

Electromagnetics 1 (ICE2003) -- Ch. 6. Capacitance --

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Chapter Outline

- Capacitance definition
- Capacitance calculation
 - Parallel plate, Coaxial transmission line, Spherical capacitor
 - Isolated sphere with a dielectric coating, Parallel plate with two-layer dielectric
 - Two-wire line, one-wire line and a plane
- Energy stored in a capacitor
- Poisson / Laplace equation
 - Laplacian operator
 - Parallel plate, Coaxial transmission line, Angled plate, Concentric sphere, Cone above a conducting plane
 - Product solution

Capacitance Definition

A simple capacitor consists of two oppositely charged conductors surrounded by a uniform dielectric.

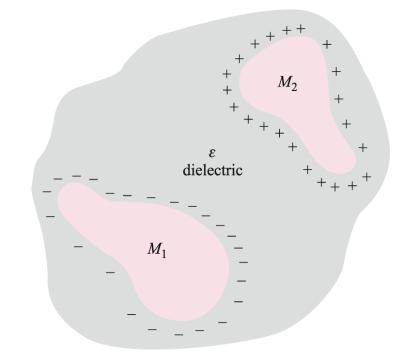
An increase in Q by some factor results in an increase in E (and in D) by the same factor.

where
$$Q = \oint_S \mathbf{D} \cdot d\mathbf{S}$$

Consequently, the potential difference between conductors:

$$V_0 = -\int_B^A \mathbf{E} \cdot d\mathbf{L}$$

will also increase by the same factor -- so the ratio of Q to V_0 is a constant. We define the *capacitance* of the structure as the ratio of the stored charge to the applied voltage, or

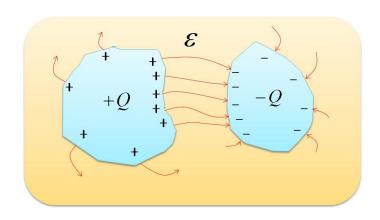


$$C = \frac{Q}{V_0}$$

Units are Coul/V or Farads

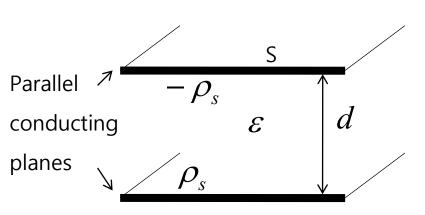
- Determined by Capacitor geometry and material
- Independent from V and Q

Calculating Capacitance



- 1. Set Coordinate system
- 2. Assume +Q, -Q to two conductors
- 3. Calculate E
- 4. Calculate V by integrating E
- 5. C=Q/V

Parallel Plate Capacitor



$$Q = \rho_s S$$

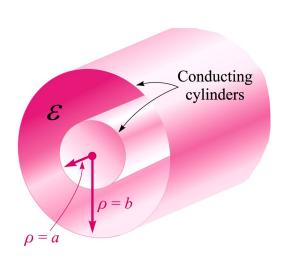
$$\mathbf{D} = \rho_{s} \mathbf{a}_{\mathbf{z}}$$

$$\mathbf{E} = \frac{\boldsymbol{\rho}_s}{\boldsymbol{\varepsilon}} \mathbf{a_z}$$

$$V = -\int_{init}^{final} \mathbf{E} \cdot d\mathbf{L} = -\int_{d}^{0} \mathbf{E} \cdot \mathbf{a}_{z} dz = -\int_{d}^{0} \frac{\rho_{s}}{\varepsilon} dz = \frac{\rho_{s} d}{\varepsilon}$$

$$C = \frac{Q}{V} = \frac{\varepsilon S}{d}$$

Coaxial Transmission Line



$$Q = 2\pi a \rho_s L$$

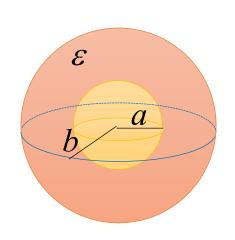
$$\mathbf{D} = \frac{a\rho_s}{\rho} \mathbf{a}_{\rho}$$

$$\mathbf{E} = \frac{a\rho_s}{\varepsilon\rho} \mathbf{a}_{\rho}$$

$$V = -\int_{init}^{final} \mathbf{E} \cdot d\mathbf{L} = -\int_{b}^{a} \mathbf{E} \cdot \mathbf{a}_{\rho} d\rho = -\int_{b}^{a} \frac{a\rho_{s}}{\varepsilon \rho} d\rho = \frac{a\rho_{s}}{\varepsilon} \ln \frac{b}{a}$$

$$C = \frac{Q}{V} = \frac{2\pi\varepsilon L}{\ln\frac{b}{a}}$$

Spherical Capacitor



$$Q = 4\pi a^2 \rho_s$$

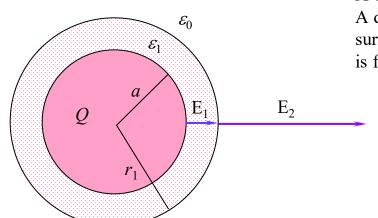
$$\mathbf{D} = \frac{a^2 \rho_s}{r^2} \mathbf{a_r}$$

