#### **Solutions Manual for**

## Fluid Mechanics: Fundamentals and Applications Third Edition

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# Chapter 12 COMPRESSIBLE FLOW

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#### **Stagnation Properties**

#### 12-1C

**Solution** We are to discuss the temperature change from an airplane's nose to far away from the aircraft.

Analysis The temperature of the air rises as it approaches the nose because of the stagnation process.

**Discussion** In the frame of reference moving with the aircraft, the air decelerates from high speed to zero at the nose (stagnation point), and this causes the air temperature to rise.

#### 12-2C

**Solution** We are to define dynamic temperature.

Analysis Dynamic temperature is the temperature rise of a fluid during a stagnation process.

**Discussion** When a gas decelerates from high speed to zero speed at a stagnation point, the temperature of the gas rises.

#### 12-3C

**Solution** We are to discuss the measurement of flowing air temperature with a probe – is there significant error?

**Analysis** No, there is not significant error, because the velocities encountered in air-conditioning applications are very low, and thus the static and the stagnation temperatures are practically identical.

**Discussion** If the air stream were supersonic, however, the error would indeed be significant.

#### **12-4**

**Solution** Air flows through a device. The stagnation temperature and pressure of air and its velocity are specified. The static pressure and temperature of air are to be determined.

**Assumptions** 1 The stagnation process is isentropic. 2 Air is an ideal gas.

**Properties** The properties of air at an anticipated average temperature of 600 K are  $c_p = 1.051$  kJ/kg·K and k = 1.376.

Analysis The static temperature and pressure of air are determined from

$$T = T_0 - \frac{V^2}{2c_p} = 673.2 - \frac{(570 \text{ m/s})^2}{2 \times 1.051 \text{ kJ/kg} \cdot \text{K}} \left( \frac{1 \text{ kJ/kg}}{1000 \text{ m}^2 / \text{s}^2} \right) = 518.6 \text{ K} \cong$$
**519 K**

and

$$P_2 = P_{02} \left(\frac{T_2}{T_{02}}\right)^{k/(k-1)} = (0.6 \text{ MPa}) \left(\frac{518.6 \text{ K}}{673.2 \text{ K}}\right)^{1.376/(1.376-1)} = \textbf{0.231 MPa}$$

**Discussion** Note that the stagnation properties can be significantly different than thermodynamic properties.

**Solution** Air at 320 K is flowing in a duct. The temperature that a stationary probe inserted into the duct will read is to be determined for different air velocities.

Assumptions The stagnation process is isentropic.

**Properties** The specific heat of air at room temperature is  $c_p = 1.005 \text{ kJ/kg·K}$ .

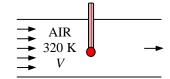
Analysis The air which strikes the probe will be brought to a complete stop, and thus it will undergo a stagnation process. The thermometer will sense the temperature of this stagnated air, which is the stagnation temperature,  $T_0$ . It is determined from  $T_0 = T + \frac{V^2}{2c_p}$ . The results for each case are calculated below:

(a) 
$$T_0 = 320 \text{ K} + \frac{(1 \text{ m/s})^2}{2 \times 1.005 \text{ kJ/kg} \cdot \text{K}} \left( \frac{1 \text{ kJ/kg}}{1000 \text{ m}^2/\text{s}^2} \right) = 320.0 \text{ K}$$

(b) 
$$T_0 = 320 \text{ K} + \frac{(10 \text{ m/s})^2}{2 \times 1.005 \text{ kJ/kg} \cdot \text{K}} \left( \frac{1 \text{ kJ/kg}}{1000 \text{ m}^2 / \text{s}^2} \right) = 320.1 \text{ K}$$

(c) 
$$T_0 = 320 \text{ K} + \frac{(100 \text{ m/s})^2}{2 \times 1.005 \text{ kJ/kg} \cdot \text{K}} \left( \frac{1 \text{ kJ/kg}}{1000 \text{ m}^2/\text{s}^2} \right) = 325.0 \text{ K}$$

(d) 
$$T_0 = 320 \text{ K} + \frac{(1000 \text{ m/s})^2}{2 \times 1.005 \text{ kJ/kg} \cdot \text{K}} \left( \frac{1 \text{ kJ/kg}}{1000 \text{ m}^2/\text{s}^2} \right) = 817.5 \text{ K}$$



**Discussion** Note that the stagnation temperature is nearly identical to the thermodynamic temperature at low velocities, but the difference between the two is significant at high velocities.

**Solution** The states of different substances and their velocities are specified. The stagnation temperature and stagnation pressures are to be determined.

Assumptions 1 The stagnation process is isentropic. 2 Helium and nitrogen are ideal gases.

**Analysis** (a) Helium can be treated as an ideal gas with  $c_p = 5.1926$  kJ/kg·K and k = 1.667. Then the stagnation temperature and pressure of helium are determined from

$$T_0 = T + \frac{V^2}{2c_p} = 50^{\circ}\text{C} + \frac{(240 \text{ m/s})^2}{2 \times 5.1926 \text{ kJ/kg} \cdot {}^{\circ}\text{C}} \left(\frac{1 \text{ kJ/kg}}{1000 \text{ m}^2/\text{s}^2}\right) = 55.5^{\circ}\text{C}$$

$$P_0 = P \left(\frac{T_0}{T}\right)^{k/(k-1)} = (0.25 \,\mathrm{MPa}) \left(\frac{328.7 \,\mathrm{K}}{323.2 \,\mathrm{K}}\right)^{1.667/(1.667-1)} = \mathbf{0.261 \,\mathrm{MPa}}$$

(b) Nitrogen can be treated as an ideal gas with  $c_p = 1.039 \text{ kJ/kg} \cdot \text{K}$  and k = 1.400. Then the stagnation temperature and pressure of nitrogen are determined from

$$T_0 = T + \frac{V^2}{2c_p} = 50^{\circ}\text{C} + \frac{(300 \text{ m/s})^2}{2 \times 1.039 \text{ kJ/kg} \cdot {}^{\circ}\text{C}} \left(\frac{1 \text{ kJ/kg}}{1000 \text{ m}^2/\text{s}^2}\right) = 93.3^{\circ}\text{C}$$

$$P_0 = P \left(\frac{T_0}{T}\right)^{k/(k-1)} = (0.15 \text{ MPa}) \left(\frac{366.5 \text{ K}}{323.2 \text{ K}}\right)^{1.4/(1.4-1)} =$$
**0.233 MPa**

(c) Steam can be treated as an ideal gas with  $c_p = 1.865 \text{ kJ/kg} \cdot \text{K}$  and k = 1.329. Then the stagnation temperature and pressure of steam are determined from

$$T_0 = T + \frac{V^2}{2c_p} = 350$$
°C +  $\frac{(480 \text{ m/s})^2}{2 \times 1.865 \text{ kJ/kg} \cdot \text{°C}} \left(\frac{1 \text{ kJ/kg}}{1000 \text{ m}^2 / \text{s}^2}\right) = 411.8$ °C = 685 K

$$P_0 = P \left(\frac{T_0}{T}\right)^{k/(k-1)} = (0.1 \,\text{MPa}) \left(\frac{685 \,\text{K}}{623.2 \,\text{K}}\right)^{1.329/(1.329-1)} =$$
**0.147 MPa**

**Discussion** Note that the stagnation properties can be significantly different than thermodynamic properties.

#### 12-7

**Solution** The state of air and its velocity are specified. The stagnation temperature and stagnation pressure of air are to be determined.

Assumptions 1 The stagnation process is isentropic. 2 Air is an ideal gas.

**Properties** The properties of air at room temperature are  $c_p = 1.005 \text{ kJ/kg} \cdot \text{K}$  and k = 1.4.

**Analysis** The stagnation temperature of air is determined from

$$T_0 = T + \frac{V^2}{2c_p} = 238 \text{ K} + \frac{(325 \text{ m/s})^2}{2 \times 1.005 \text{ kJ/kg} \cdot \text{K}} \left(\frac{1 \text{ kJ/kg}}{1000 \text{ m}^2/\text{s}^2}\right) = 290.5 \cong 291 \text{ K}$$

Other stagnation properties at the specified state are determined by considering an isentropic process between the specified state and the stagnation state,

$$P_0 = P \left(\frac{T_0}{T}\right)^{k/(k-1)} = (36 \text{ kPa}) \left(\frac{290.5 \text{ K}}{238 \text{ K}}\right)^{1.4/(1.4-1)} = 72.37 \text{ kPa} \cong 72.4 \text{ kPa}$$

**Discussion** Note that the stagnation properties can be significantly different than thermodynamic properties.

#### 12-8E

**Solution** Steam flows through a device. The stagnation temperature and pressure of steam and its velocity are specified. The static pressure and temperature of the steam are to be determined.

Assumptions 1 The stagnation process is isentropic. 2 Steam is an ideal gas.

**Properties** Steam can be treated as an ideal gas with  $c_p = 0.4455$  Btu/lbm·R and k = 1.329.

Analysis The static temperature and pressure of steam are determined from

$$T = T_0 - \frac{V^2}{2c_p} = 700^{\circ} \text{F} - \frac{(900 \text{ ft/s})^2}{2 \times 0.4455 \text{ Btu/lbm} \cdot {}^{\circ} \text{F}} \left( \frac{1 \text{ Btu/lbm}}{25,037 \text{ ft}^2 / \text{s}^2} \right) = 663.7^{\circ} \text{F}$$

$$P = P_0 \left(\frac{T}{T_0}\right)^{k/(k-1)} = (120 \text{ psia}) \left(\frac{1123.7 \text{ R}}{1160 \text{ R}}\right)^{1.329/(1.329-1)} = \textbf{105.5 psia}$$

**Discussion** Note that the stagnation properties can be significantly different than thermodynamic properties.

#### 12-9

**Solution** The inlet stagnation temperature and pressure and the exit stagnation pressure of air flowing through a compressor are specified. The power input to the compressor is to be determined.

Assumptions 1 The compressor is isentropic. 2 Air is an ideal gas.

**Properties** The properties of air at room temperature are  $c_p = 1.005 \text{ kJ/kg} \cdot \text{K}$  and k = 1.4.

*Analysis* The exit stagnation temperature of air  $T_{02}$  is determined from

$$T_{02} = T_{01} \left(\frac{P_{02}}{P_{01}}\right)^{(k-1)/k} = (308.2 \text{ K}) \left(\frac{900}{100}\right)^{(1.4-1)/1.4} = 577.4 \text{ K}$$

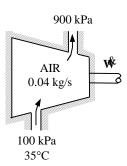
From the energy balance on the compressor,

$$W_{\text{in}}^{\text{X}} = n (h_{20} - h_{01})$$

or,

$$W_{\rm in}^{\rm SC} = n {\rm SC}_p (T_{02} - T_{01}) = (0.04 \, {\rm kg/s})(1.005 \, {\rm kJ/kg \cdot K})(577.4 - 308.2) {\rm K} = {\rm 10.8 \, kW}$$

**Discussion** Note that the stagnation properties can be used conveniently in the energy equation.



**Solution** The inlet stagnation temperature and pressure and the exit stagnation pressure of products of combustion flowing through a gas turbine are specified. The power output of the turbine is to be determined.

Assumptions 1 The expansion process is isentropic. 2 Products of combustion are ideal gases.

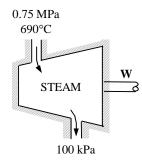
**Properties** The properties of products of combustion are  $c_p = 1.157 \text{ kJ/kg·K}$ , R = 0.287 kJ/kg·K, and k = 1.33.

**Analysis** The exit stagnation temperature  $T_{02}$  is determined to be

$$T_{02} = T_{01} \left(\frac{P_{02}}{P_{01}}\right)^{(k-1)/k} = (963.2 \text{ K}) \left(\frac{0.1}{0.75}\right)^{(1.33-1)/1.33} = 584.2 \text{ K}$$

Also,

$$c_p = kc_v = k(c_p - R) \longrightarrow c_p = \frac{kR}{k - 1}$$
$$= \frac{1.33(0.287 \text{ kJ/kg} \cdot \text{K})}{1.33 - 1}$$
$$= 1.157 \text{ kJ/kg} \cdot \text{K}$$



From the energy balance on the turbine,

$$-w_{\text{out}} = (h_{20} - h_{01})$$

or, 
$$w_{\text{out}} = c_p (T_{01} - T_{02}) = (1.157 \text{ kJ/kg} \cdot \text{K})(963.2 - 584.2)\text{K} = 438.5 \text{ kJ/kg} \cong 439 \text{ kJ/kg}$$

**Discussion** Note that the stagnation properties can be used conveniently in the energy equation.

#### **One Dimensional Isentropic Flow**

#### 12-11C

**Solution** We are to determine if it is possible to accelerate a gas to supersonic velocity in a converging nozzle.

**Analysis** No, it is not possible.

**Discussion** The only way to do it is to have first a converging nozzle, and then a diverging nozzle.

#### 12-12C

**Solution** We are to discuss what happens to several variables when a subsonic gas enters a diverging duct.

Analysis (a) The velocity decreases. (b), (c), (d) The temperature, pressure, and density of the fluid increase.

**Discussion** The velocity decrease is opposite to what happens in supersonic flow.

#### 12-13C

**Solution** We are to discuss the pressure at the throats of two different converging-diverging nozzles.

Analysis The pressures at the two throats are identical.

**Discussion** Since the gas has the same stagnation conditions, it also has the same sonic conditions at the throat.

#### 12-14C

**Solution** We are to discuss what happens to several variables when a supersonic gas enters a converging duct.

Analysis (a) The velocity decreases. (b), (c), (d) The temperature, pressure, and density of the fluid increase.

**Discussion** The velocity decrease is opposite to what happens in subsonic flow.

#### 12-15C

**Solution** We are to discuss what happens to several variables when a supersonic gas enters a diverging duct.

Analysis (a) The velocity increases. (b), (c), (d) The temperature, pressure, and density of the fluid decrease.

**Discussion** The velocity increase is opposite to what happens in subsonic flow.

#### 12-16C

**Solution** We are to discuss what happens to the exit velocity and mass flow rate through a converging nozzle at sonic exit conditions when the nozzle exit area is reduced.

Analysis (a) The exit velocity remains constant at sonic speed, (b) the mass flow rate through the nozzle decreases because of the reduced flow area.

**Discussion** Without a diverging portion of the nozzle, a converging nozzle is limited to sonic velocity at the exit.

#### 12-17C

**Solution** We are to discuss what happens to several variables when a subsonic gas enters a converging duct.

Analysis (a) The velocity increases. (b), (c), (d) The temperature, pressure, and density of the fluid decrease.

**Discussion** The velocity increase is opposite to what happens in supersonic flow.

Helium -

#### 12-18

**Solution** Helium enters a converging-diverging nozzle at specified conditions. The lowest temperature and pressure that can be obtained at the throat of the nozzle are to be determined.

**Assumptions** 1 Helium is an ideal gas with constant specific heats. 2 Flow through the nozzle is steady, one-dimensional, and isentropic.

**Properties** The properties of helium are k = 1.667 and  $c_p = 5.1926$  kJ/kg·K.

**Analysis** The lowest temperature and pressure that can be obtained at the throat are the critical temperature  $T^*$  and critical pressure  $P^*$ . First we determine the stagnation temperature  $T_0$  and stagnation pressure  $P_0$ ,

$$T_0 = T + \frac{V^2}{2c_p} = 800 \text{ K} + \frac{(100 \text{ m/s})^2}{2 \times 5.1926 \text{ kJ/kg} \cdot ^{\circ}\text{C}} \left(\frac{1 \text{ kJ/kg}}{1000 \text{ m}^2 / \text{s}^2}\right) = 801 \text{ K}$$

$$P_0 = P\left(\frac{T_0}{T}\right)^{k/(k-1)} = (0.7 \text{ MPa}) \left(\frac{801 \text{ K}}{800 \text{ K}}\right)^{1.667/(1.667-1)} = 0.702 \text{ MPa}$$

Thus,

$$T^* = T_0 \left( \frac{2}{k+1} \right) = (801 \text{ K}) \left( \frac{2}{1.667 + 1} \right) =$$
**601 K**

and

$$P^* = P_0 \left(\frac{2}{k+1}\right)^{k/(k-1)} = (0.702 \text{ MPa}) \left(\frac{2}{1.667+1}\right)^{1.667/(1.667-1)} = \textbf{0.342 MPa}$$

**Discussion** These are the temperature and pressure that will occur at the throat when the flow past the throat is supersonic.

#### 12-19

**Solution** The speed of an airplane and the air temperature are give. It is to be determined if the speed of this airplane is subsonic or supersonic.

**Assumptions** Air is an ideal gas with constant specific heats at room temperature.

**Properties** The gas constant of air is  $R = 0.287 \text{ kJ/kg} \cdot \text{K}$ . Its specific heat ratio at room temperature is k = 1.4.

Analysis The temperature is -50 + 273.15 = 223.15 K. The speed of sound is

$$c = \sqrt{kRT} = \sqrt{(1.4)(0.287 \text{ kJ/kg} \cdot \text{K})(223.15 \text{ K}) \left(\frac{1000 \text{ m}^2/\text{s}^2}{1 \text{ kJ/kg}}\right) \left(\frac{3.6 \text{ km/h}}{1 \text{ m/s}}\right)} = 1077.97 \text{ km/h}$$

and

$$Ma = \frac{V}{c} = \frac{1050 \text{ km/h}}{1077.97 \text{ km/h}} = 0.9741 \text{ km/h} \cong \mathbf{0.974}$$

The speed of the airplane is **subsonic** since the Mach number is less than 1.

**Discussion** Subsonic airplanes stay sufficiently far from the Mach number of 1 to avoid the instabilities associated with transonic flights.

**Solution** The critical temperature, pressure, and density of air and helium are to be determined at specified conditions.

**Assumptions** Air and Helium are ideal gases with constant specific heats at room temperature.

**Properties** The properties of air at room temperature are R = 0.287 kJ/kg·K, k = 1.4, and  $c_p = 1.005$  kJ/kg·K. The properties of helium at room temperature are R = 2.0769 kJ/kg·K, k = 1.667, and  $c_p = 5.1926$  kJ/kg·K.

**Analysis** (a) Before we calculate the critical temperature  $T^*$ , pressure  $P^*$ , and density  $\rho^*$ , we need to determine the stagnation temperature  $T_0$ , pressure  $P_0$ , and density  $\rho_0$ .

$$T_0 = 100^{\circ}\text{C} + \frac{V^2}{2c_p} = 100 + \frac{(250 \text{ m/s})^2}{2 \times 1.005 \text{ kJ/kg} \cdot {}^{\circ}\text{C}} \left(\frac{1 \text{ kJ/kg}}{1000 \text{ m}^2 / \text{s}^2}\right) = 131.1^{\circ}\text{C}$$

$$P_0 = P\left(\frac{T_0}{T}\right)^{k/(k-1)} = (200 \text{ kPa}) \left(\frac{404.3 \text{ K}}{373.2 \text{ K}}\right)^{1.4/(1.4-1)} = 264.7 \text{ kPa}$$

$$\rho_0 = \frac{P_0}{RT_0} = \frac{264.7 \text{ kPa}}{(0.287 \text{ kPa} \cdot \text{m}^3/\text{kg} \cdot \text{K})(404.3 \text{ K})} = 2.281 \text{ kg/m}^3$$

Thus,

$$T^* = T_0 \left(\frac{2}{k+1}\right) = (404.3 \text{ K}) \left(\frac{2}{1.4+1}\right) = \mathbf{337 \text{ K}}$$

$$P^* = P_0 \left(\frac{2}{k+1}\right)^{k/(k-1)} = (264.7 \text{ kPa}) \left(\frac{2}{1.4+1}\right)^{1.4/(1.4-1)} = \mathbf{140 \text{ kPa}}$$

$$\rho^* = \rho_0 \left(\frac{2}{k+1}\right)^{1/(k-1)} = (2.281 \text{ kg/m}^3) \left(\frac{2}{1.4+1}\right)^{1/(1.4-1)} = \mathbf{1.45 \text{ kg/m}}^3$$

(b) For helium, 
$$T_0 = T + \frac{V^2}{2c_p} = 40 + \frac{(300 \text{ m/s})^2}{2 \times 5.1926 \text{ kJ/kg} \cdot ^{\circ}\text{C}} \left(\frac{1 \text{ kJ/kg}}{1000 \text{ m}^2 / \text{s}^2}\right) = 48.7 ^{\circ}\text{C}$$

$$P_0 = P \left(\frac{T_0}{T}\right)^{k/(k-1)} = (200 \text{ kPa}) \left(\frac{321.9 \text{ K}}{313.2 \text{ K}}\right)^{1.667/(1.667-1)} = 214.2 \text{ kPa}$$

$$\rho_0 = \frac{P_0}{RT_0} = \frac{214.2 \text{ kPa}}{(2.0769 \text{ kPa} \cdot \text{m}^3/\text{kg} \cdot \text{K})(321.9 \text{ K})} = 0.320 \text{ kg/m}^3$$

Thus,

$$\begin{split} T^* &= T_0 \bigg(\frac{2}{k+1}\bigg) = (321.9 \text{ K}) \bigg(\frac{2}{1.667+1}\bigg) = \textbf{241 K} \\ P^* &= P_0 \bigg(\frac{2}{k+1}\bigg)^{k/(k-1)} = (200 \text{ kPa}) \bigg(\frac{2}{1.667+1}\bigg)^{1.667/(1.667-1)} = \textbf{97.4 kPa} \\ \rho^* &= \rho_0 \bigg(\frac{2}{k+1}\bigg)^{1/(k-1)} = (0.320 \text{ kg/m}^3) \bigg(\frac{2}{1.667+1}\bigg)^{1/(1.667-1)} = \textbf{0.208 kg/m}^3 \end{split}$$

**Discussion** These are the temperature, pressure, and density values that will occur at the throat when the flow past the throat is supersonic.

#### 12-21E

**Solution** Air flows through a duct at a specified state and Mach number. The velocity and the stagnation pressure, temperature, and density of the air are to be determined.

Assumptions Air is an ideal gas with constant specific heats at room temperature.

**Properties** The properties of air are R = 0.06855 Btu/lbm·R = 0.3704 psia·ft<sup>3</sup>/lbm·R and k = 1.4.

Analysis First, T = 320 + 459.67 = 779.67 K. The speed of sound in air at the specified conditions is

$$c = \sqrt{kRT} = \sqrt{(1.4)(0.06855 \text{ Btu/1bm} \cdot \text{R})(779.67 \text{ R}) \left(\frac{25,037 \text{ ft}^2/\text{s}^2}{1 \text{ Btu/1bm}}\right)} = 1368.72 \text{ ft/s}$$

Thus,

$$V = \text{Ma} \times c = (0.7)(1368.72 \text{ ft/s}) = 958.10 \cong 958 \text{ ft/s}$$

Also,

$$\rho = \frac{P}{RT} = \frac{25 \text{ psia}}{(0.3704 \text{ psia} \cdot \text{ft}^3/\text{lbm} \cdot \text{R})(779.67 \text{ R})} = 0.086568 \text{ 1bm/ft}^3$$

Then the stagnation properties are determined from

$$T_0 = T \left( 1 + \frac{(k-1)\text{Ma}^2}{2} \right) = (779.67 \text{ R}) \left( 1 + \frac{(1.4-1)(0.7)^2}{2} \right) = 856.08 \text{ R} \approx 856 \text{ R}$$

$$P_0 = P \left(\frac{T_0}{T}\right)^{k/(k-1)} = (25 \text{ psia}) \left(\frac{856.08 \text{ R}}{779.67 \text{ R}}\right)^{1.4/(1.4-1)} = 34.678 \text{ psia} \cong 34.7 \text{ psia}$$

$$\rho_0 = \rho \left(\frac{T_0}{T}\right)^{1/(k-1)} = (0.08656 \text{ 1bm/ft}^3) \left(\frac{856.08 \text{ R}}{779.67 \text{ R}}\right)^{1/(1.4-1)} = 0.10936 \text{ lbm/ft}^3 \cong \textbf{0.109 lbm/ft}^3$$

**Discussion** Note that the temperature, pressure, and density of a gas increases during a stagnation process.

#### 12-22

**Solution** Air enters a converging-diverging nozzle at specified conditions. The lowest pressure that can be obtained at the throat of the nozzle is to be determined.

Assumptions 1 Air is an ideal gas with constant specific heats at room temperature. 2 Flow through the nozzle is steady, one-dimensional, and isentropic.

**Properties** The specific heat ratio of air at room temperature is k = 1.4.

**Analysis** The lowest pressure that can be obtained at the throat is the critical pressure  $P^*$ , which is determined from

$$P^* = P_0 \left(\frac{2}{k+1}\right)^{k/(k-1)} = (1200 \text{ kPa}) \left(\frac{2}{1.4+1}\right)^{1.4/(1.4-1)} =$$
**634 kPa**

**Discussion** This is the pressure that occurs at the throat when the flow past the throat is supersonic.

**Solution** The Mach number of scramjet and the air temperature are given. The speed of the engine is to be determined.

**Assumptions** Air is an ideal gas with constant specific heats at room temperature.

**Properties** The gas constant of air is  $R = 0.287 \text{ kJ/kg} \cdot \text{K}$ . Its specific heat ratio at room temperature is k = 1.4.

Analysis The temperature is -20 + 273.15 = 253.15 K. The speed of sound is

$$c = \sqrt{kRT} = \sqrt{(1.4)(0.287 \text{ kJ/kg} \cdot \text{K})(253.15 \text{ K}) \left(\frac{1000 \text{ m}^2/\text{s}^2}{1 \text{ kJ/kg}}\right)} = 318.93 \text{ m/s}$$

and

$$V = c$$
Ma =  $(318.93 \text{ m/s})(7) \left( \frac{3.6 \text{ km/h}}{1 \text{ m/s}} \right) = 8037 \text{ km/h} \approx 8040 \text{ km/h}$ 

**Discussion** Note that extremely high speed can be achieved with scramjet engines. We cannot justify more than three significant digits in a problem like this.

#### 12-24E

**Solution** The Mach number of scramjet and the air temperature are given. The speed of the engine is to be determined.

Assumptions Air is an ideal gas with constant specific heats at room temperature.

**Properties** The gas constant of air is R = 0.06855 Btu/lbm·R. Its specific heat ratio at room temperature is k = 1.4.

Analysis The temperature is 0 + 459.67 = 459.67 R. The speed of sound is

$$c = \sqrt{kRT} = \sqrt{(1.4)(0.06855 \text{ Btu/lbm} \cdot \text{R})(459.67 \text{ R}) \left(\frac{25,037 \text{ ft}^2/\text{s}^2}{1 \text{ Btu/lbm}}\right)} = 1050.95 \text{ ft/s}$$

and

$$V = c$$
Ma =  $(1050.95 \text{ ft/s})(7) \left( \frac{1 \text{ mi/h}}{1.46667 \text{ ft/s}} \right) = 5015.9 \text{ mi/h} \approx 5020 \text{ mi/h}$ 

**Discussion** Note that extremely high speed can be achieved with scramjet engines. We cannot justify more than three significant digits in a problem like this.

AIR

#### 12-25

**Solution** Air flows through a duct. The state of the air and its Mach number are specified. The velocity and the stagnation pressure, temperature, and density of the air are to be determined.

Assumptions Air is an ideal gas with constant specific heats at room temperature.

**Properties** The properties of air at room temperature are  $R = 0.287 \text{ kPa.m}^3/\text{kg.K}$  and k = 1.4.

Analysis The speed of sound in air at the specified conditions is

$$c = \sqrt{kRT} = \sqrt{(1.4)(0.287 \text{ kJ/kg} \cdot \text{K})(373.2 \text{ K}) \left(\frac{1000 \text{ m}^2/\text{s}^2}{1 \text{ kJ/kg}}\right)} = 387.2 \text{ m/s}$$

Thus,

$$V = \text{Ma} \times c = (0.8)(387.2 \text{ m/s}) =$$
**310 m/s**

Also,

$$\rho = \frac{P}{RT} = \frac{200 \text{ kPa}}{(0.287 \text{ kPa} \cdot \text{m}^3/\text{kg} \cdot \text{K})(373.2 \text{ K})} = 1.867 \text{ kg/m}^3$$

Then the stagnation properties are determined from

$$T_0 = T \left( 1 + \frac{(k-1)\text{Ma}^2}{2} \right) = (373.2 \text{ K}) \left( 1 + \frac{(1.4-1)(0.8)^2}{2} \right) = 421 \text{ K}$$

$$P_0 = P \left(\frac{T_0}{T}\right)^{k/(k-1)} = (200 \text{ kPa}) \left(\frac{421.0 \text{ K}}{373.2 \text{ K}}\right)^{1.4/(1.4-1)} = 305 \text{ kPa}$$

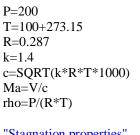
$$\rho_0 = \rho \left(\frac{T_0}{T}\right)^{1/(k-1)} = (1.867 \text{ kg/m}^3) \left(\frac{421.0 \text{ K}}{373.2 \text{ K}}\right)^{1/(1.4-1)} = \textbf{2.52 kg/m}^3$$

**Discussion** Note that both the pressure and temperature drop as the gas is accelerated as part of the internal energy of the gas is converted to kinetic energy.

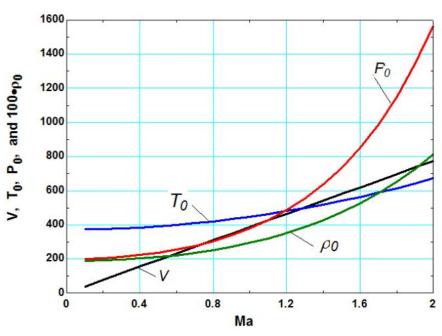


**Solution** Problem 12-41 is reconsidered. The effect of Mach number on the velocity and stagnation properties as the Ma is varied from 0.1 to 2 are to be investigated, and the results are to be plotted.

Analysis The EES Equations window is printed below, along with the tabulated and plotted results.



"Stagnation properties" T0=T\*(1+(k-1)\*Ma^2/2) P0=P\*(T0/T)^(k/(k-1)) rho0=rho\*(T0/T)^(1/(k-1))



Mach num.	Velocity,	Stag. Temp,	Stag. Press,	Stag. Density,
Ma	V, m/s	$T_0$ , K	$P_0$ , kPa	$\rho_0$ , kg/m <sup>3</sup>
0.1	38.7	373.9	201.4	1.877
0.2	77.4	376.1	205.7	1.905
0.3	116.2	379.9	212.9	1.953
0.4	154.9	385.1	223.3	2.021
0.5	193.6	391.8	237.2	2.110
0.6	232.3	400.0	255.1	2.222
0.7	271.0	409.7	277.4	2.359
0.8	309.8	420.9	304.9	2.524
0.9	348.5	433.6	338.3	2.718
1.0	387.2	447.8	378.6	2.946
1.1	425.9	463.5	427.0	3.210
1.2	464.7	480.6	485.0	3.516
1.3	503.4	499.3	554.1	3.867
1.4	542.1	519.4	636.5	4.269
1.5	580.8	541.1	734.2	4.728
1.6	619.5	564.2	850.1	5.250
1.7	658.3	588.8	987.2	5.842
1.8	697.0	615.0	1149.2	6.511
1.9	735.7	642.6	1340.1	7.267
2.0	774.4	671.7	1564.9	8.118

**Discussion** Note that as Mach number increases, so does the flow velocity and stagnation temperature, pressure, and density.

**Solution** An aircraft is designed to cruise at a given Mach number, elevation, and the atmospheric temperature. The stagnation temperature on the leading edge of the wing is to be determined.

Assumptions Air is an ideal gas with constant specific heats at room temperature.

**Properties** The properties of air are  $R = 0.287 \text{ kPa.m}^3/\text{kg.K}$ ,  $c_p = 1.005 \text{ kJ/kg·K}$ , and k = 1.4.

Analysis The speed of sound in air at the specified conditions is

$$c = \sqrt{kRT} = \sqrt{(1.4)(0.287 \text{ kJ/kg} \cdot \text{K})(236.15 \text{ K}) \left(\frac{1000 \text{ m}^2/\text{s}^2}{1 \text{ kJ/kg}}\right)} = 308.0 \text{ m/s}$$

Thus,

$$V = \text{Ma} \times c = (1.1)(308.0 \text{ m/s}) = 338.8 \text{ m/s}$$

Then,

$$T_0 = T + \frac{V^2}{2c_p} = 236.15 + \frac{(338.8 \,\mathrm{m/s})^2}{2 \times 1.005 \,\mathrm{kJ/kg \cdot K}} \left(\frac{1 \,\mathrm{kJ/kg}}{1000 \,\mathrm{m^2/s^2}}\right) = \mathbf{293 \, K}$$

**Discussion** Note that the temperature of a gas increases during a stagnation process as the kinetic energy is converted to enthalpy.

#### 12-28

**Solution** Quiescent carbon dioxide at a given state is accelerated isentropically to a specified Mach number. The temperature and pressure of the carbon dioxide after acceleration are to be determined.

**Assumptions** Carbon dioxide is an ideal gas with constant specific heats at room temperature.

**Properties** The specific heat ratio of the carbon dioxide at room temperature is k = 1.288.

Analysis The inlet temperature and pressure in this case is equivalent to the stagnation temperature and pressure since the inlet velocity of the carbon dioxide is said to be negligible. That is,  $T_0 = T_i = 400 \text{ K}$  and  $P_0 = P_i = 1200 \text{ kPa}$ . Then,

$$T = T_0 \left( \frac{2}{2 + (k - 1)\text{Ma}^2} \right) = (600 \text{ K}) \left( \frac{2}{2 + (1.288 - 1)(0.6)^2} \right) = 570.43 \text{ K} \approx 570 \text{ K}$$

and

$$P = P_0 \left(\frac{T}{T_0}\right)^{k/(k-1)} = (1200 \text{ kPa}) \left(\frac{570.43 \text{ K}}{600 \text{ K}}\right)^{1.288/(1.288-1)} = 957.23 \text{ K} \cong 957 \text{ kPa}$$

**Discussion** Note that both the pressure and temperature drop as the gas is accelerated as part of the internal energy of the gas is converted to kinetic energy.

#### **Isentropic Flow Through Nozzles**

#### 12-29C

**Solution** We are to analyze if it is possible to accelerate a fluid to supersonic speeds with a velocity that is not sonic at the throat.

**Analysis** No, if the flow in the throat is subsonic. If the velocity at the throat is subsonic, the diverging section would act like a diffuser and decelerate the flow. Yes, if the flow in the throat is already supersonic, the diverging section would accelerate the flow to even higher Mach number.

**Discussion** In duct flow, the latter situation is not possible unless a second converging-diverging portion of the duct is located upstream, and there is sufficient pressure difference to choke the flow in the upstream throat.

#### 12-30C

**Solution** We are to discuss what would happen if we add a diverging section to supersonic flow in a duct.

**Analysis** The fluid would **accelerate even further**, as desired.

**Discussion** This is the opposite of what would happen in subsonic flow.

#### 12-31C

**Solution** We are to discuss the difference between Ma\* and Ma.

Analysis Ma\* is the local velocity non-dimensionalized with respect to the sonic speed at the throat, whereas Ma is the local velocity non-dimensionalized with respect to the local sonic speed.

**Discussion** The two are identical at the throat when the flow is choked.

#### 12-32C

**Solution** We are to consider subsonic flow through a converging nozzle with critical pressure at the exit, and analyze the effect of lowering back pressure below the critical pressure.

Analysis (a) No effect on velocity. (b) No effect on pressure. (c) No effect on mass flow rate.

**Discussion** In this situation, the flow is already choked initially, so further lowering of the back pressure does not change anything upstream of the nozzle exit plane.

#### 12-33C

**Solution** We are to compare the mass flow rates through two identical converging nozzles, but with one having a diverging section.

Analysis If the back pressure is low enough so that sonic conditions exist at the throats, the mass flow rates in the two nozzles would be identical. However, if the flow is not sonic at the throat, the mass flow rate through the nozzle with the diverging section would be greater, because it acts like a subsonic diffuser.

**Discussion** Once the flow is choked at the throat, whatever happens downstream is irrelevant to the flow upstream of the throat.

#### 12-34C

**Solution** We are to discuss the hypothetical situation of hypersonic flow at the outlet of a converging nozzle.

Analysis Maximum flow rate through a converging nozzle is achieved when Ma = 1 at the exit of a nozzle. For all other Ma values the mass flow rate decreases. Therefore, the mass flow rate would decrease if hypersonic velocities were achieved at the throat of a converging nozzle.

**Discussion** Note that this is not possible unless the flow upstream of the converging nozzle is already hypersonic.

#### 12-35C

**Solution** We are to consider subsonic flow through a converging nozzle, and analyze the effect of setting back pressure to critical pressure for a converging nozzle.

Analysis (a) The exit velocity reaches the sonic speed, (b) the exit pressure equals the critical pressure, and (c) the mass flow rate reaches the maximum value.

**Discussion** In such a case, we say that the flow is *choked*.

#### 12-36C

**Solution** We are to discuss what happens to several variables in the diverging section of a subsonic converging-diverging nozzle.

Analysis (a) The velocity decreases, (b) the pressure increases, and (c) the mass flow rate remains the same.

**Discussion** Qualitatively, this is the same as what we are used to (in previous chapters) for incompressible flow.

#### 12-37C

**Solution** We are to discuss what would happen if we add a diverging section to supersonic flow in a duct.

**Analysis** The fluid would **accelerate even further** instead of decelerating.

**Discussion** This is the opposite of what would happen in subsonic flow.

**Solution** Nitrogen enters a converging-diverging nozzle at a given pressure. The critical velocity, pressure, temperature, and density in the nozzle are to be determined.

