

ACOUSTIC MODES IMPLEMENTATION

New model equations:

New parallel thermal velocity equation:

$$n_0 m_p \frac{\partial v_{\parallel, th}}{\partial t} = -\hat{b} \cdot \vec{\nabla} p - \frac{\partial \vec{B}}{|B|} \cdot \vec{\nabla} p_0$$

$$= -\frac{2\psi'}{\sqrt{g}} \frac{1}{a^2 R_0 B_0} \left(\frac{\partial}{\partial \zeta} + \tau \frac{\partial}{\partial \theta} \right) p - \frac{\vec{\nabla} \psi \wedge \hat{e}_\zeta}{R_0 B_0} \cdot \vec{\nabla} p_0$$

$$\vec{\nabla} \psi \wedge \hat{e}_\zeta = \vec{\nabla} \wedge (\psi \hat{e}_\zeta) - \cancel{\psi (\vec{\nabla} \wedge \hat{e}_\zeta)}$$

$$\Rightarrow \vec{\nabla} \wedge (\psi \hat{e}_\zeta) = \frac{1}{\sqrt{g}} \begin{vmatrix} \hat{e}_\rho & \hat{e}_\theta & \hat{e}_\zeta \\ \frac{\partial}{\partial \rho} & \frac{1}{\rho} \frac{\partial}{\partial \theta} & \frac{1}{R_0} \frac{\partial}{\partial \zeta} \\ 0 & 0 & \psi \end{vmatrix} = \frac{1}{\sqrt{g}} \left(\frac{1}{\rho} \frac{\partial \psi}{\partial \theta} \hat{e}_\rho - \frac{\partial \psi}{\partial \rho} \hat{e}_\theta \right)$$

$$\Rightarrow n_0 m_p \frac{\partial v_{\parallel, th}}{\partial t} = -\frac{2\psi'}{\sqrt{g}} \frac{1}{a^2 R_0 B_0} \left(\frac{\partial}{\partial \zeta} + \tau \frac{\partial}{\partial \theta} \right) p - \frac{1}{\sqrt{g}} \left(\frac{1}{\rho} \frac{\partial \psi}{\partial \theta} \hat{e}_\rho - \frac{\partial \psi}{\partial \rho} \hat{e}_\theta \right) \cdot \left(\frac{dp_0}{d\rho} \hat{e}_\rho \right)$$

Normalization:

$$\frac{a n_{00} \tilde{n}_0 m_p}{\tau_A^2} \frac{\partial \tilde{v}_{\parallel, th}}{\partial \tilde{t}} = -\frac{a^2 B_{00}}{\sqrt{g}} \frac{p_{00}}{a^2 R_0 B_{00} \tilde{B}_0} \left(\frac{\partial}{\partial \zeta} + \tau \frac{\partial}{\partial \theta} \right) \tilde{p} - \frac{a^2 B_{00}}{R_0 a^2 B_{00} \tilde{B}_0} \frac{p_{00}}{\sqrt{g}} \frac{1}{\tilde{\rho}} \frac{\partial \tilde{\psi}}{\partial \theta} \frac{dp_0}{d\rho}$$

$$\Rightarrow \frac{\partial \tilde{v}_{\parallel, th}}{\partial \tilde{t}} = -\frac{\beta_0}{2} \frac{\tilde{B}_0}{\tilde{n}_0 (\tilde{J} - \tilde{\mathcal{I}})} \left[\left(\frac{\partial}{\partial \zeta} + \tau \frac{\partial}{\partial \theta} \right) \tilde{p} + \frac{1}{\tilde{\rho}} \frac{\partial \tilde{\psi}}{\partial \theta} \frac{dp_0}{d\rho} \right]$$

New terms in the pressure equation, normalized:

$$\frac{\partial \tilde{p}}{\partial \tilde{t}} = -\frac{\Gamma \tilde{p}_0 \tilde{B}_0}{\tilde{J} - \tilde{\mathcal{I}}} \left(\frac{\partial}{\partial \theta} + \frac{\partial}{\partial \zeta} \right) v_{\parallel th} - \frac{\Gamma \tilde{p}_0 \tilde{B}_0 \sqrt{g}}{\tilde{J} - \tilde{\mathcal{I}}} \left[\tau \frac{\partial}{\partial \theta} \left(\frac{1}{\sqrt{g}} \right) + \frac{\partial}{\partial \zeta} \left(\frac{1}{\sqrt{g}} \right) \right] v_{\parallel th}$$

Implementation on code:

New parallel thermal velocity equation terms

! Perturbed pressure gradient term

do l=1,leqmax

$$sceq1(:,l)=bet0*r*bmod(:,l)/(2.*denseq*(feq-qqinv*cureq))$$

end do

$$sceq1 = \frac{\beta_0}{2} \frac{\rho \tilde{B}_0}{\tilde{n}_0 (\tilde{J} - \tilde{\tau})}$$

call block(sceq1,1,7,3,1,0,0,1.0_IDP)

$$Svp_1 = \frac{\beta_0}{2} \frac{\tau \tilde{B}_0}{\tilde{n}_0 (\tilde{J} - \tilde{\tau})} \frac{\partial p}{\partial \theta}$$

do l=1,leqmax

$$sceq1(:,l)=-bet0 *bmod(:,l)/(2.*denseq*(feq-qqinv*cureq))$$

end do

$$sceq1 = -\frac{\beta_0}{2} \frac{\tilde{B}_0}{\tilde{n}_0 (\tilde{J} - \tilde{\tau})}$$

call block(sceq1,1,7,3,0,0,1,1.0_IDP)

$$Svp_1 = \frac{\beta_0}{2} \frac{\tilde{B}_0}{\tilde{n}_0 (\tilde{J} - \tilde{\tau})} \frac{\partial p}{\partial \zeta}$$

! Perturbed magnetic field term

call dbydr0(sd1,preq,0.0_IDP,1.0_IDP,0)

$$sd1 = \frac{dp_0}{d\rho}$$

do l=1,leqmax

$$sceq2(:,l)=-sceq1(:,l)*sd1(:)$$

end do

$$sceq2 = -\frac{dp_0}{d\rho} \frac{\beta_0}{2} \frac{\tilde{B}_0}{\tilde{n}_0 (\tilde{J} - \tilde{\tau})}$$

call blockj(sceq2,1,7,1,1,0,0,1.0_IDP)

$$Svp_2 = -\frac{dp_0}{d\rho} \frac{\beta_0}{2} \frac{\tilde{B}_0}{\tilde{n}_0(\tilde{J}-\tilde{\mathcal{I}})} \frac{1}{\rho} \frac{\partial \psi}{\partial \theta}$$

New terms in the pressure equation:

do l=1,leqmax

$$sceq1(:,l)=preq*r*qqinv*bmod(:,l)/(feq(:)-qqinv(:)*cureq(:))$$

end do

$$sceq1 = \frac{p_0 \rho \mathcal{B}_0}{J - \mathcal{I}}$$

call blockj(sceq1,1,3,7,1,0,0,-gamma)

$$Sp_1 = -\frac{\Gamma \mathcal{P}_0 B_0}{J - \mathcal{I}} \frac{\partial v_{\parallel th}}{\partial \theta}$$

do l=1,leqmax

$$sceq1(:,l)=preq*bmod(:,l)/(feq(:)-qqinv(:)*cureq(:))$$

end do

$$sceq1 = \frac{p_0 B_0}{J - \mathcal{I}}$$

call blockj(sceq1,1,3,7,0,0,1,-gamma)

$$Sp_2 = -\frac{\Gamma p_0 B_0}{J - \mathcal{I}} \frac{\partial v_{\parallel th}}{\partial \zeta}$$

do l=1,leqmax

$$sceq1(:,l)=preq*r*qqinv*sqgibmodith(:,l)/(feq-qqinv*cureq)$$

end do

$$sceq1 = \frac{p_0 \rho \mathcal{B}_0 \sqrt{g}}{J - \mathcal{I}} \frac{\partial}{\partial \theta} \left(\frac{1}{\sqrt{g}} \right)$$

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call blockj(sceq1,1,3,7,0,0,0,-gamma)
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$$Sp_3 = -\frac{\Gamma p_0 \rho \varpi B_0 \sqrt{g}}{J - \varpi I} \frac{\partial}{\partial \theta} \left(\frac{1}{\sqrt{g}} \right) v_{\parallel, th}$$

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do l=1,leqmax
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    sceq1(:,l)=preq*sqgibmodizt(:,l)/(feq-qqinv*cureq)
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end do
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$$sceq1 = \frac{p_0 B_0 \sqrt{g}}{J - \varpi I} \frac{\partial}{\partial \zeta} \left(\frac{1}{\sqrt{g}} \right)$$

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call blockj(sceq1,1,3,7,0,0,0,-gamma)
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$$Sp_4 = -\frac{\Gamma p_0 B_0 \sqrt{g}}{J - \varpi I} \frac{\partial}{\partial \zeta} \left(\frac{1}{\sqrt{g}} \right) v_{\parallel, th}$$