$$\mathbf{D} = \frac{a^2 \rho_s}{r^2} \mathbf{a_r}$$
$$\mathbf{E} = \frac{a^2 \rho_s}{\varepsilon r^2} \mathbf{a_r}$$

$$V = -\int_{init}^{final} \mathbf{E} \cdot d\mathbf{L} = -\int_{b}^{a} \mathbf{E} \cdot \mathbf{a_r} dr = -\int_{b}^{a} \frac{a^2 \rho_s}{\varepsilon r^2} dr = \frac{a^2 \rho_s}{\varepsilon} \left(\frac{1}{a} - \frac{1}{b} \right)$$

$$C = \frac{Q}{V} = \frac{4\pi\varepsilon}{\frac{1}{a} - \frac{1}{b}}$$

Isolated Sphere with a Dielectric Coating



A conducting sphere of radius a carries charge Q. A dielectric layer of thickness r_1 - a and of permittivity ε_1 surrounds the conductor. Electric field in the two regions is found from Gauss' Law to be:

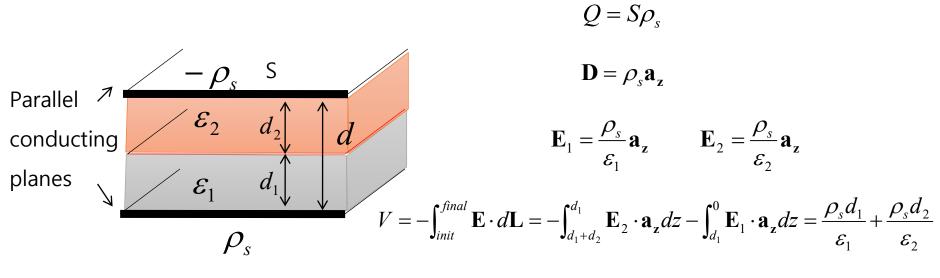
$$E_r = \frac{Q}{4\pi\epsilon_1 r^2} \qquad (a < r < r_1)$$
$$= \frac{Q}{4\pi\epsilon_0 r^2} \qquad (r_1 < r)$$

The potential at the sphere surface is (with zero reference at infinity):

$$V_a - V_{\infty} = -\int_{r_1}^a \frac{Q \, dr}{4\pi\epsilon_1 r^2} - \int_{\infty}^{r_1} \frac{Q \, dr}{4\pi\epsilon_0 r^2} = \frac{Q}{4\pi} \left[\frac{1}{\epsilon_1} \left(\frac{1}{a} - \frac{1}{r_1} \right) + \frac{1}{\epsilon_0 r_1} \right] = V_0$$

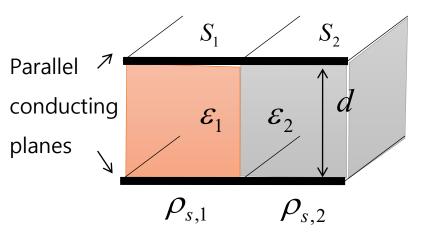
and the capacitance is:
$$C = \frac{4\pi}{\frac{1}{\epsilon_1} \left(\frac{1}{a} - \frac{1}{r_1}\right) + \frac{1}{\epsilon_0 r_1}}$$

Capacitor with a Two-Layer Dielectric



$$C = \frac{Q}{V} = \frac{1}{\frac{d_1}{\varepsilon_1 S} + \frac{d_2}{\varepsilon_2 S}}$$

The other case



$$Q = S_{1}\rho_{s,1} + S_{2}\rho_{s,2}$$

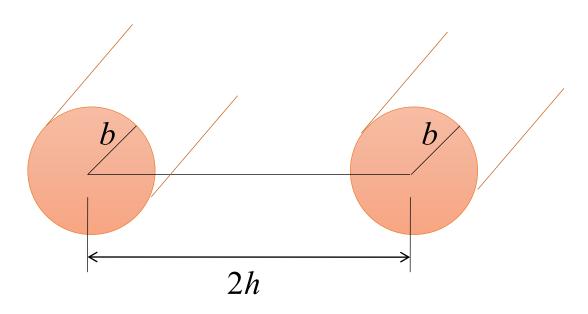
$$\mathbf{D}_{1} = \rho_{s,1}\mathbf{a}_{z} \quad \mathbf{D}_{2} = \rho_{s,2}\mathbf{a}_{z}$$

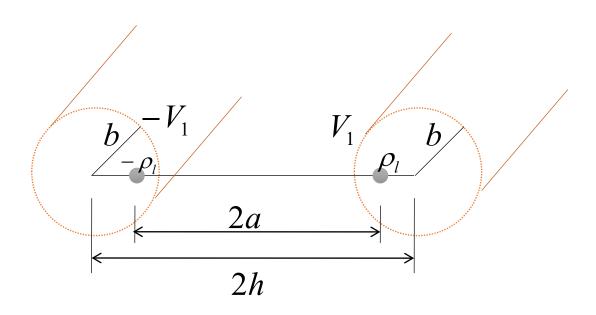
$$\mathbf{E}_{1} = \frac{\rho_{s,1}}{\varepsilon_{1}}\mathbf{a}_{z} \quad \mathbf{E}_{2} = \frac{\rho_{s,2}}{\varepsilon_{2}}\mathbf{a}_{z}$$