Assumptions 1 Nitrogen is an ideal gas. 2 Flow through the nozzle is steady, one-dimensional, and isentropic.

**Properties** The properties of nitrogen are k = 1.4 and R = 0.2968 kJ/kg·K.

**Analysis** The stagnation pressure in this case are identical to the inlet properties since the inlet velocity is negligible. They remain constant throughout the nozzle,

$$P_{0} = P_{i} = 700 \text{ kPa}$$

$$T_{0} = T_{i} = 400 \text{ K}$$

$$\rho_{0} = \frac{P_{0}}{RT_{0}} = \frac{700 \text{ kPa}}{(0.2968 \text{ kPa} \cdot \text{m}^{3}/\text{kg} \cdot \text{K})(400 \text{ K})} = 5.896 \text{ kg/m}^{3}$$
\*

Critical properties are those at a location where the Mach number is Ma = 1. From Table A-13 at Ma = 1, we read  $T/T_0$  = 0.8333,  $P/P_0$  = 0.5283, and  $\rho/\rho_0$  = 0.6339. Then the critical properties become

$$T^* = 0.8333T_0 = 0.8333(400 \text{ K}) =$$
 **333 K**  
 $P^* = 0.5283P_0 = 0.5283(700 \text{ kPa}) =$  **370 MPa**  
 $\rho^* = 0.6339\rho_0 = 0.6339(5.896 \text{ kg/m}^3) =$  **3.74 kg/m**<sup>3</sup>

Also,

$$V^* = c^* = \sqrt{kRT^*} = \sqrt{(1.4)(0.2968 \text{ kJ/kg} \cdot \text{K})(333 \text{ K}) \left(\frac{1000 \text{ m}^2/\text{s}^2}{1 \text{ kJ/kg}}\right)} = 372 \text{ m/s}$$

**Discussion** We can also solve this problem using the relations for compressible isentropic flow. The results would be identical.

#### 12-39

**Solution** For an ideal gas, an expression is to be obtained for the ratio of the speed of sound where Ma = 1 to the speed of sound based on the stagnation temperature,  $c^*/c_0$ .

**Analysis** For an ideal gas the speed of sound is expressed as  $c = \sqrt{kRT}$ . Thus,

$$\frac{c^*}{c_0} = \frac{\sqrt{kRT^*}}{\sqrt{kRT_0}} = \left(\frac{T^*}{T_0}\right)^{1/2} = \left(\frac{\mathbf{2}}{k+1}\right)^{1/2}$$

**Discussion** Note that a speed of sound changes the flow as the temperature changes.

**Solution** Air enters a converging-diverging nozzle at a specified pressure. The back pressure that will result in a specified exit Mach number is to be determined.

**Assumptions** 1 Air is an ideal gas with constant specific heats. 2 Flow through the nozzle is steady, one-dimensional, and isentropic.

**Properties** The specific heat ratio of air is k = 1.4.

**Analysis** The stagnation pressure in this case is identical to the inlet pressure since the inlet velocity is negligible. It remains constant throughout the nozzle since the flow is isentropic,

$$i$$
 AIR  $\longrightarrow$   $e$   $Ma_e = 1.8$ 

$$P_0 = P_i = 1.2 \text{ MPa}$$

From Table A-13 at Ma<sub>e</sub> = 1.8, we read  $P_e/P_0 = 0.1740$ .

Thus,  $P = 0.1740P_0 = 0.1740(1.2 \text{ MPa}) = 0.209 \text{ MPa} = 209 \text{ kPa}$ 

**Discussion** If we solve this problem using the relations for compressible isentropic flow, the results would be identical.

#### 12-41E

**Solution** Air enters a nozzle at specified temperature, pressure, and velocity. The exit pressure, exit temperature, and exit-to-inlet area ratio are to be determined for a Mach number of Ma = 1 at the exit.

Assumptions 1 Air is an ideal gas with constant specific heats at room temperature. 2 Flow through the nozzle is steady, one-dimensional, and isentropic.

**Properties** The properties of air are k = 1.4 and  $c_p = 0.240$  Btu/lbm·R (Table A-2Ea).

*Analysis* The properties of the fluid at the location where Ma = 1 are the critical properties, denoted by superscript \*. We first determine the stagnation temperature and pressure, which remain constant throughout the nozzle since the flow is isentropic.

$$T_{0} = T + \frac{V_{i}^{2}}{2c_{p}} = 630 \text{ R} + \frac{(450 \text{ ft/s})^{2}}{2 \times 0.240 \text{ Btu/lbm} \cdot \text{R}} \left( \frac{1 \text{ Btu/lbm}}{25,037 \text{ ft}^{2}/\text{s}^{2}} \right) = 646.9 \text{ R}$$

$$P_{0} = P_{i} \left( \frac{T_{0}}{T_{i}} \right)^{k/(k-1)} = (30 \text{ psia}) \left( \frac{646.9 \text{ K}}{630 \text{ K}} \right)^{1.4/(1.4-1)} = 32.9 \text{ psia}$$

$$(i) \longrightarrow \text{AIR} \longrightarrow$$

From Table A-13 (or from Eqs. 12-18 and 12-19) at Ma = 1, we read  $T/T_0 = 0.8333$ ,  $P/P_0 = 0.5283$ . Thus,

$$T = 0.8333T_0 = 0.8333(646.9 \text{ R}) =$$
**539 R** and  $P = 0.5283P_0 = 0.5283(32.9 \text{ psia}) =$ **17.4 psia**

Also,

$$c_i = \sqrt{kRT}_i = \sqrt{(1.4)(0.06855 \text{ Btu/1bm} \cdot \text{R})(630 \text{ R}) \left(\frac{25,037 \text{ ft}^2/\text{s}^2}{1 \text{ Btu/1bm}}\right)} = 1230 \text{ ft/s} \quad \text{and}$$

$$\text{Ma}_i = \frac{V_i}{c_i} = \frac{450 \text{ ft/s}}{1230 \text{ ft/s}} = 0.3657$$

From Table A-13 at this Mach number we read  $A_i/A^* = 1.7426$ . Thus the ratio of the throat area to the nozzle inlet area is

$$\frac{A^*}{A_i} = \frac{1}{1.7426} =$$
**0.574**

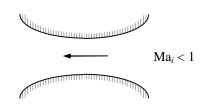
**Discussion** If we solve this problem using the relations for compressible isentropic flow, the results would be identical.

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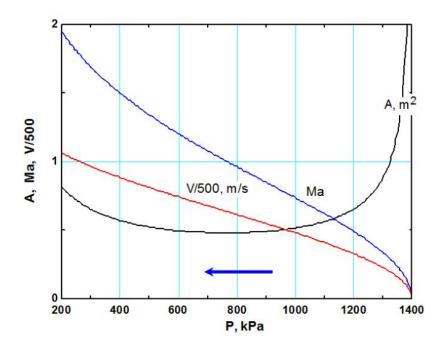
**Solution** For subsonic flow at the inlet, the variation of pressure, velocity, and Mach number along the length of the nozzle are to be sketched for an ideal gas under specified conditions.

Assumptions 1 The gas is an ideal gas. 2 Flow through the nozzle is steady, one-dimensional, and isentropic. 3 The flow is choked at the throat.

**Analysis** Using EES and  $CO_2$  as the gas, we calculate and plot flow area A, velocity V, and Mach number Ma as the pressure drops from a stagnation value of 1400 kPa to 200 kPa. Note that the curve for A is related to the shape of the nozzle, with horizontal axis serving as the centerline. The EES equation window and the plot are shown below.



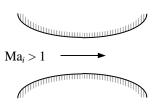
```
k=1.289
Cp=0.846 "kJ/kg.K"
R=0.1889 "kJ/kg.K"
P0=1400 "kPa"
T0=473 "K"
m=3 "kg/s"
rho_0 = P0/(R*T0)
rho=P/(R*T)
rho_norm=rho/rho_0 "Normalized density"
T=T0*(P/P0)^{(k-1)/k}
Tnorm=T/T0 "Normalized temperature"
V=SQRT(2*Cp*(T0-T)*1000)
V_norm=V/500
A=m/(rho*V)*500
C=SQRT(k*R*T*1000)
Ma=V/C
```



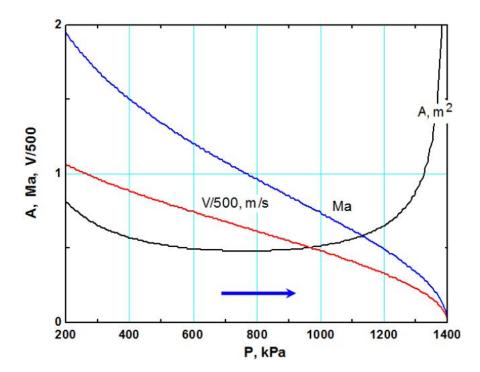
**Discussion** We are assuming that the back pressure is sufficiently low that the flow is choked at the throat, and the flow downstream of the throat is supersonic without any shock waves. Mach number and velocity continue to rise right through the throat into the diverging portion of the nozzle, since the flow becomes supersonic.

**Solution** We repeat the previous problem, but for supersonic flow at the inlet. The variation of pressure, velocity, and Mach number along the length of the nozzle are to be sketched for an ideal gas under specified conditions.

**Analysis** Using EES and  $CO_2$  as the gas, we calculate and plot flow area A, velocity V, and Mach number Ma as the pressure rises from 200 kPa at a very high velocity to the stagnation value of 1400 kPa. Note that the curve for A is related to the shape of the nozzle, with horizontal axis serving as the centerline.



```
k=1.289
Cp=0.846 "kJ/kg.K"
R=0.1889 "kJ/kg.K"
P0=1400 "kPa"
T0=473 "K"
m=3 "kg/s"
rho_0=P0/(R*T0)
rho=P/(R*T)
rho_norm=rho/rho_0 "Normalized density"
T=T0*(P/P0)^{(k-1)/k}
Tnorm=T/T0 "Normalized temperature"
V = SQRT(2*Cp*(T0-T)*1000)
V_norm=V/500
A=m/(rho*V)*500
C=SQRT(k*R*T*1000)
Ma=V/C
```



**Discussion** Note that this problem is identical to the proceeding one, except the flow direction is reversed. In fact, when plotted like this, the plots are identical.

**Solution** It is to be explained why the maximum flow rate per unit area for a given ideal gas depends only on  $P_0 / \sqrt{T_0}$ . Also for an ideal gas, a relation is to be obtained for the constant a in  $R_{\text{max}} / A^* = a \left( P_0 / \sqrt{T_0} \right)$ .

**Properties** The properties of the ideal gas considered are  $R = 0.287 \text{ kPa.m}^3/\text{kg} \cdot \text{K}$  and k = 1.4.

Analysis The maximum flow rate is given by

$$n \mathcal{K}_{\max} = A * P_0 \sqrt{k / RT_0} \left( \frac{2}{k+1} \right)^{(k+1)/2(k-1)} \quad \text{or} \quad n \mathcal{K}_{\max} / A^* = \left( P_0 / \sqrt{T_0} \right) \sqrt{k / R} \left( \frac{2}{k+1} \right)^{(k+1)/2(k-1)}$$

For a given gas, k and R are fixed, and thus the mass flow rate depends on the parameter  $P_0 / \sqrt{T_0}$ . Thus,  $n - A^*$  can be expressed as  $n - A^* = a \left( P_0 / \sqrt{T_0} \right)$  where

$$a = \sqrt{k/R} \left(\frac{2}{k+1}\right)^{(k+1)/2(k-1)} = \sqrt{\frac{1.4}{(0.287 \text{ kJ/kg.K}) \left(\frac{1000 \text{ m}^2/\text{s}^2}{1 \text{ kJ/kg}}\right)}} \left(\frac{2}{1.4+1}\right)^{2.4/0.8} = \textbf{0.0404 (m/s)}\sqrt{\textbf{K}}$$

**Discussion** Note that when sonic conditions exist at a throat of known cross-sectional area, the mass flow rate is fixed by the stagnation conditions.

#### 12-45

**Solution** An ideal gas is flowing through a nozzle. The flow area at a location where Ma = 1.8 is specified. The flow area where Ma = 0.9 is to be determined.

Assumptions Flow through the nozzle is steady, one-dimensional, and isentropic.

**Properties** The specific heat ratio is given to be k = 1.4.

Analysis The flow is assumed to be isentropic, and thus the stagnation and critical properties remain constant throughout the nozzle. The flow area at a location where  $Ma_2 = 0.9$  is determined using  $A/A^*$  data from Table A-13 to be

Ma<sub>1</sub> = 1.8: 
$$\frac{A_1}{A^*}$$
 = 1.4390  $\longrightarrow A^* = \frac{A_1}{1.4390} = \frac{36 \text{ cm}^2}{1.4390} = 25.02 \text{ cm}^2$ 

Ma<sub>2</sub> = 0.9: 
$$\frac{A_2}{A^*}$$
 = 1.0089  $\longrightarrow A_2$  = (1.0089) $A^*$  = (1.0089)(25.02 cm<sup>2</sup>) = **25.2 cm<sup>2</sup>**

**Discussion** We can also solve this problem using the relations for compressible isentropic flow. The results would be identical.

**Solution** An ideal gas is flowing through a nozzle. The flow area at a location where Ma = 1.8 is specified. The flow area where Ma = 0.9 is to be determined.

Assumptions Flow through the nozzle is steady, one-dimensional, and isentropic.

**Analysis** The flow is assumed to be isentropic, and thus the stagnation and critical properties remain constant throughout the nozzle. The flow area at a location where  $Ma_2 = 0.9$  is determined using the  $A/A^*$  relation,

$$\frac{A}{A^*} = \frac{1}{Ma} \left\{ \left( \frac{2}{k+1} \right) \left( 1 + \frac{k-1}{2} Ma^2 \right) \right\}^{(k+1)/2(k-1)}$$

For k = 1.33 and  $Ma_1 = 1.8$ :

$$\frac{A_1}{A^*} = \frac{1}{1.8} \left\{ \left( \frac{2}{1.33 + 1} \right) \left( 1 + \frac{1.33 - 1}{2} 1.8^2 \right) \right\}^{2.33/2 \times 0.33} = 1.4696$$

and,

$$A^* = \frac{A_1}{2.570} = \frac{36 \text{ cm}^2}{1.4696} = 24.50 \text{ cm}^2$$

For k = 1.33 and  $Ma_2 = 0.9$ :

$$\frac{A_2}{A^*} = \frac{1}{0.9} \left\{ \left( \frac{2}{1.33 + 1} \right) \left( 1 + \frac{1.33 - 1}{2} 0.9^2 \right) \right\}^{2.33/2 \times 0.33} = 1.0091$$

and  $A_2 = (1.0091)A^* = (1.0091)(24.50 \text{ cm}^2) = 24.7 \text{ cm}^2$ 

**Discussion** Note that the compressible flow functions in Table A-13 are prepared for k = 1.4, and thus they cannot be used to solve this problem.

#### 12-47E

**Solution** Air enters a converging-diverging nozzle at a specified temperature and pressure with low velocity. The pressure, temperature, velocity, and mass flow rate are to be calculated in the specified test section.

**Assumptions** 1 Air is an ideal gas. 2 Flow through the nozzle is steady, one-dimensional, and isentropic.

**Properties** The properties of air are k = 1.4 and R = 0.06855 Btu/lbm·R = 0.3704 psia·ft<sup>3</sup>/lbm·R.

**Analysis** The stagnation properties in this case are identical to the inlet properties since the inlet velocity is negligible. They remain constant throughout the nozzle since the flow is isentropic.

$$P_0 = P_i = 150 \text{ psia}$$
 and  $T_0 = T_i = 100^{\circ} \text{F} \approx 560 \text{ R}$ 

Then,

$$T_{e} = T_{0} \left(\frac{2}{2 + (k - 1)\text{Ma}^{2}}\right) = (560 \,\text{R}) \left(\frac{2}{2 + (1.4 - 1)2^{2}}\right) = \mathbf{311}\text{R}$$

$$P_{e} = P_{0} \left(\frac{T}{T_{0}}\right)^{k/(k - 1)} = (150 \,\text{psia}) \left(\frac{311}{560}\right)^{1.4/0.4} = \mathbf{19.1 \,\text{psia}}$$

$$\rho_{e} = \frac{P_{e}}{RT_{e}} = \frac{19.1 \,\text{psia}}{(0.3704 \,\text{psia.ft}^{3}/1\text{bm} \cdot \text{R})(311 \,\text{R})} = 0.1661 \,\text{bm/ft}^{3}$$

The nozzle exit velocity can be determined from  $V_e = \text{Ma}_e c_e$ , where  $c_e$  is the speed of sound at the exit conditions,

$$V_e = \text{Ma}_e c_e = \text{Ma}_e \sqrt{kRT_e} = (2)\sqrt{(1.4)(0.06855 \text{ Btu/1bm} \cdot \text{R})(311 \text{ R}) \left(\frac{25,037 \text{ ft}^2/\text{s}^2}{1 \text{ Btu/1bm}}\right)} = 1729 \text{ ft/s} \approx \textbf{1730 ft/s}$$

Finally

$$n = \rho_e A_e V_e = (0. 166 \text{ 1bm/ft}^3)(5 \text{ ft}^2)(1729 \text{ ft/s}) = 1435 \text{ lbm/s} \cong 1440 \text{ lbm/s}$$

**Discussion** Air must be very dry in this application because the exit temperature of air is extremely low, and any moisture in the air will turn to ice particles.

**Solution** Air enters a nozzle at specified temperature, pressure, and velocity. The exit pressure, exit temperature, and exit-to-inlet area ratio are to be determined for a Mach number of Ma = 1 at the exit.

**Assumptions** 1 Air is an ideal gas with constant specific heats at room temperature. 2 Flow through the nozzle is steady, one-dimensional, and isentropic.

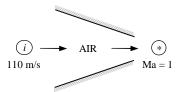
**Properties** The properties of air are k = 1.4 and  $c_p = 1.005$  kJ/kg·K.

Analysis The properties of the fluid at the location where Ma = 1 are the critical properties, denoted by superscript \*. We first determine the stagnation temperature and pressure, which remain constant throughout the nozzle since the flow is isentropic.

$$T_0 = T_i + \frac{V_i^2}{2c_p} = 420 \text{ K} + \frac{(110 \text{ m/s})^2}{2 \times 1.005 \text{ kJ/kg} \cdot \text{K}} \left(\frac{1 \text{ kJ/kg}}{1000 \text{ m}^2/\text{s}^2}\right) = 426.02$$

and

$$P_0 = P \left(\frac{T_0}{T}\right)^{k/(k-1)} = (0.5 \text{ MPa}) \left(\frac{426.02 \text{ K}}{420 \text{ K}}\right)^{1.4/(1.4-1)} = 0.52554 \text{ MPa}$$



From Table A-13 (or from Eqs. 12-18 and 12-19) at Ma = 1, we read  $T/T_0 = 0.8333$ ,  $P/P_0 = 0.5283$ . Thus,

$$T = 0.8333T_0 = 0.8333(426.02 \text{ K}) = 355.00 \text{ K} \approx 355 \text{ K}$$

and

$$P = 0.5283P_0 = 0.5283(0.52554 \text{ MPa}) = 0.27764 \text{ MPa} \approx 0.278 \text{ MPa} = 278 \text{ kPa}$$

Also,

$$c_i = \sqrt{kRT_i} = \sqrt{(1.4)(0.287 \text{ kJ/kg} \cdot \text{K})(420 \text{ K}) \left(\frac{1000 \text{ m}^2/\text{s}^2}{1 \text{ kJ/kg}}\right)} = 410.799 \text{ m/s}$$

and

$$Ma_i = \frac{V_i}{c_i} = \frac{110 \text{ m/s}}{410.799 \text{ m/s}} = 0.2678$$

$$Ma_i = \frac{V_i}{c_i} = \frac{150 \text{ m/s}}{410.799 \text{ m/s}} = 0.3651$$

From Table A-13 at this Mach number we read  $A_i/A^* = 2.3343$ . Thus the ratio of the throat area to the nozzle inlet area is

$$\frac{A*}{A} = \frac{1}{2.3343} = 0.42839 \cong \mathbf{0.428}$$

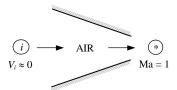
**Discussion** We can also solve this problem using the relations for compressible isentropic flow. The results would be identical.

**Solution** Air enters a nozzle at specified temperature and pressure with low velocity. The exit pressure, exit temperature, and exit-to-inlet area ratio are to be determined for a Mach number of Ma = 1 at the exit.

Assumptions 1 Air is an ideal gas. 2 Flow through the nozzle is steady, one-dimensional, and isentropic.

**Properties** The specific heat ratio of air is k = 1.4.

**Analysis** The properties of the fluid at the location where Ma = 1 are the critical properties, denoted by superscript \*. The stagnation temperature and pressure in this case are identical to the inlet temperature and pressure since the inlet velocity is negligible. They remain constant throughout the nozzle since the flow is isentropic.



$$T_0 = T_1 = 420 \text{ K}$$

$$P_0 = P_i = 0.5 \text{ MPa}$$

From Table A-13 (or from Eqs. 12-18 and 12-19) at Ma = 1, we read  $T/T_0 = 0.8333$ ,  $P/P_0 = 0.5283$ . Thus,

$$T = 0.8333T_0 = 0.8333(420 \text{ K}) = 350 \text{ K}$$

$$P = 0.5283P_0 = 0.5283(0.5 \text{ MPa}) =$$
**0.264 MPa**

The Mach number at the nozzle inlet is Ma = 0 since  $V_i \cong 0$ . From Table A-13 at this Mach number we read  $A_i/A^* = \infty$ . Thus the ratio of the throat area to the nozzle inlet area is  $\frac{A^*}{A_i} = \frac{1}{\infty} = \mathbf{0}$ .

**Discussion** If we solve this problem using the relations for compressible isentropic flow, the results would be identical.



Air enters a converging nozzle at a specified temperature and pressure with low velocity. The exit pressure, Solution the exit velocity, and the mass flow rate versus the back pressure are to be calculated and plotted.

Assumptions 1 Air is an ideal gas with constant specific heats at room temperature. 2 Flow through the nozzle is steady, one-dimensional, and isentropic.

The properties of air are k = 1.4, R = 0.287 kJ/kg·K, and  $c_p = 1.005$  kJ/kg·K. **Properties** 

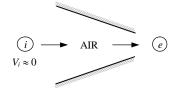
Analysis The stagnation properties in this case are identical to the inlet properties since the inlet velocity is negligible. They remain constant throughout the nozzle since the flow is isentropic.,

$$P_0 = P_i = 900 \text{ kPa}$$

$$T_0 = T_i = 400 \text{ K}$$

The critical pressure is determined to be

$$P^* = P_0 \left(\frac{2}{k+1}\right)^{k/(k-1)} = (900 \text{ kPa}) \left(\frac{2}{1.4+1}\right)^{1.4/0.4} = 475.5 \text{ kPa}$$



Then the pressure at the exit plane (throat) will be

$$P_{\rm e} = P_{\rm b}$$

$$P_{\rm b} \geq 475.5 \text{ kPa}$$

$$P_{\rm e} = P^* = 475.5 \text{ kP}$$

$$P_{\rm e} = P^* = 475.5 \text{ kPa}$$
 for  $P_{\rm b} < 475.5 \text{ kPa}$  (choked flow)

Thus the back pressure will not affect the flow when  $100 < P_b < 475.5$  kPa. For a specified exit pressure  $P_e$ , the temperature, the velocity and the mass flow rate can be determined from

Temperature

$$T_e = T_0 \left(\frac{P_e}{P_0}\right)^{(k-1)/k} = (400 \text{ K}) \left(\frac{P_e}{900}\right)^{0.4/1.4}$$

$$V = \sqrt{2c_p (T_0 - T_e)} = \sqrt{2(1.005 \text{ kJ/kg} \cdot \text{K})(400 - \text{T}_e) \left(\frac{1000 \text{ m}^2/\text{s}^2}{1 \text{ kJ/kg}}\right)}$$

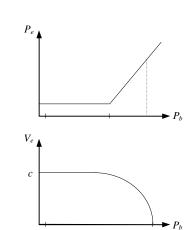
$$\rho_e = \frac{P_e}{RT_e} = \frac{P_e}{(0.287 \text{ kPa} \cdot \text{m}^3 / kg \cdot K)T_e}$$

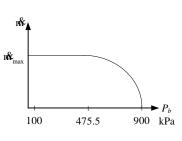
Mass flow rate

$$n = \rho_e V_e A_e = \rho_e V_e (0.001 \,\mathrm{m}^2)$$

The results of the calculations are tabulated as

$P_{\rm b}$ , kPa	P <sub>e</sub> , kPa	$T_{\rm e}$ , K	$V_{\rm e}$ , m/s	$\rho_{\rm e}$ , kg/m <sup>3</sup>	m%, kg/s
900	900	400	0	7.840	0
800	800	386.8	162.9	7.206	1.174
700	700	372.3	236.0	6.551	1.546
600	600	356.2	296.7	5.869	1.741
500	500	338.2	352.4	5.151	1.815
475.5	475.5	333.3	366.2	4.971	1.820
400	475.5	333.3	366.2	4.971	1.820
300	475.5	333.3	366.2	4.971	1.820
200	475.5	333.3	366.2	4.971	1,820
100	475.5	333.3	366.2	4.971	1.820





Discussion We see from the plots that once the flow is choked at a back pressure of 475.5 kPa, the mass flow rate remains constant regardless of how low the back pressure gets.



**Solution** We are to reconsider the previous problem. Using EES (or other) software, we are to solve the problem for the inlet conditions of 0.8 MPa and 1200 K.

Analysis Air at 800 kPa, 1200 K enters a converging nozzle with a negligible velocity. The throat area of the nozzle is 10 cm2. Assuming isentropic flow, calculate and plot the exit pressure, the exit velocity, and the mass flow rate versus the back pressure  $P_b$  for  $0.8 >= P_b >= 0.1$  MPa.

```
Procedure ExitPress(P_back,P_crit : P_exit, Condition$)
If (P_back>=P_crit) then
   P exit:=P back
                                                            "Unchoked Flow Condition"
   Condition$:='unchoked'
else
   P_exit:=P_crit
                                                            "Choked Flow Condition"
   Condition$:='choked'
Endif
End
Gas$='Air'
A cm2=10 "Throat area, cm2"
P inlet =800"kPa"
T_inlet= 1200"K"
"P back =422.7" "kPa"
A_{exit} = A_{cm2}*Convert(cm^2,m^2)
C_p=specheat(Gas$,T=T_inlet)
C_p-C_v=R
k=C_p/C_v
M=MOLARMASS(Gas$)
                                                            "Molar mass of Gas$"
R = 8.314/M
                                                            "Gas constant for Gas$"
"Since the inlet velocity is negligible, the stagnation temperature = T inlet;
and, since the nozzle is isentropic, the stagnation pressure = P_inlet."
P o=P inlet
                                                            "Stagnation pressure"
T o=T inlet
                                                            "Stagnation temperature"
P crit /P o=(2/(k+1))^{k}(k/(k-1))
                                                            "Critical pressure from Eq. 16-22"
Call ExitPress(P back, P crit: P exit, Condition$)
                                                            "Exit temperature for isentopic flow, K"
T exit/T o=(P \text{ exit/P o})^{(k-1)/k}
V_{exit} ^2/2 = C_p*(T_o-T_{exit})*1000
                                                            "Exit velocity, m/s"
Rho_exit=P_exit/(R*T_exit)
                                                            "Exit density, kg/m3"
m_dot=Rho_exit*V_exit*A_exit
                                                            "Nozzle mass flow rate, kg/s"
"If you wish to redo the plots, hide the diagram window and remove the { } from
the first 4 variables just under the procedure. Next set the desired range of
```

back pressure in the parametric table. Finally, solve the table (F3). "

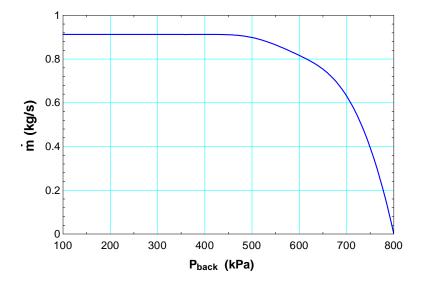
The table of results and the corresponding plot are provided below.

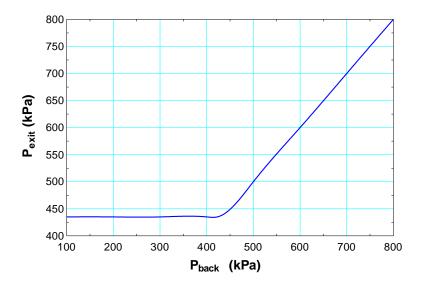
12-27

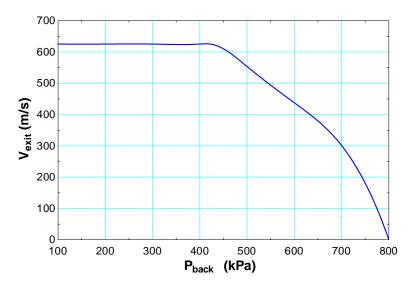
#### **EES SOLUTION**

A\_cm2=10 A\_exit=0.001 Condition\$='choked' C\_p=1.208 C\_v=0.9211 Gas\$='Air' k=1.312 M=28.97 m\_dot=0.9124 P\_back=422.7 P\_crit=434.9 P\_exit=434.9 P\_inlet=800 P\_o=800 R=0.287 Rho\_exit=1.459 T\_exit=1038 T\_inlet=1200 T\_o=1200 V\_exit=625.2

P <sub>back</sub> [kPa]	P <sub>exit</sub> [kPa]	V <sub>exit</sub> [m/s]	m [kg/s]	T <sub>exit</sub> [K]	ρ <sub>exit</sub> [kg/m <sup>3</sup> ]
100	434.9	625.2	0.9124	1038	1.459
200	434.9	625.2	0.9124	1038	1.459
300	434.9	625.2	0.9124	1038	1.459
400	434.9	625.2	0.9124	1038	1.459
422.7	434.9	625.2	0.9124	1038	1.459
500	500	553.5	0.8984	1073	1.623
600	600	437.7	0.8164	1121	1.865
700	700	300.9	0.6313	1163	2.098
800	800	0.001523	0.000003538	1200	2.323







**Discussion** We see from the plot that once the flow is choked at a back pressure of 422.7 kPa, the mass flow rate remains constant regardless of how low the back pressure gets.

#### **Shock Waves and Expansion Waves**

#### 12-52C

**Solution** We are to discuss the applicability of the isentropic flow relations across shocks and expansion waves.

Analysis The isentropic relations of ideal gases are *not* applicable for flows across (a) normal shock waves and (b) oblique shock waves, but they *are* applicable for flows across (c) Prandtl-Meyer expansion waves.

**Discussion** Flow across any kind of shock wave involves irreversible losses – hence, it cannot be isentropic.

#### 12-53C

**Solution** We are to discuss the states on the Fanno and Rayleigh lines.

Analysis The Fanno line represents the states that satisfy the conservation of mass and energy equations. The Rayleigh line represents the states that satisfy the conservation of mass and momentum equations. The intersections points of these lines represent the states that satisfy the conservation of mass, energy, and momentum equations.

**Discussion** T-s diagrams are quite helpful in understanding these kinds of flows.

#### 12-54C

**Solution** We are to analyze a claim about oblique shock analysis.

Analysis Yes, the claim is correct. Conversely, normal shocks can be thought of as special oblique shocks in which the shock angle is  $\beta = \pi/2$ , or 90°.

**Discussion** The component of flow in the direction normal to the oblique shock acts exactly like a normal shock. We can think of the flow parallel to the oblique shock as "going along for the ride" – it does not affect anything.

#### 12-55C

**Solution** We are to discuss the effect of a normal shock wave on several properties.

Analysis (a) velocity decreases, (b) static temperature increases, (c) stagnation temperature remains the same, (d) static pressure increases, and (e) stagnation pressure decreases.

**Discussion** In addition, the Mach number goes from supersonic (Ma > 1) to subsonic (Ma < 1).

#### 12-56C

**Solution** We are to discuss the formation of oblique shocks and how they differ from normal shocks.

Analysis Oblique shocks occur when a gas flowing at supersonic speeds strikes a flat or inclined surface. Normal shock waves are perpendicular to flow whereas inclined shock waves, as the name implies, are typically inclined relative to the flow direction. Also, normal shocks form a straight line whereas oblique shocks can be straight or curved, depending on the surface geometry.

**Discussion** In addition, while a normal shock must go from supersonic (Ma > 1) to subsonic (Ma < 1), the Mach number downstream of an oblique shock can be either supersonic or subsonic.

#### 12-57C

**Solution** We are to discuss whether the flow upstream and downstream of an oblique shock needs to be supersonic.

**Analysis** Yes, the *upstream flow* has to be supersonic for an oblique shock to occur. No, the flow *downstream* of an oblique shock can be subsonic, sonic, and even supersonic.

**Discussion** The latter is not true for normal shocks. For a normal shock, the flow must always go from supersonic (Ma > 1) to subsonic (Ma < 1).

#### 12-58C

**Solution** We are to determine if Ma downstream of a normal shock can be supersonic.

**Analysis** No, the second law of thermodynamics requires the flow after the shock to be subsonic.

**Discussion** A normal shock wave always goes from supersonic to subsonic in the flow direction.

#### 12-59C

**Solution** We are to discuss shock detachment at the nose of a 2-D wedge-shaped body.

Analysis When the wedge half-angle  $\delta$  is greater than the maximum deflection angle  $\theta_{max}$ , the shock becomes curved and detaches from the nose of the wedge, forming what is called a *detached oblique shock* or a *bow wave*. The numerical value of the shock angle at the nose is  $\beta = 90^{\circ}$ .

**Discussion** When  $\delta$  is less than  $\theta_{max}$ , the oblique shock is attached to the nose.

#### 12-60C

**Solution** We are to discuss the shock at the nose of a rounded body in supersonic flow.

Analysis When supersonic flow impinges on a blunt body like the rounded nose of an aircraft, the wedge half-angle  $\delta$  at the nose is 90°, and an attached oblique shock cannot exist, regardless of Mach number. Therefore, a detached oblique shock must occur in front of all such blunt-nosed bodies, whether two-dimensional, axisymmetric, or fully three-dimensional.

**Discussion** Since  $\delta = 90^{\circ}$  at the nose,  $\delta$  is always greater than  $\theta_{\text{max}}$ , regardless of Ma or the shape of the rest of the body.

#### 12-61C

**Solution** We are to discuss if a shock wave can develop in the converging section of a C-V nozzle.

**Analysis** No, because the flow must be supersonic before a shock wave can occur. The flow in the converging section of a nozzle is always subsonic.

**Discussion** A normal shock (if it is to occur) would occur in the supersonic (diverging) section of the nozzle.

#### 12-62

**Solution** Air flowing through a nozzle experiences a normal shock. Various properties are to be calculated before and after the shock.

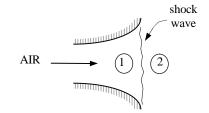
**Assumptions** 1 Air is an ideal gas with constant specific heats. 2 Flow through the nozzle is steady, one-dimensional, and isentropic before the shock occurs.

**Properties** The properties of air at room temperature are k=1.4, R=0.287 kJ/kg·K, and  $c_p=1.005$  kJ/kg·K.

Analysis The stagnation temperature and pressure before the shock are

$$T_{01} = T_1 + \frac{V_1^2}{2c_p} = 230 + \frac{(815 \text{ m/s})^2}{2(1.005 \text{ kJ/kg} \cdot \text{K})} \left( \frac{1 \text{ kJ/kg}}{1000 \text{ m}^2/\text{s}^2} \right) = 560.5 \text{ K}$$

$$P_{01} = P_1 \left(\frac{T_{01}}{T_1}\right)^{k/(k-1)} = (26 \text{ kPa}) \left(\frac{560.5 \text{ K}}{230 \text{ K}}\right)^{1.4/(1.4-1)} = 587.3 \text{ kPa}$$



The velocity and the Mach number before the shock are determined from

$$c_1 = \sqrt{kRT_1} = \sqrt{(1.4)(0.287 \text{ kJ/kg} \cdot \text{K})(230 \text{ K}) \left(\frac{1000 \text{ m}^2/\text{s}^2}{1 \text{ kJ/kg}}\right)} =$$
**304.0 m/s**

and

$$Ma_1 = \frac{V_1}{c_1} = \frac{815 \text{ m/s}}{304.0 \text{ m/s}} = 2.681$$

The fluid properties after the shock (denoted by subscript 2) are related to those before the shock through the functions listed in Table A-14. For  $Ma_1 = 2.681$  we read

Ma<sub>2</sub> = **0.4972**, 
$$\frac{P_{02}}{P_1}$$
 = 9.7330,  $\frac{P_2}{P_1}$  = 8.2208, and  $\frac{T_2}{T_1}$  = 2.3230

Then the stagnation pressure  $P_{02}$ , static pressure  $P_2$ , and static temperature  $T_2$ , are determined to be

$$P_{02} = 9.7330P_1 = (9.7330)(26 \text{ kPa}) = 253.1 \text{ kPa}$$

$$P_2 = 8.2208P_1 = (8.2208)(26 \text{ kPa}) = 213.7 \text{ kPa}$$

$$T_2 = 2.3230T_1 = (2.3230)(230 \text{ K}) =$$
**534.3 K**

The air velocity after the shock can be determined from  $V_2 = \text{Ma}_2 c_2$ , where  $c_2$  is the speed of sound at the exit conditions after the shock,

$$V_2 = \text{Ma}_2 c_2 = \text{Ma}_2 \sqrt{kRT_2} = (0.4972) \sqrt{(1.4)(0.287 \text{ kJ/kg.K})(534.3 \text{ K}) \left(\frac{1000 \text{ m}^2 / \text{s}^2}{1 \text{ kJ/kg}}\right)} = \textbf{230.4 m/s}$$

**Discussion** This problem could also be solved using the relations for compressible flow and normal shock functions. The results would be identical.

