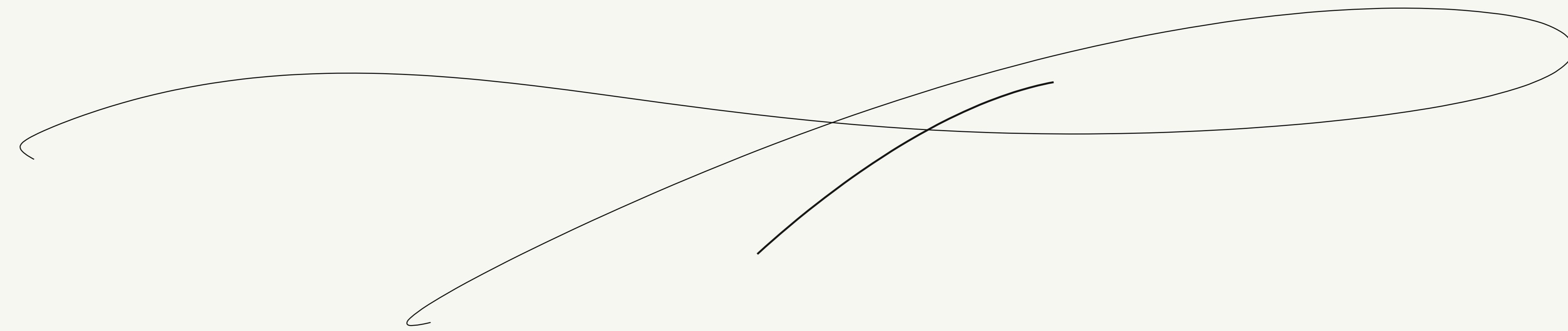


Towards a Geometry and Analysis
for Bayesian Mechanics



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Formal structures in physics:

Geometry - study of spaces and their properties,
such as the state spaces of dynamical systems

Analysis - study of functions and their
properties, such as the equations of motion of
dynamical systems

Motivation: relate gauge symmetries and gauge forces to inference and complex systems, especially in the context of the free energy principle
geometric analysis Bayesian mech.

Approach:

- ① Relate the FEP to maximum entropy
- ② Show max ent is actually a gauge theory
- ③ Profit?

Overview of some recent results

slides can be found later at darsakshi.github.io/talks ; preprints to follow (soon)

$$\text{Find } p \underset{p}{\arg \max} \left(- \int \ln p \, p \right)$$

subject to

$$\int_{\mathcal{X}} j(x) p(x) dx = C$$



$$E[j(X)] = C$$

Find θ

$$-\int \ln p p - \lambda \left(\int j p + c \right)$$

is max

$$\frac{\partial}{\partial p} \left(-\ln p p - \lambda j p \right) = 0$$

$$-\lambda p - \lambda j = 0$$

~~$p = \exp \{-\lambda j\}$~~

$$\lambda j - \lambda j = 0$$

$$-\beta V(x)$$

ℓ

$$\begin{aligned} & \partial_x (\partial_x V(p)) \\ & - \beta \partial_{xx} p \\ & = 0 \end{aligned}$$

$V(x)$: a potential

$$V(x) = f(x)$$

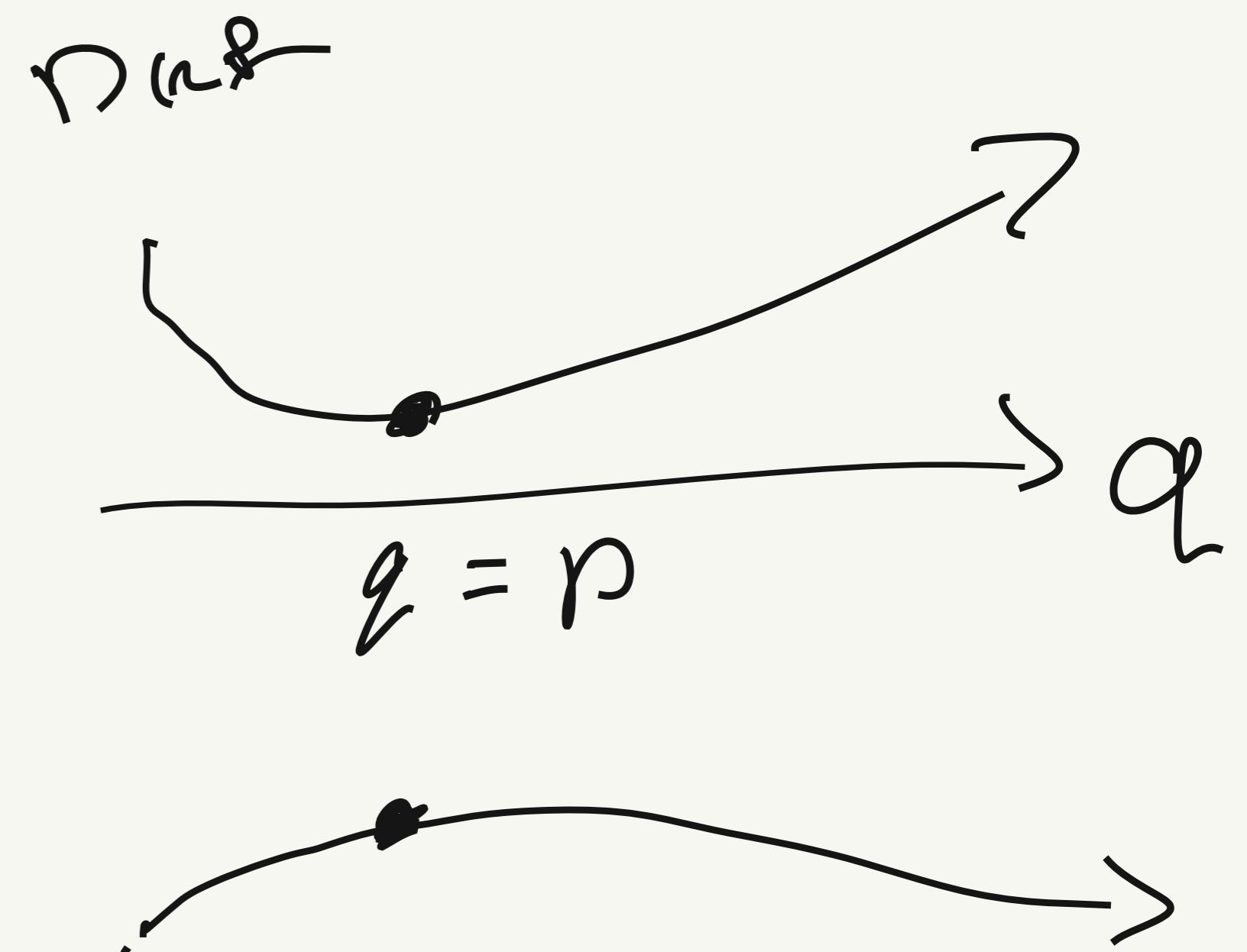
$$\hookrightarrow D_{kh}(q||p) = \mathbb{E}_q [\ln q] - \mathbb{E}_q [\ln p]$$

