

DOING PHYSICS WITH PYTHON

DYNAMICAL SYSTEMS [1D] INTRODUCTION

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ds25L1.py Damped driven mass-spring system

Jason Bramburger

History and Preliminaries - Dynamical Systems | Lecture 1

<https://www.youtube.com/watch?v=Gw19VCtHYcs>

INTRODUCTION

As an introduction to dynamical system, an oscillating mass-spring system will be considered. A dynamical system can be modelled by a set of differential equations and the behaviour of the system is completely described by the solutions of the differential equations.

The mass-spring system is an example of an initial value problem

since you need to know the equations and the initial conditions. The equations are expressed in a set of state variables x_1, x_2, x_3, \dots .

Mass – Spring system

Simple Harmonic Motion (SHM)

Damped Harmonic Motion (DHM)

Forced Harmonic Motion (FHM)

The governing equation for the oscillating mass-spring system (no external forcing) is

$$(1) \quad m\ddot{x} + b\dot{x} + kx = 0$$

where m is the mass of the oscillating object, b is the damping constant and k is the spring constant. The state variable x is the displacement from the equilibrium position ($x = 0$) along the X-axis. The equilibrium (steady-state) position is the only fixed point where $x_{ss} = 0$. This is **autonomous** equation since it does have any terms which are an explicit function of time t .

Equation 1 is a 2nd order ODE and to solve it, the equation is expressed as a set of two 1st order ODEs where x_1 and x_2 are the state variables

$$(2) \quad \begin{aligned} \dot{x}_1 &= x_2 \\ \dot{x}_2 &= -(b/m)x_2 - (k/m)x_1 \end{aligned}$$

and the initial conditions are $x_1(0)$ and $x_2(0)$. The state variable x_1 is the displacement x of the oscillating object from the Origin ($x = 0$), x_2 is velocity v of the object, and \dot{x}_2 is the objects acceleration a .

The natural frequencies and period for simple harmonic motion (SHM) $b = 0$ are

$$(3) \quad \omega_0 = \sqrt{\frac{k}{m}} \quad f_0 = \left(\frac{1}{2\pi} \right) \sqrt{\frac{k}{m}} \quad T = 2\pi \sqrt{\frac{m}{k}}$$

The ODEs can be solved very easily using the Python function **odeint**. The output of the Code is displayed graphically.

```
# ds25L1.py
```

```
import numpy as np
from scipy.integrate import odeint
import matplotlib.pyplot as plt
from scipy.signal import find_peaks
import time
from numpy import pi, sin, cos, linspace, sqrt

plt.close('all')

tStart = time.time()

# %% SOLVE ODE x
def lorenz(t, state):
    x1, x2 = state
    dx1 = x2
    dx2 = -(b/m)*x2 - (k/m)*x1
    return dx1, dx2

# %% SETUP
u0 = [0.1,0]
tMax = 10; N = 9999
m = 2
b = 0.5
k = 9
```

$$w = \sqrt{k/m}$$

$$T = 2\pi/w$$

$$f = 1/T$$

#%% SOLVE ODE

t = linspace(0,tMax,N)

sol = odeint(lorenz, u0, t, tfirst=True)

