

MATH 345 Homework 1 Solutions

Preliminary Def: Thermal Energy (T.E.), Heat Energy (H.E.), Heat Flux (H.F.), Thermal Energy Density (T.E.D.), Heat Source (H.S.), Specific Heat (S.H.)

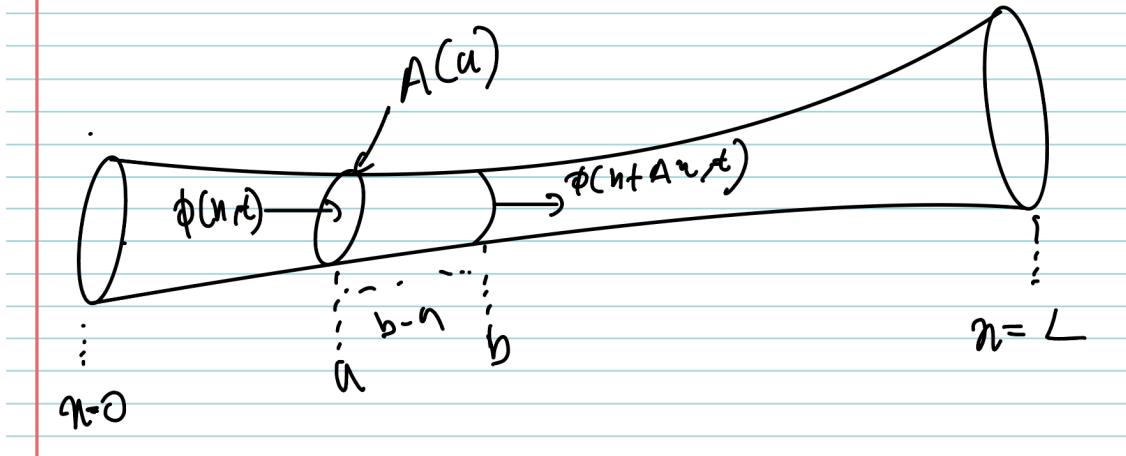
Problem 1.2.4

Figure 1: Rod

We define our surface area of the rod as $A(x)$.

T.E.D.:

$$e(x, t)$$

H.E.:

$$e(x, t)A(x) \cdot (b - a)$$

where the volume of a slice is $A(x) \cdot (b - a)$ and we define $(b - a) = \Delta x$.

$$\frac{\partial}{\partial t} [e(x, t) \cdot A(x) \Delta x] \approx \phi(x, t)A(x) - \phi(x + \Delta x, t)A(x + \Delta x)$$

(L: time rate of change of energy in slice. R: inflow at x minus outflow at $x + \Delta x$ plus sources.)

In the limit as $\Delta x \rightarrow 0$, divide by Δx and we use

$$\phi(x + \Delta x, t)A(x + \Delta x) = \phi(x, t)A(x) + \partial_x[\phi(x, t)A(x)] \Delta x + o(\Delta x)$$

which yields

$$\frac{\partial}{\partial t}[e(x, t)A(x)] = - \frac{\partial}{\partial x}[\phi(x, t)A(x)]$$

$$e(x, t) = \rho c T(x, t), \quad \phi(x, t) = -k \frac{\partial T}{\partial x}(x, t).$$

$$\frac{\partial}{\partial t}[\rho c A(x) T(x, t)] = - \frac{\partial}{\partial x}[-k A(x) T_x(x, t)].$$

Since $A = A(x)$ is time-independent and k, ρ, c are constants,

$$\rho c A(x) T_t(x, t) = \frac{\partial}{\partial x}(k A(x) T_x(x, t)).$$

Hence the heat eq. for a rod with $A(x)$ and $Q = 0$:

$$(\rho c) A(x) T_t = (k A(x) T_x)_x \iff T_t = \alpha \frac{1}{A(x)} \frac{\partial}{\partial x} (A(x) T_x), \quad \alpha = \frac{k}{\rho c}.$$

One way we can confirm we have the right answer is to check if $A(x) \equiv A_0$, then $(kA_0 T_x)_x = kA_0 T_{xx}$ and we get $T_t = \alpha T_{xx}$.