$$\arg \min_q (D_{kh}) \Rightarrow q = p$$

L \longrightarrow

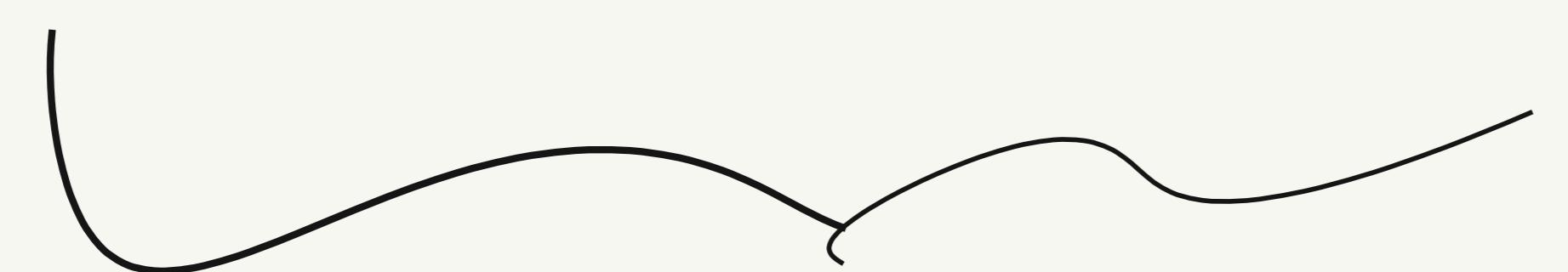
$$\arg \max_q (-D_{kh}) \Rightarrow q = p$$

$$-\mathbb{E}_q \ln q + \mathbb{E}_q \ln p$$

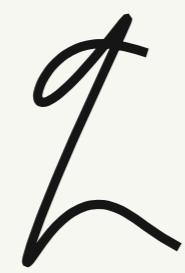


$$-\int \ln q \, q + \int \ln p \, q = 0$$

$$-\int \ln q \, q - \int -\ln p \, q + 0$$



entropy



$$E[-\ln p] = 0$$

$$-\int \ln q \, q = - \left(\int -\ln p(r|\mu, s) q + \ln p(\mu, s) \right)$$

$$\begin{aligned} \mathbb{E}[-\ln p(r|\mu, s)] \\ = -\ln p(\mu, s) \end{aligned}$$

Thm. For an appropriate set of constraints,

FEP $\Leftarrow \Rightarrow$ Max Ent.

Proof.

$$\arg \max_{\mathbf{q}} (-D_{KL}) = \arg \min_{\mathbf{q}} (D_{KL})$$

$$Z^* = P$$

$$Z^* = Z \exp \left\{ -(-\ln p) \right\}$$
$$q_{ij}^* = p$$

Symmetry of Markov blanket

a hidden assumption

$$\left\{ \begin{array}{l} \sigma(\bar{\mu}) = \bar{\mu} \\ \vartheta^{-1}(\bar{\mu}) = \bar{\mu} \end{array} \right.$$

