

# Exam 1 Cheatsheet (Methods & Templates)

## Conservation & Physical Meaning

**Heat eq (1D rod):**  $u_t = \alpha u_{xx}$ ,  $\alpha = K/(\rho c)$ .

**Flux:**  $\phi = -K u_x$ . **Energy density:**  $e = \rho c u$ .

**Integral conservation:**  $\frac{d}{dt} \int_0^L u \, dx = \frac{1}{\rho c K} \int_0^L Q \, dx$ .

**Equilibrium:**  $\Delta u = 0$ . With *pure Neumann*, steady state exists iff net flux = 0; solution unique up to constant (fix mean or value at a point).

**Divergence Thm** (2D/3D):  $\iint_{\Omega} \nabla \cdot F \, dA = \int_{\partial\Omega} F \cdot n \, ds$ ,  $\iiint_{\Omega} \nabla \cdot F \, dV = \iint_{\partial\Omega} F \cdot n \, dS$ .

## Boundary Conditions (heat)

**Dirichlet:** fixed temperature  $u = \text{given}$ . **Neumann:** insulated/flux given,  $u_x$  given.

**Robin (Newton cooling):**  $-K u_x = h(u - u_{\infty})$ .

**Periodic:**  $u(0, t) = u(L, t)$ ,  $u_x(0, t) = u_x(L, t)$ .

## Orthogonality you *don't* have on your sheet

**Half-integer cosine/sine** on  $[0, H]$ :

$$\int_0^H \cos \frac{(n + \frac{1}{2})\pi y}{H} \cos \frac{(m + \frac{1}{2})\pi y}{H} \, dy = \frac{H}{2} \delta_{mn},$$

$$\int_0^H \sin \frac{(n + \frac{1}{2})\pi y}{H} \sin \frac{(m + \frac{1}{2})\pi y}{H} \, dy = \frac{H}{2} \delta_{mn}.$$

**Fourier coeffs for these bases:**

$$C_n = \frac{2}{H} \int_0^H g(y) \cos \frac{(n + \frac{1}{2})\pi y}{H} \, dy, \quad S_n = \frac{2}{H} \int_0^H g(y) \sin \frac{(n + \frac{1}{2})\pi y}{H} \, dy.$$

## Heat Equation Templates (1D)

Domain  $0 < x < L$ ,  $u_t = \alpha u_{xx}$ .

**D–N (Dirichlet at  $x = 0$ , Neumann at  $x = L$ ):**

$$u(0, t) = 0, \quad u_x(L, t) = 0, \quad u(x, 0) = f(x).$$

Use eigenfunctions  $\sin \frac{(n + \frac{1}{2})\pi x}{L}$ .

$$u(x, t) = \sum_{n=0}^{\infty} b_n e^{-\alpha \left((n + \frac{1}{2})\pi/L\right)^2 t} \sin \frac{(n + \frac{1}{2})\pi x}{L},$$

$$b_n = \frac{2}{L} \int_0^L f(x) \sin \frac{(n + \frac{1}{2})\pi x}{L} \, dx.$$

(Satisfies  $u(0, t) = 0$  since  $\sin 0 = 0$  and  $u_x(L, t) = 0$  since  $\cos((n + \frac{1}{2})\pi) = 0$ .)

**N–D (Neumann at  $x = 0$ , Dirichlet at  $x = L$ ):**

$$u_x(0, t) = 0, \quad u(L, t) = 0, \quad u(x, 0) = f(x).$$

Use eigenfunctions  $\cos \frac{(n + \frac{1}{2})\pi x}{L}$ .

$$u(x, t) = \sum_{n=0}^{\infty} a_n e^{-\alpha \left((n + \frac{1}{2})\pi/L\right)^2 t} \cos \frac{(n + \frac{1}{2})\pi x}{L},$$

$$a_n = \frac{2}{L} \int_0^L f(x) \cos \frac{(n + \frac{1}{2})\pi x}{L} \, dx.$$

(Satisfies  $u_x(0, t) = 0$  since  $\sin 0 = 0$ , and  $u(L, t) = 0$  since  $\cos((n + \frac{1}{2})\pi) = 0$ .)

