

Modeling of 2D and 3D Gantry Systems

DPMA

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1 Dynamic Model of the 2D Gantry System

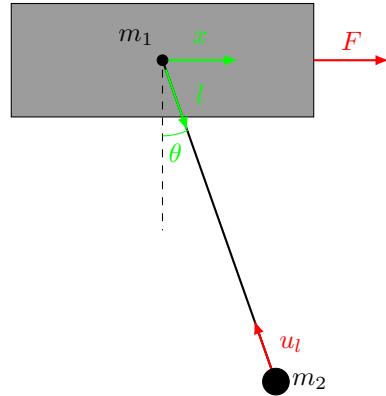


Figure 1: Diagram of the 2D-system

1.1 Kinematics

The generalized coordinates and external forces are $q^T = [x \ l \ \theta]$, $Q^T = [F, u_l, 0]$. The position of the load can then be represented as $p(q)$:

$$p(q) = \begin{bmatrix} x_2 \\ z_2 \end{bmatrix} = \begin{bmatrix} x + l \sin \theta \\ -l \cos \theta \end{bmatrix}.$$

Such that the Jacobian $J(q)$ is:

$$J(q) = \begin{bmatrix} \frac{\partial x_2}{\partial x} & \frac{\partial x_2}{\partial l} & \frac{\partial x_2}{\partial \theta} \\ \frac{\partial z_2}{\partial x} & \frac{\partial z_2}{\partial l} & \frac{\partial z_2}{\partial \theta} \end{bmatrix} = \begin{bmatrix} 1 & \sin \theta & l \cos \theta \\ 0 & -\cos \theta & l \sin \theta \end{bmatrix}.$$

The payload velocity $\dot{p}(q)$ is given by:

$$\dot{p} = J\dot{q} = \begin{bmatrix} \dot{x}_2 \\ \dot{z}_2 \end{bmatrix} = \begin{bmatrix} \dot{x} + \dot{l} \sin \theta + l \dot{\theta} \cos \theta \\ -\dot{l} \cos \theta + \dot{\theta} l \sin \theta \end{bmatrix}.$$

1.2 Energies

The cart with mass m_1 only moves in x :

$$T_1 = \frac{1}{2}m_1\dot{x}^2.$$

While the payload moves with velocity \dot{p} :

$$T_2 = \frac{1}{2}m_2\dot{p}^T\dot{p} = \frac{1}{2}m_2(J\dot{q})^T(J\dot{q}) = \frac{1}{2}m_2\dot{q}^T(J^T J)\dot{q}.$$

Total kinetic energy:

$$T = \frac{1}{2}m_1\dot{x}^2 + \frac{1}{2}m_2\dot{q}^T(J^T J)\dot{q}.$$

The only potential energy in this system comes from the payload:

$$V = m_2g(-l \cos \theta) = -m_2gl \cos \theta.$$

1.3 Compact Matrix Setup for Euler-Lagrange

$$L = T - V$$

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}_i} \right) - \frac{\partial L}{\partial q_i} = Q_i$$

Since V has no dependence on \dot{q} :

$$\begin{aligned}\frac{\partial L}{\partial \dot{q}} &= \frac{\partial T}{\partial \dot{q}} \\ \frac{\partial T_1}{\partial \dot{q}} &= \begin{bmatrix} m_1 \dot{x} \\ 0 \\ 0 \end{bmatrix} \\ \frac{\partial T_2}{\partial \dot{q}} &= m_2 J^T J \dot{q} = m_2 J^T \dot{p} \\ \frac{\partial T}{\partial \dot{q}} &= \begin{bmatrix} m_1 \dot{x} \\ 0 \\ 0 \end{bmatrix} + m_2 J^T \dot{p}\end{aligned}$$

Taking the time derivative:

$$\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}} \right) = \begin{bmatrix} m_1 \ddot{x} \\ 0 \\ 0 \end{bmatrix} + m_2 (J^T \dot{p} + J^T \ddot{p}) = \begin{bmatrix} m_1 \ddot{x} \\ 0 \\ 0 \end{bmatrix} + m_2 (J^T J \dot{q} + J^T \ddot{p})$$

Using $\ddot{p} = \frac{d}{dt}(J \dot{q}) = J \dot{q} + J \ddot{q}$:

$$\begin{aligned}\frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}} \right) &= \begin{bmatrix} m_1 \ddot{x} \\ 0 \\ 0 \end{bmatrix} + m_2 (J^T J \dot{q} + J^T (J \dot{q} + J \ddot{q})) \\ &= \begin{bmatrix} m_1 \ddot{x} \\ 0 \\ 0 \end{bmatrix} + m_2 (J^T J \dot{q} + J^T J \dot{q} + J^T J \ddot{q})\end{aligned}$$

For computing $\frac{\partial L}{\partial q}$, only T_2 and V depend on q :

$$\frac{\partial T_2}{\partial q} = \frac{1}{2} m_2 \frac{\partial}{\partial q} (\dot{p}^T \dot{p}).$$

Using $\dot{p}^T \dot{p} = \dot{q}^T J^T J \dot{q}$:

$$\frac{\partial T_2}{\partial q} = \frac{1}{2} m_2 \left(\dot{q}^T \frac{\partial}{\partial q} (J^T J) \dot{q} \right) = \frac{1}{2} m_2 \dot{q}^T \left[\left(\frac{\partial J}{\partial q} \right)^T J + J^T \left(\frac{\partial J}{\partial q} \right) \right] \dot{q}.$$