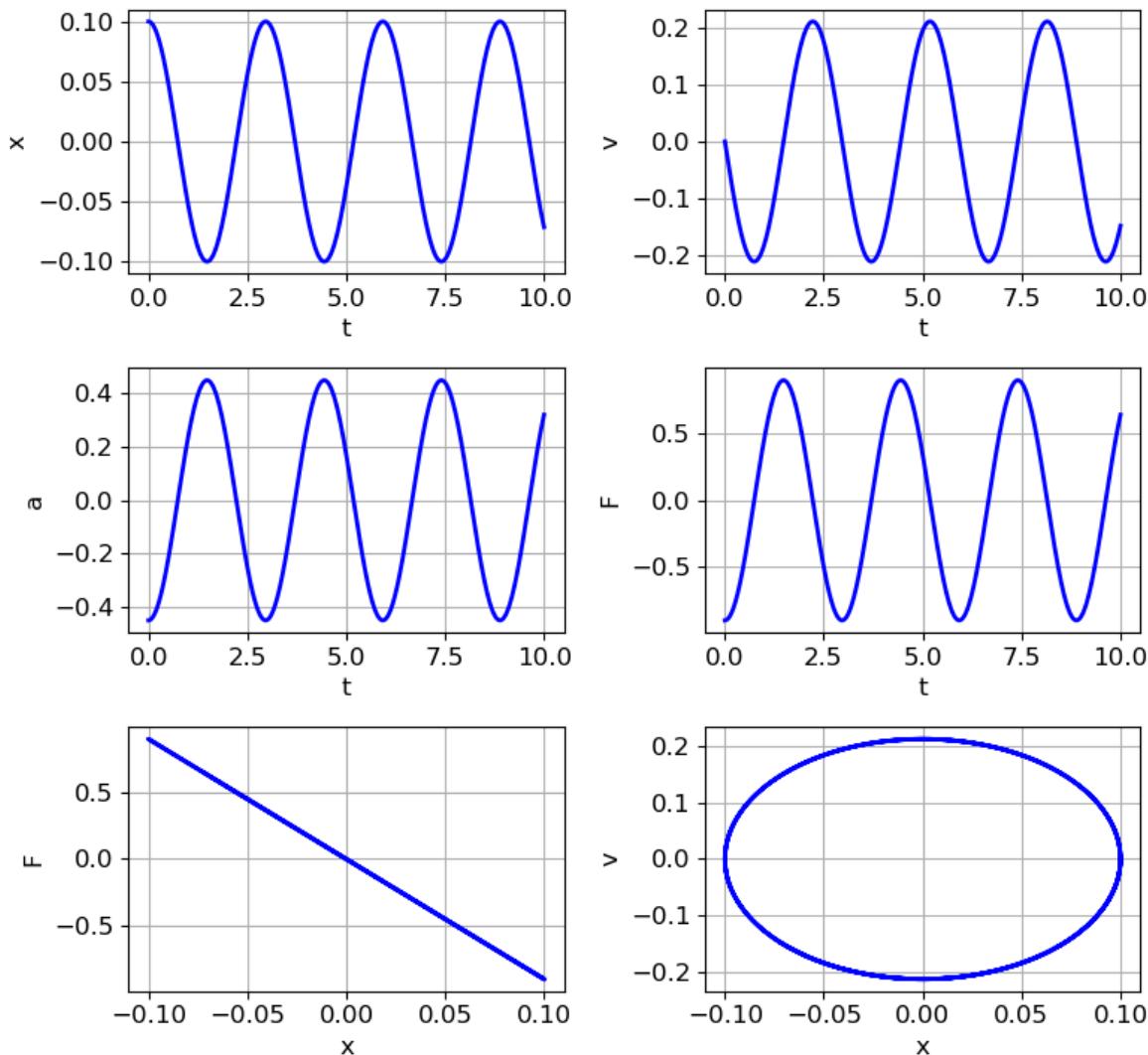
x = sol[:,0]

v = sol[:,1]