$$V = -\int_{d}^{0}\mathbf{E}_{1} \cdot \mathbf{a}_{z} dz = \frac{\rho_{s,1}d}{\varepsilon_{1}} = -\int_{d_{1}}^{0}\mathbf{E}_{2} \cdot \mathbf{a}_{z} dz = \frac{\rho_{s,2}d}{\varepsilon_{2}}$$

$$\therefore \frac{\rho_{s,1}}{\varepsilon_{1}} = \frac{\rho_{s,2}}{\varepsilon_{2}}$$

$$C = \frac{Q}{V} = \frac{\varepsilon_1 S_1}{d} + \frac{\varepsilon_2 S_2}{d}$$

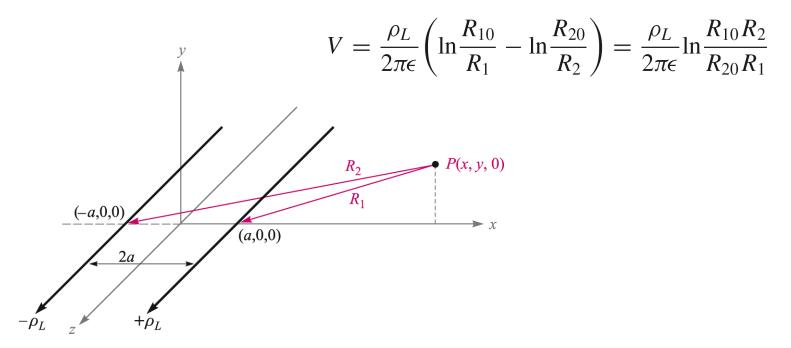




We begin with the potential field of a single line charge on the z axis, with a zero reference at $\rho = R_0$

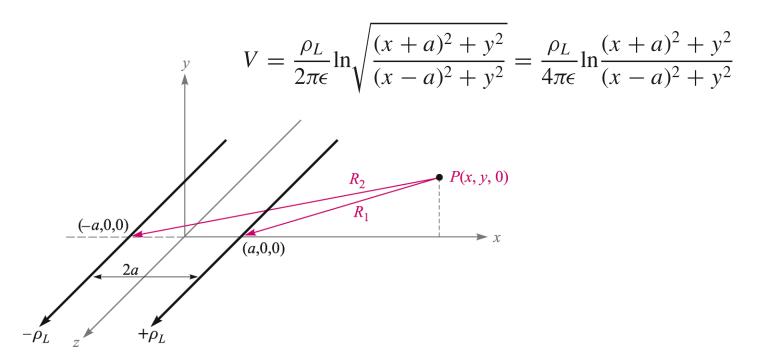
$$V = \frac{\rho_L}{2\pi\epsilon} \ln \frac{R_0}{R}$$

We can use this result to write the potential at point P from two line charges of opposite sign, postioned as shown:



We now have:
$$V = \frac{\rho_L}{2\pi\epsilon} \left(\ln \frac{R_{10}}{R_1} - \ln \frac{R_{20}}{R_2} \right) = \frac{\rho_L}{2\pi\epsilon} \ln \frac{R_{10}R_2}{R_{20}R_1}$$

We choose $R_{10} = R_{20}$, thus placing the zero reference at equal distances from each line. This surface is the x = 0 plane. Expressing R_1 and R_2 in terms of x and y,



We now have:
$$V = \frac{\rho_L}{2\pi\epsilon} \ln \sqrt{\frac{(x+a)^2 + y^2}{(x-a)^2 + y^2}} = \frac{\rho_L}{4\pi\epsilon} \ln \frac{(x+a)^2 + y^2}{(x-a)^2 + y^2}$$

Choose an equipotential surface on which $V = V_1$, and define the dimensionless parameter:

$$K_1 = e^{4\pi\epsilon V_1/\rho_L}$$

from which we identify:
$$K_1 = \frac{(x+a)^2 + y^2}{(x-a)^2 + y^2}$$

This is the equation of the equipotential surface on which the voltage is V_1

To better identify the surface, expand the squares, and collect terms:

$$x^{2} - 2ax\frac{K_{1} + 1}{K_{1} - 1} + y^{2} + a^{2} = 0 \qquad \qquad \left(x - a\frac{K_{1} + 1}{K_{1} - 1}\right)^{2} + y^{2} = \left(\frac{2a\sqrt{K_{1}}}{K_{1} - 1}\right)^{2}$$

This is the equation of a circle (actually a cylinder), displaced along the x axis, and of radius b, where

$$b = \frac{2a\sqrt{K_1}}{K_1 - 1}$$

It is all independent of z, as you might expect!