Knowing that $\frac{\partial T_2}{\partial q_i}$ is a scalar quadratic form in \dot{q} , symmetry $a^T = a$ can be used:

$$\begin{aligned}\left[\dot{q}^T \left(\frac{\partial J}{\partial q} \right)^T J \dot{q} \right]^T &= \dot{q}^T J^T \left(\frac{\partial J}{\partial q} \right) \dot{q}, \\ \frac{\partial T_2}{\partial q} &= \frac{1}{2} m_2 2 \begin{bmatrix} \dot{q}^T \left(\frac{\partial J}{\partial x} \right)^T J \dot{q} \\ \dot{q}^T \left(\frac{\partial J}{\partial l} \right)^T J \dot{q} \\ \dot{q}^T \left(\frac{\partial J}{\partial \theta} \right)^T J \dot{q} \end{bmatrix} = m_2 \begin{bmatrix} \dot{q}^T \left(\frac{\partial J}{\partial x} \right)^T J \dot{q} \\ \dot{q}^T \left(\frac{\partial J}{\partial l} \right)^T J \dot{q} \\ \dot{q}^T \left(\frac{\partial J}{\partial \theta} \right)^T J \dot{q} \end{bmatrix}.\end{aligned}$$

Finally for the potential energy component:

$$\frac{\partial V}{\partial q} = \begin{bmatrix} 0 \\ -m_2 g \cos \theta \\ m_2 g l \sin \theta \end{bmatrix}.$$

So:

$$\frac{\partial L}{\partial q} = m_2 \begin{bmatrix} \dot{q}^T \left(\frac{\partial J}{\partial x} \right)^T J \dot{q} \\ \dot{q}^T \left(\frac{\partial J}{\partial l} \right)^T J \dot{q} \\ \dot{q}^T \left(\frac{\partial J}{\partial \theta} \right)^T J \dot{q} \end{bmatrix} - \begin{bmatrix} 0 \\ -m_2 g \cos \theta \\ m_2 g l \sin \theta \end{bmatrix}$$

1.4 Explicit Expressions for the system

To write the equations in their explicit form with the addition of a damping term Q_D :

$$M(q)\ddot{q} + C(q, \dot{q})\dot{q} + G(q) = Q - Q_d\dot{q}, \quad (1)$$

Definitions for the following are required:

$$\begin{aligned} J &= \begin{bmatrix} 0 & \dot{\theta} \cos \theta & \dot{l} \cos \theta - \dot{\theta} l \sin \theta \\ 0 & \dot{\theta} \sin \theta & \dot{l} \sin \theta + \dot{\theta} l \cos \theta \end{bmatrix}, \\ M_1 &= \begin{bmatrix} m_1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}, \\ G(q) &= \begin{bmatrix} 0 \\ -m_2 g \cos \theta \\ m_2 g l \sin \theta \end{bmatrix}. \end{aligned}$$

As well as:

$$\begin{aligned} A_1 &= \frac{\partial J}{\partial x} = \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}, \\ A_2 &= \frac{\partial J}{\partial l} = \begin{bmatrix} 0 & 0 & \cos \theta \\ 0 & 0 & \sin \theta \end{bmatrix}, \\ A_3 &= \frac{\partial J}{\partial \theta} = \begin{bmatrix} 0 & \cos \theta & -l \sin \theta \\ 0 & \sin \theta & l \cos \theta \end{bmatrix}. \end{aligned}$$

Now defining:

$$\Gamma(q, \dot{q}) = m_2 \begin{bmatrix} \dot{q}^T A_1^T J \dot{q} \\ \dot{q}^T A_2^T J \dot{q} \\ \dot{q}^T A_3^T J \dot{q} \end{bmatrix},$$

and forming the matrices

$$\begin{aligned} M_1 &= \begin{bmatrix} m_1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{bmatrix}, \\ Q_d &= \begin{bmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & c \end{bmatrix}. \end{aligned}$$

Writing it out fully gives:

$$M_1 \ddot{q} + m_2 J^T J \ddot{q} + m_2 (J^T J + J^T J) \dot{q} - (\Gamma - G) = Q - Q_d \dot{q},$$

and finally:

$$\underbrace{(M_1 + m_2 J^T J)}_{M(q)} \ddot{q} + \underbrace{(m_2 J^T J + m_2 J^T J) \dot{q}}_{C(q, \dot{q}) \dot{q}} - \Gamma + G = Q - Q_d \dot{q}.$$

2 2D Gantry Optimal Control Problem

2.1 Reformulating Dynamics

For usage in optimal control, a nonlinear state space model should be formed with the chosen state space and controls vector: $x^T = [x, l, \theta, \dot{x}, \dot{l}, \dot{\theta}]$, $u^T = [F, u_l]$

$$\dot{x} = \begin{pmatrix} \dot{q} \\ \ddot{q} \end{pmatrix} = \begin{pmatrix} \dot{q} \\ M^{-1}(Q - (Q_d - C(q, \dot{q}))\dot{q} - G(q)) \end{pmatrix} = f(x, u)$$