a = -(b/m)*v - (k/m)*x

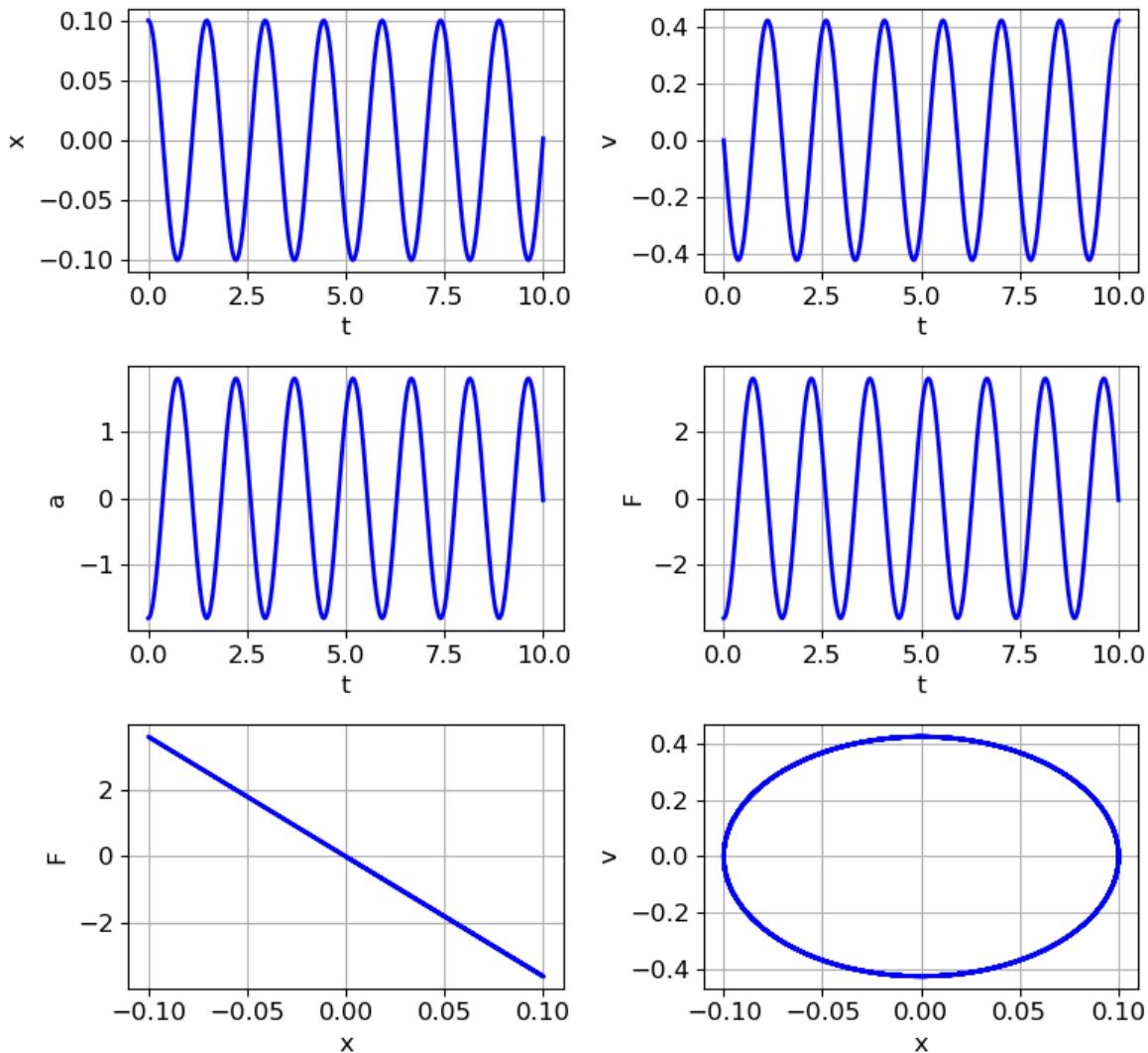
F = m*a

$$\begin{aligned} m &= 2.0 \text{ kg} & b &= 0.0 \text{ kg/s} & k &= 9.0 \text{ N/m} \\ T_0 &= 2.96 \text{ s} & f_0 &= 0.34 \text{ Hz} \end{aligned}$$