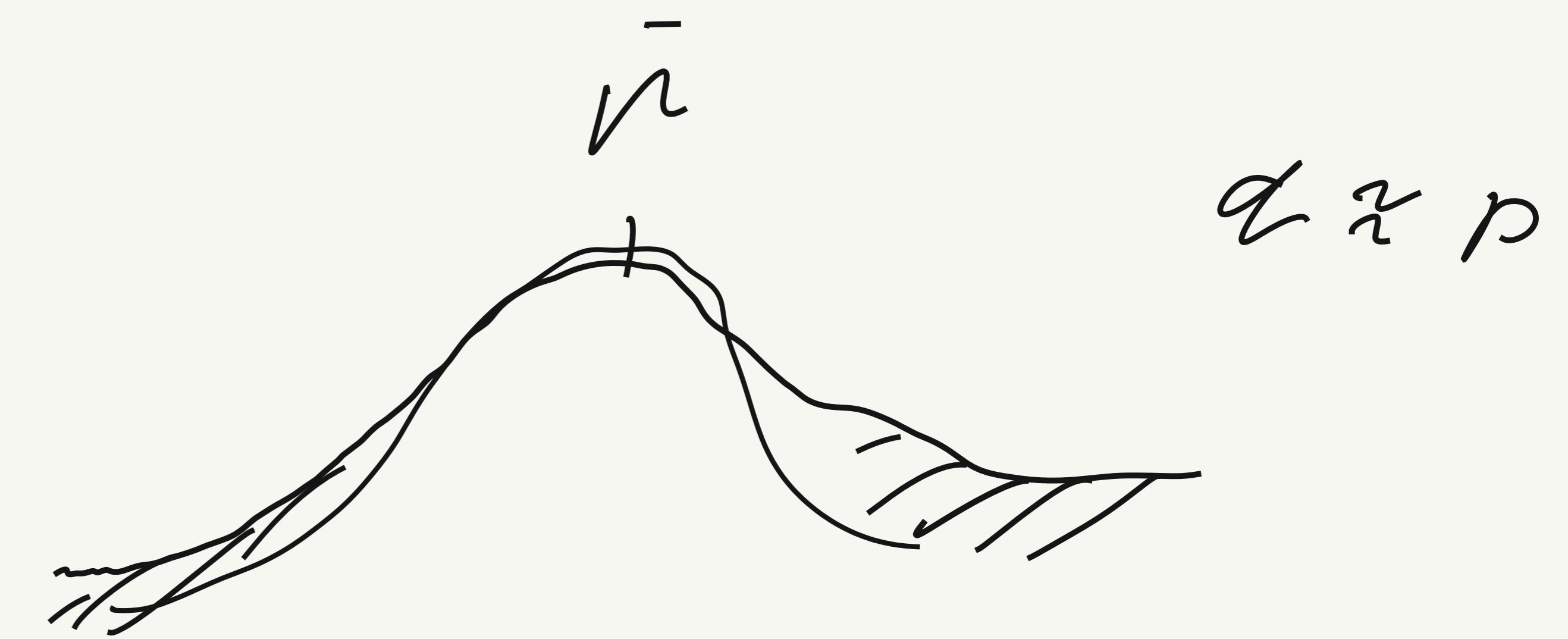
If $q = q(n; \sigma(\bar{\mu}))$ then q minimizes D_{KL} .

Proof. Suppose P has \bar{n} as a r.v. s.t.

then $p = 2e^n$ and $\mathbb{E}_P[n] = \bar{n}$.

Since $p = p(n; \bar{n})$ and $\sigma(\bar{\mu})$

$$= \bar{n} \quad q = P$$



$$J = \mu \quad \text{and} \quad \mathbb{E} [\underset{\sim}{\mu}] = \sigma(\bar{\mu})$$