Problem 1.2.7

Suppose that S.H. is $c = c(x, u)$

(a)

For a small temperature to raise $u \rightarrow u + du$ at fixed x requires per unit mass heat

$$dq = c(x, u) du.$$

$$q(x, t) = \int_0^{u(x, t)} c(x, \bar{u}) d\bar{u},$$

(b)

T.E. per unit volume is

$$e(x, t) = \rho(x) q(x, t) = \rho(x) \int_0^{u(x, t)} c(x, \bar{u}) d\bar{u}.$$

Conservation of thermal energy

$$\frac{\partial e}{\partial t} = - \frac{\partial \phi}{\partial x} + Q,$$

Leibniz rule Since \bar{u} has no explicit t -dependence,

$$\frac{\partial}{\partial t} \left[\int_0^{u(x, t)} c(x, \bar{u}) d\bar{u} \right] = c(x, u(x, t)) u_t(x, t)$$

Hence

$$\frac{\partial e}{\partial t} = \rho(x) c(x, u) u_t(x, t)$$

$$\rho(x) c(x, u) u_t = -\phi_x + Q$$

With Fourier's law

$$\begin{aligned} \phi(x, t) &= -K_0(x) u_x \implies -\phi_x = \partial_x(K_0(x) u_x) \\ \rho(x) c(x, u) u_t &= \frac{\partial}{\partial x} (K_0(x) u_x) + Q(x, t) \end{aligned}$$

If K_0, ρ are constants and $Q = 0$ this is

$$(\rho c(x, u)) u_t = K_0 u_{xx}, \quad u_t = \frac{K_0}{\rho} \frac{u_{xx}}{c(x, u)}$$

Problem 1.4.2

Equilibrium $\Rightarrow u_t = 0$. With constant properties and Fourier's law $\phi = -K_0 u_x$

$$\frac{\partial e}{\partial t} = 0 = -\phi_x + Q \implies (K_0 u_x)_x + Q = 0 \implies K_0 u_{xx} + Q = 0$$

(a)

$$Q = 0, \quad u_x(0) = 0, \quad u(L) = T.$$

$$u_{xx} = 0 \implies u(x) = ax + b, \quad u_x = a.$$

From $u_x(0) = 0$ we get $a = 0$, hence $u(x) \equiv b$. Using $u(L) = T$ gives

$$u(x) \equiv T$$

(b)

$$\frac{Q}{K_0} = 1 \implies Q = K_0, \quad u(0) = T_1, \quad u(L) = T_2.$$

$$K_0 u_{xx} + K_0 = 0 \implies u_{xx} = -1.$$

$$u_x(x) = \int u_{xx}(x) dx = \int (-1) dx = -x + C_1$$

$$u(x) = \int u_x(x) dx = \int (-x + C_1) dx = -\frac{x^2}{2} + C_1 x + C_2$$

Apply BCs:

$$u(0) = T_1 \implies C_2 = T_1, \quad u(L) = T_2 \implies -\frac{L^2}{2} + C_1 L + T_1 = T_2 \implies C_1 = \frac{T_2 - T_1}{L} + \frac{L}{2}$$

Therefore

$$u(x) = -\frac{x^2}{2} + \left(\frac{T_2 - T_1}{L} + \frac{L}{2}\right)x + T_1 = T_1 + \frac{T_2 - T_1}{L} x + \frac{x(L-x)}{2}$$

$$u(x) = T_1 + \frac{T_2 - T_1}{L} x + \frac{x(L-x)}{2}$$

Problem 1.4.8

$$u_{xx} = 0 \implies u(x) = ax + b, \quad u_x(x) \equiv a$$

BCs give

$$u_x(0, t) = 1 \Rightarrow a = 1, \quad u_x(L, t) = \beta \Rightarrow a = \beta$$