$$m = 2.0 \text{ kg} \quad b = 0.0 \text{ kg/s} \quad k = 36.0 \text{ N/m}$$

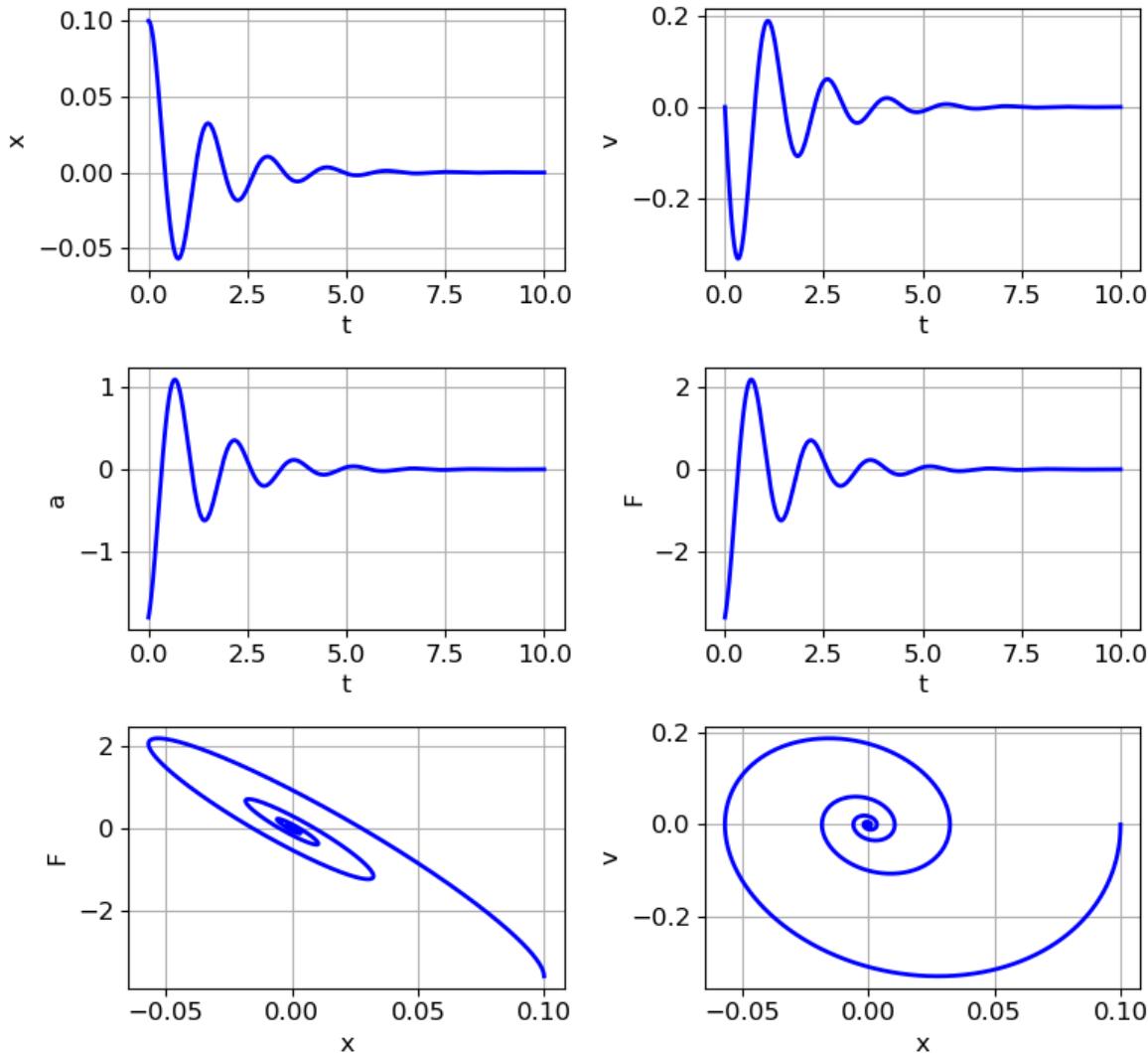
$$T_0 = 1.48 \text{ s} \quad f_0 = 0.68 \text{ Hz}$$



Increasing the value of k from 9 to 36 results in an increase in the frequency and a decrease in the period of oscillation.

$$m = 2.0 \text{ kg} \quad b = 3.0 \text{ kg/s} \quad k = 36.0 \text{ N/m}$$

$$T_0 = 1.48 \text{ s} \quad f_0 = 0.68 \text{ Hz}$$



Damping: The oscillations decay until the object comes to rest at the Origin $x = 0$.

By examining the x vs F graph, we see that the Origin acts as a basin of attraction because the force acting on the object is always pulling it towards the Origin.