$$\mu \mapsto \mu \Rightarrow \mathbb{E} \mu = \sigma^{-1}(\bar{\mu})$$

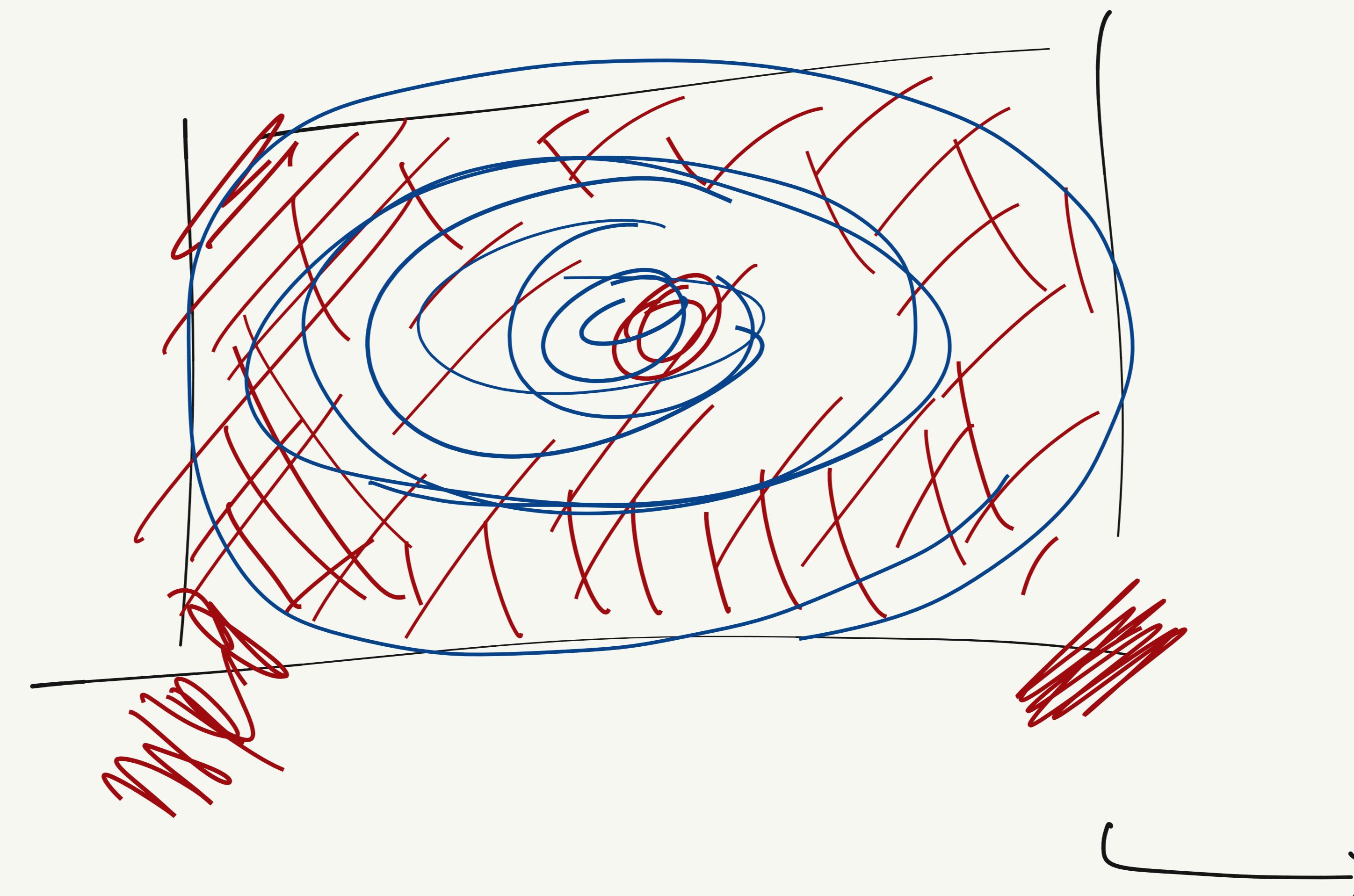
mark (S) Subject +

\sim model μ
on avg.

Introducing constraints decreases entropy

BUT does not minimise it.

$J = x^2 + y^2$
2D Gaussian
 P_J



$J =$
Constraint
 $P = \text{prob.}$

Indeed $J = x^2 + y^2$ $\|\vec{x}\|$

$$\mu = [x, y]$$

$$q(\mu) = \exp\{-\lambda J\}$$
$$l - x^2 - y^2 = \mathcal{N}(\varrho)$$

$p(\text{state}) \propto -\text{Constraints}$

V_{in} max entropy / min \bar{F}_E

Constraints on states of a system are like preferences about what states are likely to be occupied



which are like potential functions for sampling dynamics.

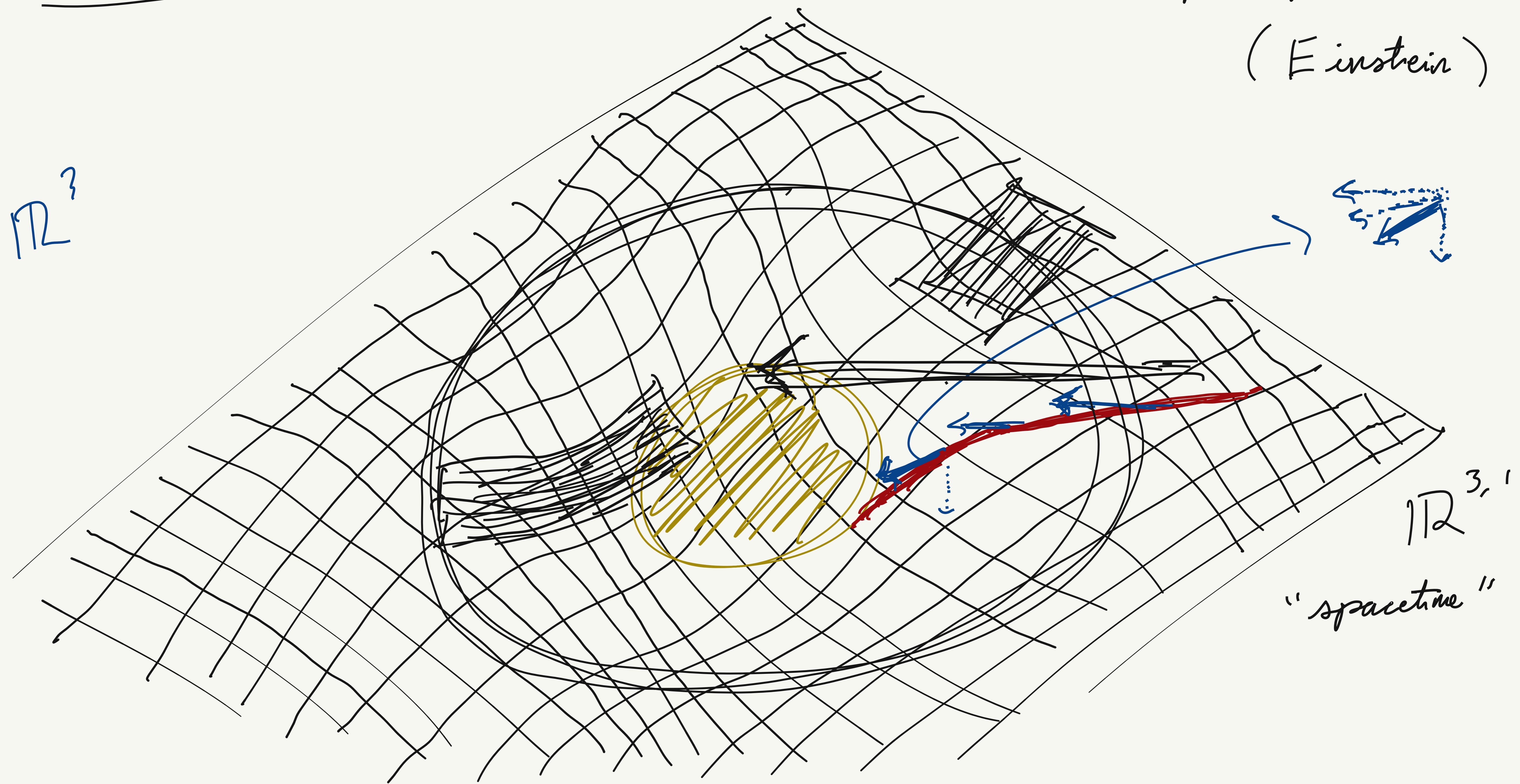
So what ?