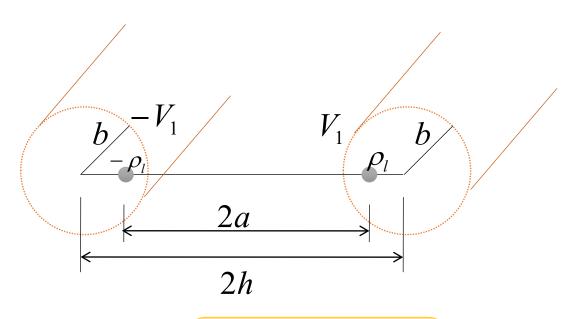
The equation of an equipotential surface is:

$$\left(x - a\frac{K_1 + 1}{K_1 - 1}\right)^2 + y^2 = \left(\frac{2a\sqrt{K_1}}{K_1 - 1}\right)^2$$

This is the equation of a circle (actually a cylinder), displaced along the x axis by distance h, and having radius b, where

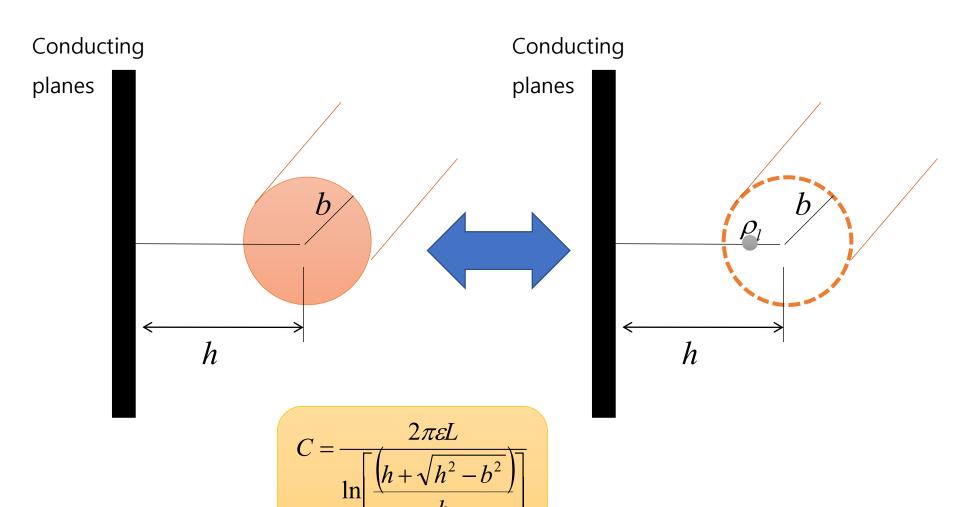
$$h = a \frac{K_1 + 1}{K_1 - 1}$$
 and $b = \frac{2a\sqrt{K_1}}{K_1 - 1}$

It is all independent of z, as you might expect!

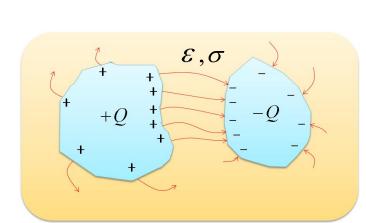


$$C = \frac{\pi \varepsilon L}{\ln \left[\frac{\left(h + \sqrt{h^2 - b^2}\right)}{b}\right]}$$

One-wire line and a conducting plane



Current Analogy



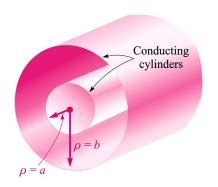
$$C = \frac{Q}{V} = \frac{\int_{v} \rho_{v} dv}{-\int \mathbf{E} \cdot d\mathbf{L}} = \frac{\int_{v} \nabla \cdot \mathbf{D} dv}{-\int \mathbf{E} \cdot d\mathbf{L}} = \frac{\oint_{S} \mathbf{D} \cdot d\mathbf{s}}{-\int \mathbf{E} \cdot d\mathbf{L}} = \frac{\oint_{S} \mathbf{E} \cdot d\mathbf{s}}{-\int \mathbf{E} \cdot d\mathbf{L}}$$

$$R = \frac{V}{I} = \frac{-\int \mathbf{E} \cdot d\mathbf{L}}{\oint_{S} \mathbf{J} \cdot d\mathbf{s}} = \frac{-\int \mathbf{E} \cdot d\mathbf{L}}{\oint_{S} \sigma \mathbf{E} \cdot d\mathbf{s}}$$

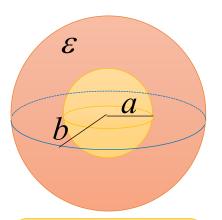
$$RC = \frac{\varepsilon}{\sigma}$$

Linear, homogeneous, isotropic material only

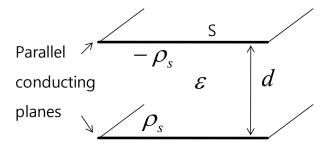
Examples



$$C = \frac{Q}{V} = \frac{2\pi\varepsilon L}{\ln\frac{b}{a}}$$

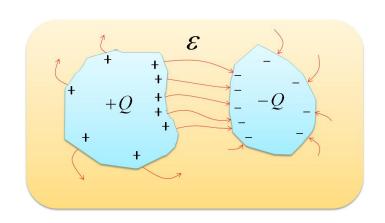


$$C = \frac{Q}{V} = \frac{2\pi\varepsilon L}{\ln\frac{b}{a}}$$



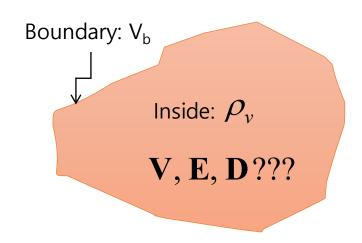
$$C = \frac{Q}{V} = \frac{\varepsilon S}{d}$$

