



TAREA 4

1. Problema 1

Dada la varianza $\langle(\Delta E)^2\rangle = \langle(E_i - \varepsilon)^2\rangle = \langle E_i^2\rangle - \varepsilon^2$, entonces, encontramos el segundo momento

$$\frac{d^2 \mathfrak{z}}{d\beta^2} = \sum E_i^2 e^{-\beta E_i} = \underbrace{\left(\sum p_i E_i^2 \right)}_{\langle E_i^2 \rangle} \mathfrak{z}. \quad (1)$$

Ahora, tomando la segunda derivada de $\ln \mathfrak{z}$, se tiene

$$\frac{d^2 \ln \mathfrak{z}}{d\beta^2} = \frac{d}{d\beta} \left(\frac{1}{\mathfrak{z}} \frac{d\mathfrak{z}}{d\beta} \right) = -\frac{1}{\mathfrak{z}^2} \left(\frac{d\mathfrak{z}}{d\beta} \right)^2 + \frac{1}{\mathfrak{z}} \frac{d^2 \mathfrak{z}}{d\beta^2},$$

despejando y reemplazando en la fórmula de la varianza

$$\langle(\Delta E)^2\rangle = \frac{d^2 \ln \mathfrak{z}}{d\beta^2} + \frac{1}{\mathfrak{z}^2} \left(\frac{d\mathfrak{z}}{d\beta} \right)^2 - \left(-\frac{1}{\mathfrak{z}} \frac{d\mathfrak{z}}{d\beta} \right)^2$$

$$\boxed{\langle(\Delta E)^2\rangle = \frac{d^2 \ln \mathfrak{z}}{d\beta^2}.$$

2. Problema 2

Dada la definición de tercer momento, se realiza la expación

$$\langle(\Delta E)^3\rangle = \sum_{E_i^3 - 3E_i^2\varepsilon + 3E_i\varepsilon^2 - \varepsilon^3} \underbrace{(E_i - \varepsilon)^3}_{p_i} = \langle E_i^3 \rangle - 3\varepsilon \langle E_i^2 \rangle + 3\varepsilon^3 - \varepsilon^3 = \langle E_i^3 \rangle - 3\varepsilon \langle E_i^2 \rangle + 2\varepsilon^3. \quad (2)$$

Ahora, siguiendo la idea del problema anterior

$$-\frac{1}{\mathfrak{z}} \frac{d^3 \mathfrak{z}}{d\beta^3} = \langle E_i^3 \rangle,$$

entontrando la tercera derivada de $\ln \mathfrak{z}$,

$$\frac{d^3 \ln \mathfrak{z}}{d\beta^3} = \frac{2}{\mathfrak{z}^3} \left(\frac{d\mathfrak{z}}{d\beta} \right)^3 - \frac{3}{\mathfrak{z}^2} \left(\frac{d\mathfrak{z}}{d\beta} \right) \frac{d^2 \mathfrak{z}}{d\beta^2} + \frac{1}{\mathfrak{z}} \frac{d^3 \mathfrak{z}}{d\beta^3}. \quad (3)$$

Sustituyendo (1), (3) y ε en (2)

$$\langle(\Delta E)^3\rangle = -\frac{d^3 \ln \mathfrak{z}}{d\beta^3} + \frac{2}{\mathfrak{z}^3} \left(\frac{d\mathfrak{z}}{d\beta} \right)^3 - \frac{3}{\mathfrak{z}^2} \left(\frac{d\mathfrak{z}}{d\beta} \right) \frac{d^2 \mathfrak{z}}{d\beta^2} - 3 \left(-\frac{1}{\mathfrak{z}} \frac{d\mathfrak{z}}{d\beta} \right) \left(\frac{1}{\mathfrak{z}} \frac{d^2 \mathfrak{z}}{d\beta^2} \right) + 2 \left(-\frac{1}{\mathfrak{z}} \frac{d\mathfrak{z}}{d\beta} \right)^3,$$

$$\boxed{\langle(\Delta E)^3\rangle = -\frac{d^3 \ln \mathfrak{z}}{d\beta^3}.$$

3. Problema 3

Dada la definición de calor específico

$$c_V = \frac{\partial \varepsilon}{\partial T},$$

entonces, tomando la energía promedio

$$\varepsilon = -\frac{\partial \ln \mathfrak{z}}{\partial \beta} = k_B T^2 \frac{\partial \ln \mathfrak{z}}{\partial T}.$$

Reemplazando en la definición de calor específico

$$c_V = \frac{\partial}{\partial T} \left(k_B T^2 \frac{\partial \ln \mathfrak{z}}{\partial T} \right),$$

$$c_V = 2k_B T \frac{\partial \ln \mathfrak{z}}{\partial T} + k_B T \frac{\partial^2 \ln \mathfrak{z}}{\partial T^2}.$$

4. Problema 4

Sabiendo que la presión la podemos escribir en términos de la energía libre de Helmholtz

$$P = - \left(\frac{\partial A}{\partial V} \right)_T,$$

tomando a A como

$$A = -\frac{\ln Z}{\beta}, \quad Z = \mathfrak{z}^n.$$

Reemplazando A en la presión, se tiene

$$P = \left(\ln \mathfrak{z} \frac{\partial}{\partial V} \frac{\partial}{\partial T} \frac{1}{\beta} + \frac{1}{\beta} \frac{\partial \ln Z}{\partial V} \right)_T,$$

entonces

$$P = k_B T \left(\frac{\partial \ln Z}{\partial V} \right)_T = n k_B T \left(\frac{\partial \ln \mathfrak{z}}{\partial V} \right)_T.$$