Potential functions are geometric
features of a system

Example : general relativity

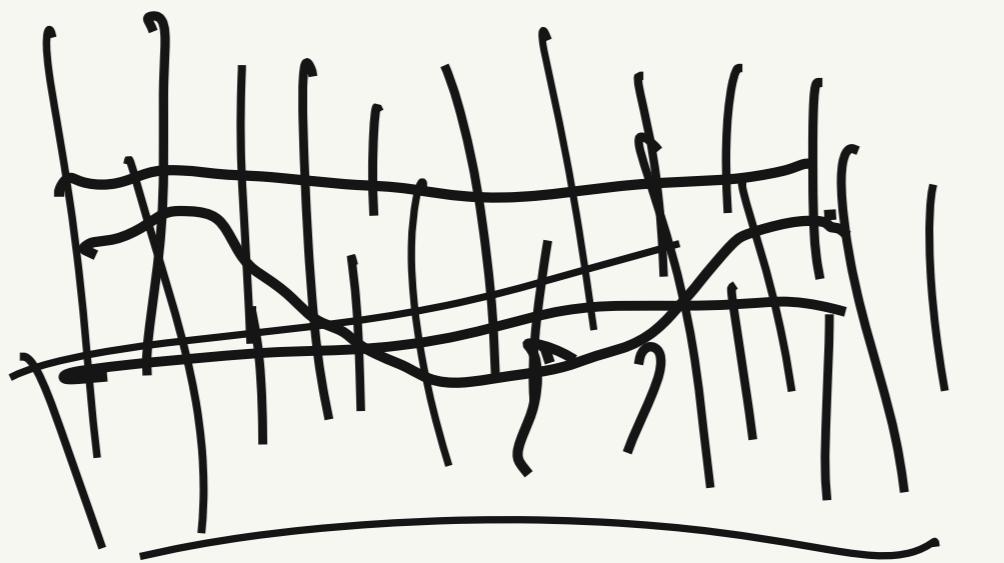
Some geometry)

Equivalence
principle
(Einstein)