**Periodic (for reference):**

$$u(0, t) = u(L, t), \quad u_x(0, t) = u_x(L, t), \quad u(x, t) = A_0 + \sum_{n \neq 0} A_n e^{-\alpha(2\pi n/L)^2 t} e^{i2\pi n x/L}.$$

Real form:  $u(x, t) = a_0 + \sum_{n \geq 1} e^{-\alpha(2\pi n/L)^2 t} [a_n \cos(2\pi n x/L) + b_n \sin(2\pi n x/L)]$ .

## Laplace on Rectangles: Ready Forms

Let  $\Delta u = 0$  on  $0 < x < L$ ,  $0 < y < H$ .

**(D at  $y = 0, H$ ;  $u(0, y) = g$ ,  $u(L, y) = 0$ ):**

$$u = \sum_{n \geq 1} \frac{G_n}{\sinh(\alpha_n L)} \sinh(\alpha_n(L - x)) \sin(\alpha_n y), \quad \alpha_n = \frac{n\pi}{H}, \quad G_n = \frac{2}{H} \int_0^H g \sin(\alpha_n y) \, dy.$$

**(N at  $y = 0$ ; D at  $y = H$ ;  $u(0, y) = g$ ,  $u(L, y) = 0$ ):**

$$u = \sum_{n \geq 0} \frac{G_n}{\sinh(\alpha_n L)} \sinh(\alpha_n(L - x)) \cos(\alpha_n y), \quad \alpha_n = \frac{(n + \frac{1}{2})\pi}{H}.$$

**(N at  $x = 0, L$ ; N at  $y = H$ ;  $u(x, 0) = F$ ):**

$$u(x, y) = A_0 + \sum_{n \geq 1} A_n \cos \frac{n\pi x}{L} \frac{\cosh\left(\frac{n\pi}{L}(H - y)\right)}{\cosh\left(\frac{n\pi}{L}H\right)},$$

$$A_0 = \frac{1}{L} \int_0^L F, \quad A_n = \frac{2}{L} \int_0^L F \cos \frac{n\pi x}{L} dx.$$

**Trick:** write  $\sinh(\alpha(L-x))$  or  $\cosh(\alpha(L-x))$  to auto-satisfy a boundary at  $x = L$  (Dirichlet or Neumann).

## Neumann Compatibility (use this!)

For steady Laplace with Neumann data:

$$\int_{\partial\Omega} \partial_n u \, ds = \iint_{\Omega} \Delta u \, dA = 0.$$

Example: if  $u_x(0, y) = u_x(L, y) = u_y(x, 0) = 0$  and  $u_y(x, H) = f(x)$ , then *must have*  $\int_0^L f(x) \, dx = 0$ . Unique up to additive constant.

## Polar Geometry: Practical Forms

**Exterior circle** ( $r \geq a$ ), **finite at**  $\infty$ , **Dirichlet at**  $r = a$ :

$$u(r, \theta) = A_0 + \sum_{n \geq 1} \left(\frac{a}{r}\right)^n (A_n \cos n\theta + B_n \sin n\theta),$$

$A_0, A_n, B_n$  are the  $2\pi$  Fourier coefficients of  $u(a, \theta)$ . If  $u \rightarrow 0$  at  $\infty$ , set  $A_0 = 0$ .