**Solution** Air flowing through a nozzle experiences a normal shock. The entropy change of air across the normal shock wave is to be determined.

**Assumptions** 1 Air is an ideal gas with constant specific heats. 2 Flow through the nozzle is steady, one-dimensional, and isentropic before the shock occurs.

**Properties** The properties of air at room temperature are  $R = 0.287 \text{ kJ/kg} \cdot \text{K}$  and  $c_p = 1.005 \text{ kJ/kg} \cdot \text{K}$ .

Analysis The entropy change across the shock is determined to be

$$s_2 - s_1 = c_p \ln \frac{T_2}{T_1} - R \ln \frac{P_2}{P_1}$$
= (1.005 kJ/kg · K)ln(2.3230) - (0.287 kJ/kg · K)ln(8.2208)
= **0.242 kJ/kg** · **K**

**Discussion** A shock wave is a highly dissipative process, and the entropy generation is large during shock waves.

#### **12-64**

**Solution** For an ideal gas flowing through a normal shock, a relation for  $V_2/V_1$  in terms of k, Ma<sub>1</sub>, and Ma<sub>2</sub> is to be developed.

Analysis The conservation of mass relation across the shock is  $\rho_1 V_1 = \rho_2 V_2$  and it can be expressed as

$$\frac{V_2}{V_1} = \frac{\rho_1}{\rho_2} = \frac{P_1 / RT_1}{P_2 / RT_2} = \left(\frac{P_1}{P_2}\right) \left(\frac{T_2}{T_1}\right)$$

From Eqs. 12-35 and 12-38,

$$\frac{V_2}{V_1} = \left(\frac{1 + k \text{Ma}_2^2}{1 + k \text{Ma}_1^2}\right) \left(\frac{1 + \text{Ma}_1^2(k-1)/2}{1 + \text{Ma}_2^2(k-1)/2}\right)$$

**Discussion** This is an important relation as it enables us to determine the velocity ratio across a normal shock when the Mach numbers before and after the shock are known.

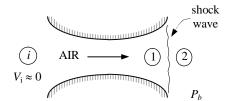
**Solution** Air enters a converging-diverging nozzle at a specified state. The required back pressure that produces a normal shock at the exit plane is to be determined for the specified nozzle geometry.

**Assumptions** 1 Air is an ideal gas. 2 Flow through the nozzle is steady, one-dimensional, and isentropic before the shock occurs. 3 The shock wave occurs at the exit plane.

Analysis The inlet stagnation pressure in this case is identical to the inlet pressure since the inlet velocity is negligible. Since the flow before the shock to be isentropic,

$$P_{01} = P_i = 2 \text{ MPa}$$

It is specified that  $A/A^* = 3.5$ . From Table A-13, Mach number and the pressure ratio which corresponds to this area ratio are the Ma<sub>1</sub> = 2.80 and  $P_1/P_{01} = 0.0368$ . The pressure ratio across the shock for this Ma<sub>1</sub> value is, from Table A-14,  $P_2/P_1 = 8.98$ . Thus the back pressure, which is equal to the static pressure at the nozzle exit, must be



$$P_2 = 8.98P_1 = 8.98 \times 0.0368P_{01} = 8.98 \times 0.0368 \times (2 \text{ MPa}) = 0.661 \text{ MPa}$$

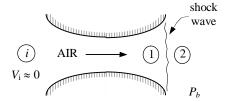
**Discussion** We can also solve this problem using the relations for compressible flow and normal shock functions. The results would be identical.

#### **12-66**

**Solution** Air enters a converging-diverging nozzle at a specified state. The required back pressure that produces a normal shock at the exit plane is to be determined for the specified nozzle geometry.

**Assumptions** 1 Air is an ideal gas. 2 Flow through the nozzle is steady, one-dimensional, and isentropic before the shock occurs.

**Analysis** The inlet stagnation pressure in this case is identical to the inlet pressure since the inlet velocity is negligible. Since the flow before the shock to be isentropic,



$$P_{0x} = P_i = 2 \text{ MPa}$$

It is specified that  $A/A^* = 2$ . From Table A-13, the Mach number and the pressure ratio which corresponds to this area ratio are the Ma<sub>1</sub> =2.20 and  $P_1/P_{01} = 0.0935$ . The pressure ratio across the shock for this M<sub>1</sub> value is, from Table A-14,  $P_2/P_1 = 5.48$ . Thus the back pressure, which is equal to the static pressure at the nozzle exit, must be

$$P_2 = 5.48P_1 = 5.48 \times 0.0935 P_{01} = 5.48 \times 0.0935 \times (2 \text{ MPa}) = 1.02 \text{ MPa}$$

**Discussion** We can also solve this problem using the relations for compressible flow and normal shock functions. The results would be identical.



**Solution** Air flowing through a nozzle experiences a normal shock. Effect of the shock wave on various properties is to be determined. Analysis is to be repeated for helium

Assumptions 1 Air and helium are ideal gases with constant specific heats. 2 Flow through the nozzle is steady, one-dimensional, and isentropic before the shock occurs.

**Properties** The properties of air are k = 1.4 and R = 0.06855 Btu/lbm·R, and the properties of helium are k = 1.667 and R = 0.4961 Btu/lbm·R.

Analysis The air properties upstream the shock are

$$Ma_1 = 2.5$$
,  $P_1 = 10$  psia, and  $T_1 = 440.5$  R

Fluid properties after the shock (denoted by subscript 2) are related to those before the shock through the functions listed in Table A-14. For  $Ma_1 = 2.5$ ,

Ma<sub>2</sub> = **0.513**, 
$$\frac{P_{02}}{P_1}$$
 = 8.5262,  $\frac{P_2}{P_1}$  = 7.125, and  $\frac{T_2}{T_1}$  = 2.1375



Then the stagnation pressure  $P_{02}$ , static pressure  $P_2$ , and static temperature  $T_2$ , are determined to be

$$P_{02} = 8.5262P_1 = (8.5262)(10 \text{ psia}) = 85.3 \text{ psia}$$

$$P_2 = 7.125P_1 = (7.125)(10 \text{ psia}) = 71.3 \text{ psia}$$

$$T_2 = 2.1375T_1 = (2.1375)(440.5 \text{ R}) = 942 \text{ R}$$

The air velocity after the shock can be determined from  $V_2 = \text{Ma}_2 c_2$ , where  $c_2$  is the speed of sound at the exit conditions after the shock,

$$V_2 = \text{Ma}_2 c_2 = \text{Ma}_2 \sqrt{kRT_2} = (0.513) \sqrt{(1.4)(0.06855 \,\text{Btu/1bm} \cdot \text{R})(941.6 \,\text{R}) \left(\frac{25,037 \,\text{ft}^2 / \text{s}^2}{1 \,\text{Btu/1bm}}\right)} = \textbf{772 ft/s}$$

We now repeat the analysis for helium. This time we cannot use the tabulated values in Table A-14 since k is not 1.4. Therefore, we have to calculate the desired quantities using the analytical relations,

$$\begin{aligned} &\text{Ma}_{2} = \left(\frac{\text{Ma}_{1}^{2} + 2/(k-1)}{2\text{Ma}_{1}^{2}k/(k-1) - 1}\right)^{1/2} = \left(\frac{2.5^{2} + 2/(1.667 - 1)}{2 \times 2.5^{2} \times 1.667/(1.667 - 1) - 1}\right)^{1/2} = \textbf{0.553} \\ &\frac{P_{2}}{P_{1}} = \frac{1 + k\text{Ma}_{1}^{2}}{1 + k\text{Ma}_{2}^{2}} = \frac{1 + 1.667 \times 2.5^{2}}{1 + 1.667 \times 0.553^{2}} = 7.5632 \\ &\frac{T_{2}}{T_{1}} = \frac{1 + \text{Ma}_{1}^{2}(k-1)/2}{1 + \text{Ma}_{2}^{2}(k-1)/2} = \frac{1 + 2.5^{2}(1.667 - 1)/2}{1 + 0.553^{2}(1.667 - 1)/2} = 2.7989 \\ &\frac{P_{02}}{P_{1}} = \left(\frac{1 + k\text{Ma}_{1}^{2}}{1 + k\text{Ma}_{2}^{2}}\right) \left(1 + (k-1)\text{Ma}_{2}^{2}/2\right)^{k/(k-1)} \\ &= \left(\frac{1 + 1.667 \times 2.5^{2}}{1 + 1.667 \times 0.553^{2}}\right) \left(1 + (1.667 - 1) \times 0.553^{2}/2\right)^{1.667/0.667} = 9.641 \end{aligned}$$

Thus,  $P_{02} = 11.546P_1 = (11.546)(10 \text{ psia}) = 115 \text{ psia}$ 

$$P_2 = 7.5632P_1 = (7.5632)(10 \text{ psia}) = 75.6 \text{ psia}$$

$$T_2 = 2.7989T_1 = (2.7989)(440.5 \text{ R}) = 1233 \text{ R}$$

$$V_2 = \text{Ma}_2 c_2 = \text{Ma}_2 \sqrt{kRT_2} = (0.553) \sqrt{(1.667)(0.4961 \,\text{Btu/1bm.R})(1232.9 \,\text{R}) \left(\frac{25,037 \,\text{ft}^2 / \text{s}^2}{1 \,\text{Btu/1bm}}\right)} = \textbf{2794} \quad \textbf{ft/s}$$

**Discussion** This problem could also be solved using the relations for compressible flow and normal shock functions. The results would be identical.



**Solution** We are to reconsider Prob. 12-67E. Using EES (or other) software, we are to study the effects of both air and helium flowing steadily in a nozzle when there is a normal shock at a Mach number in the range 2 < Mx < 3.5. In addition to the required information, we are to calculate the entropy change of the air and helium across the normal shock, and tabulate the results in a parametric table.

Analysis We use EES to calculate the entropy change of the air and helium across the normal shock. The results are given in the Parametric Table for  $2 < M_x < 3.5$ .

```
Procedure NormalShock(M_x,k:M_y,PyOPx, TyOTx,RhoyORhox, PoyOPox, PoyOPx)
    If M \times < 1 Then
              M_y = -1000; PyOPx = -1000; TyOTx = -1000; RhoyORhox = -1000
              PoyOPox=-1000;PoyOPx=-1000
    else
              M_y=sqrt((M_x^2+2/(k-1))/(2*M_x^2*k/(k-1)-1))
             PyOPx = (1+k*M_x^2)/(1+k*M_y^2)
              TyOTx=(1+M_x^2*(k-1)/2)/(1+M_y^2*(k-1)/2)
             RhoyORhox=PyOPx/TyOTx
              PoyOPox = M_x/M_y*( \ (1+M_y^2*(k-1)/2)/\ (1+M_x^2*(k-1)/2) \ )^*((k+1)/(2*(k-1)))
              PoyOPx = (1 + k*M_x^2)*(1 + M_y^2*(k-1)/2)^(k/(k-1))/(1 + k*M_y^2)
    Endif
End
Function ExitPress(P_back,P_crit)
If P_back>=P_crit then ExitPress:=P_back
                                           "Unchoked Flow Condition"
If P_back<P_crit then ExitPress:=P_crit
                                           "Choked Flow Condition"
Procedure GetProp(Gas$:Cp,k,R) "Cp and k data are from Text Table A.2E"
   M=MOLARMASS(Gas$)
                                      "Molar mass of Gas$"
   R = 1545/M
                                      "Particular gas constant for Gas$, ft-lbf/lbm-R"
                                      "k = Ratio of Cp to Cv"
                                      "Cp = Specific heat at constant pressure"
   if Gas$='Air' then
                                      Cp=0.24"Btu/lbm-R"; k=1.4
   endif
   if Gas$='CO2' then
                                      Cp=0.203"Btu/lbm_R"; k=1.289
   endif
   if Gas$='Helium' then
                                      Cp=1.25"Btu/lbm-R"; k=1.667
   endif
End
"Variable Definitions:"
"M = flow Mach Number"
"P_ratio = P/P_o for compressible, isentropic flow"
"T_ratio = T/T_o for compressible, isentropic flow"
"Rho_ratio= Rho/Rho_o for compressible, isentropic flow"
"A_ratio=A/A* for compressible, isentropic flow"
"Fluid properties before the shock are denoted with a subscript x"
"Fluid properties after the shock are denoted with a subscript y"
"M_y = Mach Number down stream of normal shock"
"PyOverPx= P_y/P_x Pressue ratio across normal shock"
"TyOverTx =T_y/T_x Temperature ratio across normal shock"
"RhoyOverRhox=Rho_y/Rho_x Density ratio across normal shock"
"PoyOverPox = P_oy/P_ox Stagantion pressure ratio across normal shock"
"PoyOverPx = P oy/P x Stagnation pressure after normal shock ratioed to pressure before shock"
"Input Data"
\{P_x = 10 \text{ "psia"}\}
                                      "Values of P_x, T_x, and M_x are set in the Parametric Table"
\{T_x = 440.5 \text{ "R"}\}
```

#### $\{M_x = 2.5\}$

Gas\$='Air' "This program has been written for the gases Air, CO2, and Helium"

Call GetProp(Gas\$:Cp,k,R)

Call NormalShock(M\_x,k:M\_y,PyOverPx, TyOverTx,RhoyOverRhox, PoyOverPox, PoyOverPx)

P\_oy\_air=P\_x\*PoyOverPx "Stagnation pressure after the shock"

P\_y\_air=P\_x\*PyOverPx "Pressure after the shock"
T\_y\_air=T\_x\*TyOverTx "Temperature after the shock"
M\_y\_air=M\_y "Mach number after the shock"

"The velocity after the shock can be found from the product of the Mach number and speed of sound after the shock."

 $C_y_air = sqrt(k*R"ft-lbf/lbm_R"*T_y_air"R"*32.2 "lbm-ft/lbf-s^2")$ 

V\_y\_air=M\_y\_air\*C\_y\_air

DELTAs\_air=entropy(air,T=T\_y\_air, P=P\_y\_air) -entropy(air,T=T\_x,P=P\_x)

Gas2\$='Helium' "Gas2\$ can be either Helium or CO2"

Call GetProp(Gas2\$:Cp\_2,k\_2,R\_2)

Call NormalShock(M\_x,k\_2:M\_y2,PyOverPx2, TyOverTx2,RhoyOverRhox2, PoyOverPox2, PoyOverPx2)

P\_oy\_he=P\_x\*PoyOverPx2 "Stagnation pressure after the shock"

P\_y\_he=P\_x\*PyOverPx2 "Pressure after the shock"
T\_y\_he=T\_x\*TyOverTx2 "Temperature after the shock"
M\_y\_he=M\_y2 "Mach number after the shock"

"The velocity after the shock can be found from the product of the Mach number and speed of sound after the shock."

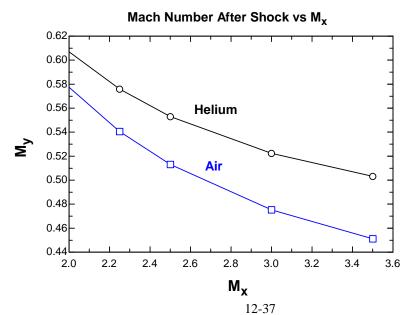
 $C_yhe = sqrt(k_2*R_2"ft-lbf/lbm_R"*T_yhe"R"*32.2 "lbm-ft/lbf-s^2")$ 

 $V_y_he=M_y_he*C_y_he$ 

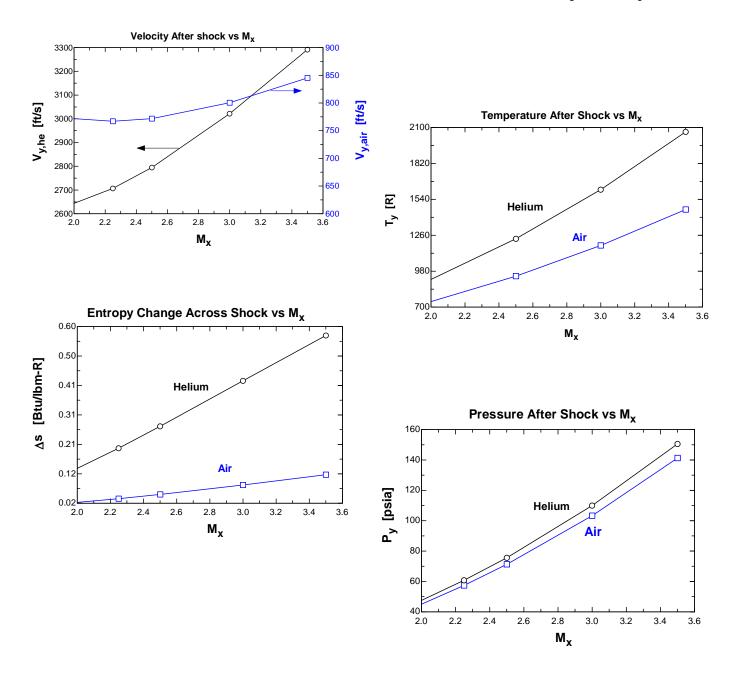
DELTAs\_he=entropy(helium,T=T\_y\_he, P=P\_y\_he) -entropy(helium,T=T\_x,P=P\_x)

The parametric table and the corresponding plots are shown below.

$V_{y,he}$	$V_{y,air}$	T <sub>y,he</sub>	T <sub>y,air</sub>	T <sub>x</sub>	P <sub>y,he</sub>	P <sub>y,air</sub>	P <sub>x</sub>	P <sub>oy,he</sub>	P <sub>oy,air</sub>	$M_{y,he}$	$M_{y,air}$	$M_x$	$\Delta s_{he}$	$\Delta s_{air}$
[ft/s]	[ft/s]	[R]	[R]	[R]	[psia]	[psia]	[psia]	[psia]	[psia]				[Btu/lbm-R]	[Btu/lbm-R]
2644	771.9	915.6	743.3	440.5	47.5	45	10	63.46	56.4	0.607	0.5774	2	0.1345	0.0228
2707	767.1	1066	837.6	440.5	60.79	57.4	10	79.01	70.02	0.5759	0.5406	2.25	0.2011	0.0351
2795	771.9	1233	941.6	440.5	75.63	71.25	10	96.41	85.26	0.553	0.513	2.5	0.2728	0.04899
3022	800.4	1616	1180	440.5	110	103.3	10	136.7	120.6	0.5223	0.4752	3	0.4223	0.08
3292	845.4	2066	1460	440.5	150.6	141.3	10	184.5	162.4	0.5032	0.4512	3.5	0.5711	0.1136



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**Discussion** In all cases, regardless of the fluid or the Mach number, entropy increases across a shock wave. This is because a shock wave involves irreversibilities.

**Solution** Air flowing through a converging-diverging nozzle experiences a normal shock at the exit. The effect of the shock wave on various properties is to be determined.

**Assumptions** 1 Air is an ideal gas with constant specific heats. 2 Flow through the nozzle is steady, one-dimensional, and isentropic before the shock occurs. 3 The shock wave occurs at the exit plane.

**Properties** The properties of air are k = 1.4 and R = 0.287 kJ/kg·K.

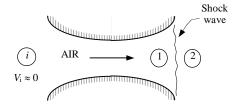
**Analysis** The inlet stagnation properties in this case are identical to the inlet properties since the inlet velocity is negligible. Then,

$$P_{01} = P_i = 1 \text{ MPa}$$

$$T_{01} = T_i = 300 \text{ K}$$

Then,

$$T_1 = T_{01} \left( \frac{2}{2 + (k-1)\text{Ma}_1^2} \right) = (300 \text{ K}) \left( \frac{2}{2 + (1.4 - 1)2.4^2} \right) = 139.4 \text{ K}$$



and

$$P_1 = P_{01} \left(\frac{T_1}{T_0}\right)^{k/(k-1)} = (1 \text{ MPa}) \left(\frac{139.4}{300}\right)^{1.4/0.4} = 0.06840 \text{ MPa}$$

The fluid properties after the shock (denoted by subscript 2) are related to those before the shock through the functions listed in Table A-14. For  $Ma_1 = 2.4$  we read

$$Ma_2 = 0.5231 \cong \mathbf{0.523}, \quad \frac{P_{02}}{P_{01}} = 0.5401, \quad \frac{P_2}{P_1} = 6.5533, \text{ and } \frac{T_2}{T_1} = 2.0403$$

Then the stagnation pressure  $P_{02}$ , static pressure  $P_2$ , and static temperature  $T_2$ , are determined to be

$$P_{02} = 0.5401P_{01} = (0.5401)(1.0 \text{ MPa}) = 0.540 \text{ MPa} = 540 \text{ kPa}$$

$$P_2 = 6.5533P_1 = (6.5533)(0.06840 \text{ MPa}) = 0.448 \text{ MPa} = 448 \text{ kPa}$$

$$T_2 = 2.0403T_1 = (2.0403)(139.4 \text{ K}) = 284 \text{ K}$$

The air velocity after the shock can be determined from  $V_2 = \text{Ma}_2 c_2$ , where  $c_2$  is the speed of sound at the exit conditions after the shock,

$$V_2 = \text{Ma}_2 c_2 = \text{Ma}_2 \sqrt{kRT_2} = (0.5231) \sqrt{(1.4)(0.287 \text{ kJ/kg} \cdot \text{K})(284 \text{ K}) \left(\frac{1000 \text{ m}^2/\text{s}^2}{1 \text{ kJ/kg}}\right)} = \textbf{177 m/s}$$

**Discussion** We can also solve this problem using the relations for normal shock functions. The results would be identical.

► Ma

0

**Solution** The entropy change of air across the shock for upstream Mach numbers between 0.5 and 1.5 is to be determined and plotted.

Assumptions 1 Air is an ideal gas. 2 Flow through the nozzle is steady, one-dimensional, and isentropic before the shock occurs.

**Properties** The properties of air are k = 1.4, R = 0.287 kJ/kg·K, and  $c_p = 1.005$  kJ/kg·K.

Analysis The entropy change across the shock is determined to be

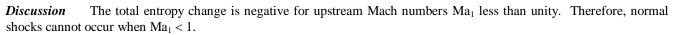
$$s_2 - s_1 = c_p \ln \frac{T_2}{T_1} - R \ln \frac{P_2}{P_1}$$

where

$$\operatorname{Ma}_{2} = \left(\frac{\operatorname{Ma}_{1}^{2} + 2/(k-1)}{2\operatorname{Ma}_{1}^{2}k/(k-1) - 1}\right)^{1/2}, \quad \frac{P_{2}}{P_{1}} = \frac{1 + k\operatorname{Ma}_{1}^{2}}{1 + k\operatorname{Ma}_{2}^{2}}, \text{ and } \frac{T_{2}}{T_{1}} = \frac{1 + \operatorname{Ma}_{1}^{2}(k-1)/2}{1 + \operatorname{Ma}_{2}^{2}(k-1)/2}$$

The results of the calculations can be tabulated as

$Ma_1$	$Ma_2$	$T_2/T_1$	$P_2/P_1$	$s_2 - s_1$
0.5	2.6458	0.1250	0.4375	-1.853
0.6	1.8778	0.2533	0.6287	-1.247
0.7	1.5031	0.4050	0.7563	-0.828
0.8	1.2731	0.5800	0.8519	-0.501
0.9	1.1154	0.7783	0.9305	-0.231
1.0	1.0000	1.0000	1.0000	0.0
1.1	0.9118	1.0649	1.2450	0.0003
1.2	0.8422	1.1280	1.5133	0.0021
1.3	0.7860	1.1909	1.8050	0.0061
1.4	0.7397	1.2547	2.1200	0.0124
1.5	0.7011	1.3202	2.4583	0.0210



## 12-71

**Solution** Supersonic airflow approaches the nose of a two-dimensional wedge and undergoes a straight oblique shock. For a specified Mach number, the minimum shock angle and the maximum deflection angle are to be determined.

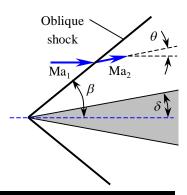
**Assumptions** Air is an ideal gas with a constant specific heat ratio of k = 1.4 (so that Fig. 12-41 is applicable).

**Analysis** For Ma = 5, we read from Fig. 12-41

Minimum shock (or wave) angle:  $\beta_{min} = 12^{\circ}$ 

Maximum deflection (or turning) angle:  $\theta_{\text{max}} = 41.5^{\circ}$ 

**Discussion** Note that the minimum shock angle decreases and the maximum deflection angle increases with increasing Mach number Ma<sub>1</sub>.

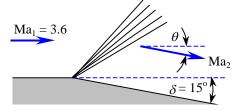


**Solution** Air flowing at a specified supersonic Mach number undergoes an expansion turn. The Mach number, pressure, and temperature downstream of the sudden expansion along a wall are to be determined.

**Assumptions** 1 The flow is steady. 2 The boundary layer on the wedge is very thin. 3 Air is an ideal gas with constant specific heats.

**Properties** The specific heat ratio of air is k = 1.4.

**Analysis** On the basis of Assumption #2, we take the deflection angle as equal to the wedge half-angle, i.e.,  $\theta \approx \delta = 15^{\circ}$ . Then the upstream and downstream Prandtl-Meyer functions are determined to be



$$v(Ma) = \sqrt{\frac{k+1}{k-1}} \tan^{-1} \left( \sqrt{\frac{k-1}{k+1} (Ma^2 - 1)} \right) - \tan^{-1} \left( \sqrt{Ma^2 - 1} \right)$$

Upstream:

$$v(\text{Ma}_1) = \sqrt{\frac{1.4+1}{1.4-1}} \tan^{-1} \left( \sqrt{\frac{1.4-1}{1.4+1} (3.6^2 - 1)} \right) - \tan^{-1} \left( \sqrt{3.6^2 - 1} \right) = 60.09^{\circ}$$

Then the downstream Prandtl-Meyer function becomes

$$v(\text{Ma}_2) = \theta + v(\text{Ma}_1) = 15^\circ + 60.09^\circ = 75.09^\circ$$

Ma<sub>2</sub> is found from the Prandtl-Meyer relation, which is now implicit:

Downstream: 
$$v(\text{Ma}_2) = \sqrt{\frac{1.4+1}{1.4-1}} \tan^{-1} \left( \sqrt{\frac{1.4-1}{1.4+1} \text{Ma}_2^2 - 1} \right) - \tan^{-1} \left( \sqrt{\text{Ma}_2^2 - 1} \right) = 75.09^{\circ}$$

Solution of this implicit equation gives  $Ma_2 = 4.81$ . Then the downstream pressure and temperature are determined from the isentropic flow relations:

$$\begin{split} P_2 &= \frac{P_2 \ / \ P_0}{P_1 \ / \ P_0} \ P_1 = \frac{\left[1 + \mathrm{Ma}_2^{\,2}(k-1) \ / \ 2\right]^{-k \ / (k-1)}}{\left[1 + \mathrm{Ma}_1^{\,2}(k-1) \ / \ 2\right]^{-k \ / (k-1)}} P_1 = \frac{\left[1 + 4.81^2 \left(1.4 - 1\right) \ / \ 2\right]^{-1.4 \ / 0.4}}{\left[1 + 3.6^2 \left(1.4 - 1\right) \ / \ 2\right]^{-1.4 \ / 0.4}} (32 \ \mathrm{kPa}) = \mathbf{6.65} \ \mathrm{kPa} \\ T_2 &= \frac{T_2 \ / \ T_0}{T_1 \ / \ T_0} T_1 = \frac{\left[1 + \mathrm{Ma}_2^{\,2}(k-1) \ / \ 2\right]^{-1}}{\left[1 + \mathrm{Ma}_1^{\,2}(k-1) \ / \ 2\right]^{-1}} T_1 = \frac{\left[1 + 4.81^2 \left(1.4 - 1\right) \ / \ 2\right]^{-1}}{\left[1 + 3.6^2 \left(1.4 - 1\right) \ / \ 2\right]^{-1}} (240 \ \mathrm{K}) = \mathbf{153} \ \mathrm{K} \end{split}$$

Note that this is an expansion, and Mach number increases while pressure and temperature decrease, as expected.

**Discussion** There are compressible flow calculators on the Internet that solve these implicit equations that arise in the analysis of compressible flow, along with both normal and oblique shock equations; e.g., see <a href="https://www.aoe.vt.edu/~devenpor/aoe3114/calc.html">www.aoe.vt.edu/~devenpor/aoe3114/calc.html</a>.

**Solution** Air flowing at a specified supersonic Mach number undergoes an expansion turn over a tilted wedge. The Mach number, pressure, and temperature downstream of the sudden expansion above the wedge are to be determined.

Assumptions 1 The flow is steady. 2 The boundary layer on the wedge is very thin. 3 Air is an ideal gas with constant specific heats.

**Properties** The specific heat ratio of air is k = 1.4.

**Analysis** On the basis of Assumption #2, the deflection angle is determined to be  $\theta \approx \delta = 25^{\circ} - 10^{\circ} = 15^{\circ}$ . Then the upstream and downstream Prandtl-Meyer functions are determined to be

$$v(Ma) = \sqrt{\frac{k+1}{k-1}} \tan^{-1} \left( \sqrt{\frac{k-1}{k+1} (Ma^2 - 1)} \right) - \tan^{-1} \left( \sqrt{Ma^2 - 1} \right)$$

Upstream:

$$v(\mathrm{Ma}_1) = \sqrt{\frac{1.4 + 1}{1.4 - 1}} \tan^{-1} \left( \sqrt{\frac{1.4 - 1}{1.4 + 1} (2.4^2 - 1)} \right) - \tan^{-1} \left( \sqrt{2.4^2 - 1} \right) = 36.75^{\circ}$$

Then the downstream Prandtl-Meyer function becomes

$$v(\text{Ma}_2) = \theta + v(\text{Ma}_1) = 15^\circ + 36.75^\circ = 51.75^\circ$$

Now Ma<sub>2</sub> is found from the Prandtl-Meyer relation, which is now implicit:

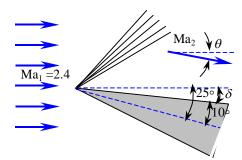
Downstream: 
$$v(\text{Ma}_2) = \sqrt{\frac{1.4+1}{1.4-1}} \tan^{-1} \left( \sqrt{\frac{1.4-1}{1.4+1}} (\text{Ma}_2^2 - 1) \right) - \tan^{-1} \left( \sqrt{\text{Ma}_2^2 - 1} \right) = 51.75^{\circ}$$

It gives  $Ma_2 = 3.105$ . Then the downstream pressure and temperature are determined from the isentropic flow relations

$$P_2 = \frac{P_2 \ / \ P_0}{P_1 \ / \ P_0} \ P_1 = \frac{[1 + \mathrm{Ma}_2^2 (k-1) \ / \ 2]^{-k \ / (k-1)}}{[1 + \mathrm{Ma}_1^2 (k-1) \ / \ 2]^{-k \ / (k-1)}} \ P_1 = \frac{[1 + 3.105^2 (1.4-1) \ / \ 2]^{-1.4 \ / 0.4}}{[1 + 2.4^2 (1.4-1) \ / \ 2]^{-1.4 \ / 0.4}} \ (70 \ \mathrm{kPa}) = \mathbf{23.8 \ kPa}$$
 
$$T_2 = \frac{T_2 \ / \ T_0}{T_1 \ / \ T_0} T_1 = \frac{[1 + \mathrm{Ma}_2^2 (k-1) \ / \ 2]^{-1}}{[1 + \mathrm{Ma}_1^2 (k-1) \ / \ 2]^{-1}} T_1 = \frac{[1 + 3.105^2 (1.4-1) \ / \ 2]^{-1}}{[1 + 2.4^2 (1.4-1) \ / \ 2]^{-1}} (260 \ \mathrm{K}) = \mathbf{191 \ K}$$

Note that this is an expansion, and Mach number increases while pressure and temperature decrease, as expected.

**Discussion** There are compressible flow calculators on the Internet that solve these implicit equations that arise in the analysis of compressible flow, along with both normal and oblique shock equations; e.g., see <a href="https://www.aoe.vt.edu/~devenpor/aoe3114/calc.html">www.aoe.vt.edu/~devenpor/aoe3114/calc.html</a>.



**Solution** Air flowing at a specified supersonic Mach number undergoes a compression turn (an oblique shock) over a tilted wedge. The Mach number, pressure, and temperature downstream of the shock below the wedge are to be determined.

**Assumptions** 1 The flow is steady. 2 The boundary layer on the wedge is very thin. 3 Air is an ideal gas with constant specific heats.

 $Ma_1 = 5$   $\delta \qquad 25^{\circ}$   $10^{\circ}$ 

**Properties** The specific heat ratio of air is k = 1.4.

**Analysis** On the basis of Assumption #2, the deflection angle is determined to be  $\theta \approx \delta = 25^{\circ} + 10^{\circ} = 35^{\circ}$ . Then the two values of oblique shock angle  $\beta$  are determined from

$$\tan \theta = \frac{2(\text{Ma}_1^2 \sin^2 \beta - 1)/\tan \beta}{\text{Ma}_1^2 (k + \cos 2\beta) + 2} \rightarrow \tan 12^\circ = \frac{2(3.4^2 \sin^2 \beta - 1)/\tan \beta}{3.4^2 (1.4 + \cos 2\beta) + 2}$$

which is implicit in  $\beta$ . Therefore, we solve it by an iterative approach or with an equation solver such as EES. It gives  $\beta_{\text{weak}} = 49.86^{\circ}$  and  $\beta_{\text{strong}} = 77.66^{\circ}$ . Then for the case of strong oblique shock, the upstream "normal" Mach number  $Ma_{1,n}$  becomes

$$Ma_{1n} = Ma_1 \sin \beta = 5 \sin 77.66^\circ = 4.884$$

Also, the downstream normal Mach numbers Ma<sub>2,n</sub> become

$$Ma_{2,n} = \sqrt{\frac{(k-1)Ma_{1,n}^2 + 2}{2kMa_{1,n}^2 - k + 1}} = \sqrt{\frac{(1.4-1)(4.884)^2 + 2}{2(1.4)(4.884)^2 - 1.4 + 1}} = 0.4169$$

The downstream pressure and temperature are determined to be

$$P_2 = P_1 \frac{2k \text{Ma}_{1,n}^2 - k + 1}{k + 1} = (70 \text{ kPa}) \frac{2(1.4)(4.884)^2 - 1.4 + 1}{1.4 + 1} = 1940 \text{ kPa}$$

$$T_2 = T_1 \frac{P_2}{P_1} \frac{\rho_1}{\rho_2} = T_1 \frac{P_2}{P_1} \frac{2 + (k - 1) \text{Ma}_{1,n}^2}{(k + 1) \text{Ma}_{1,n}^2} = (260 \text{ K}) \frac{1940 \text{ kPia}}{70 \text{ kPa}} \frac{2 + (1.4 - 1)(4.884)^2}{(1.4 + 1)(4.884)^2} = \textbf{1450 K}$$

The downstream Mach number is determined to be

$$Ma_2 = \frac{Ma_{2,n}}{\sin(\beta - \theta)} = \frac{0.4169}{\sin(77.66^\circ - 35^\circ)} =$$
**0.615**

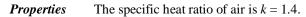
**Discussion** Note that Ma<sub>1,n</sub> is supersonic and Ma<sub>2,n</sub> and Ma<sub>2</sub> are subsonic. Also note the huge rise in temperature and pressure across the strong oblique shock, and the challenges they present for spacecraft during reentering the earth's atmosphere.

shock

## 12-75E

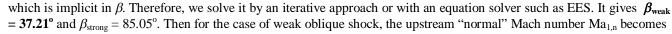
**Solution** Air flowing at a specified supersonic Mach number is forced to turn upward by a ramp, and weak oblique shock forms. The wave angle, Mach number, pressure, and temperature after the shock are to be determined.

**Assumptions** 1 The flow is steady. 2 The boundary layer on the wedge is very thin. 3 Air is an ideal gas with constant specific heats.