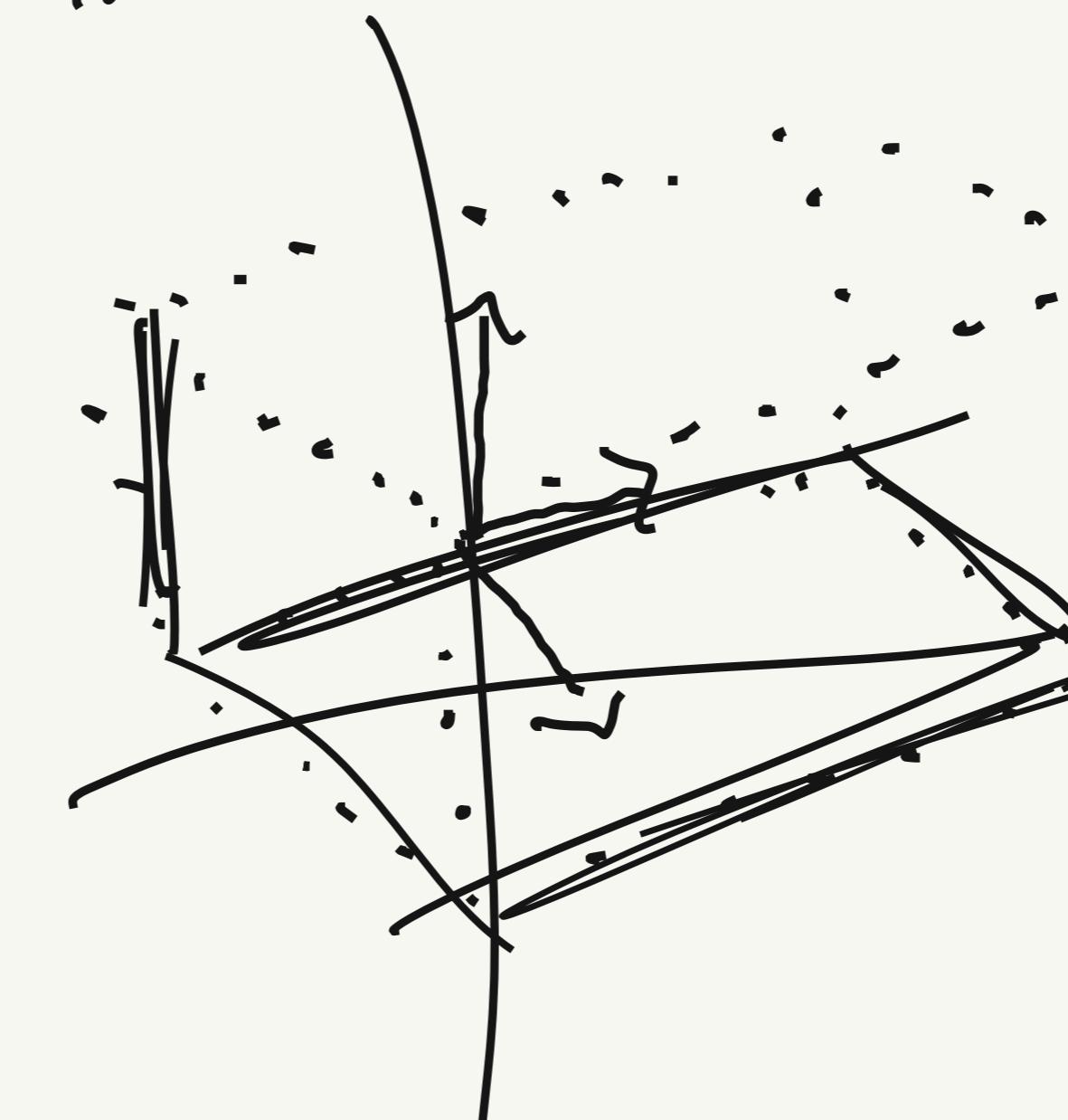


dynamics
of gauge
forced
particles

split into particle
dynamics and gauge
force

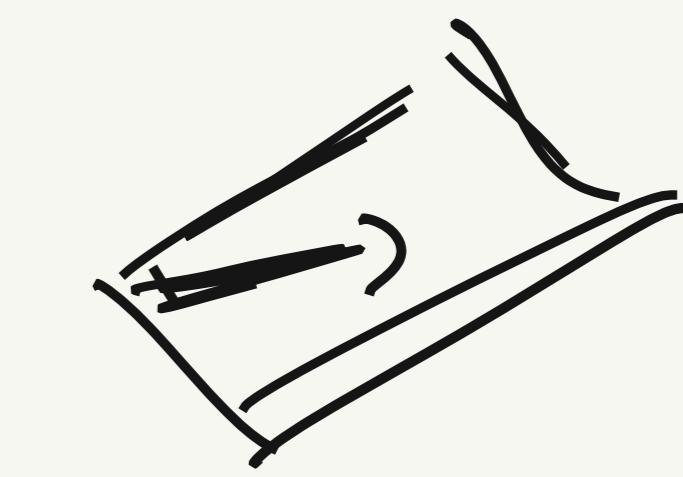


$$\mathbb{R}^2 + \mathbb{R} = \mathbb{R}^3$$



Spacetime

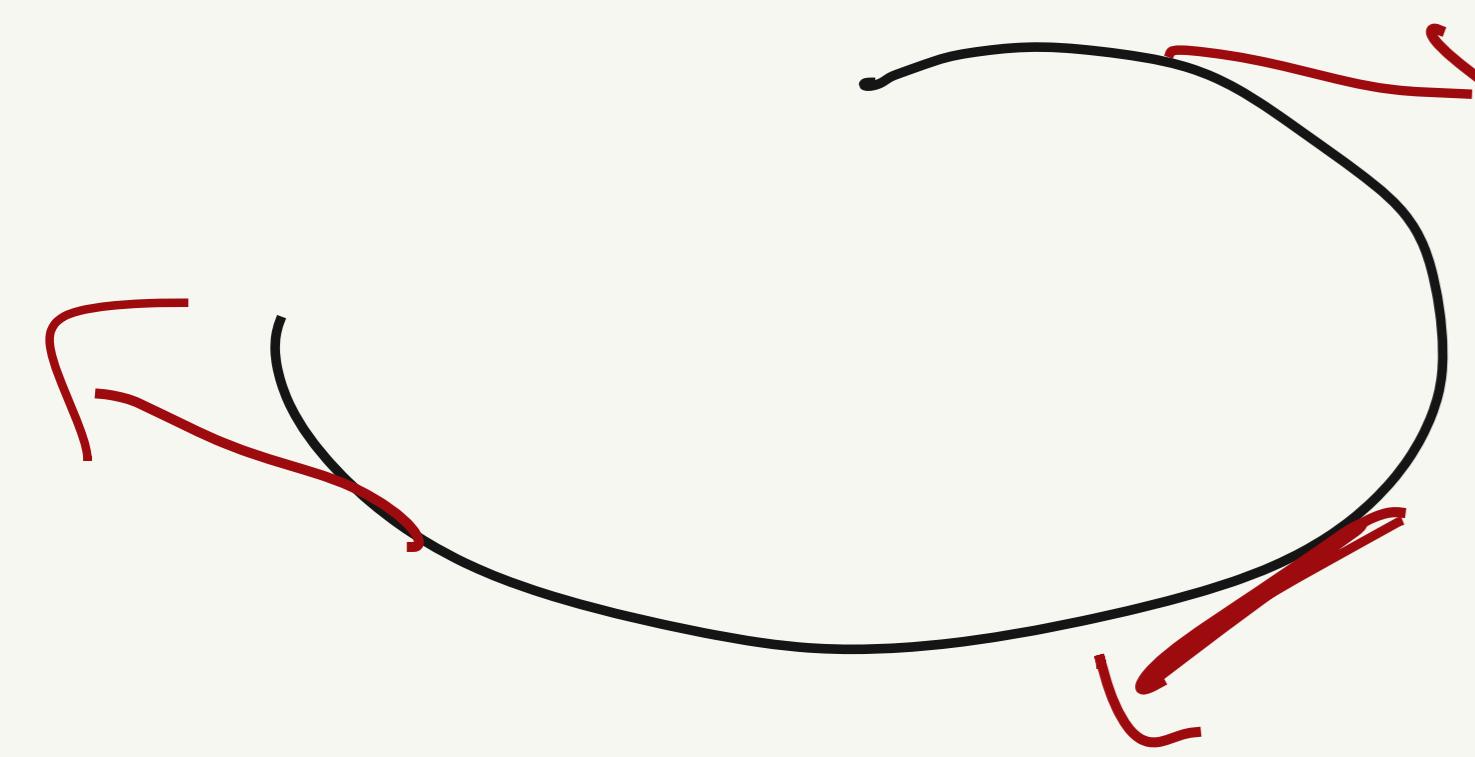
↑ force
(gravity)