Energy stored in a capacitor



$$W_E = \frac{1}{2} \int_{vol} \rho_v V dv = \frac{1}{2} \int_{vol} \mathbf{D} \cdot \mathbf{E} dv = \frac{1}{2} \int_{vol} \mathcal{E} E^2 dv$$
$$= \frac{1}{2} CV^2 = \frac{1}{2} QV = \frac{Q^2}{2C}$$

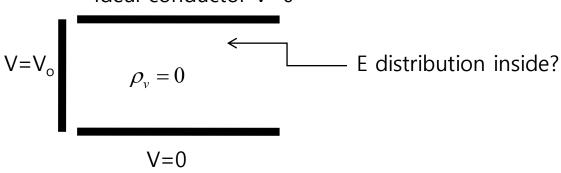
Poisson's / Laplace's Equation



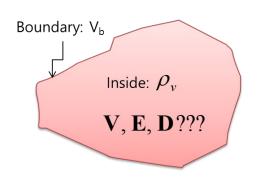
- When (1) internal charge distribution and (2) boundary condition are known,
- Calculate V, E, and D inside the boundary

Example

Ideal conductor V=0



Poisson's / Laplace's Equation



$$\nabla \cdot \mathbf{D} = \rho_{v}$$

$$\nabla \cdot \varepsilon \mathbf{E} = \rho_{v}$$

$$\nabla \cdot \varepsilon (-\nabla V) = \rho_{v}$$

$$\nabla \cdot \nabla V = -\frac{\rho_{v}}{\varepsilon}$$

$$\nabla^2 V = -\frac{\rho_v}{\varepsilon}$$
 Poisson's Equation

When no charge,

$$\nabla^2 V = 0$$
 Laplace's Equation

• E can be obtained by solving Poisson's or Laplace equation with boundary conditions

Laplacian

$$\nabla^2 V = \frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} + \frac{\partial^2 V}{\partial z^2}$$



Rectangular coordinates

$$\nabla^{2}V = \frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial V}{\partial \rho} \right) + \frac{1}{\rho^{2}} \left(\frac{\partial^{2}V}{\partial \phi^{2}} \right) + \frac{\partial^{2}V}{\partial z^{2}}$$



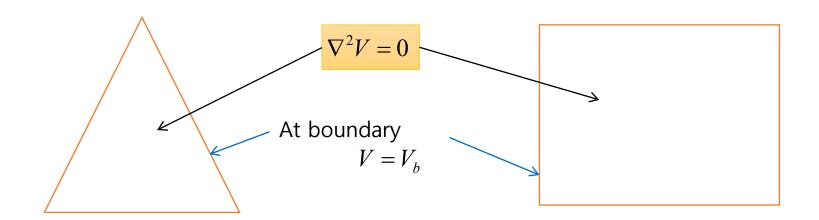
Cylindrical coordinates

$$\nabla^2 V = \frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial V}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial V}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 V}{\partial \phi^2} \implies \text{Spherical coordinates}$$

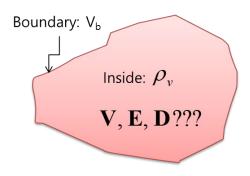
Uniqueness Theorem

Is the solution of Poisson's or Laplace's equation unique for given boundary condition?





Using Poisson's / Laplace's Equation



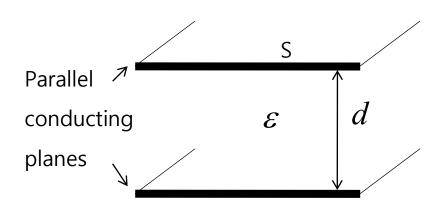
$$\nabla^2 V = -\frac{\rho_v}{\mathcal{E}}$$

- 1. Set coordinate system
- 2. Obtain V by solving Poisson's / Laplace equation with given boundary condition
- 3. Obtain E,D,Q using V

$$\mathbf{E} = -\nabla V$$

$$\mathbf{D} = \varepsilon \mathbf{E}$$

Surface charge $\rho_s = |\mathbf{D}|$



$$\mathbf{E} = -\nabla V = \frac{V_o}{d} \mathbf{a_z} \quad \leqslant \quad$$

$$\mathbf{D} = \frac{\varepsilon V_o}{\sigma} \mathbf{a_z} \quad \leqslant \quad$$

$$\mathbf{D} = \frac{c r_o}{d} \mathbf{a_z}$$

$$\rho_s = \frac{\varepsilon V_o}{d}$$

$$C = \frac{Q}{V} = \frac{\frac{\varepsilon V_o}{d} S}{V_o} = \frac{\varepsilon S}{d}$$

Laplace equation
$$\nabla^2 V = 0$$

BC:
$$V=V_0$$
 (at $z=0$), $V=0$ (at $z=d$)