**Analysis** On the basis of Assumption #2, we take the deflection angle as equal to the ramp, i.e.,  $\theta \approx \delta = 8^{\circ}$ . Then the two values of oblique shock angle  $\beta$  are determined from

$$\tan \theta = \frac{2(\text{Ma}_1^2 \sin^2 \beta - 1)/\tan \beta}{\text{Ma}_1^2 (k + \cos 2\beta) + 2} \rightarrow \tan 8^\circ = \frac{2(2^2 \sin^2 \beta - 1)/\tan \beta}{2^2 (1.4 + \cos 2\beta) + 2}$$



$$Ma_{1.n} = Ma_1 \sin \beta = 2 \sin 37.21^\circ = 1.209$$

Also, the downstream normal Mach numbers Ma2,n become

$$Ma_{2,n} = \sqrt{\frac{(k-1)Ma_{1,n}^2 + 2}{2kMa_{1,n}^2 - k + 1}} = \sqrt{\frac{(1.4-1)(1.209)^2 + 2}{2(1.4)(1.209)^2 - 1.4 + 1}} = 0.8363$$

The downstream pressure and temperature are determined to be

$$P_{2} = P_{1} \frac{2k \operatorname{Ma}_{1,n}^{2} - k + 1}{k + 1} = (12 \operatorname{psia}) \frac{2(1.4)(1.209)^{2} - 1.4 + 1}{1.4 + 1} = \mathbf{18.5 \operatorname{psia}}$$

$$T_{2} = T_{1} \frac{P_{2}}{P_{1}} \frac{\rho_{1}}{\rho_{2}} = T_{1} \frac{P_{2}}{P_{1}} \frac{2 + (k - 1)\operatorname{Ma}_{1,n}^{2}}{(k + 1)\operatorname{Ma}_{1,n}^{2}} = (490 \operatorname{R}) \frac{18.5 \operatorname{psia}}{12 \operatorname{psia}} \frac{2 + (1.4 - 1)(1.209)^{2}}{(1.4 + 1)(1.209)^{2}} = \mathbf{556 \operatorname{R}}$$

The downstream Mach number is determined to be

$$Ma_2 = \frac{Ma_{2,n}}{\sin(\beta - \theta)} = \frac{0.8363}{\sin(37.21^\circ - 8^\circ)} = 1.71$$

**Discussion** Note that  $Ma_{1,n}$  is supersonic and  $Ma_{2,n}$  is subsonic. However,  $Ma_2$  is *supersonic* across the weak oblique shock (it is *subsonic* across the strong oblique shock).

### 12-76E

**Solution** Air flowing at a specified supersonic Mach number is forced to undergo a compression turn (an oblique shock)., The Mach number, pressure, and temperature downstream of the oblique shock are to be determined.

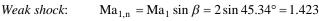
Assumptions 1 The flow is steady. 2 The boundary layer on the wedge is very thin. 3 Air is an ideal gas with constant specific heats.

**Properties** The specific heat ratio of air is k = 1.4.

**Analysis** On the basis of Assumption #2, we take the deflection angle as equal to the wedge half-angle, i.e.,  $\theta \approx \delta = 15^{\circ}$ . Then the two values of oblique shock angle  $\beta$  are determined from

$$\tan \theta = \frac{2(\text{Ma}_1^2 \sin^2 \beta - 1) / \tan \beta}{\text{Ma}_1^2 (k + \cos 2\beta) + 2} \quad \rightarrow \quad \tan 15^\circ = \frac{2(2^2 \sin^2 \beta - 1) / \tan \beta}{2^2 (1.4 + \cos 2\beta) + 2}$$

which is implicit in  $\beta$ . Therefore, we solve it by an iterative approach or with an equation solver such as EES. It gives  $\beta_{\text{weak}} = 45.34^{\circ}$  and  $\beta_{\text{strong}} = 79.83^{\circ}$ . Then the upstream "normal" Mach number Ma<sub>1,n</sub> becomes



Strong shock: 
$$Ma_{1.n} = Ma_1 \sin \beta = 2 \sin 79.83^\circ = 1.969$$

Also, the downstream normal Mach numbers Ma<sub>2,n</sub> become

Weak shock: 
$$\operatorname{Ma}_{2,n} = \sqrt{\frac{(k-1)\operatorname{Ma}_{1,n}^2 + 2}{2k\operatorname{Ma}_{1,n}^2 - k + 1}} = \sqrt{\frac{(1.4-1)(1.423)^2 + 2}{2(1.4)(1.423)^2 - 1.4 + 1}} = 0.7304$$

Strong shock: 
$$\operatorname{Ma}_{2,n} = \sqrt{\frac{(k-1)\operatorname{Ma}_{1,n}^2 + 2}{2k\operatorname{Ma}_{1,n}^2 - k + 1}} = \sqrt{\frac{(1.4-1)(1.969)^2 + 2}{2(1.4)(1.969)^2 - 1.4 + 1}} = 0.5828$$

The downstream pressure and temperature for each case are determined to be

Weak shock: 
$$P_2 = P_1 \frac{2k \text{Ma}_{1,n}^2 - k + 1}{k + 1} = (6 \text{ psia}) \frac{2(1.4)(1.423)^2 - 1.4 + 1}{1.4 + 1} = 17.57 \cong 17.6 \text{ psia}$$

$$T_2 = T_1 \frac{P_2}{P_1} \frac{\rho_1}{\rho_2} = T_1 \frac{P_2}{P_1} \frac{2 + (k - 1) \text{Ma}_{1,n}^2}{(k + 1) \text{Ma}_{1,n}^2} = (480 \text{ R}) \frac{17.57 \text{ psia}}{8 \text{ psia}} \frac{2 + (1.4 - 1)(1.423)^2}{(1.4 + 1)(1.423)^2} = 609.5 \text{ R} \approx 610 \text{ R}$$

Strong shock: 
$$P_2 = P_1 \frac{2k \text{Ma}_{1,n}^2 - k + 1}{k + 1} = (8 \text{ psia}) \frac{2(1.4)(1.969)^2 - 1.4 + 1}{1.4 + 1} = 34.85 \cong 34.9 \text{ psia}$$

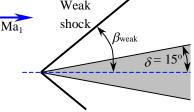
$$T_2 = T_1 \frac{P_2}{P_1} \frac{\rho_1}{\rho_2} = T_1 \frac{P_2}{P_1} \frac{2 + (k - 1)\text{Ma}_{1,n}^2}{(k + 1)\text{Ma}_{1,n}^2} = (480 \text{ R}) \frac{34.85 \text{ psia}}{8 \text{ psia}} \frac{2 + (1.4 - 1)(1.969)^2}{(1.4 + 1)(1.969)^2} = 797.9 \text{ R} \cong \textbf{798 R}$$

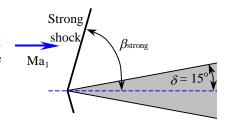
The downstream Mach number is determined to be

Weak shock: 
$$\operatorname{Ma}_2 = \frac{\operatorname{Ma}_{2,n}}{\sin(\beta - \theta)} = \frac{0.7304}{\sin(45.34^\circ - 15^\circ)} = 1.45$$

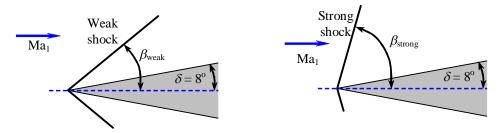
Strong shock: 
$$\operatorname{Ma}_{2} = \frac{\operatorname{Ma}_{2,n}}{\sin(\beta - \theta)} = \frac{0.5828}{\sin(79.83^{\circ} - 15^{\circ})} = \mathbf{0.644}$$

**Discussion** Note that the change in Mach number, pressure, temperature across the *strong shock* are much greater than the changes across the *weak shock*, as expected. For both the weak and strong oblique shock cases,  $Ma_{1,n}$  is supersonic and  $Ma_{2,n}$  is subsonic. However,  $Ma_2$  is *supersonic* across the weak oblique shock, but *subsonic* across the strong oblique shock.





**Solution** Air flowing at a specified supersonic Mach number impinges on a two-dimensional wedge, The shock angle, Mach number, and pressure downstream of the weak and strong oblique shock formed by a wedge are to be determined.



Assumptions 1 The flow is steady. 2 The boundary layer on the wedge is very thin. 3 Air is an ideal gas with constant specific heats.

**Properties** The specific heat ratio of air is k = 1.4.

**Analysis** On the basis of Assumption #2, we take the deflection angle as equal to the wedge half-angle, i.e.,  $\theta \approx \delta = 8^{\circ}$ . Then the two values of oblique shock angle  $\beta$  are determined from

$$\tan \theta = \frac{2(\text{Ma}_1^2 \sin^2 \beta - 1)/\tan \beta}{\text{Ma}_1^2 (k + \cos 2\beta) + 2} \rightarrow \tan 8^\circ = \frac{2(3.4^2 \sin^2 \beta - 1)/\tan \beta}{3.4^2 (1.4 + \cos 2\beta) + 2}$$

which is implicit in  $\beta$ . Therefore, we solve it by an iterative approach or with an equation solver such as EES. It gives  $\beta_{\text{weak}} = 23.15^{\circ}$  and  $\beta_{\text{strong}} = 87.45^{\circ}$ . Then the upstream "normal" Mach number  $Ma_{1,n}$  becomes

Weak shock:  $Ma_{1,n} = Ma_1 \sin \beta = 3.4 \sin 23.15^{\circ} = 1.336$ 

Strong shock:  $Ma_{1n} = Ma_1 \sin \beta = 3.4 \sin 87.45^{\circ} = 3.397$ 

Also, the downstream normal Mach numbers Ma2,n become

Weak shock: 
$$\operatorname{Ma}_{2,n} = \sqrt{\frac{(k-1)\operatorname{Ma}_{1,n}^2 + 2}{2k\operatorname{Ma}_{1,n}^2 - k + 1}} = \sqrt{\frac{(1.4-1)(1.336)^2 + 2}{2(1.4)(1.336)^2 - 1.4 + 1}} = 0.7681$$

Strong shock: 
$$\operatorname{Ma}_{2,n} = \sqrt{\frac{(k-1)\operatorname{Ma}_{1,n}^2 + 2}{2k\operatorname{Ma}_{1,n}^2 - k + 1}} = \sqrt{\frac{(1.4-1)(3.397)^2 + 2}{2(1.4)(3.397)^2 - 1.4 + 1}} = 0.4553$$

The downstream pressure for each case is determined to be

Weak shock: 
$$P_2 = P_1 \frac{2k \text{Ma}_{1,n}^2 - k + 1}{k + 1} = (60 \text{ kPa}) \frac{2(1.4)(1.336)^2 - 1.4 + 1}{1.4 + 1} = 115.0 \text{ kPa}$$

Strong shock: 
$$P_2 = P_1 \frac{2k\text{Ma}_{1,n}^2 - k + 1}{k + 1} = (60 \text{ kPa}) \frac{2(1.4)(3.397)^2 - 1.4 + 1}{1.4 + 1} = 797.6 \text{ kPa}$$

The downstream Mach number is determined to be

Weak shock: 
$$\operatorname{Ma}_{2} = \frac{\operatorname{Ma}_{2,n}}{\sin(\beta - \theta)} = \frac{0.7681}{\sin(23.15^{\circ} - 8^{\circ})} = 2.94$$

Strong shock: 
$$\operatorname{Ma}_2 = \frac{\operatorname{Ma}_{2,n}}{\sin(\beta - \theta)} = \frac{0.4553}{\sin(87.45^\circ - 8^\circ)} = \mathbf{0.463}$$

**Discussion** Note that the change in Mach number and pressure across the *strong shock* are much greater than the changes across the *weak shock*, as expected. For both the weak and strong oblique shock cases,  $Ma_{1,n}$  is supersonic and  $Ma_{2,n}$  is subsonic. However,  $Ma_2$  is *supersonic* across the weak oblique shock, but *subsonic* across the strong oblique shock.

**Solution** Air flowing through a nozzle experiences a normal shock. The effect of the shock wave on various properties is to be determined. Analysis is to be repeated for helium under the same conditions.

**Assumptions** 1 Air and helium are ideal gases with constant specific heats. 2 Flow through the nozzle is steady, one-dimensional, and isentropic before the shock occurs.

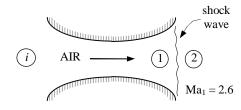
**Properties** The properties of air are k = 1.4 and R = 0.287 kJ/kg·K, and the properties of helium are k = 1.667 and R = 2.0769 kJ/kg·K.

Analysis The air properties upstream the shock are

$$Ma_1 = 2.6$$
,  $P_1 = 58$  kPa, and  $T_1 = 270$  K

Fluid properties after the shock (denoted by subscript 2) are related to those before the shock through the functions in Table A-14. For  $Ma_1 = 2.6$ ,

Ma<sub>2</sub> = **0.5039**, 
$$\frac{P_{02}}{P_1}$$
 = 9.1813,  $\frac{P_2}{P_1}$  = 7.7200, and  $\frac{T_2}{T_1}$  = 2.2383



We obtained these values using analytical relations in Table 14. Then the stagnation pressure  $P_{02}$ , static pressure  $P_2$ , and static temperature  $T_2$ , are determined to be

$$P_{02} = 9.1813P_1 = (9.1813)(58 \text{ kPa}) = 532.5 \text{ kPa}$$

$$P_2 = 7.7200P_1 = (7.7200)(58 \text{ kPa}) = 447.8 \text{ kPa}$$

$$T_2 = 2.9220T_1 = (2.2383)(270 \text{ K}) =$$
**604.3 K**

The air velocity after the shock can be determined from  $V_2 = \text{Ma}_2 c_2$ , where  $c_2$  is the speed of sound at the exit conditions after the shock,

$$V_2 = \text{Ma}_2 c_2 = \text{Ma}_2 \sqrt{kRT_2} = (0.5039) \sqrt{(1.4)(0.287 \text{ kJ/kg} \cdot \text{K})(604.3 \text{ K}) \left(\frac{1000 \text{ m}^2 / \text{s}^2}{1 \text{ kJ/kg}}\right)} = \textbf{248.3 m/s}$$

We now repeat the analysis for helium. This time we cannot use the tabulated values in Table A-14 since k is not 1.4. Therefore, we have to calculate the desired quantities using the analytical relations,

$$Ma_2 = \left(\frac{Ma_1^2 + 2/(k-1)}{2Ma_1^2k/(k-1) - 1}\right)^{1/2} = \left(\frac{2.6^2 + 2/(1.667 - 1)}{2 \times 2.6^2 \times 1.667/(1.667 - 1) - 1}\right)^{1/2} = \mathbf{0.5455}$$

$$\frac{P_2}{P_1} = \frac{1 + k \text{Ma}_1^2}{1 + k \text{Ma}_2^2} = \frac{1 + 1.667 \times 2.6^2}{1 + 1.667 \times 0.5455^2} = 8.2009$$

$$\frac{T_2}{T_1} = \frac{1 + \text{Ma}_1^2(k-1)/2}{1 + \text{Ma}_2^2(k-1)/2} = \frac{1 + 2.6^2(1.667 - 1)/2}{1 + 0.5455^2(1.667 - 1)/2} = 2.9606$$

$$\frac{P_{02}}{P_1} = \left(\frac{1 + k \text{Ma}_1^2}{1 + k \text{Ma}_2^2}\right) \left(1 + (k - 1) \text{Ma}_2^2 / 2\right)^{k/(k-1)}$$

$$= \left(\frac{1 + 1.667 \times 2.6^2}{1 + 1.667 \times 0.5455^2}\right) \left(1 + (1.667 - 1) \times 0.5455^2 / 2\right)^{1.667 / 0.667} = 10.389$$

Thus,  $P_{02} = 10.389P_1 = (10.389)(58 \text{ kPa}) = 602.5 \text{ kPa}$ 

$$P_2 = 8.2009P_1 = (8.2009)(58 \text{ kPa}) = 475.7 \text{ kPa}$$

$$T_2 = 2.9606T_1 = (2.9606)(270 \text{ K}) = 799.4 \text{ K}$$

$$V_2 = \text{Ma}_2 c_2 = \text{Ma}_2 \sqrt{kRT_2} = (0.5455) \sqrt{(1.667)(2.0769 \, \text{kJ/kg} \cdot \text{K})(799.4 \, \text{K}) \left(\frac{1000 \, \text{m}^2 \, / \, \text{s}^2}{1 \, \text{kJ/kg}}\right)} = \textbf{907.5 m/s}$$

**Discussion** The velocity and Mach number are higher for helium than for air due to the different values of k and R.

**Solution** Air flowing through a nozzle experiences a normal shock. The entropy change of air across the normal shock wave is to be determined.

**Assumptions** 1 Air and helium are ideal gases with constant specific heats. 2 Flow through the nozzle is steady, one-dimensional, and isentropic before the shock occurs.

**Properties** The properties of air are  $R = 0.287 \text{ kJ/kg} \cdot \text{K}$  and  $c_p = 1.005 \text{ kJ/kg} \cdot \text{K}$ , and the properties of helium are  $R = 2.0769 \text{ kJ/kg} \cdot \text{K}$  and  $c_p = 5.1926 \text{ kJ/kg} \cdot \text{K}$ .

Analysis For air, the entropy change across the shock is determined to be

$$s_2 - s_1 = c_p \ln \frac{T_2}{T_1} - R \ln \frac{P_2}{P_1} = (1.005 \text{ kJ/kg} \cdot \text{K}) \ln(2.2383) - (0.287 \text{ kJ/kg} \cdot \text{K}) \ln(7.7200) = \textbf{0.223 kJ/kg} \cdot \textbf{K}$$

For helium, the entropy change across the shock is determined to be

$$s_2 - s_1 = c_p \ln \frac{T_2}{T_1} - R \ln \frac{P_2}{P_1} = (5.1926 \text{ kJ/kg} \cdot \text{K}) \ln(2.9606) - (2.0769 \text{ kJ/kg} \cdot \text{K}) \ln(8.2009) = \textbf{1.27 kJ/kg} \cdot \textbf{K}$$

**Discussion** Note that shock wave is a highly dissipative process, and the entropy generation is large during shock waves.

## **Duct Flow with Heat Transfer and Negligible Friction (Rayleigh Flow)**

#### 12-80C

**Solution** We are to discuss the effect of heating on the flow velocity in subsonic Rayleigh flow.

Analysis Heating the fluid increases the flow velocity in subsonic Rayleigh flow, but decreases the flow velocity in supersonic Rayleigh flow.

**Discussion** These results are not necessarily intuitive, but must be true in order to satisfy the conservation laws.

#### 12-81C

**Solution** We are to discuss what the points on a *T-s* diagram of Rayleigh flow represent.

Analysis The points on the Rayleigh line represent the states that satisfy the conservation of mass, momentum, and energy equations as well as the property relations for a given state. Therefore, for a given inlet state, the fluid cannot exist at any downstream state outside the Rayleigh line on a *T-s* diagram.

**Discussion** The T-s diagram is quite useful, since any downstream state must lie on the Rayleigh line.

## 12-82C

**Solution** We are to discuss the effect of heat gain and heat loss on entropy during Rayleigh flow.

Analysis In Rayleigh flow, the effect of heat gain is to increase the entropy of the fluid, and the effect of heat loss is to decrease the entropy.

**Discussion** You should recall from thermodynamics that the entropy of a system can be lowered by removing heat.

### 12-83C

**Solution** We are to discuss how temperature and stagnation temperature change in subsonic Rayleigh flow.

Analysis In Rayleigh flow, the stagnation temperature  $T_0$  always increases with heat transfer to the fluid, but the temperature T decreases with heat transfer in the Mach number range of 0.845 < Ma < 1 for air. Therefore, the temperature in this case will decrease.

**Discussion** This at first seems counterintuitive, but if heat were *not* added, the temperature would drop even *more* if the air were accelerated isentropically from Ma = 0.92 to 0.95.

#### 12-84C

**Solution** We are to discuss the characteristic aspect of Rayleigh flow, and its main assumptions.

**Analysis** The characteristic aspect of Rayleigh flow is **its involvement of heat transfer**. The main assumptions associated with Rayleigh flow are: the flow is **steady**, **one-dimensional**, and **frictionless** through a constant-area duct, and the fluid is an **ideal gas with constant specific heats**.

**Discussion** Of course, there is no such thing as frictionless flow. It is better to say that frictional effects are negligible compared to the heating effects.

### 12-85C

**Solution** We are to examine the Mach number at the end of a choked duct in Rayleigh flow when more heat is added.

**Analysis** The flow is choked, and thus the flow at the duct exit **remains sonic**.

**Discussion** There is no mechanism for the flow to become supersonic in this case.

Solution Argon flowing at subsonic velocity in a constant-diameter duct is accelerated by heating. The highest rate of heat transfer without reducing the mass flow rate is to be determined.

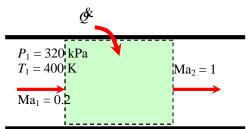
Assumptions 1 The assumptions associated with Rayleigh flow (i.e., steady one-dimensional flow of an ideal gas with constant properties through a constant cross-sectional area duct with negligible frictional effects) are valid. 2 Mass flow rate remains constant.

We take the properties of argon to be k = 1.667,  $c_p =$ **Properties**  $0.5203 \text{ kJ/kg} \cdot \text{K}$ , and  $R = 0.2081 \text{ kJ/kg} \cdot \text{K}$ .

Heat transfer stops when the flow is choked, and thus  $Ma_2 = V_2/c_2 = 1$ . The inlet stagnation temperature is

numbers are

$$T_{01} = T_1 \left( 1 + \frac{k-1}{2} \operatorname{Ma}_1^2 \right) = (400 \, \text{K}) \left( 1 + \frac{1.667 \cdot 1}{2} \cdot 0.2^2 \right) = 405.3 \, \text{K}$$
The Rayleigh flow functions corresponding to the inlet and exit Mach



$$T_{02}/T_0^* = 1$$
 (since Ma<sub>2</sub> = 1)

$$\frac{T_{01}}{T_0^*} = \frac{(k+1)\mathrm{Ma}_1^2[2 + (k-1)\mathrm{Ma}_1^2]}{(1+k\mathrm{Ma}_1^2)^2} = \frac{(1.667+1)0.2^2[2 + (1.667-1)0.2^2]}{(1+1.667\times0.2^2)^2} = 0.1900 \text{ Therefore,}$$

$$\frac{T_{02}}{T_{01}} = \frac{T_{02} / T_0^*}{T_{01} / T_0^*} = \frac{1}{0.1900} \rightarrow T_{02} = T_{01} / 0.1900 = (405.3 \text{ K}) / 0.1900 = 2133 \text{ K}$$

Then the rate of heat transfer becomes

$$\mathcal{Q}^{k} = n k_{\text{air}} c_p (T_{02} - T_{01}) = (1.2 \text{ kg/s})(0.5203 \text{ kJ/kg} \cdot \text{K})(2133 - 400) \text{ K} = 1080 \text{ kW}$$

It can also be shown that  $T_2 = 1600$  K, which is the highest thermodynamic temperature that can be attained Discussion under stated conditions. If more heat is transferred, the additional temperature rise will cause the mass flow rate to decrease. Also, in the solution of this problem, we cannot use the values of Table A-15 since they are based on k = 1.4.

 $P_2 = 270 \text{ kPa}$ 

V<sub>2</sub>⊨ 680 m/s

 $Ma_2 = 1$ 

q = 67 kJ/kg

 $P_1$ 

 $T_1$ 

 $Ma_1$ 

### 12-87

**Solution** Air is heated in a duct during subsonic flow until it is choked. For specified pressure and velocity at the exit, the temperature, pressure, and velocity at the inlet are to be determined.

**Assumptions** The assumptions associated with Rayleigh flow (i.e., steady one-dimensional flow of an ideal gas with constant properties through a constant cross-sectional area duct with negligible frictional effects) are valid.

**Properties** We take the properties of air to be k = 1.4,  $c_p = 1.005$  kJ/kg·K, and R = 0.287 kJ/kg·K.

Analysis Noting that sonic conditions exist at the exit, the exit temperature is

$$c_2 = V_2/\text{Ma}_2 = (680 \text{ m/s})/1 = 680 \text{ m/s}$$

$$c_2 = \sqrt{kRT_2} \rightarrow \sqrt{(1.4)(0.287 \text{ kJ/kg} \cdot \text{K})T_2 \left(\frac{1000 \text{ m}^2/\text{s}^2}{1 \text{ kJ/kg}}\right)} = 680 \text{ m/s}$$

It gives  $T_2 = 1151$  K. Then the exit stagnation temperature becomes

$$T_{02} = T_2 + \frac{V_2^2}{2c_p} = 1151 \,\text{K} + \frac{(680 \,\text{m/s})^2}{2 \times 1.005 \,\text{kJ/kg} \cdot \text{K}} \left(\frac{1 \,\text{kJ/kg}}{1000 \,\text{m}^2/\text{s}^2}\right) = 1381 \,\text{K}$$

The inlet stagnation temperature is, from the energy equation  $q = c_p (T_{02} - T_{01})$ ,

$$T_{01} = T_{02} - \frac{q}{c_n} = 1381 \text{ K} - \frac{67 \text{ kJ/kg}}{1.005 \text{ kJ/kg} \cdot \text{K}} = 1314 \text{ K}$$

The maximum value of stagnation temperature  $T_0^*$  occurs at Ma = 1, and its value in this case is  $T_{02}$  since the flow is choked. Therefore,  $T_0^* = T_{02} = 1381$  K. Then the stagnation temperature ratio at the inlet, and the Mach number corresponding to it are, from Table A-15,

$$\frac{T_{01}}{T_0^*} = \frac{1314 \text{ K}}{1381 \text{ K}} = 0.9516$$
  $\rightarrow$   $Ma_1 = 0.7779 \cong 0.778$ 

The Rayleigh flow functions corresponding to the inlet and exit Mach numbers are (Table A-15):

Ma<sub>1</sub> = 0.7779: 
$$T_1/T^* = 1.018$$
,  $P_1/P^* = 1.301$ ,  $V_1/V^* = 0.7852$   
Ma<sub>2</sub> = 1:  $T_2/T^* = 1$ ,  $P_2/P^* = 1$ ,  $V_2/V^* = 1$ 

Then the inlet temperature, pressure, and velocity are determined to be

$$\frac{T_2}{T_1} = \frac{T_2 / T^*}{T_1 / T^*} = \frac{1}{1.018}$$
  $\rightarrow T_1 = 1.018T_2 = 1.018(1151 \, \text{K}) = 1172 \, \text{K}$ 

$$\frac{P_2}{P_1} = \frac{P_2 / P^*}{P_1 / P^*} = \frac{1}{1.301}$$
  $\rightarrow P_1 = 1.301P_2 = 1.301(270 \, \text{kPa}) = 351.3 \, \text{kPa}$ 

$$\frac{V_2}{V_1} = \frac{V_2 / V^*}{V_1 / V^*} = \frac{1}{0.7852}$$
  $\rightarrow V_1 = 0.7852V_2 = 0.7852(680 \, \text{m/s}) = 533.9 \, \text{m/s}$ 

**Discussion** Note that the temperature and pressure decreases with heating during this subsonic Rayleigh flow while velocity increases. This problem can also be solved using appropriate relations instead of tabulated values, which can likewise be coded for convenient computer solutions.

**Solution** Air enters the combustion chamber of a gas turbine at a subsonic velocity. For a specified rate of heat transfer, the Mach number at the exit and the loss in stagnation pressure to be determined.

**Assumptions** 1 The assumptions associated with Rayleigh flow (i.e., steady one-dimensional flow of an ideal gas with constant properties through a constant cross-sectional area duct with negligible frictional effects) are valid. 2 The cross-sectional area of the combustion chamber is constant. 3 The increase in mass flow rate due to fuel injection is disregarded.

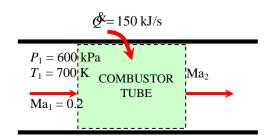
**Properties** We take the properties of air to be k = 1.4,  $c_p = 1.005$  kJ/kg·K, and R = 0.287 kJ/kg·K.

Analysis The inlet stagnation temperature and pressure are

$$T_{01} = T_1 \left( 1 + \frac{k-1}{2} \operatorname{Ma}_1^2 \right) = (700 \,\mathrm{K}) \left( 1 + \frac{1.4 - 1}{2} \,0.2^2 \right) = 705.6 \,\mathrm{K}$$

$$R_1 = R \left( 1 + \frac{k-1}{2} \operatorname{Ma}_2^2 \right)^{k/(k-1)} = (600 \,\mathrm{kPg}) \left( 1 + \frac{1.4 - 1}{2} \,0.2^2 \right)^{1.4/0.6}$$

$$P_{01} = P_1 \left( 1 + \frac{k - 1}{2} \text{ Ma}_1^2 \right)^{k/(k - 1)} = \left( 600 \text{ kPa} \right) \left( 1 + \frac{1.4 - 1}{2} \cdot 0.2^2 \right)^{1.4/0.4}$$
  
= 617.0 kPa



The exit stagnation temperature is determined from

$$Q^{k} = n^{k}_{kir} c_{p} (T_{02} - T_{01}) \rightarrow 150 \text{ kJ/s} = (0.3 \text{ kg/s})(1.005 \text{ kJ/kg} \cdot \text{K})(T_{02} - 705.6) \text{ K}$$

It gives

$$T_{02} = 1203 \text{ K}$$

At  $Ma_1 = 0.2$  we read from  $T_{01}/T_0^* = 0.1736$  (Table A-15). Therefore,

$$T_0^* = \frac{T_{01}}{0.1736} = \frac{705.6 \text{ K}}{0.1736} = 4064.5 \text{ K}$$

Then the stagnation temperature ratio at the exit and the Mach number corresponding to it are (Table A-15)

$$\frac{T_{02}}{T_0^*} = \frac{1203 \text{ K}}{4064.5 \text{ K}} = 0.2960$$
  $\rightarrow$   $Ma_2 = 0.2706 \cong 0.271$ 

Also,

$$Ma_1 = 0.2$$
  $\rightarrow P_{01}/P_0^* = 1.2346$   
 $Ma_2 = 0.2706$   $\rightarrow P_{02}/P_0^* = 1.2091$ 

Then the stagnation pressure at the exit and the pressure drop become

$$\frac{P_{02}}{P_{01}} = \frac{P_{02} / P_0^*}{P_{01} / P_0^*} = \frac{1.2091}{1.2346} = 0.9794 \quad \rightarrow \quad P_{02} = 0.9794 P_{01} = 0.9794 (617 \text{ kPa}) = 604.3 \text{ kPa}$$

and

$$\Delta P_0 = P_{01} - P_{02} = 617.0 - 604.3 =$$
**12.7 kPa**

**Discussion** This problem can also be solved using appropriate relations instead of tabulated values, which can likewise be coded for convenient computer solutions.

**Solution** Air enters the combustion chamber of a gas turbine at a subsonic velocity. For a specified rate of heat transfer, the Mach number at the exit and the loss in stagnation pressure to be determined.

**Assumptions** 1 The assumptions associated with Rayleigh flow (i.e., steady one-dimensional flow of an ideal gas with constant properties through a constant cross-sectional area duct with negligible frictional effects) are valid. 2 The cross-sectional area of the combustion chamber is constant. 3 The increase in mass flow rate due to fuel injection is disregarded.

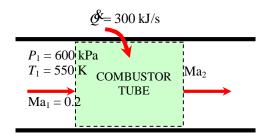
**Properties** We take the properties of air to be k = 1.4,  $c_p = 1.005$  kJ/kg·K, and R = 0.287 kJ/kg·K.

**Analysis** The inlet stagnation temperature and pressure are

$$T_{01} = T_1 \left( 1 + \frac{k - 1}{2} \operatorname{Ma}_1^2 \right) = (700 \text{ K}) \left( 1 + \frac{1.4 - 1}{2} 0.2^2 \right) = 705.6 \text{ K}$$

$$P_{01} = P_1 \left( 1 + \frac{k - 1}{2} \operatorname{Ma}_1^2 \right)^{k/(k - 1)} = \left( 600 \text{ kPa} \right) \left( 1 + \frac{1.4 - 1}{2} 0.2^2 \right)^{1.4/0.4}$$

= 617.0 kPa



The exit stagnation temperature is determined from

$$\mathcal{O}^{k} = n k_{\text{air}} c_p (T_{02} - T_{01}) \rightarrow 300 \text{ kJ/s} = (0.3 \text{ kg/s})(1.005 \text{ kJ/kg} \cdot \text{K})(T_{02} - 705.6) \text{ K}$$

It gives

$$T_{02} = 1701 \text{ K}$$

At Ma<sub>1</sub> = 0.2 we read from  $T_{01}/T_0^* = 0.1736$  (Table A-15). Therefore,

$$T_0^* = \frac{T_{01}}{0.1736} = \frac{705.6 \text{ K}}{0.1736} = 4064.5 \text{ K}$$

Then the stagnation temperature ratio at the exit and the Mach number corresponding to it are (Table A-15)

$$\frac{T_{02}}{T_0^*} = \frac{1701 \text{ K}}{4064.5 \text{ K}} = 0.4185$$
  $\rightarrow$   $Ma_2 = 0.3393 \cong$ **0.339**

Also,

$$Ma_1 = 0.2$$
  $\rightarrow P_{01}/P_0^* = 1.2346$   
 $Ma_2 = 0.3393$   $\rightarrow P_{02}/P_0^* = 1.1820$ 

Then the stagnation pressure at the exit and the pressure drop become

$$\frac{P_{02}}{P_{01}} = \frac{P_{02} / P_0^*}{P_{01} / P_0^*} = \frac{1.1820}{1.2346} = 0.9574 \quad \rightarrow \quad P_{02} = 0.9574 P_{01} = 0.9574 (617 \text{ kPa}) = 590.7 \text{ kPa}$$

and

$$\Delta P_0 = P_{01} - P_{02} = 617.0 - 590.7 =$$
**26.3 kPa**

**Discussion** This problem can also be solved using appropriate relations instead of tabulated values, which can likewise be coded for convenient computer solutions.

### 12-90E

**Solution** Air flowing with a subsonic velocity in a round duct is accelerated by heating until the flow is choked at the exit. The rate of heat transfer and the pressure drop are to be determined.

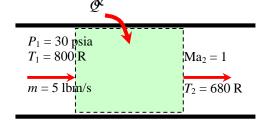
**Assumptions** 1 The assumptions associated with Rayleigh flow (i.e., steady one-dimensional flow of an ideal gas with constant properties through a constant cross-sectional area duct with negligible frictional effects) are valid. 2 The flow is choked at the duct exit. 3 Mass flow rate remains constant.

**Properties** We take the properties of air to be k = 1.4,  $c_p = 0.2400$  Btu/lbm·R, and R = 0.06855 Btu/lbm·R = 0.3704 psia·ft<sup>3</sup>/lbm·R.

Analysis The inlet density and velocity of air are

$$\rho_1 = \frac{P_1}{RT_1} = \frac{30 \text{ psia}}{(0.3704 \text{ psia} \cdot \text{ft}^3/\text{lbm} \cdot \text{R})(800 \text{ R})} = 0.1012 \text{ lbm/ft}^3$$

$$V_1 = \frac{m_{uir}^8}{\rho_1 A_{c1}} = \frac{5 \text{ lbm/s}}{(0.1012 \text{ lbm/ft}^3)[\pi (4/12 \text{ ft})^2 / 4]} = 565.9 \text{ ft/s}$$



The stagnation temperature and Mach number at the inlet are

$$T_{01} = T_1 + \frac{V_1^2}{2c_p} = 800 \text{ R} + \frac{(565.9 \text{ ft/s})^2}{2 \times 0.2400 \text{ Btu/lbm} \cdot \text{R}} \left( \frac{1 \text{ Btu/lbm}}{25,037 \text{ ft}^2/\text{s}^2} \right) = 826.7 \text{ R}$$

$$c_1 = \sqrt{kRT_1} = \sqrt{(1.4)(0.06855 \text{ Btu/lbm} \cdot \text{R})(800 \text{ R}) \left( \frac{25,037 \text{ ft}^2/\text{s}^2}{1 \text{ Btu/lbm}} \right)} = 1386 \text{ ft/s}$$

$$Ma_1 = \frac{V_1}{c_1} = \frac{565.9 \text{ ft/s}}{1386 \text{ ft/s}} = 0.4082$$

The Rayleigh flow functions corresponding to the inlet and exit Mach numbers are (Table A-15):

$$Ma_1 = 0.4082$$
:  $T_1/T^* = 0.6310$ ,  $P_1/P^* = 1.946$ ,  $T_{01}/T_0^* = 0.5434$   $Ma_2 = 1$ :  $T_2/T^* = 1$ ,  $P_2/P^* = 1$ ,  $T_{02}/T_0^* = 1$ 

Then the exit temperature, pressure, and stagnation temperature are determined to be

$$\frac{T_2}{T_1} = \frac{T_2 / T^*}{T_1 / T^*} = \frac{1}{0.6310}$$
  $\rightarrow T_2 = T_1 / 0.6310 = (800 \,\mathrm{R}) / 0.6310 = 1268 \,\mathrm{R}$ 

$$\frac{P_2}{P_1} = \frac{P_2 / P^*}{P_1 / P^*} = \frac{1}{1.946}$$
  $\rightarrow P_2 = P_1 / 2.272 = (30 \,\mathrm{psia}) / 1.946 = 15.4 \,\mathrm{psia}$ 

$$\frac{T_{02}}{T_{01}} = \frac{T_{02} / T^*}{T_{01} / T^*} = \frac{1}{0.5434}$$
  $\rightarrow T_{02} = T_{01} / 0.1743 = (826.7 \,\mathrm{R}) / 0.5434 = 1521 \,\mathrm{R}$ 

Then the rate of heat transfer and the pressure drop become

$$Q^{k} = n_{\text{Air}}^{k} c_{p} (T_{02} - T_{01}) = (5 \text{ lbm/s})(0.2400 \text{ Btu/lbm} \cdot \text{R})(1521 - 826.7) \text{ R} = 834 \text{ Btu/s}$$
  
 $\Delta P = P_{1} - P_{2} = 30 - 15.4 = 14.6 \text{ psia}$ 

**Discussion** Note that the entropy of air increases during this heating process, as expected.

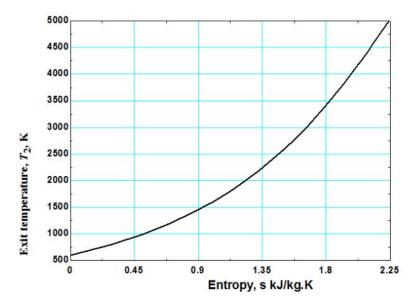


**Solution** Air flowing with a subsonic velocity in a duct. The variation of entropy with temperature is to be investigated as the exit temperature varies from 600 K to 5000 K in increments of 200 K. The results are to be tabulated and plotted.