Consistency $\Rightarrow \beta = 1$

If $\beta = 1$, the family of steady states is

$$u_{\text{eq}}(x) = x + C, \quad C \in \mathbb{R}$$

Physical explanation. Integrate the PDE over $(0, L)$:

$$\frac{d}{dt} \int_0^L u \, dx = \int_0^L u_{xx} \, dx = u_x(L, t) - u_x(0, t) = \beta - 1.$$

Thus:

$$\begin{aligned} \beta \neq 1 &\Rightarrow \text{net heat flux } \neq 0 \Rightarrow \int_0^L u \, dx \text{ drifts linearly in time } \Rightarrow \text{no equilibrium.} \\ \beta = 1 &\Rightarrow \text{net flux } = 0 \Rightarrow \text{equilibrium exists.} \end{aligned}$$

Which C ? When $\beta = 1$, the spatial mean is conserved:

$$\begin{aligned} \int_0^L u(x, t) \, dx &= \int_0^L f(x) \, dx. \\ \int_0^L (x + C) \, dx &= \int_0^L f(x) \, dx \Rightarrow C = \frac{1}{L} \int_0^L f(x) \, dx - \frac{L}{2}. \end{aligned}$$

So the steady state compatible with the initial data is

$$u_{\text{eq}}(x) = x + \left(\frac{1}{L} \int_0^L f(x) \, dx - \frac{L}{2} \right) \quad \text{exists iff } \beta = 1.$$

Problem 1.4.11

Let

$$E(t) = \int_0^L u(x, t) dx$$

$$E'(t) = \int_0^L u_t dx = \int_0^L (u_{xx} + 4) dx = [u_x]_0^L + 4L = u_x(L, t) - u_x(0, t) + 4L = 6 - 5 + 4L = 1 + 4L.$$

$$E(t) = E(0) + (1 + 4L)t = \int_0^L f(x) dx + (1 + 4L)t.$$

With constant density ρ , S.H. c , and cross-sectional area A , total T.E. is

$$\mathcal{H}(t) = \rho c A E(t) = \rho c A \left(\int_0^L f(x) dx + (1 + 4L)t \right).$$

Problem 1.5.3. $x = r \cos \theta$, $y = r \sin \theta$.

(a) Since $r^2 = x^2 + y^2$,

$$2r r_x = 2x \Rightarrow r_x = \frac{x}{r} = \cos \theta, \quad 2r r_y = 2y \Rightarrow r_y = \frac{y}{r} = \sin \theta.$$

Also $\theta = \arctan(y/x)$, so

$$\theta_x = \frac{-y}{x^2 + y^2} = -\frac{\sin \theta}{r}, \quad \theta_y = \frac{x}{x^2 + y^2} = \frac{\cos \theta}{r}.$$

(b) Position vector $\vec{r} = x \hat{i} + y \hat{j} = r(\cos \theta \hat{i} + \sin \theta \hat{j})$, so

$$\hat{r} = \cos \theta \hat{i} + \sin \theta \hat{j}, \quad \hat{\theta} = -\sin \theta \hat{i} + \cos \theta \hat{j}$$

(Obtain $\hat{\theta}$ by $+90^\circ$ rotation or by $\partial_\theta \hat{r}$.)

(c) Chain rule:

$$\partial_x = r_x \partial_r + \theta_x \partial_\theta = \cos \theta \partial_r - \frac{\sin \theta}{r} \partial_\theta, \quad \partial_y = r_y \partial_r + \theta_y \partial_\theta = \sin \theta \partial_r + \frac{\cos \theta}{r} \partial_\theta.$$

Hence

$$\nabla = \hat{i} \partial_x + \hat{j} \partial_y = (\cos \theta \hat{i} + \sin \theta \hat{j}) \partial_r + (-\sin \theta \hat{i} + \cos \theta \hat{j}) \frac{1}{r} \partial_\theta = \boxed{\hat{r} \partial_r + \hat{\theta} \frac{1}{r} \partial_\theta}.$$

Applying to u ,

$$\nabla u = \hat{r} u_r + \hat{\theta} \frac{1}{r} u_\theta$$