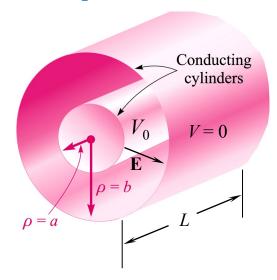
$$\nabla^2 V = 0$$

$$\frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} + \frac{\partial^2 V}{\partial z^2} = 0$$

$$\frac{d^2V}{dz^2} = 0$$
 From symmetry

$$V = Az + B$$

$$V = -\frac{V_o}{d}z + V_o$$



$$\mathbf{E} = -\nabla V = \frac{V_o}{\rho} \frac{1}{\ln(b/a)} \mathbf{a}_{\rho}$$

$$\mathbf{D} = \frac{\varepsilon V_o}{\rho} \frac{1}{\ln(b/a)} \mathbf{a}_{\rho}$$

$$\rho_s = \frac{\varepsilon V_o}{a} \frac{1}{\ln(b/a)}$$

$$\rho_s = \frac{\varepsilon V_o}{a} \frac{1}{\ln(b/a)} \qquad C = \frac{Q}{V} = \frac{2\pi \varepsilon L}{\ln(b/a)}$$

Laplace equation
$$\nabla^2 V = 0$$

BC:
$$V=V_0$$
 (at $\rho=a$), $V=0$ (at $\rho=b$)

$$\nabla^2 V = 0$$

$$\nabla^{2}V = \frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial V}{\partial \rho} \right) + \frac{1}{\rho^{2}} \left(\frac{\partial^{2}V}{\partial \phi^{2}} \right) + \frac{\partial^{2}V}{\partial z^{2}} = 0$$

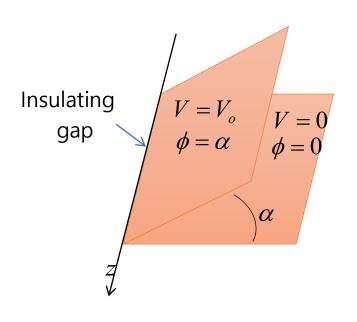
$$\nabla^{2}V = \frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial V}{\partial \rho} \right) + \frac{1}{\rho^{2}} \left(\frac{\partial^{2}V}{\partial \phi^{2}} \right) + \frac{\partial^{2}V}{\partial z^{2}} = 0$$

$$\frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial V}{\partial \rho} \right) = 0$$
From symmetry

$$V = A \ln \rho + B$$

$$V = A \ln \rho + B$$

$$V = V_o \frac{\ln(b/\rho)}{\ln(b/a)}$$
BC



$$\mathbf{E} = -\nabla V = -\frac{V_o}{\alpha \rho} \mathbf{a}_{\phi}$$

$$\mathbf{D} = -\frac{\varepsilon V_o}{\alpha \rho} \mathbf{a}_{\phi}$$

$$\rho_s = \frac{\varepsilon V_o}{\alpha \rho}$$

Laplace equation $\nabla^2 V = 0$

BC: $V=V_0$ (at $\phi=\alpha$), V=0 (at $\phi=0$)

$$\nabla^2 V = 0$$

$$\nabla^{2}V = \frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial V}{\partial \rho} \right) + \frac{1}{\rho^{2}} \left(\frac{\partial^{2}V}{\partial \phi^{2}} \right) + \frac{\partial^{2}V}{\partial z^{2}} = 0$$

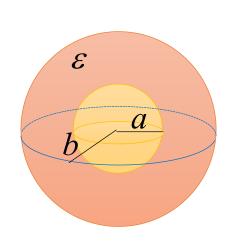
$$\nabla^{2}V = \frac{1}{\rho} \frac{\partial}{\partial \rho} \left(\rho \frac{\partial V}{\partial \rho} \right) + \frac{1}{\rho^{2}} \left(\frac{\partial^{2}V}{\partial \phi^{2}} \right) + \frac{\partial^{2}V}{\partial z^{2}} = 0$$

$$\frac{1}{\rho^{2}} \left(\frac{\partial^{2}V}{\partial \phi^{2}} \right) = 0$$
From symmetry

$$V = A\phi + B$$

$$V = V_o \frac{\phi}{\alpha}$$

$$=V_o\frac{\phi}{\alpha}$$



Laplace equation $\nabla^2 V = 0$

경계 조건: V=Vo (at r=a), V=0 (at r=b)

$$\nabla^2 V = 0$$

$$\frac{1}{r^{2}} \frac{\partial}{\partial r} \left(r^{2} \frac{\partial V}{\partial r} \right) + \frac{1}{r^{2} \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial V}{\partial \theta} \right) + \frac{1}{r^{2} \sin^{2} \theta} \frac{\partial^{2} V}{\partial \phi^{2}} = 0$$

$$\frac{1}{r^{2}} \frac{\partial}{\partial r} \left(r^{2} \frac{\partial V}{\partial r} \right) = 0$$
From symmetry

$$\mathbf{E} = -\nabla V = \frac{1}{r^2} \frac{V_o}{\frac{1}{a} - \frac{1}{b}} \mathbf{a_r}$$

$$\mathbf{D} = \frac{\varepsilon}{r^2} \frac{V_o}{\frac{1}{a} - \frac{1}{b}} \mathbf{a_r}$$

$$\rho_s = \frac{\varepsilon}{a^2} \frac{V_o}{\frac{1}{a} - \frac{1}{h}}$$

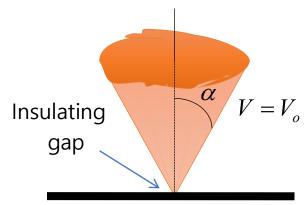
$$\rho_s = \frac{\varepsilon}{a^2} \frac{V_o}{\frac{1}{a} - \frac{1}{b}}$$