Analysis We solve this problem using EES making use of Rayleigh functions. The EES Equations window is printed below, along with the tabulated and plotted results.

k=1.4
cp=1.005
R=0.287
P1=350
T1=600
V1=70
C1 = sqrt(k*R*T1*1000)
Ma1=V1/C1
T01=T1*(1+0.5*(k-1)*Ma1^2)
P01=P1*(1+0.5*(k-1)*Ma1^2)^(k/(k-1))
F1=1+0.5*(k-1)*Ma1^2
T01Ts=2*(k+1)*Ma1^2*F1/(1+k*Ma1^2)^2
$P01Ps = ((1+k)/(1+k*Ma1^2))*(2*F1/(k+1))^(k/(k-1))$
$T1Ts = (Ma1*((1+k)/(1+k*Ma1^2)))^2$
$P1Ps=(1+k)/(1+k*Ma1^2)$
V1Vs=Ma1^2*(1+k)/(1+k*Ma1^2)
F2=1+0.5*(k-1)*Ma2^2
T02Ts=2*(k+1)*Ma2^2*F2/(1+k*Ma2^2)^2
$P02Ps = ((1+k)/(1+k*Ma2^2))*(2*F2/(k+1))^(k/(k-1))$
$T2Ts = (Ma2*((1+k)/(1+k*Ma2^2)))^2$
$P2Ps=(1+k)/(1+k*Ma2^2)$
V2Vs=Ma2^2*(1+k)/(1+k*Ma2^2)
T02=T02Ts/T01Ts*T01
P02=P02Ps/P01Ps*P01
T2=T2Ts/T1Ts*T1
P2=P2Ps/P1Ps*P1
V2=V2Vs/V1Vs*V1
$Delta\_s = cp*ln(T2/T1)-R*ln(P2/P1)$

Exit temperature	Exit Mach	Exit entropy relative to inlet,
$T_2$ , K	number, Ma <sub>2</sub>	s₂, kJ/kg⋅K
600	0.143	0.000
800	0.166	0.292
1000	0.188	0.519
1200	0.208	0.705
1400	0.227	0.863
1600	0.245	1.001
1800	0.263	1.123
2000	0.281	1.232
2200	0.299	1.331
2400	0.316	1.423
2600	0.333	1.507
2800	0.351	1.586
3000	0.369	1.660
3200	0.387	1.729
3400	0.406	1.795
3600	0.426	1.858
3800	0.446	1.918
4000	0.467	1.975
4200	0.490	2.031
4400	0.515	2.085
4600	0.541	2.138
4800	0.571	2.190
5000	0.606	2.242



**Discussion** Note that the entropy of air increases during this heating process, as expected.

## 12-92E

**Solution** Air flowing with a subsonic velocity in a square duct is accelerated by heating until the flow is choked at the exit. The rate of heat transfer and the entropy change are to be determined.

**Assumptions** 1 The assumptions associated with Rayleigh flow (i.e., steady one-dimensional flow of an ideal gas with constant properties through a constant cross-sectional area duct with negligible frictional effects) are valid. 2 The flow is choked at the duct exit. 3 Mass flow rate remains constant.

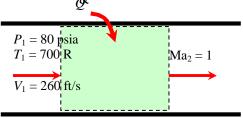
**Properties** We take the properties of air to be k = 1.4,  $c_p = 0.2400$  Btu/lbm·R, and R = 0.06855 Btu/lbm·R = 0.3704 psia·ft<sup>3</sup>/lbm·R.

Analysis The inlet density and mass flow rate of air are

$$\rho_1 = \frac{P_1}{RT_1} = \frac{80 \text{ psia}}{(0.3704 \text{ psia} \cdot \text{ft}^3/\text{lbm} \cdot \text{R})(700 \text{ R})} = 0.3085 \text{ lbm/ft}^3$$

$$m_{air}^{S} = \rho_1 A_{c1} V_1 = (0.3085 \text{ lbm/ft}^3) (6 \times 6/144 \text{ ft}^2) (260 \text{ ft/s}) = 20.06 \text{ lbm/s}$$

The stagnation temperature and Mach number at the inlet are



$$T_{01} = T_1 + \frac{V_1^2}{2c_p} = 700 \,\mathrm{R} + \frac{(260 \,\mathrm{ft/s})^2}{2 \times 0.2400 \,\mathrm{Btu/lbm \cdot R}} \left( \frac{1 \,\mathrm{Btu/lbm}}{25,037 \,\mathrm{ft}^2/\mathrm{s}^2} \right) = 705.6 \,\mathrm{R}$$

$$c_1 = \sqrt{kRT_1} = \sqrt{(1.4)(0.06855 \,\mathrm{Btu/lbm \cdot R})(700 \,\mathrm{R}) \left( \frac{25,037 \,\mathrm{ft}^2/\mathrm{s}^2}{1 \,\mathrm{Btu/lbm}} \right)} = 1297 \,\mathrm{ft/s}$$

$$Ma_1 = \frac{V_1}{c_1} = \frac{260 \,\mathrm{ft/s}}{1297 \,\mathrm{ft/s}} = 0.2005$$

The Rayleigh flow functions corresponding to the inlet and exit Mach numbers are (Table A-15):

$$Ma_1 = 0.2005$$
:  $T_1/T^* = 0.2075$ ,  $P_1/P^* = 2.272$ ,  $T_{01}/T_0^* = 0.1743$   
 $Ma_2 = 1$ :  $T_2/T^* = 1$ ,  $P_2/P^* = 1$ ,  $T_{02}/T_0^* = 1$ 

Then the exit temperature, pressure, and stagnation temperature are determined to be

$$\frac{T_2}{T_1} = \frac{T_2 / T^*}{T_1 / T^*} = \frac{1}{0.2075} \qquad \rightarrow T_2 = T_1 / 0.2075 = (700 \,\mathrm{R}) / 0.2075 = 3374 \,\mathrm{R}$$

$$\frac{P_2}{P_1} = \frac{P_2 / P^*}{P_1 / P^*} = \frac{1}{2.272} \qquad \rightarrow P_2 = P_1 / 2.272 = (80 \,\mathrm{psia}) / 2.272 = 35.2 \,\mathrm{psia}$$

$$\frac{T_{02}}{T_{01}} = \frac{T_{02} / T^*}{T_{01} / T^*} = \frac{1}{0.1743} \qquad \rightarrow T_{02} = T_{01} / 0.1743 = (705.6 \,\mathrm{R}) / 0.1743 = 4048 \,\mathrm{R}$$

Then the rate of heat transfer and entropy change become

$$\mathcal{Q} = n \mathbf{k}_{\text{air}} c_p (T_{02} - T_{01}) = (20.06 \,\text{lbm/s})(0.2400 \,\text{Btu/lbm} \cdot \text{R})(4048 - 705.6) \,\text{R} = \mathbf{16,090} \,\text{Btu/s}$$

$$\Delta s = c_p \ln \frac{T_2}{T_1} - R \ln \frac{P_2}{P_1} = (0.2400 \,\text{Btu/lbm} \cdot \text{R}) \ln \frac{3374 \,\text{R}}{700 \,\text{R}} - (0.06855 \,\text{Btu/lbm} \cdot \text{R}) \ln \frac{35.2 \,\text{psia}}{80 \,\text{psia}} = \mathbf{0.434} \,\text{Btu/lbm} \cdot \text{R}$$

**Discussion** Note that the entropy of air increases during this heating process, as expected.

**Solution** Fuel is burned in a rectangular duct with compressed air. For specified heat transfer, the exit temperature and Mach number are to be determined.

**Assumptions** The assumptions associated with Rayleigh flow (i.e., steady one-dimensional flow of an ideal gas with constant properties through a constant cross-sectional area duct with negligible frictional effects) are valid.

**Properties** We take the properties of air to be k = 1.4,  $c_p = 1.005$  kJ/kg·K, and R = 0.287 kJ/kg·K.

Analysis The stagnation temperature and Mach number at the inlet are

$$c_{1} = \sqrt{kRT_{1}} = \sqrt{(1.4)(0.287 \text{ kJ/kg} \cdot \text{K})(300 \text{ K})} \left(\frac{1000 \text{ m}^{2} / \text{s}^{2}}{1 \text{ kJ/kg}}\right) = 347.2 \text{ m/s}$$

$$Q = 55 \text{ kJ/kg}$$

$$V_{1} = \text{Ma}_{1}c_{1} = 2(347.2 \text{ m/s}) = 694.4 \text{ m/s}$$

$$T_{1} = 300 \text{ K}$$

$$T_{1} = 300 \text{ K}$$

$$T_{1} = 300 \text{ K}$$

$$T_{2}, \text{Ma}_{2}$$

$$T_{1} = 420 \text{ kPa}$$

$$T_{1} = 300 \text{ K}$$

$$T_{2} = 300 \text{ K} + \frac{(694.4 \text{ m/s})^{2}}{2 \times 1.005 \text{ kJ/kg} \cdot \text{K}} \left(\frac{1 \text{ kJ/kg}}{1000 \text{ m}^{2}/\text{s}^{2}}\right) = 539.9 \text{ K}$$

The exit stagnation temperature is, from the energy equation  $q = c_p (T_{02} - T_{01})$ ,

$$T_{02} = T_{01} + \frac{q}{c_p} = 539.9 \text{ K} + \frac{55 \text{ kJ/kg}}{1.005 \text{ kJ/kg} \cdot \text{K}} = 594.6 \text{ K}$$

The maximum value of stagnation temperature  $T_0^*$  occurs at Ma = 1, and its value can be determined from Table A-15 or from the appropriate relation. At Ma<sub>1</sub> = 2 we read  $T_{01}/T_0^* = 0.7934$ . Therefore,

$$T_0^* = \frac{T_{01}}{0.7934} = \frac{539.9 \text{ K}}{0.7934} = 680.5 \text{ K}$$

The stagnation temperature ratio at the exit and the Mach number corresponding to it are, from Table A-15,

$$\frac{T_{02}}{T_0^*} = \frac{594.6 \text{ K}}{680.5 \text{ K}} = 0.8738$$
  $\rightarrow$   $Ma_2 = 1.642 \cong 1.642$ 

Also,

$$Ma_1 = 2$$
  $\rightarrow$   $T_1/T^* = 0.5289$   
 $Ma_2 = 1.642$   $\rightarrow$   $T_2/T^* = 0.6812$ 

Then the exit temperature becomes

$$\frac{T_2}{T_1} = \frac{T_2 / T^*}{T_1 / T^*} = \frac{0.6812}{0.5289} = 1.288 \quad \rightarrow \quad T_2 = 1.288 T_1 = 1.288 (300 \text{ K}) = 386 \text{ K}$$

**Discussion** Note that the temperature increases during this supersonic Rayleigh flow with heating. This problem can also be solved using appropriate relations instead of tabulated values, which can likewise be coded for convenient computer solutions.

**Solution** Compressed air is cooled as it flows in a rectangular duct. For specified heat rejection, the exit temperature and Mach number are to be determined.

**Assumptions** The assumptions associated with Rayleigh flow (i.e., steady one-dimensional flow of an ideal gas with constant properties through a constant cross-sectional area duct with negligible frictional effects) are valid.

**Properties** We take the properties of air to be k = 1.4,  $c_p = 1.005$  kJ/kg·K, and R = 0.287 kJ/kg·K.

Analysis The stagnation temperature and Mach number at the inlet are

$$c_{1} = \sqrt{kRT_{1}} = \sqrt{(1.4)(0.287 \text{ kJ/kg} \cdot \text{K})(300 \text{ K}) \left(\frac{1000 \text{ m}^{2}/\text{s}^{2}}{1 \text{ kJ/kg}}\right)} = 347.2 \text{ m/s}$$

$$V_{1} = \text{Ma}_{1}c_{1} = 2(347.2 \text{ m/s}) = 694.4 \text{ m/s}$$

$$T_{01} = T_{1} + \frac{V_{1}^{2}}{2c_{p}} = 300 \text{ K} + \frac{(694.4 \text{ m/s})^{2}}{2 \times 1.005 \text{ kJ/kg} \cdot \text{K}} \left(\frac{1 \text{ kJ/kg}}{1000 \text{ m}^{2}/\text{s}^{2}}\right) = 539.9 \text{ K}$$

The exit stagnation temperature is, from the energy equation  $q = c_n (T_{02} - T_{01})$ ,

$$T_{02} = T_{01} + \frac{q}{c_p} = 539.9 \text{ K} + \frac{-55 \text{ kJ/kg}}{1.005 \text{ kJ/kg} \cdot \text{K}} = 485.2 \text{ K}$$

The maximum value of stagnation temperature  $T_0^*$  occurs at Ma = 1, and its value can be determined from Table A-15 or from the appropriate relation. At Ma<sub>1</sub> = 2 we read  $T_{01}/T_0^* = 0.7934$ . Therefore,

$$T_0^* = \frac{T_{01}}{0.7934} = \frac{539.9 \text{ K}}{0.7934} = 680.5 \text{ K}$$

The stagnation temperature ratio at the exit and the Mach number corresponding to it are, from Table A-15,

$$\frac{T_{02}}{T_0^*} = \frac{485.2 \text{ K}}{680.5 \text{ K}} = 0.7130$$
  $\rightarrow$  Ma<sub>2</sub> = 2.479  $\cong$  **2.48**

Also,

$$Ma_1 = 2$$
  $\rightarrow$   $T_1/T^* = 0.5289$   
 $Ma_2 = 2.479$   $\rightarrow$   $T_2/T^* = 0.3838$ 

Then the exit temperature becomes

$$\frac{T_2}{T_1} = \frac{T_2 / T^*}{T_1 / T^*} = \frac{0.3838}{0.5289} = 0.7257 \quad \rightarrow \quad T_2 = 0.7257 T_1 = 0.7257(300 \text{ K}) = \textbf{218 K}$$

**Discussion** Note that the temperature decreases and Mach number increases during this supersonic Rayleigh flow with cooling. This problem can also be solved using appropriate relations instead of tabulated values, which can likewise be coded for convenient computer solutions.

**Solution** Fuel is burned in a tubular combustion chamber with compressed air. For a specified exit Mach number, the exit temperature and the rate of fuel consumption are to be determined.

**Assumptions** 1 The assumptions associated with Rayleigh flow (i.e., steady one-dimensional flow of an ideal gas with constant properties through a constant cross-sectional area duct with negligible frictional effects) are valid. 2 Combustion is complete, and it is treated as a heat addition process, with no change in the chemical composition of flow. 3 The increase in mass flow rate due to fuel injection is disregarded.

**Properties** We take the properties of air to be k = 1.4,  $c_p = 1.005$  kJ/kg·K, and R = 0.287 kJ/kg·K.

Analysis The inlet density and mass flow rate of air are

$$\rho_1 = \frac{P_1}{RT_1} = \frac{380 \text{ kPa}}{(0.287 \text{ kJ/kgK})(450 \text{ K})} = 2.942 \text{ kg/m}^3$$

$$n_{air}^{g} = \rho_1 A_{c1} V_1 = (2.942 \text{ kg/m}^3) [\pi (0.16 \text{ m})^2 / 4] (55 \text{ m/s}) = 3.254 \text{ kg/s}$$

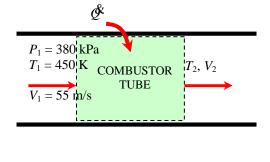
The stagnation temperature and Mach number at the inlet are

$$T_{01} = T_1 + \frac{V_1^2}{2c_p} = 450 \text{ K} + \frac{(55 \text{ m/s})^2}{2 \times 1.005 \text{ kJ/kg} \cdot \text{K}} \left(\frac{1 \text{ kJ/kg}}{1000 \text{ m}^2/\text{s}^2}\right) = 451.5 \text{ K}$$

$$c_1 = \sqrt{kRT_1} = \sqrt{(1.4)(0.287 \text{ kJ/kg} \cdot \text{K})(450 \text{ K}) \left(\frac{1000 \text{ m}^2/\text{s}^2}{1 \text{ kJ/kg}}\right)} = 425.2 \text{ m/s}$$

$$\frac{c_1 - \sqrt{kRT_1} - \sqrt{(1.4)(0.287 \text{ kJ/kg} \cdot \text{K})(430 \text{ K})}}{1 \text{ kJ/kg}} = \frac{V}{1 \text{ kJ/kg}}$$

$$Ma_1 = \frac{V_1}{c_1} = \frac{55 \text{ m/s}}{425.2 \text{ m/s}} = 0.1293$$



The Rayleigh flow functions corresponding to the inlet and exit Mach numbers are (Table A-15) (We used analytical functions):

$$Ma_1 = 0.1293$$
:  $T_1/T^* = 0.09201$ ,  $T_{01}/T^* = 0.07693$ ,  $V_1/V^* = 0.03923$   
 $Ma_2 = 0.8$ :  $T_2/T^* = 1.0255$ ,  $T_{02}/T^* = 0.9639$ ,  $V_2/V^* = 0.8101$ 

The exit temperature, stagnation temperature, and velocity are determined to be

$$\frac{T_2}{T_1} = \frac{T_2 / T^*}{T_1 / T^*} = \frac{1.0255}{0.09201} = 11.146 \rightarrow T_2 = 11.146 T_1 = 11.146 (450 \text{ K}) = 5016 \text{ K}$$

$$\frac{T_{02}}{T_{01}} = \frac{T_{02} / T^*}{T_{01} / T^*} = \frac{0.9639}{0.07693} = 12.530 \rightarrow T_{02} = 12.530 (451.5 \text{ K}) = 5658 \text{ K}$$

$$\frac{V_2}{V_1} = \frac{V_2 / V^*}{V_1 / V^*} = \frac{0.8101}{0.03923} = 20.650 \rightarrow V_2 = 20.650 V_1 = 20.650 (55 \text{ m/s}) = 1136 \text{ m/s}$$

Then the mass flow rate of the fuel is determined to be

$$q = c_p (T_{02} - T_{01}) = (1.005 \text{ kJ/kg} \cdot \text{K})(5658 - 451.5) \text{ K} = 5232 \text{ kJ/kg}$$

$$\mathcal{G} = n_{\text{Air}} q = (3.254 \text{ kg/s})(5232 \text{ kJ/kg}) = 17,024 \text{ kW}$$

$$n_{\text{fuel}} = \frac{\mathcal{G}}{HV} = \frac{17,024 \text{ kJ/s}}{39,000 \text{ kJ/kg}} = \textbf{0.4365 \text{ kg/s}}$$

**Discussion** Note that both the temperature and velocity increase during this subsonic Rayleigh flow with heating, as expected. This problem can also be solved using appropriate relations instead of tabulated values, which can likewise be coded for convenient computer solutions.

**Solution** Air flowing at a supersonic velocity in a duct is decelerated by heating. The highest temperature air can be heated by heat addition and the rate of heat transfer are to be determined.

**Assumptions** 1 The assumptions associated with Rayleigh flow (i.e., steady one-dimensional flow of an ideal gas with constant properties through a constant cross-sectional area duct with negligible frictional effects) are valid. 2 Mass flow rate remains constant.

**Properties** We take the properties of air to be k = 1.4,  $c_p = 1.005$  kJ/kg·K, and R = 0.287 kJ/kg·K.

Analysis Heat transfer will stop when the flow is choked, and thus  $Ma_2 = V_2/c_2 = 1$ . Knowing stagnation properties, the static properties are determined to be

$$T_{1} = T_{01} \left( 1 + \frac{k - 1}{2} \operatorname{Ma}_{1}^{2} \right)^{-1} = (600 \,\mathrm{K}) \left( 1 + \frac{1.4 - 1}{2} 1.8^{2} \right)^{-1} = 364.1 \,\mathrm{K}$$

$$P_{1} = P_{01} \left( 1 + \frac{k - 1}{2} \operatorname{Ma}_{1}^{2} \right)^{-k/(k - 1)} = (140 \,\mathrm{K}) \left( 1 + \frac{1.4 - 1}{2} 1.8^{2} \right)^{-1.4/0.4}$$

$$= 24.37 \,\mathrm{kPa}$$

$$P_{1} = 24.37 \,\mathrm{kPa}$$

$$P_{2} = 24.37 \,\mathrm{kPa}$$

$$P_{3} = 24.37 \,\mathrm{kPa}$$

$$P_{4} = 24.37 \,\mathrm{kPa}$$

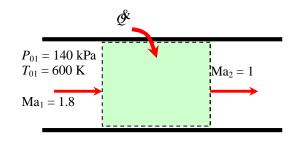
$$P_{5} = 24.37 \,\mathrm{kPa}$$

$$P_{6} = 24.37 \,\mathrm{kPa}$$

$$P_{7} = 24.37 \,\mathrm{kPa}$$

$$P_{8} = 24.37 \,\mathrm{kPa}$$

$$P_{1} = 24.37 \,\mathrm{kPa}$$



$$\rho_1 = \frac{P_1}{RT_1} = \frac{24.37 \text{ kPa}}{(0.287 \text{ kJ/kgK})(364.1 \text{ K})} = 0.2332 \text{ kg/m}^3$$

Then the inlet velocity and the mass flow rate become

$$c_1 = \sqrt{kRT_1} = \sqrt{(1.4)(0.287 \text{ kJ/kg} \cdot \text{K})(364.1 \text{ K}) \left(\frac{1000 \text{ m}^2 / \text{s}^2}{1 \text{ kJ/kg}}\right)} = 382.5 \text{ m/s}$$

$$V_1 = \text{Ma}_1 c_1 = 1.8(382.5 \text{ m/s}) = 688.5 \text{ m/s}$$

$$m_{air}^{2} = \rho_{1}A_{c1}V_{1} = (0.2332 \text{ kg/m}^{3})[\pi(0.07 \text{ m})^{2}/4](688.5 \text{ m/s}) = 0.6179 \text{ kg/s}$$

The Rayleigh flow functions corresponding to the inlet and exit Mach numbers are (Table A-15):

$$Ma_1 = 1.8$$
:  $T_1/T^* = 0.6089$ ,  $T_{01}/T_0^* = 0.8363$   
 $Ma_2 = 1$ :  $T_2/T^* = 1$ ,  $T_{02}/T_0^* = 1$ 

Then the exit temperature and stagnation temperature are determined to be

$$\frac{T_2}{T_1} = \frac{T_2 / T^*}{T_1 / T^*} = \frac{1}{0.6089} \rightarrow T_2 = T_1 / 0.6089 = (364.1 \,\mathrm{K}) / 0.6089 = \mathbf{598} \,\mathrm{K}$$

$$\frac{T_{02}}{T_{01}} = \frac{T_{02} / T_0^*}{T_{01} / T_0^*} = \frac{1}{0.8363} \rightarrow T_{02} = T_{01} / 0.8363 = (600 \,\mathrm{K}) / 0.8363 = 717.4 \,\mathrm{K} \cong \mathbf{717} \,\mathrm{K}$$

Finally, the rate of heat transfer is

$$\mathcal{Q}^{k} = n k_{air} c_p (T_{02} - T_{01}) = (0.6179 \text{ kg/s})(1.005 \text{ kJ/kg} \cdot \text{K})(717.4 - 600) \text{ K} = 72.9 \text{ kW}$$

**Discussion** Note that this is the highest temperature that can be attained under stated conditions. If more heat is transferred, the additional temperature will cause the mass flow rate to decrease. Also, once the sonic conditions are reached, the thermodynamic temperature can be increased further by cooling the fluid and reducing the velocity (see the *T-s* diagram for Rayleigh flow).

## **Adiabatic Duct Flow with Friction (Fanno Flow)**

### 12-97C

**Solution** We are to discuss the effect of friction on velocity in Fanno flow.

Analysis Friction increases the flow velocity in subsonic Fanno flow, but decreases the flow velocity in supersonic flow.

**Discussion** These results may not be intuitive, but they come from following the Fanno line, which satisfies the conservation equations.

### 12-98C

**Solution** We are to discuss the T-s diagram for Fanno flow.

Analysis The points on the Fanno line on a *T-s* diagram represent the **states that satisfy the conservation of mass, momentum, and energy equations as well as the property relations for a given inlet state**. Therefore, for a given initial state, the fluid cannot exist at any downstream state outside the Fanno line on a *T-s* diagram.

**Discussion** The *T-s* diagram is quite useful, since any downstream state must lie on the Fanno line.

#### 12-99C

**Solution** We are to discuss the effect of friction on the entropy during Fanno flow.

**Analysis** In Fanno flow, the effect of friction is **always to increase the entropy of the fluid**. Therefore Fanno flow always proceeds in the direction of increasing entropy.

**Discussion** To do otherwise would violate the second law of thermodynamics.

#### 12-100C

**Solution** We are to discuss what happens to supersonic Fanno flow, initially sonic at the exit, when the duct is extended.

Analysis The flow at the duct exit remains sonic. The mass flow rate must remain constant since upstream conditions are not affected by the added duct length.

**Discussion** The mass flow rate is fixed by the upstream stagnation conditions and the size of the throat – therefore, the mass flow rate does not change by extending the duct. However, a *shock wave* appears in the duct when it is extended.

## 12-101C

**Solution** We are to examine what happens when the Mach number of air decreases in supersonic Fanno flow.

Analysis During supersonic Fanno flow, the stagnation temperature  $T_{\theta}$  remains constant, stagnation pressure  $P_{\theta}$  decreases, and entropy s increases.

**Discussion** Friction leads to irreversible losses, which are felt as a loss of stagnation pressure and an increase of entropy. However, since the flow is adiabatic, the stagnation temperature does not change downstream.

### 12-102C

**Solution** We are to discuss the characteristic aspect of Fanno flow and its main assumptions.

**Analysis** The characteristic aspect of Fanno flow is its **consideration of friction**. The main assumptions associated with Fanno flow are: the flow is **steady**, **one-dimensional**, and **adiabatic** through a **constant-area duct**, and the fluid is an **ideal gas** with **constant specific heats**.

**Discussion** Compared to Rayleigh flow, Fanno flow accounts for friction but neglects heat transfer effects, whereas Rayleigh flow accounts for heat transfer but neglects frictional effects.

### 12-103C

**Solution** We are to discuss what happens to choked subsonic Fanno flow when the duct is extended.

Analysis The flow is choked, and thus the flow at the duct exit must remain sonic. The mass flow rate has to decrease as a result of extending the duct length in order to compensate.

**Discussion** Since there is no way for the flow to become supersonic (e.g., there is no throat), the upstream flow must adjust itself such that the flow at the exit plan remains sonic.

#### 12-104C

**Solution** We are to examine what happens when the Mach number of air increases in subsonic Fanno flow.

Analysis During subsonic Fanno flow, the stagnation temperature  $T_{\theta}$  remains constant, stagnation pressure  $P_{\theta}$  decreases, and entropy s increases.

**Discussion** Friction leads to irreversible losses, which are felt as a loss of stagnation pressure and an increase of entropy. However, since the flow is adiabatic, the stagnation temperature does not change downstream.

**Solution** Subsonic airflow in a constant cross-sectional area adiabatic duct is considered. For a specified exit Mach number, the duct length, temperature, pressure, and velocity at the duct exit are to be determined.

**Assumptions** 1 The assumptions associated with Fanno flow (i.e., steady, frictional flow of an ideal gas with constant properties through a constant cross-sectional area adiabatic duct) are valid. 2 The friction factor remains constant along the duct.

**Properties** We take the properties of air to be k = 1.4,  $c_p = 1.005$  kJ/kg·K, and R = 0.287 kJ/kg·K. The average friction factor is given to be f = 0.021.

**Analysis** 

The inlet velocity is

$$c_1 = \sqrt{kRT_1} = \sqrt{(1.4)(0.287 \text{ kJ/kg} \cdot \text{K})(550 \text{ K}) \left(\frac{1000 \text{ m}^2/\text{s}^2}{1 \text{ kJ/kg}}\right)} = 470.1 \text{ m/s}$$

$$V_1 = \text{Ma}_1 c_1 = 0.4(470.1 \,\text{m/s}) = 188.0 \,\text{m/s}$$

The Fanno flow functions corresponding to the inlet and exit Mach numbers are, from Table A-16,

$$Ma_1 = 0.4$$
:  $(fL^*/D_h)_1 = 2.3085$   $T_1/T^* = 1.1628$ ,  $P_1/P^* = 2.6958$ ,  $V_1/V^* = 0.4313$   $Ma_2 = 0.8$ :  $(fL^*/D_h)_2 = 0.0723$   $T_2/T^* = 1.0638$ ,  $P_2/P^* = 1.2893$ ,  $V_2/V^* = 0.8251$ 

Then the temperature, pressure, and velocity at the duct exit are determined to be

$$\frac{T_2}{T_1} = \frac{T_2 / T^*}{T_1 / T^*} = \frac{1.0638}{1.1628} = 0.9149 \qquad \rightarrow T_2 = 0.9149 T_1 = 0.9149 (550 \text{ K}) = \textbf{503.2 K}$$
 
$$\frac{P_2}{P_1} = \frac{P_2 / P^*}{P_1 / P^*} = \frac{1.2893}{2.6958} = 0.4783 \qquad \rightarrow P_2 = 0.4783 P_1 = 0.4783 (200 \text{ kPa}) = \textbf{95.65 kPa}$$
 
$$\frac{V_2}{V_1} = \frac{V_2 / V^*}{V_1 / V^*} = \frac{0.8251}{0.4313} = 1.9131 \qquad \rightarrow V_2 = 1.9131 V_1 = 1.9131 (188.0 \text{ m/s}) = \textbf{359.7 m/s}$$

Finally, the actual duct length is determined to be

$$L = L_1^* - L_2^* = \left(\frac{fL_1^*}{D_h} - \frac{fL_2^*}{D_h}\right) \frac{D_h}{f} = (2.3085 - 0.0723) \frac{0.12 \text{ m}}{0.021} = \mathbf{12.8 m}$$

**Discussion** Note that it takes a duct length of 12.8 m for the Mach number to increase from 0.4 to 0.8. The Mach number rises at a much higher rate as sonic conditions are approached. The maximum (or sonic) duct lengths at the inlet and exit states in this case are  $L_1^* = 13.2$  m and  $L_2^* = 0.413$  m. Therefore, the flow would reach sonic conditions if a 0.413-m long section were added to the existing duct.

**Solution** Air enters a constant-area adiabatic duct of given length at a specified state. The exit Mach number, exit velocity, and the mass flow rate are to be determined.

**Assumptions** 1 The assumptions associated with Fanno flow (i.e., steady, frictional flow of an ideal gas with constant properties through a constant cross-sectional area adiabatic duct) are valid. 2 The friction factor is constant along the duct.

**Properties** We take the properties of air to be k = 1.4,  $c_p = 1.005$  kJ/kg·K, and R = 0.287 kJ/kg·K. The friction factor is given to be f = 0.023.

**Analysis** The first thing we need to know is whether the flow is choked at the exit or not. Therefore, we first determine the inlet Mach number and the corresponding value of the function  $fL^*/D_h$ ,

$$L = 15 \text{ m}$$

$$L_2*$$

$$P_1 = 300 \text{ kPa}$$

$$T_1 = 500 \text{ K}$$

$$V_1 = 70 \text{ m/s}$$

$$L_2*$$

$$M\bar{a}*=1$$

$$T^*$$

$$P^*$$

$$V$$

$$c_1 = \sqrt{kRT_1} = \sqrt{(1.4)(0.287 \text{ kJ/kg} \cdot \text{K})(500 \text{ K}) \left(\frac{1000 \text{ m}^2/\text{s}^2}{1 \text{ kJ/kg}}\right)} = 448.2 \text{ m/s}$$

$$Ma_1 = \frac{V_1}{c_1} = \frac{70 \text{ m/s}}{448.2 \text{ m/s}} = 0.1562$$

Corresponding to this Mach number we calculate (or read) from Table A-16),  $(fL^*/D_h)_1 = 25.540$ . Also, using the actual duct length L, we have

$$\frac{fL}{D_h} = \frac{(0.023)(15 \text{ m})}{0.04 \text{ m}} = 8.625 < 25.540$$

Therefore, flow is *not* choked and exit Mach number is less than 1. Noting that  $L = L_1^* - L_2^*$ , the function  $fL^*/D_h$  at the exit state is calculated from

$$\left(\frac{fL^*}{D_h}\right)_2 = \left(\frac{fL^*}{D_h}\right)_1 - \frac{fL}{D_h} = 25.540 - 8.625 = 16.915$$

The Mach number corresponding to this value of  $fL^*/D$  is obtained from Table A-16 to be

$$Ma_2 = 0.187$$

which is the Mach number at the duct exit. The mass flow rate of air is determined from the inlet conditions to be

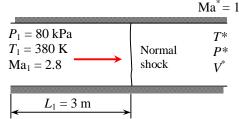
$$\rho_1 = \frac{P_1}{RT_1} = \frac{300 \text{ kPa}}{(0.287 \text{ kJ/kgK})(500 \text{ K})} \left(\frac{1 \text{ kJ}}{1 \text{ kPa} \cdot \text{m}^3}\right) = 2.091 \text{ kg/m}^3$$

$$m_{air}^{S} = \rho_1 A_{c1} V_1 = (2.091 \,\text{kg/m}^3) [\pi (0.04 \,\text{m})^2 \,/\, 4] (70 \,\text{m/s}) =$$
**0.184 kg/s**

**Discussion** It can be shown that  $L_2^* = 29.4$  m, indicating that it takes a duct length of 15 m for the Mach number to increase from 0.156 to 0.187, but only 29.4 m to increase from 0.187 to 1. Therefore, the Mach number rises at a much higher rate as sonic conditions are approached.

**Solution** Air enters a constant-area adiabatic duct at a specified state, and undergoes a normal shock at a specified location. The exit velocity, temperature, and pressure are to be determined.

**Assumptions** 1 The assumptions associated with Fanno flow (i.e., steady, frictional flow of an ideal gas with constant properties through a constant cross-sectional area adiabatic duct) are valid. 2 The friction factor is constant along the duct.



**Properties** We take the properties of air to be k = 1.4,  $c_p = 1.005$  kJ/kg·K, and R = 0.287 kJ/kg·K. The friction factor is given to be f = 0.007.

Analysis The Fanno flow functions corresponding to the inlet Mach number of 2.8 are, from Table A-16,

$$Ma_1 = 2.8$$
:  $(fL^*/D_h)_1 = 0.4898$   $T_1/T^* = 0.4673$ ,  $P_1/P^* = 0.2441$ 

First we check to make sure that the flow everywhere upstream the shock is supersonic. The required duct length from the inlet  $L_1^*$  for the flow to reach sonic conditions is

$$L_1^* = 0.4898 \frac{D}{f} = 0.4898 \frac{0.05 \text{ m}}{0.007} = 3.50 \text{ m}$$

which is greater than the actual length 3 m. Therefore, the flow is indeed supersonic when the normal shock occurs at the indicated location. Also, using the actual duct length  $L_1$ , we have  $\frac{fL_1}{D_h} = \frac{(0.007)(3 \text{ m})}{0.05 \text{ m}} = 0.4200$ . Noting that  $L_1 = L_1^* - L_2^*$ ,

the function  $fL^*/D_h$  at the exit state and the corresponding Mach number are

$$\left(\frac{fL^*}{D_h}\right)_2 = \left(\frac{fL^*}{D_h}\right)_1 - \frac{fL_1}{D_h} = 0.4898 - 0.4200 = 0.0698 \rightarrow Ma_2 = 1.315$$

From Table A-16, at Ma<sub>2</sub> = 1.315:  $T_2/T^* = 0.8918$  and  $P_2/P^* = 0.7183$ . Then the temperature, pressure, and velocity before the shock are determined to be

$$\frac{T_2}{T_1} = \frac{T_2 / T^*}{T_1 / T^*} = \frac{0.8918}{0.4673} = 1.9084$$
  $\rightarrow T_2 = 1.9084(380 \text{ K}) = 725.2 \text{ K}$ 

$$\frac{P_2}{P_1} = \frac{P_2 / P^*}{P_1 / P^*} = \frac{0.7183}{0.2441} = 2.9426$$
  $\rightarrow P_2 = 2.9426 P_1 = 2.9426(80 \text{ kPa}) = 235.4 \text{ kPa}$ 

The normal shock functions corresponding to a Mach number of 1.315 are, from Table A-14,

$$Ma_2 = 1.315$$
:  $Ma_3 = 0.7786$ ,  $T_3/T_2 = 1.2001$ ,  $P_3/P_2 = 1.8495$ 

Then the temperature and pressure after the shock become

$$T_3 = 1.2001T_2 = 1.2001(725.2 \text{ K}) = 870.3 \text{ K}$$
 and  $P_3 = 1.8495P_2 = 1.8495(235.4 \text{ kPa}) = 435.4 \text{ kPa}$ 

Sonic conditions exist at the duct exit, and the flow downstream the shock is still Fanno flow. From Table A-16,

Ma<sub>3</sub> = 0.7786: 
$$T_3/T^* = 1.0702$$
,  $P_3/P^* = 1.3286$   
Ma<sub>4</sub> = 1:  $T_4/T^* = 1$ ,  $P_4/P^* = 1$ 

Then the temperature, pressure, and velocity at the duct exit are determined to be

$$\frac{T_4}{T_3} = \frac{T_4 / T^*}{T_3 / T^*} = \frac{1}{1.0702}$$

$$\Rightarrow T_4 = T_3 / 1.0702 = (870.3 \text{ K}) / 1.0702 = 813 K$$

$$\frac{P_4}{P_3} = \frac{P_4 / P^*}{P_3 / P^*} = \frac{1}{1.3286}$$

$$\Rightarrow P_4 = P_3 / 1.3286 = (435.4 \text{ kPa}) / 1.3286 = 328 kPa$$

$$V_4 = \text{Ma}_4 c_4 = (1)\sqrt{kRT_4} = \sqrt{(1.4)(0.287 \text{ kJ/kg} \cdot \text{K})(813 \text{ K}) \left(\frac{1000 \text{ m}^2/\text{s}^2}{1 \text{ kJ/kg}}\right)} =$$
**572 m/s**

**Discussion** It can be shown that  $L_3^* = 0.67$  m, and thus the total length of this duct is 3.67 m. If the duct is extended, the normal shock will move further upstream, and eventually to the inlet of the duct.

