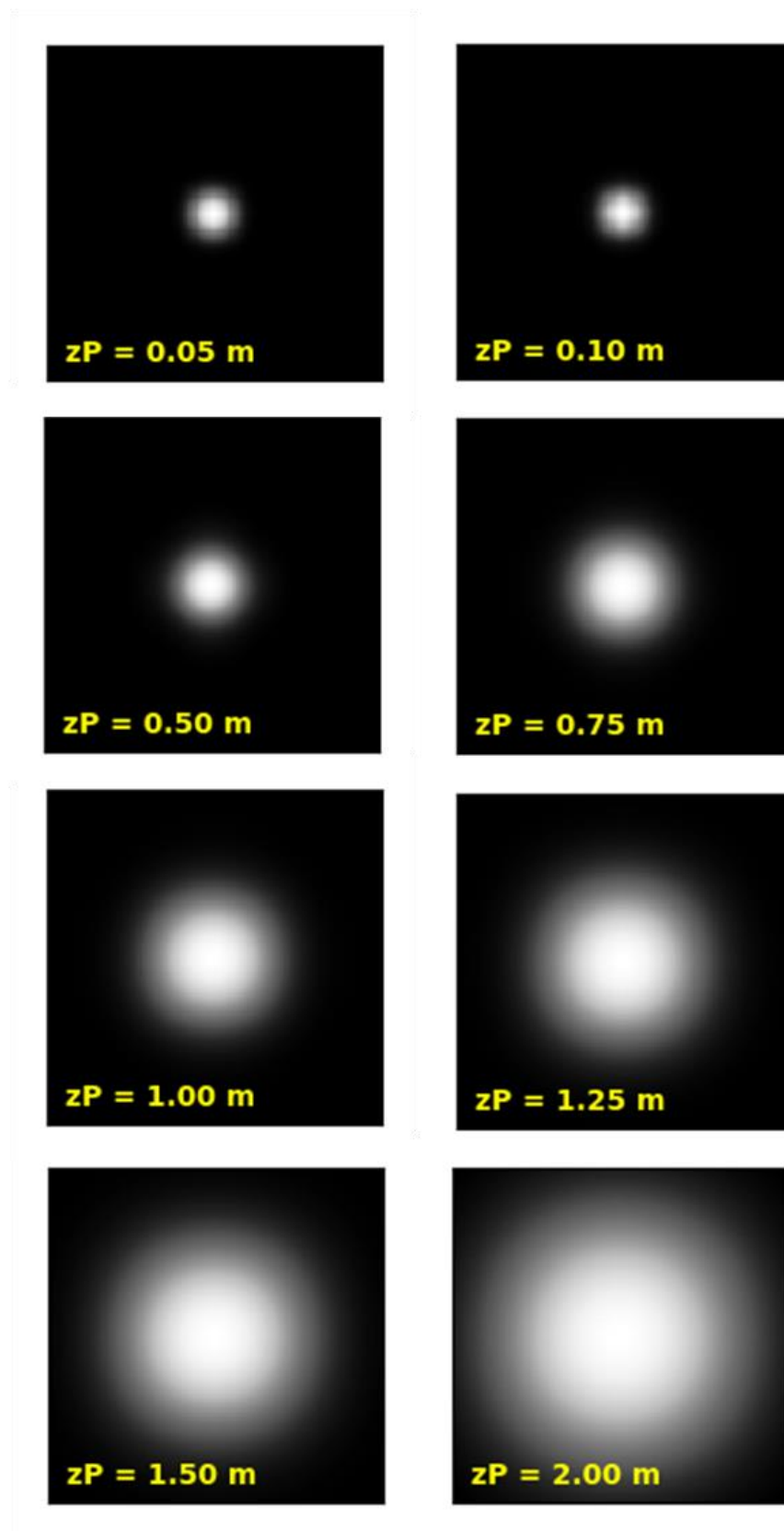


Fig. 5. The irradiance in XY planes for increasing values of zP .