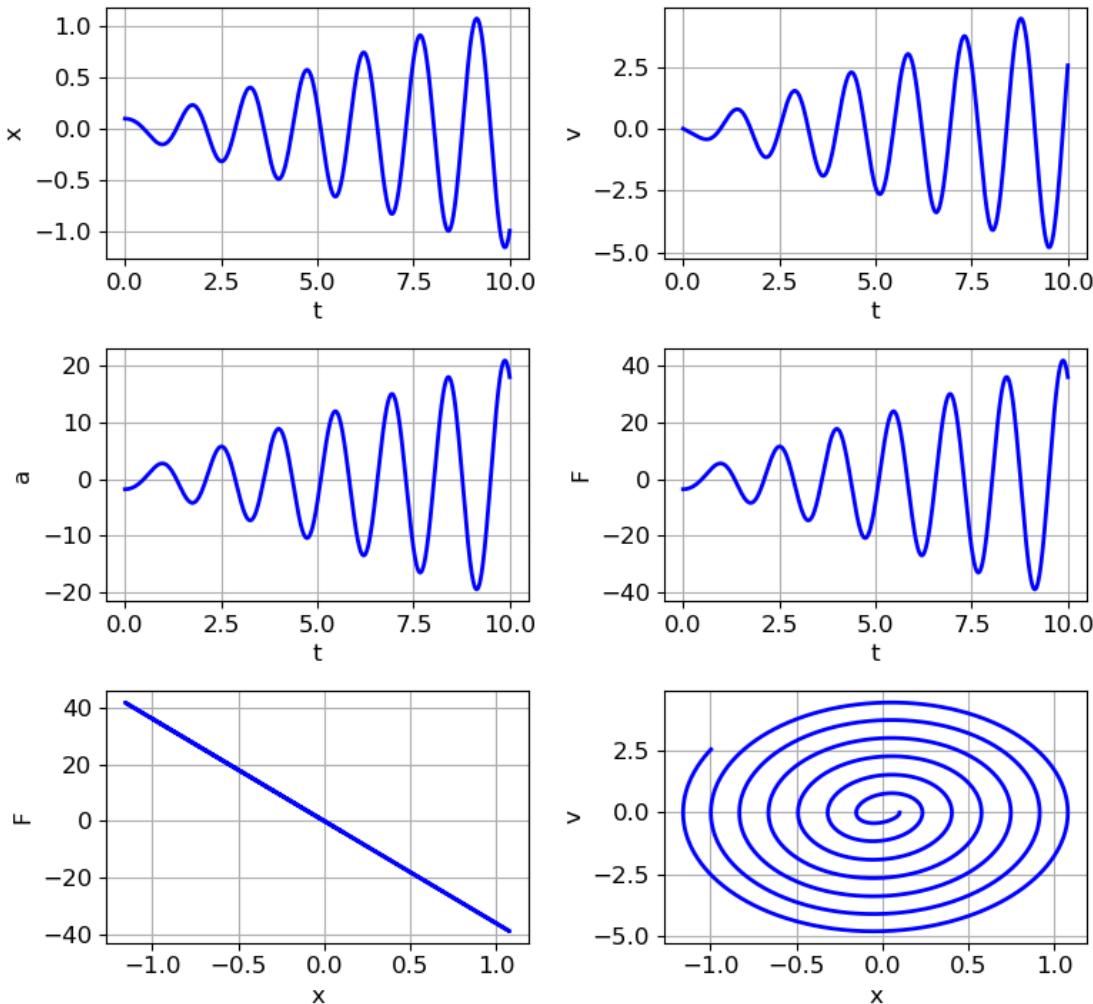
When an external sinusoidal force acts on the mass-spring system, then the equation of motion must include a time dependent term and the equation is now a **nonautonomous** equation

$$(3) \quad m \ddot{x} + b \dot{x} + k x = A_D \cos(\omega_D t)$$

where A_D is the strength of the driving force and the driving frequency is f_D ($\omega_D = 2\pi f_D$). The set of 1st order ODEs are