D=6 in

 $Ma_2 = 1$ 

 $T^*$ 

 $P^*$ 

V

### 12-108E

Solution Helium enters a constant-area adiabatic duct at a specified state, and leaves at sonic state. The maximum duct length that will not cause the mass flow rate to be reduced is to be determined.

 $P_1 = 60 \text{ psia}$ 

 $T_1 = 600 \text{ R}$ 

 $Ma_1 = 0.2$ 

Assumptions 1 The assumptions associated with Fanno flow (i.e., steady, frictional flow of an ideal gas with constant properties through a constant cross-sectional area adiabatic duct) are valid. 2 The friction factor is constant along the duct.

We take the properties of helium to be k = 1.667,  $c_p =$ **Properties** 1.2403 Btu/lbm·R, and R = 0.4961 Btu/lbm·R. The friction factor is given to be f = 0.025.

Analysis Mach number of 0.2 is (Table A-16)

Btu/lbm·R, and 
$$R = 0.4961$$
 Btu/lbm·R. The friction factor is to be  $f = 0.025$ .

is The Fanno flow function  $fL^*/D$  corresponding to the inlet number of 0.2 is (Table A-16)

$$\frac{fL_1^*}{D} = 14.5333$$

Noting that \* denotes sonic conditions, which exist at the exit state, the duct length is determined to be

$$L_1^* = 14.5333D/f = 14.5333(6/12 \text{ ft})/0.025 = 291 \text{ ft}$$

Thus, for the given friction factor, the duct length must be 291 ft for the Mach number to reach Ma = 1 at the duct exit.

This problem can also be solved using equations instead of tabulated values for the Fanno functions.

**Solution** Subsonic airflow in a constant cross-sectional area adiabatic duct is considered. The duct length from the inlet where the inlet velocity doubles and the pressure drop in that section are to be determined.

**Assumptions** 1 The assumptions associated with Fanno flow (i.e., steady, frictional flow of an ideal gas with constant properties through a constant cross-sectional area adiabatic duct) are valid. 2 The friction factor remains constant along the duct.

**Properties** We take the properties of air to be k = 1.4,  $c_p = 1.005$  kJ/kg·K, and R = 0.287 kJ/kg·K. The average friction factor is given to be f = 0.014.

Analysis

The inlet Mach number is

$$c_1 = \sqrt{kRT_1} = \sqrt{(1.4)(0.287 \text{ kJ/kg} \cdot \text{K})(500 \text{ K}) \left(\frac{1000 \text{ m}^2/\text{s}^2}{1 \text{ kJ/kg}}\right)} = 448.2 \text{ m/s} \rightarrow \text{Ma}_1 = \frac{V_1}{c_1} = \frac{150 \text{ m/s}}{448.2 \text{ m/s}} = 0.3347 \text{ m/s}$$

The Fanno flow functions corresponding to the inlet and exit Mach numbers are, from Table A-16,

$$Ma_1 = 0.3347$$
:

$$(fL^*/D_h)_1 = 3.924$$

$$P_1/P^* = 3.2373, V_1/V^* = 0.3626$$

Therefore, 
$$V_1 = 0.3626V^*$$
. Then the Fanno function  $V_2/V^*$  becomes  $\frac{V_2}{V^*} = \frac{2V_1}{V^*} = \frac{2 \times 0.3626V^*}{V^*} = 0.7252$ .

The corresponding Mach number and Fanno flow functions are, from Table A-16,

$$Ma_2 = 0.693, (fL^*/D_h)_1 = 0.2220, \text{ and } P_2/P^* = 1.5099.$$

Then the duct length where the velocity doubles, the exit pressure, and the pressure drop become

$$L = L_1^* - L_2^* = \left(\frac{fL_1^*}{D_h} - \frac{fL_2^*}{D_h}\right) \frac{D_h}{f} = (3.924 - 0.2220) \frac{0.15 \text{ m}}{0.014} = 39.7 \text{ m}$$

$$\frac{P_2}{P_1} = \frac{P_2 / P^*}{P_1 / P^*} = \frac{1.5099}{3.2373} = 0.4664 \qquad \rightarrow P_2 = 0.4664 P_1 = 0.4664 (200 \text{ kPa}) = 93.3 \text{ kPa}$$

$$\Delta P = P_1 - P_2 = 200 - 93.3 = 106.7 \text{ kPa} \cong 107 \text{ kPa}$$

**Discussion** Note that it takes a duct length of 39.7 m for the velocity to double, and the Mach number to increase from 0.3347 to 0.693. The maximum (or sonic) duct lengths at the inlet and exit states in this case are  $L_1^* = 42.1$  m and  $L_2^* = 2.38$  m. Therefore, the flow would reach sonic conditions if there is an additional 2.38 m of duct length.

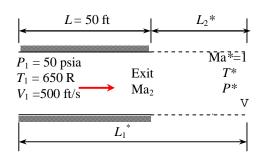
### 12-110E

**Solution** Air enters a constant-area adiabatic duct of given length at a specified state. The velocity, temperature, and pressure at the duct exit are to be determined.

**Assumptions** 1 The assumptions associated with Fanno flow (i.e., steady, frictional flow of an ideal gas with constant properties through a constant cross-sectional area adiabatic duct) are valid. 2 The friction factor is constant along the duct.

**Properties** We take the properties of helium to be k = 1.4,  $c_p = 0.2400 \text{ Btu/lbm} \cdot \text{R}$ , and  $R = 0.06855 \text{ Btu/lbm} \cdot \text{R} = 0.3704 \text{ psia} \cdot \text{ft}^3/\text{lbm} \cdot \text{R}$ . The friction factor is given to be f = 0.025.

**Analysis** The first thing we need to know is whether the flow is choked at the exit or not. Therefore, we first determine the inlet Mach number and the corresponding value of the function  $fL^*/D_h$ ,



$$T_1 = T_{01} - \frac{{V_1}^2}{2c_p} = 650 \,\mathrm{R} - \frac{(500 \,\mathrm{ft/s})^2}{2 \times 0.2400 \,\mathrm{Btu/lbm} \cdot \mathrm{R}} \left( \frac{1 \,\mathrm{Btu/lbm}}{25,037 \,\mathrm{ft}^2 /\mathrm{s}^2} \right) = 629.2 \,\mathrm{R}$$

$$c_1 = \sqrt{kRT_1} = \sqrt{(1.4)(0.06855 \text{ Btu/lbm} \cdot \text{R})(629.2 \text{ R}) \left(\frac{25,037 \text{ ft}^2/\text{s}^2}{1 \text{ Btu/lbm}}\right)} = 1230 \text{ ft/s}$$

$$Ma_1 = \frac{V_1}{c_1} = \frac{500 \text{ m/s}}{1230 \text{ ft/s}} = 0.4066$$

Corresponding to this Mach number we calculate (or read) from Table A-16),  $(fL^*/D_h)_1 = 2.1911$ . Also, using the actual duct length L, we have

$$\frac{fL}{D_h} = \frac{(0.02)(50 \text{ ft})}{6/12 \text{ ft}} = 2 < 2.1911$$

Therefore, the flow is not choked and exit Mach number is less than 1. Noting that

 $L = L_1^* - L_2^*$ , the function  $fL^*/D_h$  at the exit state is calculated from

$$\left(\frac{fL^*}{D_h}\right)_2 = \left(\frac{fL^*}{D_h}\right)_1 - \frac{fL}{D_h} = 2.1911 - 2 = 0.1911$$

The Mach number corresponding to this value of  $fL^*/D$  is obtained from Table A-16 to be Ma<sub>2</sub> = 0.7091. The Fanno flow functions corresponding to the inlet and exit Mach numbers are, from Table A-16,

$$Ma_1 = 0.4066$$
:  $T_1/T^* = 1.1616$ ,  $P_1/P^* = 2.6504$ ,  $V_1/V^* = 0.4383$   
 $Ma_2 = 0.7091$ :  $T_2/T^* = 1.0903$ ,  $P_2/P^* = 1.4726$ ,  $V_2/V^* = 0.7404$ 

Then the temperature, pressure, and velocity at the duct exit are determined to be

$$\frac{T_2}{T_1} = \frac{T_2 / T^*}{T_1 / T^*} = \frac{1.0903}{1.1616} = 0.9386 \qquad \rightarrow T_2 = 0.9386 (629.2 \text{ R}) = \mathbf{591R}$$

$$\frac{P_2}{P_1} = \frac{P_2 / P^*}{P_1 / P^*} = \frac{1.4726}{2.6504} = 0.5556 \qquad \rightarrow P_2 = 0.5556 (50 \text{ psia}) = \mathbf{27.8 \text{ psia}}$$

$$\frac{V_2}{V_1} = \frac{V_2 / V^*}{V_1 / V^*} = \frac{0.7404}{0.4383} = 1.6893 \qquad \rightarrow V_2 = 1.6893 (500 \text{ ft/s}) = \mathbf{845 \text{ ft/s}}$$

**Discussion** It can be shown that  $L_2^* = 4.8$  ft, indicating that it takes a duct length of 50 ft for the Mach number to increase from 0.4066 to 0.7091, but only 4.8 ft to increase from 0.7091 to 1. Therefore, the Mach number rises at a much higher rate as sonic conditions are approached.



**Solution** Choked subsonic airflow in a constant cross-sectional area adiabatic duct is considered. The variation of duct length with Mach number is to be investigated, and the results are to be plotted.

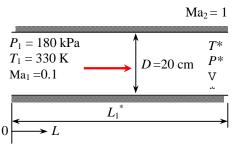
Assumptions 1 The assumptions associated with Fanno flow (i.e., steady, frictional flow of an ideal gas with constant properties through a constant cross-sectional area adiabatic duct) are valid. 2 The friction factor remains constant along the duct.

**Properties** We take the properties of air to be k = 1.4,  $c_p = 1.005$  kJ/kg·K, and R = 0.287 kJ/kg·K. The average friction factor is given to be f = 0.02.

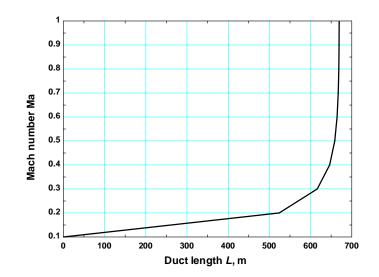
*Analysis* The flow is choked, and thus  $Ma_2 = 1$ . Corresponding to the inlet Mach number of  $Ma_1 = 0.1$  we have, from Table A-16,  $fL^*/D_h = 66.922$ , Therefore, the original duct length is

$$L_1^* = 66.922 \frac{D}{f} = 66.922 \frac{0.20 \text{ m}}{0.02} = 669 \text{ m}$$

Repeating the calculations for different Ma<sub>2</sub> as it varies from 0.1 to 1 results in the following table for the location on the duct from the inlet. The EES *Equations* window is printed below, along with the plotted results.



Mach	Duct length
number, Ma	L, m
0.10	0
0.20	524
0.30	616
0.40	646
0.50	659
0.60	664
0.70	667
0.80	668
0.90	669
1.00	669



## **EES program:**

k=1.4

cp=1.005

R=0.287

P1=180

T1=330

Ma1 = 0.1

"Ma2=1"

f=0.02

D=0.2

C1 = sqrt(k\*R\*T1\*1000)

Ma1=V1/C1

T01=T02

 $T01=T1*(1+0.5*(k-1)*Ma1^2)$ 

T02=T2\*(1+0.5\*(k-1)\*Ma2^2)

P01=P1\*(1+0.5\*(k-1)\*Ma1^2)^(k/(k-1))

rho1=P1/(R\*T1)Ac=pi\*D^2/4 mair=rho1\*Ac\*V1 P01Ps=((2+(k-1)\*Ma1^2)/(k+1))^(0.5\*(k+1)/(k-1))/Ma1 P1Ps=((k+1)/(2+(k-1)\*Ma1^2))^0.5/Ma1  $T1Ts=(k+1)/(2+(k-1)*Ma1^2)$ R1Rs=((2+(k-1)\*Ma1^2)/(k+1))^0.5/Ma1 V1Vs=1/R1Rs $fLs1 = (1-Ma1^2)/(k*Ma1^2) + (k+1)/(2*k)*ln((k+1)*Ma1^2/(2+(k-1)*Ma1^2))$ Ls1=fLs1\*D/f $P02Ps = ((2+(k-1)*Ma2^2)/(k+1))^(0.5*(k+1)/(k-1))/Ma2$ P2Ps=((k+1)/(2+(k-1)\*Ma2^2))^0.5/Ma2  $T2Ts=(k+1)/(2+(k-1)*Ma2^2)$ R2Rs=((2+(k-1)\*Ma2^2)/(k+1))^0.5/Ma2 V2Vs=1/R2Rs $fLs2 = (1-Ma2^2)/(k*Ma2^2) + (k+1)/(2*k)*ln((k+1)*Ma2^2/(2+(k-1)*Ma2^2))$ Ls2=fLs2\*D/f L=Ls1-Ls2 P02=P02Ps/P01Ps\*P01 P2=P2Ps/P1Ps\*P1 V2=V2Vs/V1Vs\*V1

**Discussion** Note that the Mach number increases very mildly at the beginning, and then rapidly near the duct outlet. It takes 262 m of duct length for Mach number to increase from 0.1 to 0.2, but only 1 m to increase from 0.7 to 1.



**Solution** Choked subsonic airflow in a constant cross-sectional area adiabatic duct is considered. The variation of duct length with Mach number is to be investigated, and the results are to be plotted.

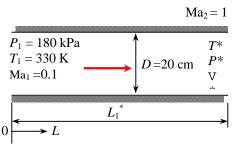
Assumptions 1 The assumptions associated with Fanno flow (i.e., steady, frictional flow of an ideal gas with constant properties through a constant cross-sectional area adiabatic duct) are valid. 2 The friction factor remains constant along the duct.

**Properties** We take the properties of helium to be k = 1.667,  $c_p = 5.193$  kJ/kg·K, and R = 2.077 kJ/kg·K. The average friction factor is given to be f = 0.02.

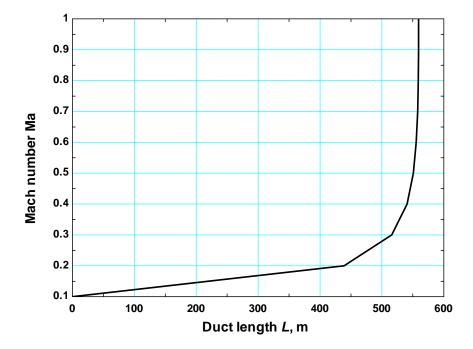
*Analysis* The flow is choked, and thus  $Ma_2 = 1$ . Corresponding to the inlet Mach number of  $Ma_1 = 0.1$  we have, from Table A-16,  $fL^*/D_h = 66.922$ , Therefore, the original duct length is

$$L_1^* = 66.922 \frac{D}{f} = 66.922 \frac{0.20 \text{ m}}{0.02} = 669 \text{ m}$$

Repeating the calculations for different  $Ma_2$  as it varies from 0.1 to 1 results in the following table for the location on the duct from the inlet. The EES *Equations* window is printed below, along with the plotted results.



Mach	Duct length
number, Ma	<i>L</i> , m
0.10	0
0.20	439
0.30	516
0.40	541
0.50	551
0.60	555
0.70	558
0.80	559
0.90	559
1.00	559



# **EES program:**

k=1.667

cp=5.193

R=2.077

P1=180

T1=330

Ma1=0.1

"Ma2=1"

f=0.02 D=0.2

12-71

```
C1=sqrt(k*R*T1*1000)
Ma1=V1/C1
T01=T02
T01=T1*(1+0.5*(k-1)*Ma1^2)
T02=T2*(1+0.5*(k-1)*Ma2^2)
P01=P1*(1+0.5*(k-1)*Ma1^2)^(k/(k-1))
rho1=P1/(R*T1)
Ac=pi*D^2/4
mair=rho1*Ac*V1
P01Ps = ((2+(k-1)*Ma1^2)/(k+1))^(0.5*(k+1)/(k-1))/Ma1
P1Ps=((k+1)/(2+(k-1)*Ma1^2))^0.5/Ma1
T1Ts=(k+1)/(2+(k-1)*Ma1^2)
R1Rs=((2+(k-1)*Ma1^2)/(k+1))^0.5/Ma1
V1Vs=1/R1Rs
fLs1 = (1-Ma1^2)/(k*Ma1^2) + (k+1)/(2*k)*ln((k+1)*Ma1^2/(2+(k-1)*Ma1^2))
Ls1=fLs1*D/f
P02Ps = ((2+(k-1)*Ma2^2)/(k+1))^{(0.5*(k+1)/(k-1))/Ma2
P2Ps=((k+1)/(2+(k-1)*Ma2^2))^0.5/Ma2
T2Ts=(k+1)/(2+(k-1)*Ma2^2)
R2Rs = ((2+(k-1)*Ma2^2)/(k+1))^0.5/Ma2
V2Vs=1/R2Rs
fLs2 = (1-Ma2^2)/(k*Ma2^2) + (k+1)/(2*k)*ln((k+1)*Ma2^2/(2+(k-1)*Ma2^2))
Ls2=fLs2*D/f
L=Ls1-Ls2
P02=P02Ps/P01Ps*P01
P2=P2Ps/P1Ps*P1
V2=V2Vs/V1Vs*V1
```

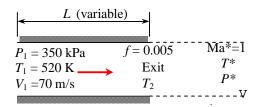
**Discussion** Note that the Mach number increases very mildly at the beginning, and then rapidly near the duct outlet. It takes 262 m of duct length for Mach number to increase from 0.1 to 0.2, but only 1 m to increase from 0.7 to 1.



**Solution** The flow of argon gas in a constant cross-sectional area adiabatic duct is considered. The variation of entropy change with exit temperature is to be investigated, and the calculated results are to be plotted on a *T-s* diagram.

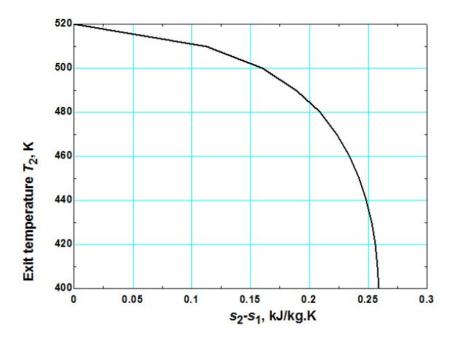
**Assumptions** 1 The assumptions associated with Fanno flow (i.e., steady, frictional flow of an ideal gas with constant properties through a constant cross-sectional area adiabatic duct) are valid. 2 The friction factor remains constant along the duct.

**Properties** The properties of argon are given to be k = 1.667,  $c_p = 0.5203$  kJ/kg·K, and R = 0.2081 kJ/kg·K. The average friction factor is given to be f = 0.005.



**Analysis** Using EES, we determine the entropy change and tabulate and plot the results as follows:

Exit temp.	Mach umber	Entropy change
$T_2$ , K	$Ma_2$	∆s, kg/kg·K
520	0.165	0.000
510	0.294	0.112
500	0.385	0.160
490	0.461	0.189
480	0.528	0.209
470	0.591	0.224
460	0.649	0.234
450	0.706	0.242
440	0.760	0.248
430	0.813	0.253
420	0.865	0.256
410	0.916	0.258
400	0.967	0.259

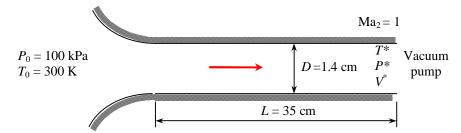


# **EES Program:**

```
k=1.667
cp=0.5203
R=0.2081
P1 = 350
T1 = 520
V1=70
"T2=400"
f=0.005
D=0.08
C1=sqrt(k*R*T1*1000)
Ma1=V1/C1
T01=T02
T01=T1*(1+0.5*(k-1)*Ma1^2)
P01=P1*(1+0.5*(k-1)*Ma1^2)^(k/(k-1))
rho1=P1/(R*T1)
Ac=pi*D^2/4
mair=rho1*Ac*V1
P01Ps = ((2+(k-1)*Ma1^2)/(k+1))^(0.5*(k+1)/(k-1))/Ma1
P1Ps=((k+1)/(2+(k-1)*Ma1^2))^0.5/Ma1
T1Ts=(k+1)/(2+(k-1)*Ma1^2)
R1Rs = ((2+(k-1)*Ma1^2)/(k+1))^0.5/Ma1
V1Vs=1/R1Rs
fLs1 = (1-Ma1^2)/(k*Ma1^2) + (k+1)/(2*k)*ln((k+1)*Ma1^2/(2+(k-1)*Ma1^2))
Ls1=fLs1*D/f
P02Ps = ((2+(k-1)*Ma2^2)/(k+1))^(0.5*(k+1)/(k-1))/Ma2
P2Ps = ((k+1)/(2+(k-1)*Ma2^2))^0.5/Ma2
T2Ts=(k+1)/(2+(k-1)*Ma2^2)
R2Rs = ((2+(k-1)*Ma2^2)/(k+1))^0.5/Ma2
V2Vs=1/R2Rs
fLs2 = (1-Ma2^2)/(k*Ma2^2) + (k+1)/(2*k)*ln((k+1)*Ma2^2/(2+(k-1)*Ma2^2))
Ls2=fLs2*D/f
L=Ls1-Ls2
P02=P02Ps/P01Ps*P01
P2=P2Ps/P1Ps*P1
T2=T2Ts/T1Ts*T1
V2=V2Vs/V1Vs*V1
Del_s=cp*ln(T2/T1)-R*ln(P2/P1)
```

**Discussion** Note that entropy increases with increasing duct length and Mach number (and thus decreasing temperature). It reached a maximum value of 0.259 kJ/kg.K when the Mach number reaches  $Ma_2 = 1$  and thus the flow is choked.

**Solution** Air enters a constant-area adiabatic duct at a specified state, and leaves at sonic state. The maximum duct length that will not cause the mass flow rate to be reduced is to be determined.



Assumptions 1 The assumptions associated with Fanno flow (i.e., steady, frictional flow of an ideal gas with constant properties through a constant cross-sectional area adiabatic duct) are valid. 2 The friction factor is constant along the duct.

**Properties** We take the properties of air to be k = 1.4,  $c_p = 1.005$  kJ/kg·K, and R = 0.287 kJ/kg·K. The friction factor is given to be f = 0.018.

Analysis The mass flow rate will be maximum when the flow is choked, and thus the exit Mach number is  $Ma_2 = 1$ . In that case we have

$$\frac{fL_1^*}{D} = \frac{fL_1}{D} = \frac{(0.018)(0.35 \text{ m})}{0.014 \text{ m}} = 0.45$$

The Mach number corresponding to this value of  $fL^*/D$  at the tube inlet is obtained from Table A-16 to be Ma<sub>1</sub> = 0.6107  $\approx$  **0.611**. This value is obtained using the analytical relation. An interpolation on Table 16 gives 0.614. Noting that the flow in the nozzle section is isentropic, the thermodynamic temperature, pressure, and density at the tube inlet become

$$T_{1} = T_{01} \left( 1 + \frac{k-1}{2} \operatorname{Ma}_{1}^{2} \right)^{-1} = (300 \,\mathrm{K}) \left( 1 + \frac{1.4-1}{2} (0.6107)^{2} \right)^{-1} = 279.2 \,\mathrm{K}$$

$$P_{1} = P_{01} \left( 1 + \frac{k-1}{2} \operatorname{Ma}_{1}^{2} \right)^{-k/(k-1)} = (100 \,\mathrm{kPa}) \left( 1 + \frac{1.4-1}{2} (0.6107)^{2} \right)^{-1.4/0.4} = 77.74 \,\mathrm{kPa}$$

$$\rho_{1} = \frac{P_{1}}{RT_{1}} = \frac{77.74 \,\mathrm{kPa}}{(0.287 \,\mathrm{kJ/kgK})(279.2 \,\mathrm{K})} = 0.9702 \,\mathrm{kg/m}^{3}$$

Then the inlet velocity and the mass flow rate become

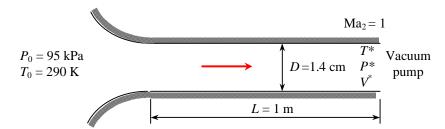
$$c_1 = \sqrt{kRT_1} = \sqrt{(1.4)(0.287 \text{ kJ/kg} \cdot \text{K})(279.2 \text{ K}) \left(\frac{1000 \text{ m}^2/\text{s}^2}{1 \text{ kJ/kg}}\right)} = 334.9 \text{ m/s}$$

$$V_1 = \text{Ma}_1 c_1 = 0.6107(334.9 \text{ m/s}) = 204.5 \text{ m/s}$$

$$m_{air}^{\text{K}} = \rho_1 A_{cl} V_1 = (0.9702 \text{ kg/m}^3) [\pi (0.014 \text{ m})^2/4](204.5 \text{ m/s}) = \textbf{0.0305 kg/s}$$

**Discussion** This is the maximum mass flow rate through the tube for the specified stagnation conditions at the inlet. The flow rate will remain at this level even if the vacuum pump drops the pressure even further.

**Solution** Air enters a constant-area adiabatic duct at a specified state, and leaves at sonic state. The maximum duct length that will not cause the mass flow rate to be reduced is to be determined.



Assumptions 1 The assumptions associated with Fanno flow (i.e., steady, frictional flow of an ideal gas with constant properties through a constant cross-sectional area adiabatic duct) are valid. 2 The friction factor is constant along the duct.

**Properties** We take the properties of air to be k = 1.4,  $c_p = 1.005$  kJ/kg·K, and R = 0.287 kJ/kg·K. The friction factor is given to be f = 0.025.

**Analysis** The mass flow rate will be maximum when the flow is choked, and thus the exit Mach number is  $Ma_2 = 1$ . In that case we have

$$\frac{fL_1^*}{D} = \frac{fL_1}{D} = \frac{(0.025)(1\,\mathrm{m})}{0.014\,\mathrm{m}} = 1.786$$

The Mach number corresponding to this value of  $fL^*/D$  at the tube inlet is obtained from Table A-16 to be Ma<sub>1</sub> = **0.4422**. Noting that the flow in the nozzle section is isentropic, the thermodynamic temperature, pressure, and density at the tube inlet become

$$T_{1} = T_{01} \left( 1 + \frac{k-1}{2} \operatorname{Ma}_{1}^{2} \right)^{-1} = (290 \,\mathrm{K}) \left( 1 + \frac{1.4-1}{2} (0.4422)^{2} \right)^{-1} = 279.1 \,\mathrm{K}$$

$$P_{1} = P_{01} \left( 1 + \frac{k-1}{2} \operatorname{Ma}_{1}^{2} \right)^{-k/(k-1)} = (95 \,\mathrm{kPa}) \left( 1 + \frac{1.4-1}{2} (0.4422)^{2} \right)^{-1.4/0.4} = 83.06 \,\mathrm{kPa}$$

$$\rho_{1} = \frac{P_{1}}{RT_{1}} = \frac{83.06 \,\mathrm{kPa}}{(0.287 \,\mathrm{kJ/kgK})(279.1 \,\mathrm{K})} = 1.037 \,\mathrm{kg/m}^{3}$$

Then the inlet velocity and the mass flow rate become

$$c_1 = \sqrt{kRT_1} = \sqrt{(1.4)(0.287 \text{ kJ/kg} \cdot \text{K})(279.1 \text{ K}) \left(\frac{1000 \text{ m}^2/\text{s}^2}{1 \text{ kJ/kg}}\right)} = 334.9 \text{ m/s}$$

$$V_1 = \text{Ma}_1 c_1 = 0.4422(334.9 \text{ m/s}) = 148.1 \text{ m/s}$$

$$m_{air}^{2} = \rho_{1}A_{c1}V_{1} = (1.037 \text{ kg/m}^{3})[\pi(0.014 \text{ m})^{2}/4](148.1 \text{ m/s}) =$$
**0.0236 kg/s**

**Discussion** This is the maximum mass flow rate through the tube for the specified stagnation conditions at the inlet. The flow rate will remain at this level even if the vacuum pump drops the pressure even further.

#### **Review Problems**

# **12-116**

**Solution** The thrust developed by the engine of a Boeing 777 is about 380 kN. The mass flow rate of gases through the nozzle is to be determined.

Assumptions 1 Air is an ideal gas with constant specific heats. 2 Flow of combustion gases through the nozzle is isentropic. 3 Choked flow conditions exist at the nozzle exit. 4 The velocity of gases at the nozzle inlet is negligible.

**Properties** The gas constant of air is  $R = 0.287 \text{ kPa} \cdot \text{m}^3/\text{kg} \cdot \text{K}$ , and it can also be used for combustion gases. The specific heat ratio of combustion gases is k = 1.33.

Analysis The velocity at the nozzle exit is the sonic speed, which is determined to be

$$V = c = \sqrt{kRT} = \sqrt{(1.33)(0.287 \text{ kJ/kg} \cdot \text{K}) \left(\frac{1000 \text{ m}^2/\text{s}^2}{1 \text{ kJ/kg}}\right) (220 \text{ K})} = 289.8 \text{ m/s}$$

Noting that thrust F is related to velocity by F = n N V, the mass flow rate of combustion gases is determined to be

$$m = \frac{F}{V} = \frac{380,000 \text{ N}}{289.8 \text{ m/s}} \left( \frac{1 \text{ kg} \cdot \text{m/s}^2}{1 \text{ N}} \right) = 1311 \text{ kg/s} \approx 1310 \text{ kg/s}$$

**Discussion** The combustion gases are mostly nitrogen (due to the 78% of  $N_2$  in air), and thus they can be treated as air with a good degree of approximation.

# 12-117

**Solution** A stationary temperature probe is inserted into an air duct reads 85°C. The actual temperature of air is to be determined.

Assumptions 1 Air is an ideal gas with constant specific heats at room temperature. 2 The stagnation process is isentropic.

**Properties** The specific heat of air at room temperature is  $c_p = 1.005 \text{ kJ/kg·K}$ .

Analysis The air that strikes the probe will be brought to a complete stop, and thus it will undergo a stagnation process. The thermometer will sense the temperature of this stagnated air, which is the stagnation temperature. The actual air temperature is determined from

$$T = T_0 - \frac{V^2}{2c_p} = 85^{\circ}\text{C} - \frac{(190 \text{ m/s})^2}{2 \times 1.005 \text{ kJ/kg} \cdot \text{K}} \left( \frac{1 \text{ kJ/kg}}{1000 \text{ m}^2 / \text{s}^2} \right) = 67.0^{\circ}\text{C}$$

**Discussion** Temperature rise due to stagnation is very significant in high-speed flows, and should always be considered when compressibility effects are not negligible.

**Solution** Nitrogen flows through a heat exchanger. The stagnation pressure and temperature of the nitrogen at the inlet and the exit states are to be determined.

Assumptions 1 Nitrogen is an ideal gas with constant specific heats. 2 Flow of nitrogen through the heat exchanger is isentropic.

**Properties** The properties of nitrogen are  $c_p = 1.039 \text{ kJ/kg} \cdot \text{K}$  and k = 1.4.

Analysis The stagnation temperature and pressure of nitrogen at the inlet and the exit states are determined from

$$T_{01} = T_1 + \frac{V_1^2}{2c_p} = 10^{\circ}\text{C} + \frac{(100 \text{ m/s})^2}{2 \times 1.039 \text{ kJ/kg} \cdot ^{\circ}\text{C}} \left( \frac{1 \text{ kJ/kg}}{1000 \text{ m}^2 \text{ /s}^2} \right) = 14.8^{\circ}\text{C}$$

$$P_{01} = P_1 \left( \frac{T_{01}}{T_1} \right)^{k/(k-1)} = (150 \text{ kPa}) \left( \frac{288.0 \text{ K}}{283.2 \text{ K}} \right)^{1.4/(1.4-1)} = 159 \text{ kPa}$$

$$150 \text{ kPa}$$

$$10^{\circ}\text{C}$$

$$100 \text{ m/s}$$
Nitrogen
$$100 \text{ kPa}$$

$$100 \text{ kPa}$$

From the energy balance relation  $E_{\rm in} - E_{\rm out} = \Delta E_{\rm system}$  with w = 0

$$q_{\text{in}} = c_p (T_2 - T_1) + \frac{{V_2}^2 - {V_1}^2}{2} + \Delta p e^{70}$$

$$150 \text{ kJ/kg} = (1.039 \text{ kJ/kg} \cdot ^{\circ}\text{C})(T_2 - 10^{\circ}\text{C}) + \frac{(200 \text{ m/s})^2 - (100 \text{ m/s})^2}{2} \left(\frac{1 \text{ kJ/kg}}{1000 \text{ m}^2 / \text{s}^2}\right)$$

$$T_2 = 139.9^{\circ}\text{C}$$

and

$$T_{02} = T_2 + \frac{{V_2}^2}{2c_p} = 139.9 \text{°C} + \frac{(200 \text{ m/s})^2}{2 \times 1.039 \text{ kJ/kg} \cdot \text{°C}} \left( \frac{1 \text{ kJ/kg}}{1000 \text{ m}^2 / \text{s}^2} \right) = \mathbf{159} \text{°C}$$

$$P_{02} = P_2 \left( \frac{T_{02}}{T_2} \right)^{k/(k-1)} = (100 \text{ kPa}) \left( \frac{432.3 \text{ K}}{413.1 \text{ K}} \right)^{1.4/(1.4-1)} = \mathbf{117} \text{ kPa}$$

**Discussion** Note that the stagnation temperature and pressure can be very different than their thermodynamic counterparts when dealing with compressible flow.

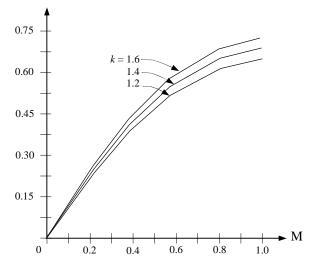
**Solution** The mass flow parameter  $n \sqrt[4]{RT_0} / (AP_0)$  versus the Mach number for k = 1.2, 1.4, and 1.6 in the range of  $0 \le \text{Ma} \le 1$  is to be plotted.

**Analysis** The mass flow rate parameter  $(R\sqrt[4]{RT_0})/P_0A$  can be expressed as

$$\frac{n \& \sqrt{RT_0}}{P_0 A} = \text{Ma}\sqrt{k} \left(\frac{2}{2 + (k-1)M^2}\right)^{(k+1)/2(k-1)}$$

Thus.

inus,			
Ma	k = 1.2	k = 1.4	k = 1.6
0.0	0	0	0
0.1	0.1089	0.1176	0.1257
0.2	0.2143	0.2311	0.2465
0.3	0.3128	0.3365	0.3582
0.4	0.4015	0.4306	0.4571
0.5	0.4782	0.5111	0.5407
0.6	0.5411	0.5763	0.6077
0.7	0.5894	0.6257	0.6578
0.8	0.6230	0.6595	0.6916
0.9	0.6424	0.6787	0.7106
1.0	0.6485	0.6847	0.7164



**Discussion** Note that the mass flow rate increases with increasing Mach number and specific heat ratio. It levels off at Ma = 1, and remains constant (choked flow).

**Solution** The equivalent relation for the speed of sound is to be verified using thermodynamic relations.

**Analysis** The two relations are 
$$c^2 = \left(\frac{\partial P}{\partial \rho}\right)_s$$
 and  $c^2 = k \left(\frac{\partial P}{\partial \rho}\right)_T$ 

From 
$$r = 1/v$$
  $\longrightarrow$   $dr = -dv/v^2$ . Thus,  $c^2 = \left(\frac{\partial P}{\partial r}\right)_s = -v^2 \left(\frac{\partial P}{\partial v}\right)_s = -v^2 \left(\frac{\partial P}{\partial T}\frac{\partial T}{\partial v}\right)_s = -v^2 \left(\frac{\partial P}{\partial T}\right)_s \left(\frac{\partial T}{\partial v}\right)_s$ 

From the cyclic rule,

$$(P,T,s): \left(\frac{\partial P}{\partial T}\right)_{s} \left(\frac{\partial T}{\partial s}\right)_{P} \left(\frac{\partial s}{\partial P}\right)_{T} = -1 \longrightarrow \left(\frac{\partial P}{\partial T}\right)_{s} = -\left(\frac{\partial s}{\partial T}\right)_{P} \left(\frac{\partial P}{\partial s}\right)_{T}$$

$$(T, v, s) : \left(\frac{\partial T}{\partial v}\right)_{s} \left(\frac{\partial v}{\partial s}\right)_{T} \left(\frac{\partial s}{\partial T}\right)_{v} = -1 \longrightarrow \left(\frac{\partial T}{\partial v}\right)_{s} = -\left(\frac{\partial s}{\partial v}\right)_{T} \left(\frac{\partial T}{\partial s}\right)_{v}$$

Substituting.

$$c^{2} = -v^{2} \left( \frac{\partial s}{\partial T} \right)_{P} \left( \frac{\partial P}{\partial s} \right)_{T} \left( \frac{\partial s}{\partial v} \right)_{T} \left( \frac{\partial T}{\partial s} \right)_{v} = -v^{2} \left( \frac{\partial s}{\partial T} \right)_{P} \left( \frac{\partial T}{\partial s} \right)_{v} \left( \frac{\partial P}{\partial v} \right)_{T}$$

Recall that 
$$\frac{c_p}{T} = \left(\frac{\partial s}{\partial T}\right)_P$$
 and  $\frac{c_v}{T} = \left(\frac{\partial s}{\partial T}\right)_v$ . Substituting,

$$c^{2} = -v^{2} \left( \frac{c_{p}}{T} \right) \left( \frac{T}{c_{v}} \right) \left( \frac{\partial P}{\partial v} \right)_{T} = -v^{2} k \left( \frac{\partial P}{\partial v} \right)_{T}$$

Replacing  $-dv/v^2$  by  $d\rho$ , we get  $c^2 = k \left(\frac{\partial P}{\partial \rho}\right)_T$ , which is the desired expression

**Discussion** Note that the differential thermodynamic property relations are very useful in the derivation of other property relations in differential form.