$$C = \frac{Q}{V} = \frac{4\pi\varepsilon}{\frac{1}{a} - \frac{1}{b}}$$

$$\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial V}{\partial r} \right) = 0$$
 From symmetry

$$V = \frac{A}{r} + B$$

$$V = V_o \frac{\frac{1}{r} - \frac{1}{b}}{\frac{1}{r} - \frac{1}{b}}$$
BO



Laplace equation
$$\nabla^2 V = 0$$

BC:
$$V=V_0$$
 (at $\theta=\alpha$), $V=0$ (at $\theta=\pi/2$)

$$\nabla^2 V = 0$$

$$V = 0$$

$$\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial V}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial V}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 V}{\partial \phi^2} = 0$$

$$E = -\nabla V = \frac{V_o}{r \sin \theta \ln \left(\cot \frac{\alpha}{2} \right)} \mathbf{a}_{\theta}$$

$$\frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial V}{\partial \theta} \right) = 0$$
From symmetry

$$\mathbf{E} = -\nabla V = \frac{V_o}{r \sin \theta \ln \left(\cot \frac{\alpha}{2}\right)} \mathbf{a_{\theta}}$$

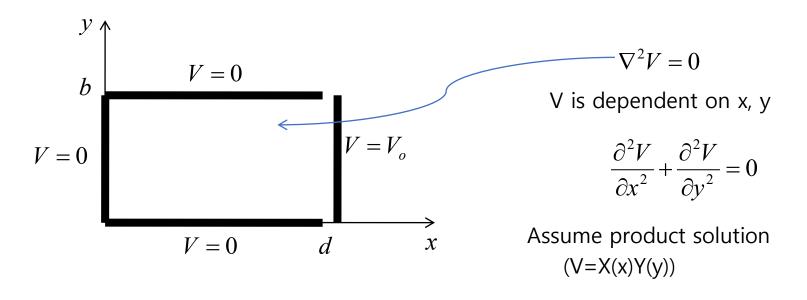
$$\mathbf{D} = \frac{\varepsilon V_o}{r \sin \theta \ln \left(\cot \frac{\alpha}{2}\right)} \mathbf{a}_{\theta}$$

$$\rho_s = \frac{\varepsilon V_o}{r \sin \alpha \ln \left(\cot \frac{\alpha}{2}\right)}$$

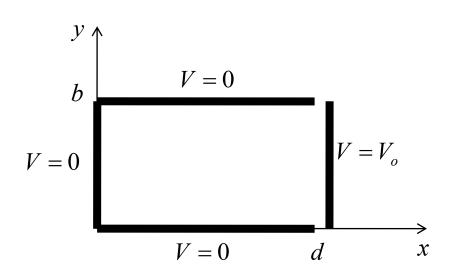
$$\frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial V}{\partial \theta} \right) = 0$$
 From symmetr

$$V = A \ln\left(\cot\frac{\theta}{2}\right) + B$$

$$V = V_o \frac{\ln\left(\cot\frac{\theta}{2}\right)}{\ln\left(\cot\frac{\alpha}{2}\right)}$$
BC



$$Y\frac{\partial^2 X}{\partial x^2} + X\frac{\partial^2 Y}{\partial y^2} = 0$$



$$Y \frac{\partial^2 X}{\partial x^2} + X \frac{\partial^2 Y}{\partial y^2} = 0$$

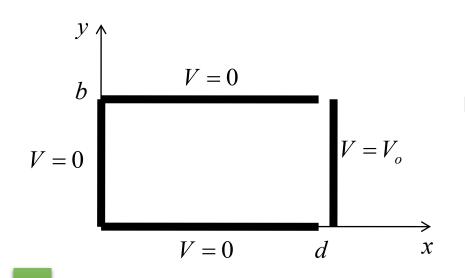
$$\frac{1}{X} \frac{\partial^2 X}{\partial x^2} = -\frac{1}{Y} \frac{\partial^2 Y}{\partial y^2} = k^2$$

$$\frac{\partial^2 X}{\partial x^2} = k^2 X$$

$$\frac{\partial^2 Y}{\partial y^2} = -k^2 Y$$

$$X = A \cosh kx + B \sinh kx = A'e^{kx} + B'e^{-kx}$$
$$Y = C \cos ky + D \sin ky = C'e^{iky} + D'e^{-iky}$$