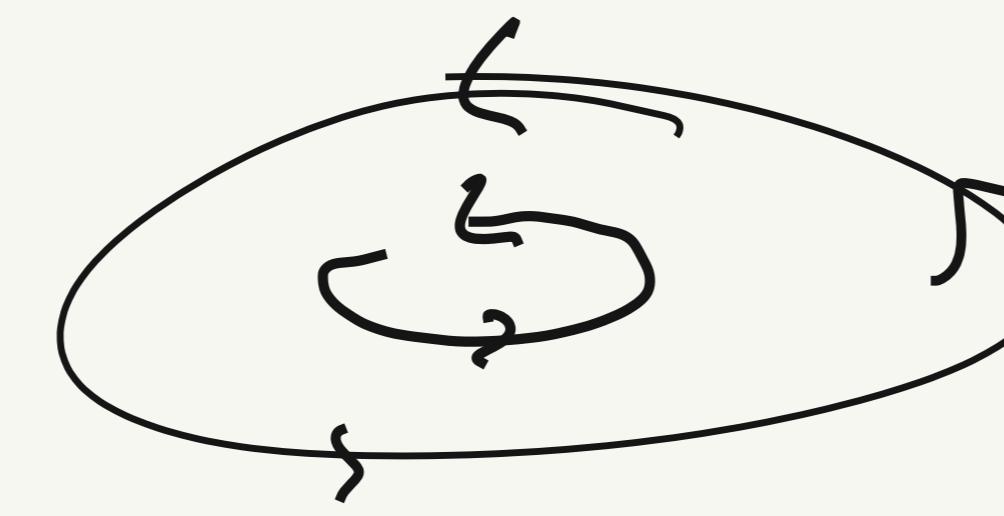
$H_x T$

$V_x T$

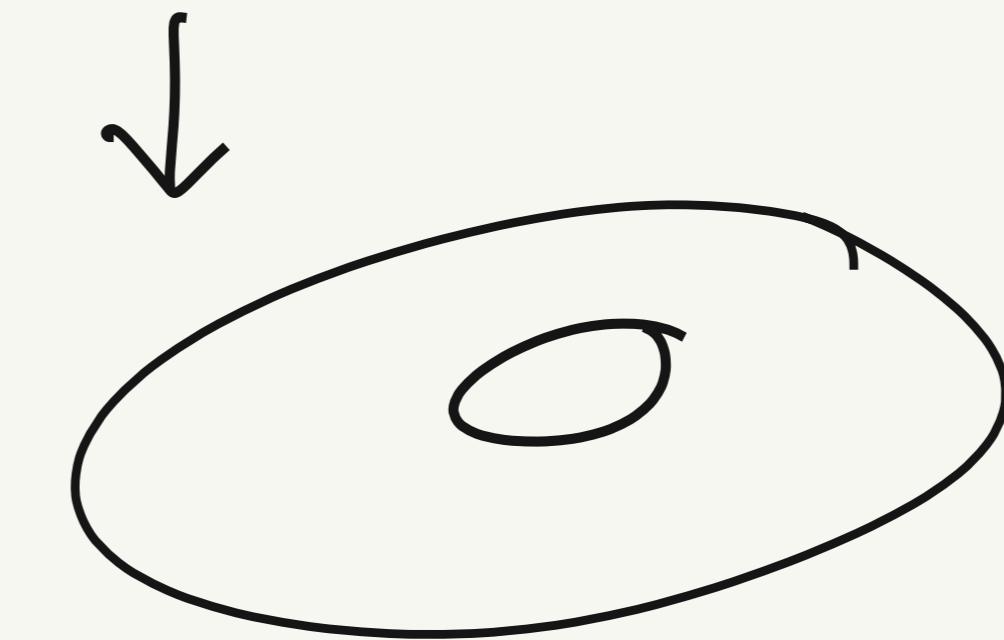
Example:



H/P



VP

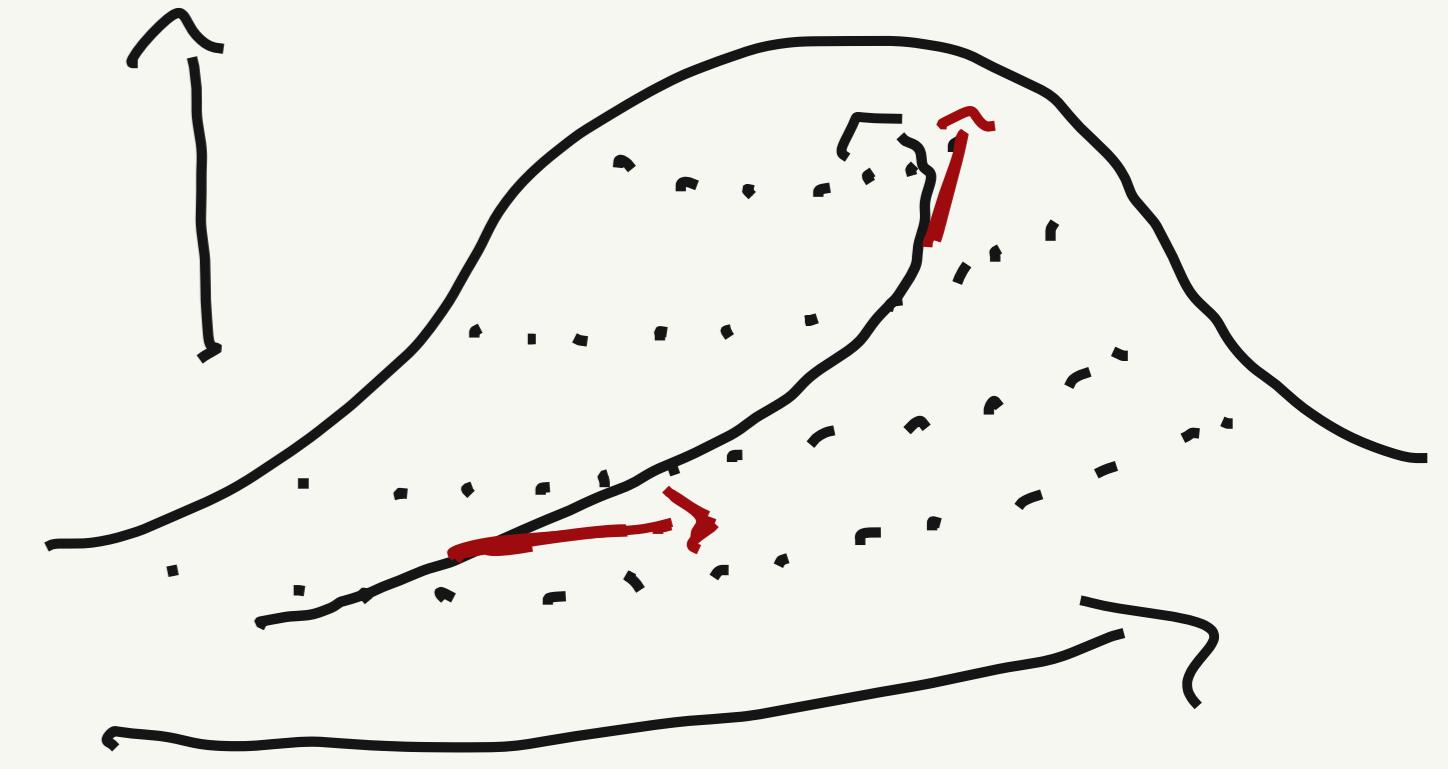


R^3



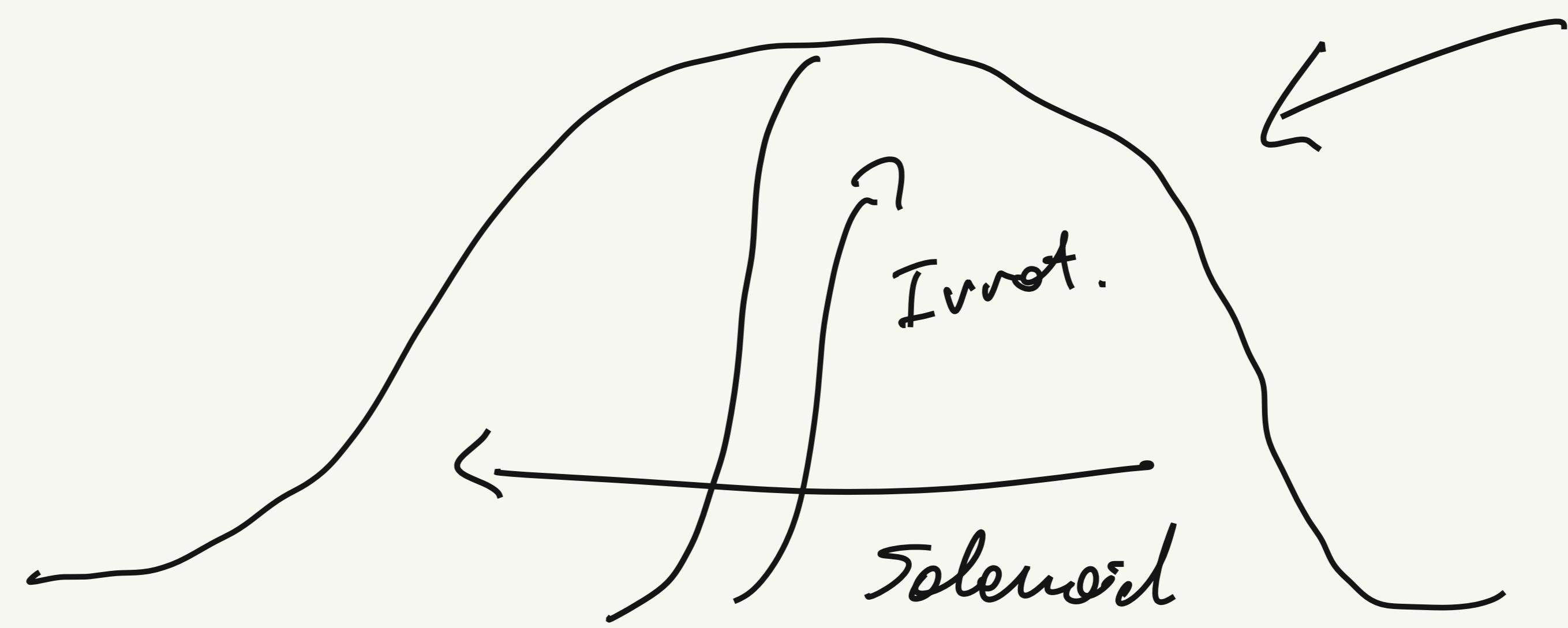
along isocont
⇒
max int.

We have
seen this
before



up
isocont.
⇒

MAP



$$\varphi = \mathcal{J}(x)$$

$$\exp\{-j(x)\}$$

Taking \mathcal{J} as $-\ln \{ p(\mu) \}$,

$$\ln p = -\mathcal{J}$$

$$p = e^{-\mathcal{J}}$$

this flow is $-\nabla \ln \{ p(\mu) \}$.

$$-\mathcal{J} = -\mathcal{I}$$

Proof.

$$\partial_{x_i} p(x_i) = -\partial_{x_i} \mathcal{J}(x_i) p(x_i)$$

$$\frac{\nabla p}{p} = \nabla \ln p$$

$$\nabla p(x_i) = -\nabla \mathcal{J}(x_i) p(x_i)$$

$$\nabla \ln p = -\nabla \mathcal{J}$$

What we have shown:

- $\min \text{FE} \Leftrightarrow \max \text{ent}$
- $\max \text{ent} \Leftrightarrow$ potential on the state space

↳ Consequence:

flow towards or mode
is a sort of gauge
force.

Appendix : gauge symmetry

Frederic Schuller

Valery Rubakov

Lorenzo A Lessa

"Notes on Gauge theories"

2019