# 12-121

**Solution** For ideal gases undergoing isentropic flows, expressions for  $P/P^*$ ,  $T/T^*$ , and  $\rho/\rho^*$  as functions of k and Ma are to be obtained.

Analysis Equations 12-18 and 12-21 are given to be  $\frac{T_0}{T} = \frac{2 + (k-1)\text{Ma}^2}{2}$  and  $\frac{T^*}{T_0} = \frac{2}{k+1}$ 

Multiplying the two,  $\left(\frac{T_0}{T} \frac{T^*}{T_0}\right) = \left(\frac{2 + (k-1)\text{Ma}^2}{2}\right) \left(\frac{2}{k+1}\right)$ 

Simplifying and inverting,  $\frac{T}{T^*} = \frac{k+1}{2 + (k-1)\text{Ma}^2}$  (1)

From  $\frac{P}{P^*} = \left(\frac{T}{T^*}\right)^{k/(k-1)} \longrightarrow \boxed{\frac{P}{P^*} = \left(\frac{k+1}{2 + (k-1)\text{Ma}^2}\right)^{k/(k-1)}}$  (2)

From  $\frac{\rho}{\rho^*} = \left(\frac{\rho}{\rho^*}\right)^{k/(k-1)} \longrightarrow \left[\frac{\rho}{\rho^*} = \left(\frac{k+1}{2+(k-1)\text{Ma}^2}\right)^{k/(k-1)}\right]$  (3)

**Discussion** Note that some very useful relations can be obtained by very simple manipulations.

**Solution** It is to be verified that for the steady flow of ideal gases  $dT_0/T = dA/A + (1-Ma^2) dV/V$ . The effect of heating and area changes on the velocity of an ideal gas in steady flow for subsonic flow and supersonic flow are to be explained.

**Analysis** We start with the relation 
$$\frac{V^2}{2} = c_p (T_0 - T)$$
 (1)

Differentiating, 
$$VdV = c_p (dT_0 - dT)$$
 (2)

We also have 
$$\frac{d\rho}{\rho} + \frac{dA}{A} + \frac{dV}{V} = 0$$
 (3)

and 
$$\frac{dP}{\rho} + V dV = 0 \tag{4}$$

Differentiating the ideal gas relation 
$$P = \rho RT$$
,  $\frac{dP}{P} = \frac{d\rho}{\rho} + \frac{dT}{T} = 0$  (5)

From the speed of sound relation, 
$$c^2 = kRT = (k-1)c_nT = kP/\rho$$
 (6)

Combining Eqs. (3) and (5), 
$$\frac{dP}{P} - \frac{dT}{T} + \frac{dA}{A} + \frac{dV}{V} = 0$$
 (7)

Combining Eqs. (4) and (6), 
$$\frac{dP}{\rho} = \frac{dP}{kP/c^2} = -VdV$$

or, 
$$\frac{dP}{P} = -\frac{k}{C^2}VdV = -k\frac{V^2}{C^2}\frac{dV}{V} = -k\text{Ma}^2\frac{dV}{V}$$
 (8)

Combining Eqs. (2) and (6), 
$$dT = dT_0 - V \frac{dV}{c_p}$$

or, 
$$\frac{dT}{T} = \frac{dT_0}{T} - \frac{V^2}{C_p T} \frac{dV}{V} = \frac{dT}{T} = \frac{dT_0}{T} - \frac{V^2}{C^2 / (k-1)} \frac{dV}{V} = \frac{dT_0}{T} - (k-1) \text{Ma}^2 \frac{dV}{V}$$
(9)

Combining Eqs. (7), (8), and (9), 
$$-(k-1)\text{Ma}^2 \frac{dV}{V} - \frac{dT_0}{T} + (k-1)\text{Ma}^2 \frac{dV}{V} + \frac{dA}{A} + \frac{dV}{V} = 0$$

or, 
$$\frac{dT_0}{T} = \frac{dA}{A} + \left[ -k\text{Ma}^2 + (k-1)\text{Ma}^2 + 1 \right] \frac{dV}{V}$$
Thus, 
$$\frac{dT_0}{T} = \frac{dA}{A} + (1 - \text{Ma}^2) \frac{dV}{V}$$
(10)

Differentiating the steady-flow energy equation  $q = h_{02} - h_{01} = c_p (T_{02} - T_{01})$ 

$$\delta q = c_n dT_0 \tag{11}$$

Eq. (11) relates the stagnation temperature change  $dT_0$  to the net heat transferred to the fluid. Eq. (10) relates the velocity changes to area changes dA, and the stagnation temperature change  $dT_0$  or the heat transferred.

(a) When Ma < 1 (subsonic flow), the fluid accelerates if the duct converges (dA < 0) or the fluid is heated  $(dT_0 > 0)$  or  $\delta q > 0$ . The fluid decelerates if the duct converges (dA < 0) or the fluid is cooled  $(dT_0 < 0)$  or  $\delta q < 0$ .

(b) When Ma > 1 (supersonic flow), the fluid accelerates if the duct diverges (dA > 0) or the fluid is cooled  $(dT_0 < 0)$  or  $\delta q < 0$ . The fluid decelerates if the duct converges (dA < 0) or the fluid is heated  $(dT_0 > 0)$  or  $\delta q > 0$ .

**Discussion** Some of these results are not intuitively obvious, but come about by satisfying the conservation equations.

**Solution** A Pitot-static probe measures the difference between the static and stagnation pressures for a subsonic airplane. The speed of the airplane and the flight Mach number are to be determined.

**Assumptions** 1 Air is an ideal gas with a constant specific heat ratio. 2 The stagnation process is isentropic.

**Properties** The properties of air are  $R = 0.287 \text{ kJ/kg} \cdot \text{K}$  and k = 1.4.

Analysis The stagnation pressure of air at the specified conditions is

$$P_0 = P + \Delta P = 54 + 16 = 70 \text{ kPa}$$

Then,

$$\frac{P_0}{P} = \left(1 + \frac{(k-1)\text{Ma}^2}{2}\right)^{k/k-1} \longrightarrow \frac{70}{54} = \left(1 + \frac{(1.4-1)\text{Ma}^2}{2}\right)^{1.4/0.4}$$

It yields

$$Ma = 0.620$$

The speed of sound in air at the specified conditions is

$$c = \sqrt{kRT} = \sqrt{(1.4)(0.287 \text{ kJ/kg} \cdot \text{K})(256 \text{ K}) \left(\frac{1000 \text{ m}^2/\text{s}^2}{1 \text{ kJ/kg}}\right)} = 320.7 \text{ m/s}$$

Thus,

$$V = \text{Ma} \times c = (0.620)(320.7 \text{ m/s}) = 199 \text{ m/s}$$

**Discussion** Note that the flow velocity can be measured in a simple and accurate way by simply measuring pressure.

**Solution** An expression for the speed of sound based on van der Waals equation of state is to be derived. Using this relation, the speed of sound in carbon dioxide is to be determined and compared to that obtained by ideal gas behavior.

**Properties** The properties of  $CO_2$  are R = 0.1889 kJ/kg·K and k = 1.279 at T = 50°C = 323.2 K.

Analysis Van der Waals equation of state can be expressed as  $P = \frac{RT}{v-b} - \frac{a}{v^2}$ .

Differentiating,  $\left(\frac{\partial P}{\partial v}\right)_T = \frac{RT}{(v-b)^2} + \frac{2a}{v^3}$ 

Noting that  $\rho = 1/v \longrightarrow d\rho = -dv/v^2$ , the speed of sound relation becomes

Substituting,  $c^{2} = k \left(\frac{\partial P}{\partial r}\right)_{T} = v^{2} k \left(\frac{\partial P}{\partial v}\right)_{T}$  $c^{2} = \frac{v^{2} k RT}{(v - b)^{2}} - \frac{2ak}{v}$ 

Using the molar mass of  $CO_2$  (M = 44 kg/kmol), the constant a and b can be expressed per unit mass as

$$a = 0.1882 \text{ kPa} \cdot \text{m}^6/\text{kg}^2$$
 and  $b = 9.70 \times 10^{-4} \text{ m}^3/\text{kg}$ 

The specific volume of CO<sub>2</sub> is determined to be

$$200 \text{ kPa} = \frac{(0.1889 \text{ kPa} \cdot \text{m}^3 / \text{kg} \cdot \text{K})(323.2 \text{ K})}{\text{v} - 0.000970 \text{ m}^3/\text{kg}} - \frac{2 \times 0.1882 \text{ kPa} \cdot \text{m}^6 / \text{kg}^2}{\text{v}^2} \rightarrow \text{v} = 0.300 \text{ m}^3 / \text{kg}$$

Substituting,

$$c = \left( \left( \frac{(0.300 \, \text{m}^3 \, / \, \text{kg})^2 \, (1.279) (0.1889 \, \text{kJ/kg} \cdot \text{K}) (323.2 \, \text{K})}{(0.300 - 0.000970 \, \text{m}^3 \, / \, \text{kg})^2} - \frac{2 (0.1882 \, \text{kPa} \cdot \text{m}^6 \, / \, \text{kg}^3) (1.279)}{(0.300 \, \text{m}^3 \, / \, \text{kg})^2} \right) \left( \frac{1000 \, \text{m}^2 \, / \, \text{s}^2}{1 \, \text{kPa} \cdot \text{m}^3 / \text{kg}} \right) \right)^{1/2}$$

If we treat CO<sub>2</sub> as an ideal gas, the speed of sound becomes

$$c = \sqrt{kRT} = \sqrt{(1.279)(0.1889 \text{ kJ/kg} \cdot \text{K})(323.2 \text{ K}) \left(\frac{1000 \text{ m}^2 / \text{s}^2}{1 \text{ kJ/kg}}\right)} = 279 \text{ m/s}$$

**Discussion** Note that the ideal gas relation is the simplest equation of state, and it is very accurate for most gases encountered in practice. At high pressures and/or low temperatures, however, the gases deviate from ideal gas behavior, and it becomes necessary to use more complicated equations of state.

(i) He \* (\*)
120 m/s Ma = 1

# 12-125

**Solution** Helium gas is accelerated in a nozzle. The pressure and temperature of helium at the location where Ma = 1 and the ratio of the flow area at this location to the inlet flow area are to be determined.

**Assumptions** 1 Helium is an ideal gas with constant specific heats. 2 Flow through the nozzle is steady, one-dimensional, and isentropic.

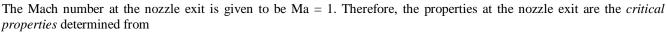
**Properties** The properties of helium are  $R = 2.0769 \text{ kJ/kg} \cdot \text{K}$ ,  $c_p = 5.1926 \text{ kJ/kg} \cdot \text{K}$ , and k = 1.667.

**Analysis** The properties of the fluid at the location where Ma = 1 are the critical properties, denoted by superscript \*. We first determine the stagnation temperature and pressure, which remain constant throughout the nozzle since the flow is isentropic.

$$T_0 = T_i + \frac{V_i^2}{2c_p} = 560 \text{ K} + \frac{(120 \text{ m/s})^2}{2 \times 5.1926 \text{ kJ/kg} \cdot \text{K}} \left(\frac{1 \text{ kJ/kg}}{1000 \text{ m}^2/\text{s}^2}\right) = 561.4 \text{ K}$$

and

$$P_0 = P_i \left(\frac{T_0}{T_i}\right)^{k/(k-1)} = (0.6 \text{ MPa}) \left(\frac{561.4 \text{ K}}{560 \text{ K}}\right)^{1.667/(1.667-1)} = 0.6037 \text{ MPa}$$



$$T^* = T_0 \left(\frac{2}{k+1}\right) = (561.4 \text{ K}) \left(\frac{2}{1.667+1}\right) = 421.0 \text{ K} = 421 \text{ K}$$

$$P^* = P_0 \left(\frac{2}{k+1}\right)^{k/(k-1)} = (0.6037 \text{ MPa}) \left(\frac{2}{1.667+1}\right)^{1.667/(1.667-1)} = 0.2941 \text{ MPa} \cong 0.294 \text{ MPa}$$

The speed of sound and the Mach number at the nozzle inlet are

$$c_i = \sqrt{kRT_i} = \sqrt{(1.667)(2.0769 \text{ kJ/kg} \cdot \text{K})(560 \text{ K}) \left(\frac{1000 \text{ m}^2/\text{s}^2}{1 \text{ kJ/kg}}\right)} = 1392 \text{ m/s}$$

$$Ma_i = \frac{V_i}{c_i} = \frac{120 \text{ m/s}}{1392 \text{ m/s}} = 0.08618$$

The ratio of the entrance-to-throat area is

$$\frac{A_i}{A^*} = \frac{1}{\text{Ma}_i} \left[ \left( \frac{2}{k+1} \right) \left( 1 + \frac{k-1}{2} \text{Ma}_i^2 \right) \right]^{(k+1)/[2(k-1)]} \\
= \frac{1}{0.08618} \left[ \left( \frac{2}{1.667 + 1} \right) \left( 1 + \frac{1.667 - 1}{2} (0.08618)^2 \right) \right]^{2.667/(2 \times 0.667)} \\
= 8.745$$

Then the ratio of the throat area to the entrance area becomes

$$\frac{A^*}{A_i} = \frac{1}{8.745} = 0.1144 \cong \mathbf{0.1144}$$

**Discussion** The compressible flow functions are essential tools when determining the proper shape of the compressible flow duct.

**Solution** Helium gas enters a nozzle with negligible velocity, and is accelerated in a nozzle. The pressure and temperature of helium at the location where Ma = 1 and the ratio of the flow area at this location to the inlet flow area are to be determined.

**Assumptions** 1 Helium is an ideal gas with constant specific heats. 2 Flow through the nozzle is steady, one-dimensional, and isentropic. 3 The entrance velocity is negligible.

**Properties** The properties of helium are  $R = 2.0769 \text{ kJ/kg} \cdot \text{K}$ ,  $c_p = 5.1926 \text{ kJ/kg} \cdot \text{K}$ , and k = 1.667.

**Analysis** We treat helium as an ideal gas with k = 1.667. The properties of the fluid at the location where Ma = 1 are the critical properties, denoted by superscript \*.

The stagnation temperature and pressure in this case are identical to the inlet temperature and pressure since the inlet velocity is negligible. They remain constant throughout the nozzle since the flow is isentropic.

$$T_0 = T_i = 560 \text{ K}$$

$$P_0 = P_i = 0.6 \text{ MPa}$$

The Mach number at the nozzle exit is given to be Ma = 1. Therefore, the properties at the nozzle exit are the *critical properties* determined from

$$T^* = T_0 \left(\frac{2}{k+1}\right) = (560 \text{ K}) \left(\frac{2}{1.667+1}\right) = 420 \text{ K}$$

$$P^* = P_0 \left(\frac{2}{k+1}\right)^{k/(k-1)} = (0.6 \text{ MPa}) \left(\frac{2}{1.667+1}\right)^{1.667/(1.667-1)} = \textbf{0.292 MPa}$$

$$V_i \approx 0$$

$$W_{a} = 0$$

The ratio of the nozzle inlet area to the throat area is determined from

$$\frac{A_i}{A^*} = \frac{1}{\text{Ma}_i} \left[ \left( \frac{2}{k+1} \right) \left( 1 + \frac{k-1}{2} \text{Ma}_i^2 \right) \right]^{(k+1)/[2(k-1)]}$$

But the Mach number at the nozzle inlet is Ma = 0 since  $V_i \cong 0$ . Thus the ratio of the throat area to the nozzle inlet area is

$$\frac{A^*}{A_i} = \frac{1}{\infty} = \mathbf{0}$$

**Discussion** The compressible flow functions are essential tools when determining the proper shape of the compressible flow duct.



**Solution** Air enters a converging nozzle. The mass flow rate, the exit velocity, the exit Mach number, and the exit pressure-stagnation pressure ratio versus the back pressure-stagnation pressure ratio for a specified back pressure range are to be calculated and plotted.

Assumptions 1 Air is an ideal gas with constant specific heats at room temperature. 2 Flow through the nozzle is steady, one-dimensional, and isentropic.

**Properties** The properties of air at room temperature are R = 0.287 kJ/kg·K,  $c_p = 1.005 \text{ kJ/kg·K}$ , and k = 1.4.

**Analysis** We use EES to tabulate and plot the results. The stagnation properties remain constant throughout the nozzle since the flow is isentropic. They are determined from

$$T_0 = T_i + \frac{{V_i}^2}{2c_p} = 400 \text{ K} + \frac{(180 \text{ m/s})^2}{2 \times 1.005 \text{ kJ/kg} \cdot \text{K}} \left(\frac{1 \text{ kJ/kg}}{1000 \text{ m}^2 / \text{s}^2}\right) = 416.1 \text{ K}$$

$$P_0 = P_i \left(\frac{T_0}{T_i}\right)^{k/(k-1)} = (900 \text{ kPa}) \left(\frac{416.1 \text{ K}}{400 \text{ K}}\right)^{1.4/(1.4-1)} = 1033.3 \text{ kPa}$$

The critical pressure is determined to be

$$P^* = P_0 \left(\frac{2}{k+1}\right)^{k/(k-1)} = (1033.3 \text{ kPa}) \left(\frac{2}{1.4+1}\right)^{1.4/0.4} = 545.9 \text{ kPa}$$

Then the pressure at the exit plane (throat) is

$$P_e = P_b$$
 for  $P_e = P^* = 545.9 \text{ kPa}$  for

for 
$$P_b < 545.9 \text{ kPa}$$
 (choked flow)

 $P_b \ge 545.9 \text{ kPa}$ 

Thus the back pressure does not affect the flow when  $100 < P_b < 545.9$  kPa. For a specified exit pressure  $P_e$ , the temperature, velocity, and mass flow rate are

Temperature 
$$T_e = T_0 \left(\frac{P_e}{P_0}\right)^{(k-1)/k} = (416.1 \text{ K}) \left(\frac{P_e}{1033.3}\right)^{0.4/1.2}$$

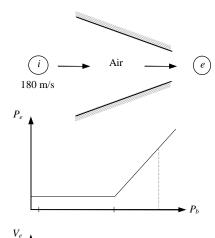
Velocity 
$$V = \sqrt{2c_p(T_0 - T_e)} = \sqrt{2(1.005 \text{ kJ/kg} \cdot \text{K})(416.1 - T_e) \left(\frac{1000 \text{ m}^2/\text{s}^2}{1 \text{ kJ/kg}}\right)}$$

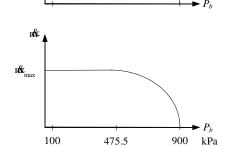
Speed of sound 
$$c_e = \sqrt{kRT_e} = \sqrt{(1.4)(0.287 \text{ kJ/kg} \cdot \text{K}) \left(\frac{1000 \text{ m}^2/\text{s}^2}{1 \text{ kJ/kg}}\right)}$$

Mach number  $Ma_e = V_e / c_e$ 

Density 
$$\rho_e = \frac{P_e}{RT_e} = \frac{P_e}{(0.287 \text{ kPa} \cdot \text{m}^3 / \text{kg} \cdot \text{K})T_e}$$

Mass flow rate  $n = \rho_e V_e A_e = \rho_e V_e (0.001 \,\mathrm{m}^2)$ 





$P_b$ , kPa	$P_b, P_0$	$P_e$ , kPa	$P_b$ , $P_0$	$T_e$ , K	$V_e$ , m/s	Ma	$\rho_{\rm e}$ , kg/m <sup>3</sup>	m%, kg/s
900	0.871	900	0.871	400.0	180.0	0.45	7.840	0
800	0.774	800	0.774	386.8	162.9	0.41	7.206	1.174
700	0.677	700	0.677	372.3	236.0	0.61	6.551	1.546
600	0.581	600	0.581	356.2	296.7	0.78	5.869	1.741
545.9	0.528	545.9	0.528	333.3	366.2	1.00	4.971	1.820
500	0.484	545.9	0.528	333.2	366.2	1.00	4.971	1.820
400	0.387	545.9	0.528	333.3	366.2	1.00	4.971	1.820
300	0.290	545.9	0.528	333.3	366.2	1.00	4.971	1.820
200	0.194	545.9	0.528	333.3	366.2	1.00	4.971	1.820
100	0.097	545.9	0.528	333.3	366.2	1.00	4.971	1.820

**Discussion** Once the back pressure drops below 545.0 kPa, the flow is choked, and not remain constant from then on.

**Solution** Nitrogen gas enters a converging nozzle. The properties at the nozzle exit are to be determined.

**Assumptions** 1 Nitrogen is an ideal gas with k = 1.4. 2 Flow through the nozzle is steady, one-dimensional, and isentropic.

Analysis The schematic of the duct is shown in Fig. 12–25. For isentropic flow through a duct, the area ratio  $A/A^*$  (the flow area over the area of the throat where Ma = 1) is also listed in Table A–13. At the initial Mach number of Ma = 0.3, we read

$$\frac{A_1}{A^*} = 2.0351,$$
  $\frac{T_1}{T_0} = 0.9823,$  and  $\frac{P_1}{P_0} = 0.9395$ 

With a 20 percent reduction in flow area,  $A_2 = 0.8A_1$ , and

$$\frac{A_2}{A^*} = \frac{A_2}{A_1} \frac{A_1}{A^*} = (0.8)(2.0351) = 1.6281$$

For this value of  $A_2/A^*$  from Table A–13, we read

$$\frac{T_2}{T_0} = 0.9791$$
,  $\frac{P_2}{P_0} = 0.8993$ , and Ma<sub>2</sub> =

Here we chose the subsonic Mach number for the calculated  $A_2/A^*$  instead of the supersonic one because the duct is converging in the flow direction and the initial flow is subsonic. Since the stagnation properties are constant for isentropic flow, we can write

$$\frac{T_2}{T_1} = \frac{T_2 / T_0}{T_1 / T_0} \longrightarrow T_2 = T_1 \left(\frac{T_2 / T_0}{T_1 / T_0}\right) = (400 \text{ K}) \left(\frac{0.9791}{0.9823}\right) = 399 \text{ K}$$

$$\frac{P_2}{P_1} = \frac{P_2 / P_0}{P_1 / P_0} \longrightarrow P_2 = P_1 \left(\frac{P_2 / P_0}{P_1 / P_0}\right) = (100 \text{ kPa}) \left(\frac{0.8993}{0.9395}\right) = \mathbf{95.7 K}$$

which are the temperature and pressure at the desired location.

**Discussion** Note that the temperature and pressure drop as the fluid accelerates in a converging nozzle.

**Solution** Nitrogen gas enters a converging nozzle. The properties at the nozzle exit are to be determined.

**Assumptions** 1 Nitrogen is an ideal gas with k = 1.4. 2 Flow through the nozzle is steady, one-dimensional, and isentropic.

Analysis The schematic of the duct is shown in Fig. 12–25. For isentropic flow through a duct, the area ratio  $A/A^*$  (the flow area over the area of the throat where Ma = 1) is also listed in Table A–13. At the initial Mach number of Ma = 0.5, we read

$$\frac{A_1}{A^*} = 1.3398$$
,  $\frac{T_1}{T_0} = 0.9524$ , and  $\frac{P_1}{P_0} = 0.8430$ 

With a 20 percent reduction in flow area,  $A_2 = 0.8A_1$ , and

$$\frac{A_2}{A^*} = \frac{A_2}{A_1} \frac{A_1}{A^*} = (0.8)(1.3398) = 1.0718$$

For this value of  $A_2/A^*$  from Table A–13, we read

$$\frac{T_2}{T_0} = 0.9010$$
,  $\frac{P_2}{P_0} = 0.6948$ , and Ma<sub>2</sub> = **0.740**

Here we chose the subsonic Mach number for the calculated  $A_2/A^*$  instead of the supersonic one because the duct is converging in the flow direction and the initial flow is subsonic. Since the stagnation properties are constant for isentropic flow, we can write

$$\frac{T_2}{T_1} = \frac{T_2 / T_0}{T_1 / T_0} \longrightarrow T_2 = T_1 \left(\frac{T_2 / T_0}{T_1 / T_0}\right) = (400 \text{ K}) \left(\frac{0.9010}{0.9524}\right) = 378 \text{ K}$$

$$\frac{P_2}{P_1} = \frac{P_2 / P_0}{P_1 / P_0} \longrightarrow P_2 = P_1 \left(\frac{P_2 / P_0}{P_1 / P_0}\right) = (100 \text{ kPa}) \left(\frac{0.6948}{0.8430}\right) = 82.4 \text{ K}$$

which are the temperature and pressure at the desired location.

**Discussion** Note that the temperature and pressure drop as the fluid accelerates in a converging nozzle.

**Solution** Nitrogen entering a converging-diverging nozzle experiences a normal shock. The pressure, temperature, velocity, Mach number, and stagnation pressure downstream of the shock are to be determined. The results are to be compared to those of air under the same conditions.

Assumptions 1 Nitrogen is an ideal gas with constant specific heats. 2 Flow through the nozzle is steady, one-dimensional, and isentropic. 3 The nozzle is adiabatic.

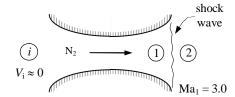
**Properties** The properties of nitrogen are  $R = 0.297 \text{ kJ/kg} \cdot \text{K}$  and k = 1.4.

**Analysis** The inlet stagnation properties in this case are identical to the inlet properties since the inlet velocity is negligible. Assuming the flow before the shock to be isentropic,

$$P_{01} = P_i = 620 \text{ kPa}$$
  
 $T_{01} = T_i = 310 \text{ K}$ 

Then,

$$T_1 = T_{01} \left( \frac{2}{2 + (k - 1)\text{Ma}_1^2} \right) = (310 \text{ K}) \left( \frac{2}{2 + (1.4 - 1)3^2} \right) = 110.7 \text{ K}$$



and

$$P_1 = P_{01} \left(\frac{T_1}{T_{01}}\right)^{k/(k-1)} = (620 \text{ kPa}) \left(\frac{110.7}{310}\right)^{1.4/0.4} = 16.88 \text{ kPa}$$

The fluid properties after the shock (denoted by subscript 2) are related to those before the shock through the functions listed in Table A-14. For  $Ma_1 = 3.0$  we read

$$\text{Ma}_2 = 0.4752 \cong \textbf{0.475}, \qquad \frac{P_{02}}{P_{01}} = 0.32834, \qquad \frac{P_2}{P_1} = 10.333, \quad \text{and} \quad \frac{T_2}{T_1} = 2.679$$

Then the stagnation pressure  $P_{02}$ , static pressure  $P_2$ , and static temperature  $T_2$ , are determined to be

$$P_{02}=0.32834P_{01}=(0.32834)(620~\mathrm{kPa})=203.6~\mathrm{kPa}\cong$$
 **204 kPa**  $P_2=10.333P_1=(10.333)(16.88~\mathrm{kPa})=174.4~\mathrm{kPa}\cong$  **174 kPa**  $T_2=2.679T_1=(2.679)(110.7~\mathrm{K})=296.6~\mathrm{K}\cong$  **297 K**

The velocity after the shock can be determined from  $V_2 = \text{Ma}_2 c_2$ , where  $c_2$  is the speed of sound at the exit conditions after the shock,

$$V_2 = \text{Ma}_2 c_2 = \text{Ma}_2 \sqrt{kRT_2} = (0.4752) \sqrt{(1.4)(0.297 \text{ kJ/kg} \cdot \text{K})(296.6 \text{ K}) \left(\frac{1000 \text{ m}^2/\text{s}^2}{1 \text{ kJ/kg}}\right)} = 166.9 \text{ m/s} \approx \textbf{167 m/s}$$

**Discussion** For air at specified conditions k = 1.4 (same as nitrogen) and R = 0.287 kJ/kg·K. Thus the only quantity which will be different in the case of air is the velocity after the normal shock, which happens to be 164.0 m/s.

**Solution** The diffuser of an aircraft is considered. The static pressure rise across the diffuser and the exit area are to be determined.

**Assumptions** 1 Air is an ideal gas with constant specific heats at room temperature. 2 Flow through the diffuser is steady, one-dimensional, and isentropic. 3 The diffuser is adiabatic.

**Properties** Air properties at room temperature are R = 0.287 kJ/kg·K,  $c_p = 1.005 \text{ kJ/kg·K}$ , and k = 1.4.

Analysis The inlet velocity is

$$V_1 = \text{Ma}_1 c_1 = M_1 \sqrt{kRT_1} = (0.9) \sqrt{(1.4)(0.287 \text{ kJ/kg} \cdot \text{K})(242.7 \text{ K}) \left(\frac{1000 \text{ m}^2 / \text{s}^2}{1 \text{ kJ/kg}}\right)} = 281.0 \text{ m/s}$$

Then the stagnation temperature and pressure at the diffuser inlet become

$$T_{01} = T_1 + \frac{V_1^2}{2c_p} = 242.7 + \frac{(281.0 \text{ m/s})^2}{2(1.005 \text{ kJ/kg} \cdot \text{K})} \left(\frac{1 \text{ kJ/kg}}{1000 \text{ m}^2/\text{s}^2}\right) = 282.0 \text{ K}$$

$$P_{01} = P_1 \left(\frac{T_{01}}{T_1}\right)^{k/(k-1)} = (41.1 \text{ kPa}) \left(\frac{282.0 \text{ K}}{242.7 \text{ K}}\right)^{1.4/(1.4-1)} = 69.50 \text{ kPa}$$

For an adiabatic diffuser, the energy equation reduces to  $h_{01} = h_{02}$ . Noting that  $h = c_p T$  and the specific heats are assumed to be constant, we have

$$T_{01} = T_{02} = T_0 = 282.0 \text{ K}$$

The isentropic relation between states 1 and 02 gives

$$P_{02} = P_{02} = P_1 \left(\frac{T_{02}}{T_1}\right)^{k/(k-1)} = (41.1 \,\text{kPa}) \left(\frac{282.0 \,\text{K}}{242.7 \,\text{K}}\right)^{1.4/(1.4-1)} = 69.50 \,\text{kPa}$$

The exit velocity can be expressed as

$$V_2 = \text{Ma}_2 c_2 = \text{Ma}_2 \sqrt{kRT_2} = (0.3)\sqrt{(1.4)(0.287 \text{ kJ/kg} \cdot \text{K}) T_2 \left(\frac{1000 \text{ m}^2 / \text{s}^2}{1 \text{ kJ/kg}}\right)} = 6.01\sqrt{T_2}$$

Thus 
$$T_2 = T_{02} - \frac{V_2^2}{2c_p} = (282.0) - \frac{6.01^2 \text{ T}_2 \text{ m}^2/\text{s}^2}{2(1.005 \text{ kJ/kg} \cdot \text{K})} \left( \frac{1 \text{ kJ/kg}}{1000 \text{ m}^2/\text{s}^2} \right) = 277.0 \text{ K}$$

Then the static exit pressure becomes

$$P_2 = P_{02} \left(\frac{T_2}{T_{02}}\right)^{k/(k-1)} = (69.50 \text{ kPa}) \left(\frac{277.0 \text{ K}}{282.0 \text{ K}}\right)^{1.4/(1.4-1)} = 65.28 \text{ kPa}$$

Thus the static pressure rise across the diffuser is

$$\Delta P = P_2 - P_1 = 65.28 - 41.1 =$$
**24.2 kPa**

Also, 
$$\rho_2 = \frac{P_2}{RT_2} = \frac{65.28 \,\text{kPa}}{(0.287 \,\text{kPa} \cdot \text{m}^3/\text{kg} \cdot \text{K})(277.0 \,\text{K})} = 0.8211 \,\text{kg/m}^3$$

$$V_2 = 6.01\sqrt{T_2} = 6.01\sqrt{277.0} = 100.0 \text{ m/s}$$

Thus 
$$A_2 = \frac{n k}{\rho_2 V_2} = \frac{38 \text{ kg/s}}{(0.8211 \text{ kg/m}^3)(100.0 \text{ m/s})} = \textbf{0.463 m}^2$$

**Discussion** The pressure rise in actual diffusers will be lower because of the irreversibilities. However, flow through well-designed diffusers is very nearly isentropic.

**Solution** The critical temperature, pressure, and density of an equimolar mixture of oxygen and nitrogen for specified stagnation properties are to be determined.

**Assumptions** Both oxygen and nitrogen are ideal gases with constant specific heats at room temperature.

**Properties** The specific heat ratio and molar mass are k = 1.395 and M = 32 kg/kmol for oxygen, and k = 1.4 and M = 28 kg/kmol for nitrogen.

**Analysis** The gas constant of the mixture is

$$M_m = y_{O_2} M_{O_2} + y_{N_2} M_{N_2} = 0.5 \times 32 + 0.5 \times 28 = 30 \text{ kg/kmol}$$

$$R_m = \frac{R_u}{M_m} = \frac{8.314 \text{ kJ/kmol} \cdot \text{K}}{30 \text{ kg/kmol}} = 0.2771 \text{ kJ/kg} \cdot \text{K}$$

The specific heat ratio is 1.4 for nitrogen, and nearly 1.4 for oxygen. Therefore, the specific heat ratio of the mixture is also 1.4. Then the critical temperature, pressure, and density of the mixture become

$$T^* = T_0 \left(\frac{2}{k+1}\right) = (550 \text{ K}) \left(\frac{2}{1.4+1}\right) = 458.3 \text{ K} \cong \textbf{458 K}$$

$$P^* = P_0 \left(\frac{2}{k+1}\right)^{k/(k-1)} = (350 \text{ kPa}) \left(\frac{2}{1.4+1}\right)^{1.4/(1.4-1)} = 184.9 \text{ kPa} \cong \textbf{185 kPa}$$

$$\rho^* = \frac{P^*}{RT^*} = \frac{184.9 \text{ kPa}}{(0.2771 \text{ kPa} \cdot \text{m}^3/\text{kg} \cdot \text{K})(458.3 \text{ K})} = \textbf{1.46 kg/m}^3$$

**Discussion** If the specific heat ratios k of the two gases were different, then we would need to determine the k of the mixture from  $k = c_{p,m}/c_{v,m}$  where the specific heats of the mixture are determined from

$$\begin{split} &C_{p,m} = \mathrm{mf}_{O_2} c_{p,O_2} + \mathrm{mf}_{N_2} c_{p,N_2} = (y_{O_2} M_{O_2} / M_m) c_{p,O_2} + (y_{N_2} M_{N_2} / M_m) c_{p,N_2} \\ &C_{v,m} = \mathrm{mf}_{O_2} c_{v,O_2} + \mathrm{mf}_{N_2} c_{v,N_2} = (y_{O_2} M_{O_2} / M_m) c_{v,O_2} + (y_{N_2} M_{N_2} / M_m) c_{v,N_2} \end{split}$$

where mf is the mass fraction and y is the mole fraction. In this case it would give

$$c_{nm} = (0.5 \times 32/30) \times 0.918 + (0.5 \times 28/30) \times 1.039 = 0.974 \text{ kJ/kg.K}$$

$$c_{p,m} = (0.5 \times 32/30) \times 0.658 + (0.5 \times 28/30) \times 0.743 = 0.698 \text{ kJ/kg.K}$$

and

$$k = 0.974/0.698 = 1.40$$

# 12-133E

Solution Helium gas is accelerated in a nozzle. For a specified mass flow rate, the throat and exit areas of the nozzle are to be determined for the case of isentropic nozzle.

Assumptions 1 Helium is an ideal gas with constant specific heats. 2 Flow through the nozzle is steady, one-dimensional, and isentropic. 3 The nozzle is adiabatic.

The properties of helium are R = 0.4961 Btu/lbm·R = 2.6809 psia·ft<sup>3</sup>/lbm·R,  $c_p = 1.25$  Btu/lbm·R, and k = 1.25 Btu/lbm·R, and an analysis Btu/lbm·R, an an analysis Btu/lbm·R, an a **Properties** 1.667.