A, B, C, D, k obtained from BC



$$X = A \cosh kx + B \sinh kx = A'e^{kx} + B'e^{-kx}$$

$$Y = C \cos ky + D \sin ky = C'e^{iky} + D'e^{-iky}$$

$$BC$$

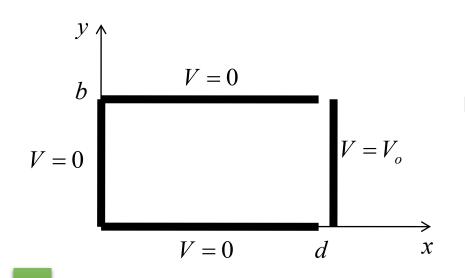
$$V = V_o \qquad at \qquad x = d$$

$$V = 0 \qquad at \qquad y = 0, \ y = d, \ x = 0$$

V = 0 at x = 0

 $V = (A\cosh kx + B\sinh kx)(C\cos ky + D\sin ky) = A(C\cos ky + D\sin ky) = 0$

$$\therefore A = 0$$



$$X = A \cosh kx + B \sinh kx = A'e^{kx} + B'e^{-kx}$$

$$Y = C \cos ky + D \sin ky = C'e^{iky} + D'e^{-iky}$$

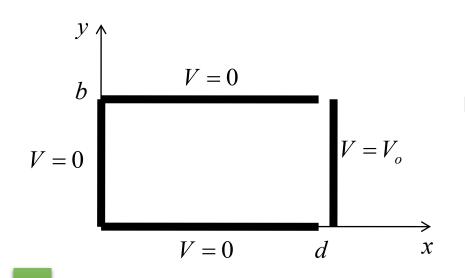
$$V = V_o$$

$$V = V_o \quad at \quad x = d$$

$$V = 0 \quad at \quad y = 0, y = d, x = 0$$

V = 0 at v = 0

 $V = B \sinh kx (C \cos ky + D \sin ky) = BC \sinh kx = 0$ $\therefore C = 0$



$$X = A \cosh kx + B \sinh kx = A'e^{kx} + B'e^{-kx}$$

$$Y = C \cos ky + D \sin ky = C'e^{iky} + D'e^{-iky}$$

$$V = V_o$$

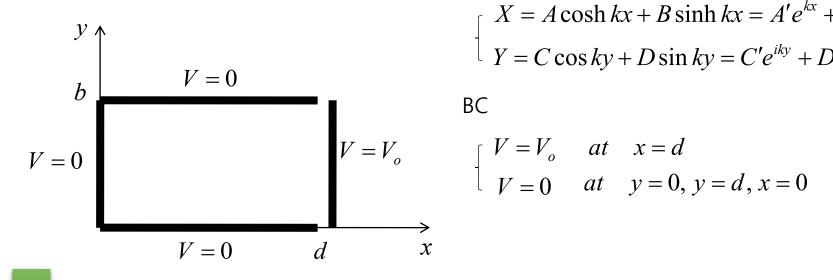
$$V = V_o \quad at \quad x = d$$

$$V = 0 \quad at \quad y = 0, y = d, x = 0$$

V = 0 at y = b

 $V = BD \sinh kx \sin ky = BD \sinh kx \sin kb = 0$

$$\therefore k = \frac{m\pi}{b}, \quad m = 0,1,\cdots$$



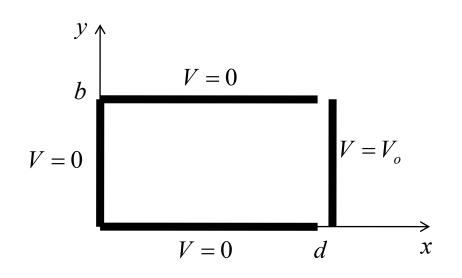
$$X = A \cosh kx + B \sinh kx = A'e^{kx} + B'e^{-kx}$$

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BC
$$V = V_o \quad at \quad x = d$$

$$V = 0 \quad at \quad y = 0, y = d, x = 0$$

$$V = \sum_{m=0}^{\infty} V_m \sinh \frac{m \pi x}{b} \sin \frac{m \pi y}{b} = \sum_{m=0}^{\infty} V_m \sinh \frac{m \pi d}{b} \sin \frac{m \pi y}{b} = \sum_{m=0}^{\infty} c_m \sin \frac{m \pi y}{b} = V_o$$

$$c_{m} = \frac{1}{b} \left[\int_{0}^{b} V_{o} \sin \frac{m\pi y}{b} dy + \int_{b}^{2b} -V_{o} \sin \frac{m\pi y}{b} dy \right] = \begin{cases} \frac{4V_{o}}{m\pi} & (m \text{ odd}) \\ 0 & (m \text{ even}) \end{cases}$$



$$X = A \cosh kx + B \sinh kx = A'e^{kx} + B'e^{-kx}$$

$$Y = C \cos ky + D \sin ky = C'e^{iky} + D'e^{-iky}$$

$$V = V_o$$

$$V = V_o \quad at \quad x = d$$

$$V = 0 \quad at \quad y = 0, y = d, x = 0$$

$$V = \frac{4V_o}{\pi} \sum_{1,odd}^{\infty} \frac{1}{m} \frac{\sinh(m\pi x/b)}{\sinh(m\pi d/b)} \sin\frac{m\pi y}{b}$$

Chapter Summary

- Capacitance definition
- Capacitance calculation
 - Parallel plate, Coaxial transmission line, Spherical capacitor
 - Isolated sphere with a dielectric coating, Parallel plate with two-layer dielectric
 - Two-wire line, one-wire line and a plane
- Energy stored in a capacitor
- Poisson / Laplace equation
 - Laplacian operator
 - Parallel plate, Coaxial transmission line, Angled plate, Concentric sphere, Cone above a conducting plane
 - Product solution