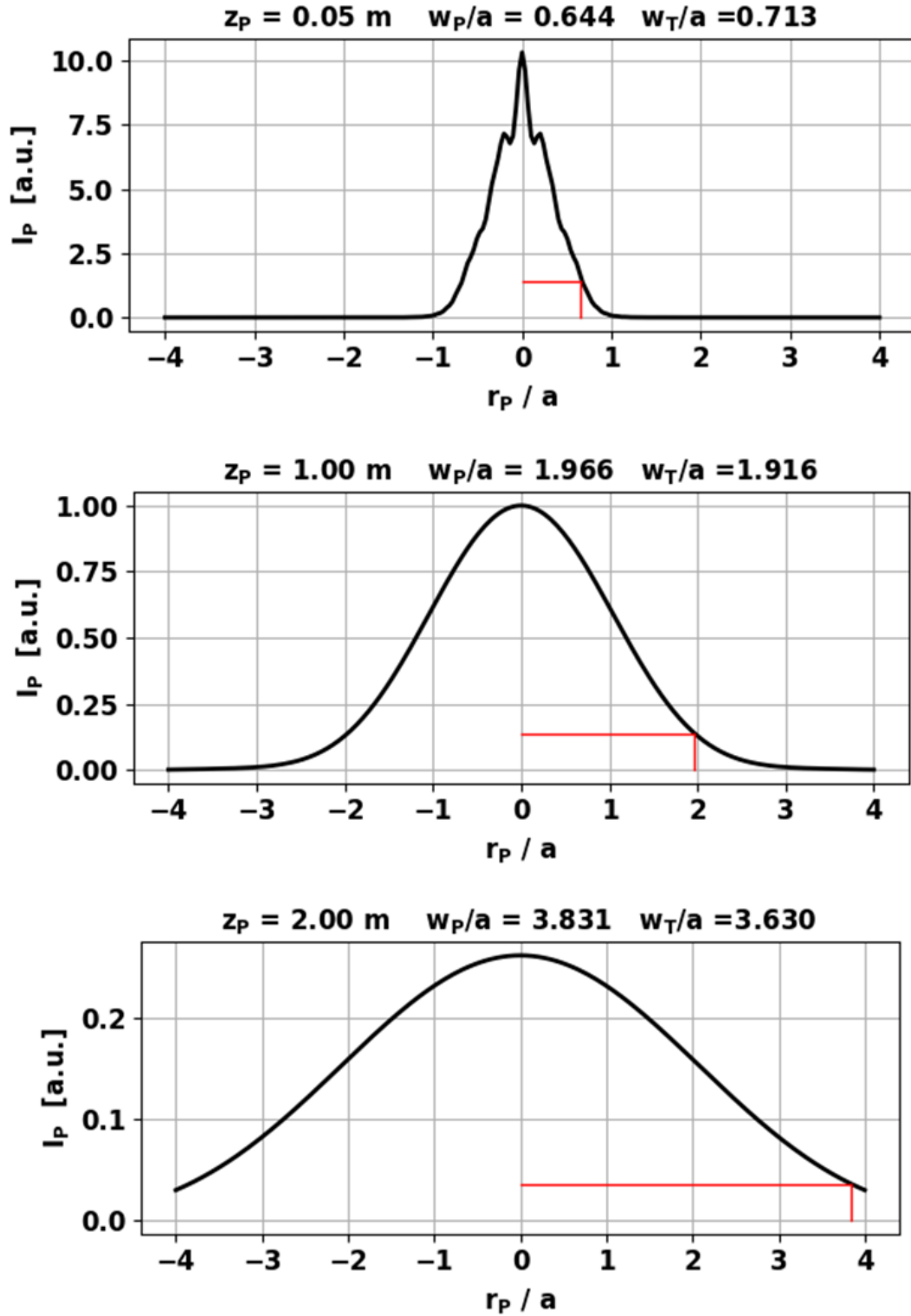


Fig. 6. Irradiance for three z_P values. As the distance from the aperture increases, the beam spot increases. w_T is the beam spot calculated from equation 1. In the far-field ($z_P > 1 \text{ m}$), the peak intensity obeys the inverse square law to a good approximation.