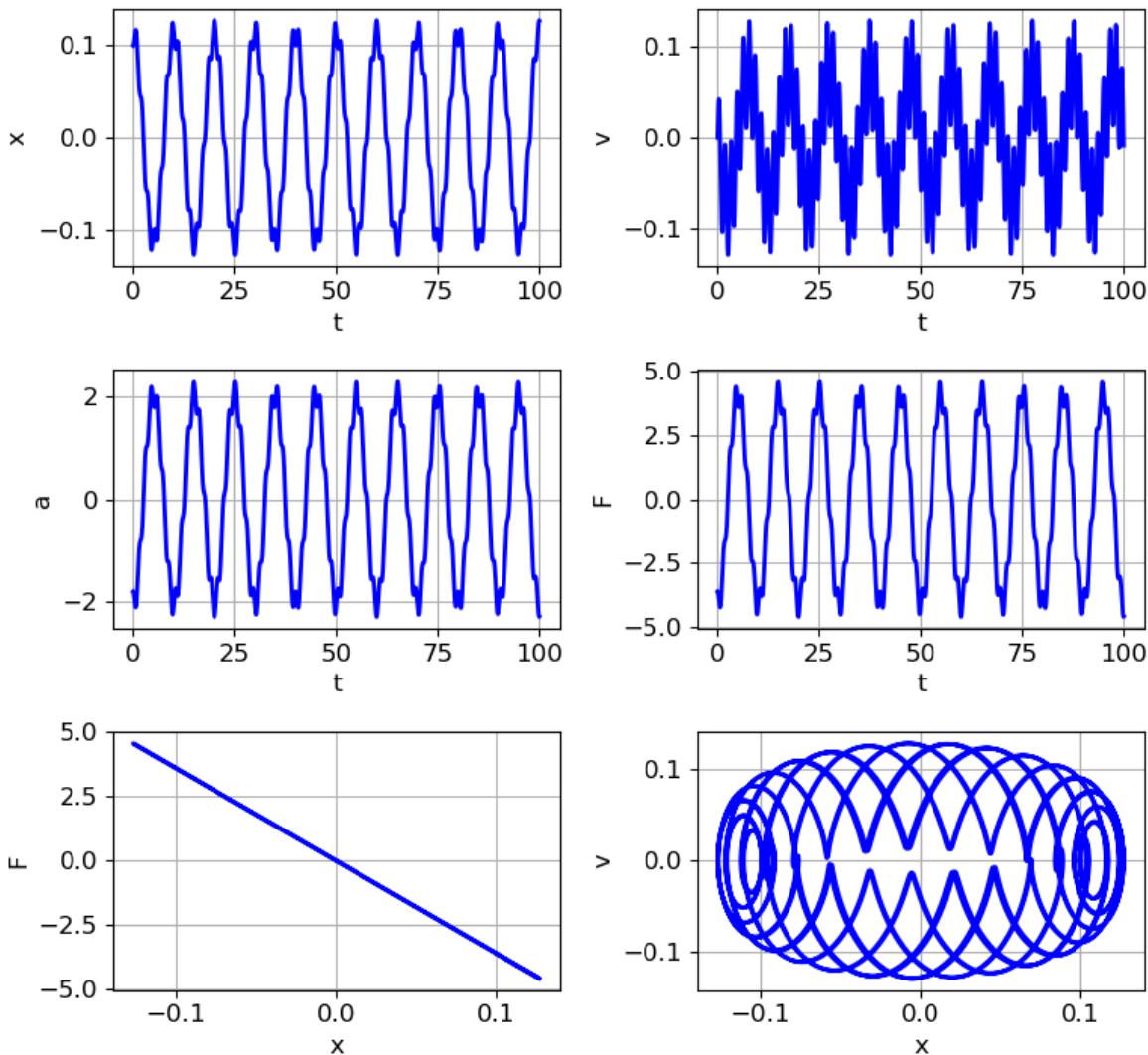
$$(4) \quad \begin{aligned} \dot{x}_1 &= x_2 \\ \dot{x}_2 &= -(b/m)x_2 - (k/m)x_1 + A_D \cos(\omega t) \end{aligned}$$

$$\begin{aligned} m &= 2.0 \text{ kg} & b &= 0.0 \text{ kg/s} & k &= 36.0 \text{ N/m} \\ T_0 &= 1.48 \text{ s} & f_0 &= 0.68 \text{ Hz} & T_D &= 1.45 \text{ s} & A_D &= 1.00 \text{ m/s} \end{aligned}$$



$b = 0 \quad f_D \approx f_0 \Rightarrow$ oscillations grow and grow (**resonance**)

$$\begin{aligned}
m &= 2.0 \text{ kg} & b &= 0.0 \text{ kg/s} & k &= 36.0 \text{ N/m} \\
T_0 &= 1.48 \text{ s} & f_0 &= 0.68 \text{ Hz} & T_D &= 10.00 \text{ s} & A_D &= 2.00 \text{ m/s}
\end{aligned}$$



When the driving frequency f_D is not near the natural frequency f_0 , then oscillations do not grow and the system oscillates near the driving frequency.

View

[Chaotic dynamical systems: a damped driven pendulum](#)