**Quarter-disk** ( $0 \leq \theta \leq \pi/2$ ) with  $u_\theta(r, 0) = 0$ ,  $u(r, \frac{\pi}{2}) = 0$ ,  $u(1, \theta) = f(\theta)$ :

$$u(r, \theta) = \sum_{n=0}^{\infty} A_n r^{2n+1} \cos((2n+1)\theta), \quad A_n = \frac{4}{\pi} \int_0^{\pi/2} f(\varphi) \cos((2n+1)\varphi) \, d\varphi.$$

## Fourier Tips You'll Use

**Step on half-interval:**  $F(x) = \mathbf{1}_{[0, L/2)}(x)$  has cosine coeffs

$$A_n = \frac{2}{n\pi} \sin \frac{n\pi}{2} \Rightarrow A_{2k} = 0, \quad A_{2k+1} = \frac{2(-1)^k}{(2k+1)\pi}.$$

**Even/odd extensions** (to use pure cos/sin on  $[0, L]$ ):

Even  $\Rightarrow$  cosine series; Odd  $\Rightarrow$  sine series.

## Divergence Theorem (2D/3D Full Forms)

**General 2D (Cartesian):**

$$\iint_{\Omega} \left( \frac{\partial F_x}{\partial x} + \frac{\partial F_y}{\partial y} \right) dA = \oint_{\partial\Omega} (F_x n_x + F_y n_y) \, ds = \oint_{\partial\Omega} \mathbf{F} \cdot \mathbf{n} \, ds$$

**Rectangular region**  $\Omega = [0, L] \times [0, H]$ :

$$\iint_{\Omega} \left( \frac{\partial F_x}{\partial x} + \frac{\partial F_y}{\partial y} \right) dA = \int_0^H [F_x(L, y) - F_x(0, y)] \, dy + \int_0^L [F_y(x, H) - F_y(x, 0)] \, dx.$$

For heat flux  $\mathbf{F} = -K\nabla u$ , this equals total outward heat flow through all 4 sides.

**Polar coordinates**  $(r, \theta)$ : If  $\mathbf{F} = F_r \mathbf{e}_r + F_\theta \mathbf{e}_\theta$ ,

$$\iint_{\Omega} \left[ \frac{1}{r} \frac{\partial}{\partial r} (r F_r) + \frac{1}{r} \frac{\partial F_\theta}{\partial \theta} \right] r \, dr \, d\theta = \oint_{\partial\Omega} \mathbf{F} \cdot \mathbf{n} \, ds$$

For a circular region  $r = a$ :

$$\iint_{\Omega} \left[ \frac{1}{r} \frac{\partial}{\partial r} (r F_r) + \frac{1}{r} \frac{\partial F_\theta}{\partial \theta} \right] r \, dr \, d\theta = \int_0^{2\pi} F_r(a, \theta) a \, d\theta.$$

For  $\mathbf{F} = -K\nabla u$ , this gives total heat leaving the circle of radius  $a$ .

**3D version (for reference):**

$$\iiint_{\Omega} \left( \frac{\partial F_x}{\partial x} + \frac{\partial F_y}{\partial y} + \frac{\partial F_z}{\partial z} \right) dV = \iint_{\partial\Omega} (F_x n_x + F_y n_y + F_z n_z) \, dS.$$

## Proof Tools (one-liners)

**Dirichlet uniqueness** for  $\Delta u = g$ : difference is harmonic  $\Rightarrow$  max principle  $\Rightarrow$  zero.

**Mean value property:** harmonic  $u(x_0)$  equals average of  $u$  on any circle/sphere centered at  $x_0$  inside domain (quick sanity check).

**Strong max principle:** interior max/min  $\Rightarrow u$  constant.

## Checklists

**1D Heat:**

- Choose sin/cos basis.
- Compute Fourier coeffs of  $u(x, 0)$ .
- Multiply by  $e^{-\alpha \lambda_n t}$ .
- Add constant mode (Neumann).

**Polar:**

- Periodicity  $\Rightarrow$  angular  $\cos n\theta, \sin n\theta$ .
- Radial Euler  $\Rightarrow$  keep  $r^n$  (interior) or  $r^{-n}$  (exterior).
- Coeffs are Fourier coeffs of boundary data.