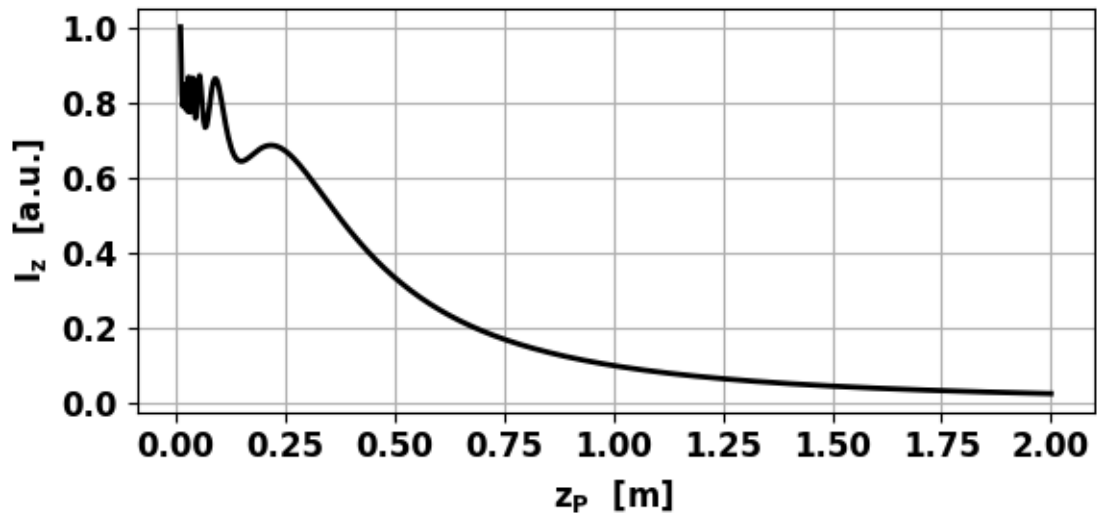
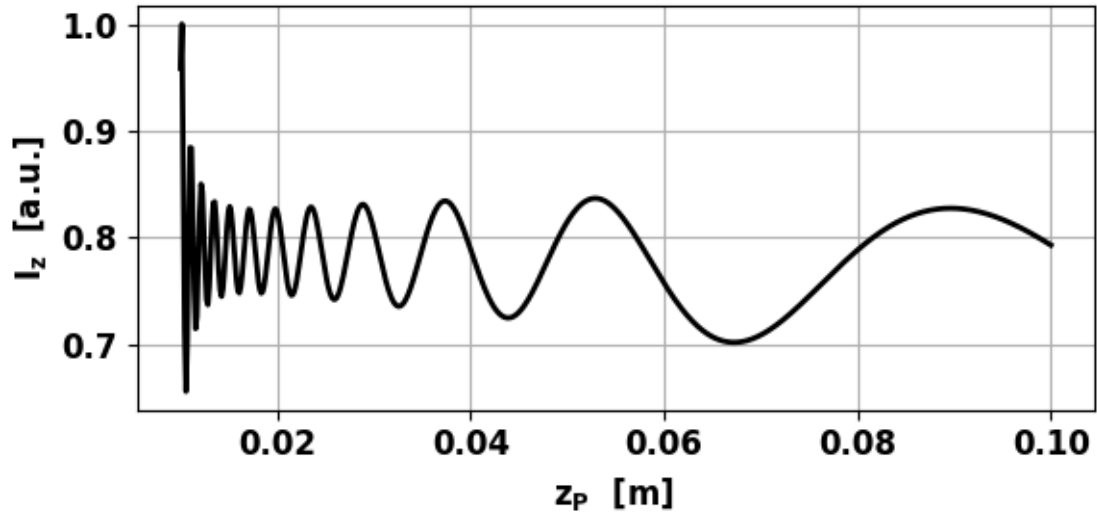


Fig. 7. The irradiance I_z along the optical axis (+Z axis). In the far-field the irradiance decreases according to the inverse square law.

$$z_p = 1.00 \text{ m} \quad I_z = 0.100 \text{ a.u.} \quad z_p = 2.00 \text{ m} \quad I_z = 0.026 \text{ a.u.}$$

emRSGb01Z.py