Analysis The inlet stagnation properties in this case are identical to the inlet properties since the inlet velocity is negligible,

$$T_{01} = T_1 = 740 \,\mathrm{R}$$
  
 $P_{01} = P_1 = 220 \,\mathrm{psia}$ 

The flow is assumed to be isentropic, thus the stagnation temperature and pressure remain constant throughout the nozzle,

$$T_{02} = T_{01} = 740 \text{ R}$$
  
 $P_{02} = P_{01} = 220 \text{ psia}$ 

The critical pressure and temperature are determined from

$$T^* = T_0 \left(\frac{2}{k+1}\right) = (740 \text{ R}) \left(\frac{2}{1.667+1}\right) = 554.9 \text{ R}$$

$$P^* = P_0 \left(\frac{2}{k+1}\right)^{k/(k-1)} = (220 \text{ psia}) \left(\frac{2}{1.667+1}\right)^{1.667/(1.667-1)} = 107.2 \text{ psia}$$

$$\rho^* = \frac{P^*}{RT^*} = \frac{107.2 \text{ psia}}{(2.6809 \text{ psia} \cdot \text{ft}^3/\text{lbm} \cdot \text{R})(554.9 \text{ R})} = 0.07203 \text{ lbm/ft}^3$$

$$V^* = c^* = \sqrt{kRT^*} = \sqrt{(1.667)(0.4961 \text{ Btu/lbm} \cdot \text{R})(554.9 \text{ R})} \left(\frac{25,037 \text{ ft}^2/\text{s}^2}{\text{s}^2}\right)$$

$$V^* = c^* = \sqrt{kRT^*} = \sqrt{(1.667)(0.4961 \,\text{Btu/lbm} \cdot \text{R})(554.9 \,\text{R}) \left(\frac{25,037 \,\text{ft}^2 / \text{s}^2}{1 \,\text{Btu/lbm}}\right)} = 3390 \,\text{ft/s}$$

$$A^* = \frac{n^2 v}{\rho^* V^*} = \frac{0.2 \text{ 1bm/s}}{(0.072031 \text{bm/ft}^3)(3390 \text{ ft/s})} = 8.19 \times 10^{-4} \text{ ft}^2$$

At the nozzle exit the pressure is 
$$P_2 = 15$$
 psia. Then the other properties at the nozzle exit are determined to be 
$$\frac{p_0}{p_2} = \left(1 + \frac{k-1}{2} \text{Ma}_2^2\right)^{k/(k-1)} \longrightarrow \frac{220 \text{ psia}}{15 \text{ psia}} = \left(1 + \frac{1.667 - 1}{2} \text{Ma}_2^2\right)^{1.667/0.667}$$

It yields  $Ma_2 = 2.405$ , which is greater than 1. Therefore, the nozzle must be converging-diverging.

$$T_2 = T_0 \left( \frac{2}{2 + (k - 1)\text{Ma}_2^2} \right) = (740 \text{ R}) \left( \frac{2}{2 + (1.667 - 1) \times 2.405^2} \right) = 252.6 \text{ R}$$

$$\rho_2 = \frac{P_2}{RT_2} = \frac{15 \text{ psia}}{(2.6809 \text{ psia} \cdot \text{ft}^3/\text{lbm} \cdot \text{R})(252.6 \text{ R})} = 0.02215 \text{ lbm/ft}^3$$

$$V_2 = \text{Ma}_2 c_2 = \text{Ma}_2 \sqrt{kRT_2} = (2.405) \sqrt{(1.667)(0.4961 \,\text{Btu/lbm} \cdot \text{R})(252.6 \,\text{R}) \left(\frac{25,037 \,\text{ft}^2 / \text{s}^2}{1 \,\text{Btu/lbm}}\right)} = 5500 \,\text{ft/s}$$

Thus the exit area is

$$A_2 = \frac{n^{8}c}{\rho_2 V_2} = \frac{0.2 \text{ lbm/s}}{(0.02215 \text{ lbm/ft}^3)(5500 \text{ ft/s})} = \textbf{0.00164 ft}^2$$

Discussion Flow areas in actual nozzles would be somewhat larger to accommodate the irreversibilities.



**Solution** Using the compressible flow relations, the one-dimensional compressible flow functions are to be evaluated and tabulated as in Table A-13 for an ideal gas with k = 1.667.

**Properties** The specific heat ratio of the ideal gas is given to be k = 1.667.

Analysis The compressible flow functions listed below are expressed in EES and the results are tabulated.

$$Ma^* = Ma \sqrt{\frac{k+1}{2 + (k-1)Ma^2}}$$

$$\frac{P}{P_0} = \left(1 + \frac{k-1}{2}Ma^2\right)^{-k/(k-1)}$$

$$\frac{P}{P_0} = \left(1 + \frac{k-1}{2}Ma^2\right)^{-k/(k-1)}$$

$$\frac{P}{P_0} = \left(1 + \frac{k-1}{2}Ma^2\right)^{-1/(k-1)}$$

$$\frac{P}{P_0} = \left(1 + \frac{k-1}{2}Ma^2\right)^{-1/(k-1)}$$

k=1.667

 $PP0=(1+(k-1)*M^2/2)^{-(-k/(k-1))}$ 

TT0=1/(1+(k-1)\*M^2/2)

 $DD0=(1+(k-1)*M^2/2)^{-1/(k-1)}$ 

 $Mcr=M*SQRT((k+1)/(2+(k-1)*M^2))$ 

 $AAcr=((2/(k+1))*(1+0.5*(k-1)*M^2))^{(0.5*(k+1)/(k-1))/M}$ 

Ma	Ma*	$A/A^*$	$P/P_0$	$ ho/ ho_0$	$T/T_0$
0.0	0	∞	1.0000	1.0000	1.0000
0.1	0.1153	5.6624	0.9917	0.9950	0.9967
0.2	0.2294	2.8879	0.9674	0.9803	0.9868
0.3	0.3413	1.9891	0.9288	0.9566	0.9709
0.4	0.4501	1.5602	0.8782	0.9250	0.9493
0.5	0.5547	1.3203	0.8186	0.8869	0.9230
0.6	0.6547	1.1760	0.7532	0.8437	0.8928
0.7	0.7494	1.0875	0.6850	0.7970	0.8595
0.8	0.8386	1.0351	0.6166	0.7482	0.8241
0.9	0.9222	1.0081	0.5501	0.6987	0.7873
1.0	1.0000	1.0000	0.4871	0.6495	0.7499
1.2	1.1390	1.0267	0.3752	0.5554	0.6756
1.4	1.2572	1.0983	0.2845	0.4704	0.6047
1.6	1.3570	1.2075	0.2138	0.3964	0.5394
1.8	1.4411	1.3519	0.1603	0.3334	0.4806
2.0	1.5117	1.5311	0.1202	0.2806	0.4284
2.2	1.5713	1.7459	0.0906	0.2368	0.3825
2.4	1.6216	1.9980	0.0686	0.2005	0.3424
2.6	1.6643	2.2893	0.0524	0.1705	0.3073
2.8	1.7007	2.6222	0.0403	0.1457	0.2767
3.0	1.7318	2.9990	0.0313	0.1251	0.2499
5.0	1.8895	9.7920	0.0038	0.0351	0.1071
$\infty$	1.9996	8	0	0	0

**Discussion** The tabulated values are useful for quick calculations, but be careful – they apply only to one specific value of k, in this case k = 1.667.



**Solution** Using the normal shock relations, the normal shock functions are to be evaluated and tabulated as in Table A-14 for an ideal gas with k = 1.667.

**Properties** The specific heat ratio of the ideal gas is given to be k = 1.667.

Analysis The normal shock relations listed below are expressed in EES and the results are tabulated.

$$\begin{split} \operatorname{Ma}_2 &= \sqrt{\frac{(k-1)\operatorname{Ma}_1^2 + 2}{2k\operatorname{Ma}_1^2 - k + 1}} & \frac{P_2}{P_1} &= \frac{1 + k\operatorname{Ma}_1^2}{1 + k\operatorname{Ma}_2^2} = \frac{2k\operatorname{Ma}_1^2 - k + 1}{k + 1} \\ \frac{T_2}{T_1} &= \frac{2 + \operatorname{Ma}_1^2(k-1)}{2 + \operatorname{Ma}_2^2(k-1)} & \frac{\rho_2}{\rho_1} &= \frac{P_2 / P_1}{T_2 / T_1} = \frac{(k+1)\operatorname{Ma}_1^2}{2 + (k-1)\operatorname{Ma}_1^2} = \frac{V_1}{V_2} \,, \\ \frac{P_{02}}{P_{01}} &= \frac{\operatorname{Ma}_1}{\operatorname{Ma}_2} \left[ \frac{1 + \operatorname{Ma}_2^2(k-1) / 2}{1 + \operatorname{Ma}_1^2(k-1) / 2} \right]^{\frac{k+1}{2(k-1)}} & \frac{P_{02}}{P_1} &= \frac{(1 + k\operatorname{Ma}_1^2)[1 + \operatorname{Ma}_2^2(k-1) / 2]^{k / (k-1)}}{1 + k\operatorname{Ma}_2^2} \end{split}$$

k=1.667

 $My=SQRT((Mx^2+2/(k-1))/(2*Mx^2*k/(k-1)-1))$ 

 $PyPx=(1+k*Mx^2)/(1+k*My^2)$ 

 $TyTx = (1+Mx^2*(k-1)/2)/(1+My^2*(k-1)/2)$ 

RyRx=PyPx/TyTx

 $P0yP0x = (Mx/My)*((1 + My^2*(k-1)/2)/(1 + Mx^2*(k-1)/2))^*(0.5*(k+1)/(k-1))$ 

 $P0yPx = (1+k*Mx^2)*(1+My^2*(k-1)/2)^(k/(k-1))/(1+k*My^2)$ 

Ma <sub>1</sub>	Ma <sub>2</sub>	$P_2/P_1$	$\rho_2/\rho_1$	$T_2/T_1$	$P_{02}/P_{01}$	$P_{02}/P_1$
1.0	1.0000	1.0000	1.0000	1.0000	1	2.0530
1.1	0.9131	1.2625	1.1496	1.0982	0.999	2.3308
1.2	0.8462	1.5500	1.2972	1.1949	0.9933	2.6473
1.3	0.7934	1.8626	1.4413	1.2923	0.9813	2.9990
1.4	0.7508	2.2001	1.5805	1.3920	0.9626	3.3838
1.5	0.7157	2.5626	1.7141	1.4950	0.938	3.8007
1.6	0.6864	2.9501	1.8415	1.6020	0.9085	4.2488
1.7	0.6618	3.3627	1.9624	1.7135	0.8752	4.7278
1.8	0.6407	3.8002	2.0766	1.8300	0.8392	5.2371
1.9	0.6227	4.2627	2.1842	1.9516	0.8016	5.7767
2.0	0.6070	4.7503	2.2853	2.0786	0.763	6.3462
2.1	0.5933	5.2628	2.3802	2.2111	0.7243	6.9457
2.2	0.5814	5.8004	2.4689	2.3493	0.6861	7.5749
2.3	0.5708	6.3629	2.5520	2.4933	0.6486	8.2339
2.4	0.5614	6.9504	2.6296	2.6432	0.6124	8.9225
2.5	0.5530	7.5630	2.7021	2.7989	0.5775	9.6407
2.6	0.5455	8.2005	2.7699	2.9606	0.5442	10.3885
2.7	0.5388	8.8631	2.8332	3.1283	0.5125	11.1659
2.8	0.5327	9.5506	2.8923	3.3021	0.4824	11.9728
2.9	0.5273	10.2632	2.9476	3.4819	0.4541	12.8091
3.0	0.5223	11.0007	2.9993	3.6678	0.4274	13.6750
4.0	0.4905	19.7514	3.3674	5.8654	0.2374	23.9530
5.0	0.4753	31.0022	3.5703	8.6834	0.1398	37.1723
$\infty$	0.4473	8	3.9985	$\infty$	0	$\infty$

**Discussion** The tabulated values are useful for quick calculations, but be careful – they apply only to one specific value of k, in this case k = 1.667.

 $\begin{array}{cccc}
1 & \text{He} & * & -\\
V_{i} \approx 0 & & & & \\
\end{array}$ 

#### 12-136

**Solution** Helium gas is accelerated in a nozzle isentropically. For a specified mass flow rate, the throat and exit areas of the nozzle are to be determined.

**Assumptions** 1 Helium is an ideal gas with constant specific heats. 2 Flow through the nozzle is steady, one-dimensional, and isentropic. 3 The nozzle is adiabatic.

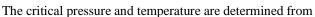
**Properties** The properties of helium are R = 2.0769 kJ/kg.K,  $c_p = 5.1926 \text{ kJ/kg.K}$ , and k = 1.667.

**Analysis** The inlet stagnation properties in this case are identical to the inlet properties since the inlet velocity is negligible,

$$T_{01} = T_1 = 500 \text{ K}$$
  
 $P_{01} = P_1 = 1.0 \text{ MPa}$ 

The flow is assumed to be isentropic, thus the stagnation temperature and pressure remain constant throughout the nozzle,

$$T_{02} = T_{01} = 500 \text{ K}$$
  
 $P_{02} = P_{01} = 1.0 \text{ MPa}$ 



$$\begin{split} T^* &= T_0 \bigg(\frac{2}{k+1}\bigg) = (500 \text{ K}) \bigg(\frac{2}{1.667+1}\bigg) = 375.0 \text{ K} \\ P^* &= P_0 \bigg(\frac{2}{k+1}\bigg)^{k/(k-1)} = (1.0 \text{ MPa}) \bigg(\frac{2}{1.667+1}\bigg)^{1.667/(1.667-1)} = 0.487 \text{ MPa} \\ \rho^* &= \frac{P^*}{RT^*} = \frac{487 \text{ kPa}}{(2.0769 \text{ kPa} \cdot \text{m}^3/\text{kg} \cdot \text{K})(375 \text{ K})} = 0.625 \text{ kg/m}^3 \\ V^* &= c^* = \sqrt{kRT^*} = \sqrt{(1.667)(2.0769 \text{ kJ/kg} \cdot \text{K})(375 \text{ K})} \bigg(\frac{1000 \text{ m}^2/\text{s}^2}{1 \text{ kJ/kg}}\bigg) = 1139.4 \text{ m/s} \end{split}$$

Thus the throat area is

$$A^* = \frac{m^2 k}{\rho * V *} = \frac{0.46 \text{ kg/s}}{(0.625 \text{ kg/m}^3)(1139.4 \text{ m/s})} = 6.460 \times 10^{-4} \text{ m}^2 = \textbf{6.460} \times 10^{-4} \text{ m}^2$$

At the nozzle exit the pressure is  $P_2 = 0.1$  MPa. Then the other properties at the nozzle exit are determined to be

$$\frac{P_0}{P_2} = \left(1 + \frac{k - 1}{2} \operatorname{Ma}_2^2\right)^{k/(k - 1)} \longrightarrow \frac{1.0 \operatorname{MPa}}{0.1 \operatorname{MPa}} = \left(1 + \frac{1.667 - 1}{2} \operatorname{Ma}_2^2\right)^{1.667/0.667}$$

It yields  $Ma_2 = 2.130$ , which is greater than 1. Therefore, the nozzle must be converging-diverging.

$$T_2 = T_0 \left(\frac{2}{2 + (k - 1)\text{Ma}_2^2}\right) = (500 \text{ K}) \left(\frac{2}{2 + (1.667 - 1) \times 2.13^2}\right) = 199.0 \text{ K}$$

$$\rho_2 = \frac{P_2}{RT_2} = \frac{100 \text{ kPa}}{(2.0769 \text{ kPa} \cdot \text{m}^3/\text{kg} \cdot \text{K})(199 \text{ K})} = 0.242 \text{ kg/m}^3$$

$$V_2 = \text{Ma}_2 c_2 = \text{Ma}_2 \sqrt{kRT_2} = (2.13) \sqrt{(1.667)(2.0769 \text{ kJ/kg} \cdot \text{K})(199 \text{ K}) \left(\frac{1000 \text{ m}^2 / \text{s}^2}{1 \text{ kJ/kg}}\right)} = 1768.0 \text{ m/s}$$

Thus the exit area is

$$A_2 = \frac{n^{\rm lk}}{\rho_2 V_2} = \frac{0.46\,{\rm kg/s}}{(0.242\,{\rm kg/m^3})(1768\,{\rm m/s})} = 0.1075 \times 10^{-3}\,{\rm m^2} \cong {\bf 10.8\,cm^2}$$

**Discussion** Flow areas in actual nozzles would be somewhat larger to accommodate the irreversibilities.

**Solution** The flow velocity of air in a channel is to be measured using a Pitot-static probe, which causes a shock wave to occur. For measured values of static pressure before the shock and stagnation pressure and temperature after the shock, the flow velocity before the shock is to be determined.

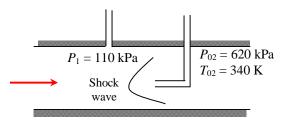
Assumptions 1 Air is an ideal gas with constant specific heats. 2 Flow through the nozzle is steady and one-dimensional.

**Properties** The specific heat ratio of air at room temperature is k=1.4.

Analysis The nose of the probe is rounded (instead of being pointed), and thus it will cause a bow shock wave to form. Bow shocks are difficult to analyze. But they are normal to the body at the nose, and thus we can approximate them as normal shocks in the vicinity of the probe. It is given that the static pressure before the shock is  $P_1 = 110$  kPa, and the stagnation pressure and temperature after the shock are  $P_{02} = 620$  kPa, and  $T_{02} = 340$  K. Noting that the stagnation temperature remains constant, we have

$$T_{01} = T_{02} = 340 \text{ K}$$
  
Also,  $\frac{P_{02}}{P_1} = \frac{620 \text{ kPa}}{110 \text{ kPa}} = 5.6364 \approx 5.64$ 

The fluid properties after the shock are related to those before the shock through the functions listed in Table A-14. For  $P_{02}$  /  $P_1$  = 5.64 we read



$$Ma_1 = 2.0$$
,  $Ma_2 = 0.5774$ ,  $\frac{P_{02}}{P_{01}} = 0.7209$ ,  $\frac{V_1}{V_2} = \frac{\rho_2}{\rho_1} = 2.6667$ ,

Then the stagnation pressure and temperature before the shock become

$$P_{01} = P_{02} / 0.7209 = (620 \text{ kPa}) / 0.7209 = 860 \text{ kPa}$$

$$T_1 = T_{01} \left(\frac{P_1}{P_{01}}\right)^{(k-1)/k} = (340 \text{ K}) \left(\frac{110 \text{ kPa}}{860 \text{ kPa}}\right)^{(1.4-1)/1.4} = 188.9 \text{ K}$$

The flow velocity before the shock can be determined from  $V_1 = \text{Ma}_1 c_1$ , where  $c_1$  is the speed of sound before the shock,

$$c_1 = \sqrt{kRT_1} = \sqrt{(1.4)(0.287 \text{ kJ/kg} \cdot \text{K})(188.9 \text{ K}) \left(\frac{1000 \text{ m}^2/\text{s}^2}{1 \text{ kJ/kg}}\right)} = 275.5 \text{ m/s}$$

$$V_1 = \text{Ma}_1 c_1 = 2(275.5 \,\text{m/s}) = 551 \,\text{m/s}$$

**Discussion** The flow velocity after the shock is  $V_2 = V_1/2.6667 = 551/2.6667 = 207$  m/s. Therefore, the velocity measured by a Pitot-static probe would be very different that the flow velocity.



**Solution** Using the normal shock relations, the normal shock functions are to be evaluated and tabulated as in Table

A-14 for air.

**Properties** The specific heat ratio is given to be k = 1.4 for air.

Analysis The normal shock relations listed below are expressed in EES and the results are tabulated.

$$\begin{split} \mathbf{Ma}_2 &= \sqrt{\frac{(k-1)\mathbf{Ma}_1^2 + 2}{2k\mathbf{Ma}_1^2 - k + 1}} & \frac{P_2}{P_1} = \frac{1 + k\mathbf{Ma}_1^2}{1 + k\mathbf{Ma}_2^2} = \frac{2k\mathbf{Ma}_1^2 - k + 1}{k + 1} \\ \frac{T_2}{T_1} &= \frac{2 + \mathbf{Ma}_1^2(k-1)}{2 + \mathbf{Ma}_2^2(k-1)} & \frac{\rho_2}{\rho_1} = \frac{P_2 / P_1}{T_2 / T_1} = \frac{(k+1)\mathbf{Ma}_1^2}{2 + (k-1)\mathbf{Ma}_1^2} = \frac{V_1}{V_2} \,, \\ \frac{P_{02}}{P_{01}} &= \frac{\mathbf{Ma}_1}{\mathbf{Ma}_2} \left[ \frac{1 + \mathbf{Ma}_2^2(k-1) / 2}{1 + \mathbf{Ma}_1^2(k-1) / 2} \right]^{\frac{k+1}{2(k-1)}} & \frac{P_{02}}{P_1} = \frac{(1 + k\mathbf{Ma}_1^2)[1 + \mathbf{Ma}_2^2(k-1) / 2]^{k/(k-1)}}{1 + k\mathbf{Ma}_2^2} \end{split}$$

Air:

k=1.4

 $My=SQRT((Mx^2+2/(k-1))/(2*Mx^2*k/(k-1)-1))$ 

 $PyPx=(1+k*Mx^2)/(1+k*My^2)$ 

 $TyTx = (1+Mx^2*(k-1)/2)/(1+My^2*(k-1)/2)$ 

RyRx=PyPx/TyTx

 $P0yP0x = (Mx/My)*((1+My^2*(k-1)/2)/(1+Mx^2*(k-1)/2))^*(0.5*(k+1)/(k-1))$ 

 $P0yPx = (1+k*Mx^2)*(1+My^2*(k-1)/2)^(k/(k-1))/(1+k*My^2)$ 

$Ma_1$	$Ma_2$	$P_2/P_1$	$ ho_2/ ho_1$	$T_2/T_1$	$P_{02}/P_{01}$	$P_{02}/P_1$
1.0	1.0000	1.0000	1.0000	1.0000	1	1.8929
1.5	0.7011	2.4583	1.8621	1.3202	0.9298	3.4133
2.0	0.5774	4.5000	2.6667	1.6875	0.7209	5.6404
2.5	0.5130	7.1250	3.3333	2.1375	0.499	8.5261
3.0	0.4752	10.3333	3.8571	2.6790	0.3283	12.0610
3.5	0.4512	14.1250	4.2609	3.3151	0.2129	16.2420
4.0	0.4350	18.5000	4.5714	4.0469	0.1388	21.0681
4.5	0.4236	23.4583	4.8119	4.8751	0.0917	26.5387
5.0	0.4152	29.0000	5.0000	5.8000	0.06172	32.6535
5.5	0.4090	35.1250	5.1489	6.8218	0.04236	39.4124
6.0	0.4042	41.8333	5.2683	7.9406	0.02965	46.8152
6.5	0.4004	49.1250	5.3651	9.1564	0.02115	54.8620
7.0	0.3974	57.0000	5.4444	10.4694	0.01535	63.5526
7.5	0.3949	65.4583	5.5102	11.8795	0.01133	72.8871
8.0	0.3929	74.5000	5.5652	13.3867	0.008488	82.8655
8.5	0.3912	84.1250	5.6117	14.9911	0.006449	93.4876
9.0	0.3898	94.3333	5.6512	16.6927	0.004964	104.7536
9.5	0.3886	105.1250	5.6850	18.4915	0.003866	116.6634
10.0	0.3876	116.5000	5.7143	20.3875	0.003045	129.2170

**Discussion** The tabulated values are useful for quick calculations, but be careful – they apply only to one specific value of k, in this case k = 1.4.



**Solution** Using the normal shock relations, the normal shock functions are to be evaluated and tabulated as in Table A-14 for methane.

**Properties** The specific heat ratio is given to be k = 1.3 for methane.

Analysis The normal shock relations listed below are expressed in EES and the results are tabulated.

$$\begin{split} \mathbf{Ma}_2 &= \sqrt{\frac{(k-1)\mathbf{Ma}_1^2 + 2}{2k\mathbf{Ma}_1^2 - k + 1}} & \frac{P_2}{P_1} = \frac{1 + k\mathbf{Ma}_1^2}{1 + k\mathbf{Ma}_2^2} = \frac{2k\mathbf{Ma}_1^2 - k + 1}{k + 1} \\ \frac{T_2}{T_1} &= \frac{2 + \mathbf{Ma}_1^2(k-1)}{2 + \mathbf{Ma}_2^2(k-1)} & \frac{\rho_2}{\rho_1} = \frac{P_2 / P_1}{T_2 / T_1} = \frac{(k+1)\mathbf{Ma}_1^2}{2 + (k-1)\mathbf{Ma}_1^2} = \frac{V_1}{V_2} \,, \\ \frac{P_{02}}{P_{01}} &= \frac{\mathbf{Ma}_1}{\mathbf{Ma}_2} \left[ \frac{1 + \mathbf{Ma}_2^2(k-1) / 2}{1 + \mathbf{Ma}_1^2(k-1) / 2} \right]^{\frac{k+1}{2(k-1)}} & \frac{P_{02}}{P_1} = \frac{(1 + k\mathbf{Ma}_1^2)[1 + \mathbf{Ma}_2^2(k-1) / 2]^{k/(k-1)}}{1 + k\mathbf{Ma}_2^2} \end{split}$$

# Methane:

k=1.3

 $My=SQRT((Mx^2+2/(k-1))/(2*Mx^2*k/(k-1)-1))$ 

 $PyPx=(1+k*Mx^2)/(1+k*My^2)$ 

 $TyTx = (1+Mx^2*(k-1)/2)/(1+My^2*(k-1)/2)$ 

RyRx=PyPx/TyTx

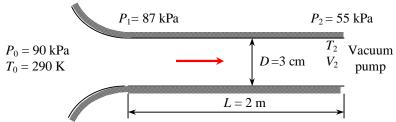
 $P0yP0x = (Mx/My)*((1+My^2*(k-1)/2)/(1+Mx^2*(k-1)/2))^*(0.5*(k+1)/(k-1))$ 

 $P0yPx = (1+k*Mx^2)*(1+My^2*(k-1)/2)^(k/(k-1))/(1+k*My^2)$ 

$Ma_1$	$Ma_2$	$P_2/P_1$	$\rho_2/\rho_1$	$T_2/T_1$	$P_{02}/P_{01}$	$P_{02}/P_1$
1.0	1.0000	1.0000	1.0000	1.0000	1	1.8324
1.5	0.6942	2.4130	1.9346	1.2473	0.9261	3.2654
2.0	0.5629	4.3913	2.8750	1.5274	0.7006	5.3700
2.5	0.4929	6.9348	3.7097	1.8694	0.461	8.0983
3.0	0.4511	10.0435	4.4043	2.2804	0.2822	11.4409
3.5	0.4241	13.7174	4.9648	2.7630	0.1677	15.3948
4.0	0.4058	17.9565	5.4118	3.3181	0.09933	19.9589
4.5	0.3927	22.7609	5.7678	3.9462	0.05939	25.1325
5.0	0.3832	28.1304	6.0526	4.6476	0.03613	30.9155
5.5	0.3760	34.0652	6.2822	5.4225	0.02243	37.3076
6.0	0.3704	40.5652	6.4688	6.2710	0.01422	44.3087
6.5	0.3660	47.6304	6.6218	7.1930	0.009218	51.9188
7.0	0.3625	55.2609	6.7485	8.1886	0.006098	60.1379
7.5	0.3596	63.4565	6.8543	9.2579	0.004114	68.9658
8.0	0.3573	72.2174	6.9434	10.4009	0.002827	78.4027
8.5	0.3553	81.5435	7.0190	11.6175	0.001977	88.4485
9.0	0.3536	91.4348	7.0837	12.9079	0.001404	99.1032
9.5	0.3522	101.8913	7.1393	14.2719	0.001012	110.367
10.0	0.3510	112.9130	7.1875	15.7096	0.000740	122.239

**Discussion** The tabulated values are useful for quick calculations, but be careful – they apply only to one specific value of k, in this case k = 1.3.

**Solution** Air enters a constant-area adiabatic duct at a specified state, and leaves at a specified pressure. The mass flow rate of air, the exit velocity, and the average friction factor are to be determined.



Assumptions 1 The assumptions associated with Fanno flow (i.e., steady, frictional flow of an ideal gas with constant properties through a constant cross-sectional area adiabatic duct) are valid. 2 The friction factor is constant along the duct.

**Properties** We take the properties of air to be k = 1.4,  $c_p = 1.005$  kJ/kg·K, and R = 0.287 kJ/kg·K. The friction factor is given to be f = 0.025.

**Analysis** Noting that the flow in the nozzle section is isentropic, the Mach number, thermodynamic temperature, and density at the tube inlet become

$$P_{1} = P_{01} \left( 1 + \frac{k-1}{2} \operatorname{Ma}_{1}^{2} \right)^{-k/(k-1)} \rightarrow 87 \, \text{kPa} = (90 \, \text{kPa}) \left( 1 + \frac{1.4-1}{2} \operatorname{Ma}_{1}^{2} \right)^{-1.4/0.4} \rightarrow \operatorname{Ma}_{1} = 0.2206$$

$$T_{1} = T_{01} \left( 1 + \frac{k-1}{2} \operatorname{Ma}_{1}^{2} \right)^{-1} = (290 \, \text{K}) \left( 1 + \frac{1.4-1}{2} (0.2206)^{2} \right)^{-1} = 287.2 \, \text{K}$$

$$\rho_{1} = \frac{P_{1}}{RT_{1}} = \frac{87 \, \text{kPa}}{(0.287 \, \text{kJ/kgK})(287.2 \, \text{K})} = 1.055 \, \text{kg/m}^{3}$$

Then the inlet velocity and the mass flow rate become

$$c_1 = \sqrt{kRT_1} = \sqrt{(1.4)(0.287 \text{ kJ/kg} \cdot \text{K})(287.2 \text{ K}) \left(\frac{1000 \text{ m}^2/\text{s}^2}{1 \text{ kJ/kg}}\right)} = 339.7 \text{ m/s}$$

$$V_1 = \text{Ma}_1 c_1 = 0.2206(339.7 \text{ m/s}) = 74.94 \text{ m/s}$$

$$m_{air}^{g} = \rho_1 A_{cl} V_1 = (1.055 \text{ kg/m}^3) [\pi (0.03 \text{ m})^2 / 4] (74.94 \text{ m/s}) = \textbf{0.0559 kg/s}$$

The Fanno flow functions corresponding to the inlet Mach number are, from Table A-16 (we used analytical relations),

$$Ma_1 = 0.2206$$
:  $(fL^*/D_h)_1 = 11.520$   $T_1/T^* = 1.1884$ ,  $P_1/P^* = 4.9417$ ,  $V_1/V^* = 0.2405$ 

Therefore,  $P_1 = 4.9417P^*$ . Then the Fanno function  $P_2/P^*$  becomes

$$\frac{P_2}{P^*} = \frac{P_2}{P_1 / 5.2173} = \frac{4.9417(55 \text{ kPa})}{87 \text{ kPa}} = 3.124$$

The corresponding Mach number and Fanno flow functions are, from Table A-16,

$$Ma_2 = 0.3465$$
,  $(fL^*/D_h)_2 = 3.5536$ , and  $V_2/V^* = 0.3751$ .

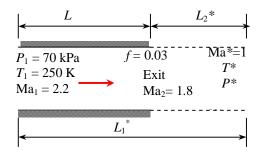
Then the air velocity at the duct exit and the average friction factor become

$$\frac{V_2}{V_1} = \frac{V_2 / V^*}{V_1 / V^*} = \frac{0.3751}{0.2405} = 1.5597 \qquad \rightarrow V_2 = 1.5597 V_1 = 1.5597 (74.94 \text{ m/s}) = 117 \text{ m/s}$$

$$L = L_1^* - L_2^* = \left(\frac{fL_1^*}{D_h} - \frac{fL_2^*}{D_h}\right) \frac{D_h}{f} \rightarrow 2 \text{ m} = (11.520 - 3.5536) \frac{0.03 \text{ m}}{f} \rightarrow f = 0.120$$

**Discussion** Note that the mass flow rate and the average friction factor can be determined by measuring static pressure, as in incompressible flow.

**Solution** Supersonic airflow in a constant cross-sectional area adiabatic duct is considered. For a specified exit Mach number, the temperature, pressure, and velocity at the duct exit are to be determined.



**Assumptions** 1 The assumptions associated with Fanno flow (i.e., steady, frictional flow of an ideal gas with constant properties through a constant cross-sectional area adiabatic duct) are valid. 2 The friction factor remains constant along the duct.

**Properties** We take the properties of air to be k = 1.4,  $c_p = 1.005$  kJ/kg·K, and R = 0.287 kJ/kg·K. The average friction factor is given to be f = 0.03.

Analysis The inlet velocity is

$$c_1 = \sqrt{kRT_1} = \sqrt{(1.4)(0.287 \text{ kJ/kg} \cdot \text{K})(250 \text{ K}) \left(\frac{1000 \text{ m}^2/\text{s}^2}{1 \text{ kJ/kg}}\right)} = 316.9 \text{ m/s}$$

$$V_1 = \text{Ma}_1 c_1 = 2.2(316.9 \text{ m/s}) = 697.3 \text{ m/s}$$

The Fanno flow functions corresponding to the inlet and exit Mach numbers are, from Table A-16,

$$Ma_1 = 2.2$$
:  $(fL^*/D_h)_1 = 0.3609$   $T_1/T^* = 0.6098$ ,  $P_1/P^* = 0.3549$ ,  $V_1/V^* = 1.7179$   $Ma_2 = 1.8$ :  $(fL^*/D_h)_2 = 0.2419$   $T_2/T^* = 0.7282$ ,  $P_2/P^* = 0.4741$ ,  $V_2/V^* = 1.5360$ 

Then the temperature, pressure, and velocity at the duct exit are determined to be

$$\frac{T_2}{T_1} = \frac{T_2 / T^*}{T_1 / T^*} = \frac{0.7282}{0.6098} = 1.1942 \qquad \rightarrow T_2 = 1.1942 (250 \text{ K}) = \mathbf{299 K}$$

$$\frac{P_2}{P_1} = \frac{P_2 / P^*}{P_1 / P^*} = \frac{0.4741}{0.3549} = 1.3359 \qquad \rightarrow P_2 = 1.3359 P_1 = 1.3359 (70 \text{ kPa}) = \mathbf{93.5 \text{ kPa}}$$

$$\frac{V_2}{V_1} = \frac{V_2 / V^*}{V_1 / V^*} = \frac{1.5360}{1.7179} = 0.8941 \qquad \rightarrow V_2 = 0.8941 V_1 = 0.8941 (697.3 \text{ m/s}) = \mathbf{623 \text{ m/s}}$$

**Discussion** The duct length is determined to be

$$L = L_1^* - L_2^* = \left(\frac{fL_1^*}{D_h} - \frac{fL_2^*}{D_h}\right) \frac{D_h}{f} = (0.3609 - 0.2419) \frac{0.055 \text{ m}}{0.03} = \textbf{0.218 m}$$

Note that it takes a duct length of only 0.218 m for the Mach number to decrease from 2.2 to 1.8. The maximum (or sonic) duct lengths at the inlet and exit states in this case are  $L_1^* = 0.662$  m and  $L_2^* = 0.443$  m. Therefore, the flow would reach sonic conditions if a 0.443-m long section were added to the existing duct.



**Solution** Choked supersonic airflow in a constant cross-sectional area adiabatic duct is considered. The variation of duct length with Mach number is to be investigated, and the results are to be plotted.

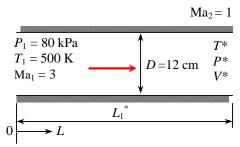
**Assumptions** 1 The assumptions associated with Fanno flow (i.e., steady, frictional flow of an ideal gas with constant properties through a constant cross-sectional area adiabatic duct) are valid. 2 The friction factor remains constant along the duct.

**Properties** We take the properties of air to be k = 1.4,  $c_p = 1.005$  kJ/kg·K, and R = 0.287 kJ/kg·K. The average friction factor is given to be f = 0.03.

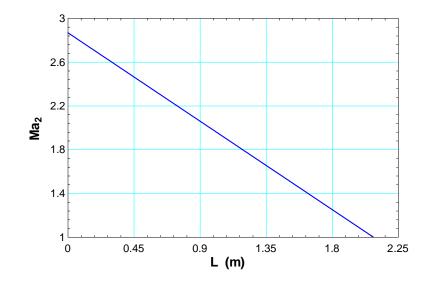
**Analysis** We use EES to solve the problem. The flow is choked, and thus Ma<sub>2</sub> =1. Corresponding to the inlet Mach number of Ma<sub>1</sub> = 3 we have, from Table A-16,  $fL^*/D_h = 0.5222$ , Therefore, the original duct length is

$$L_1^* = 0.5222 \frac{D}{f} = 0.5222 \frac{0.18 \text{ m}}{0.03} = 3.13 \text{ m}$$

Repeating the calculations for different  $Ma_2$  as it varies from 3 to 1 results in the following table for the location on the duct from the inlet:



Mach	Duct length
number, Ma	<i>L</i> , m
1.00	2.09
1.25	1.89
1.50	1.54
1.75	1.19
2.00	0.87
2.25	0.59
2.50	0.36
2.75	0.16
3.00	0.00



# **EES program:**

k=1.4

cp=1.005

R=0.287

P1=80

T1 = 500

Ma1=3
"Ma2=1"

f=0.03

D=0.12

C1=sqrt(k\*R\*T1\*1000)

Ma1=V1/C1

T01=T02

12-101

```
T01=T1*(1+0.5*(k-1)*Ma1^2)
T02=T2*(1+0.5*(k-1)*Ma2^2)
P01=P1*(1+0.5*(k-1)*Ma1^2)^(k/(k-1))
rho1=P1/(R*T1)
Ac=pi*D^2/4
mair=rho1*Ac*V1
P01Ps = ((2+(k-1)*Ma1^2)/(k+1))^(0.5*(k+1)/(k-1))/Ma1
P1Ps=((k+1)/(2+(k-1)*Ma1^2))^0.5/Ma1
T1Ts=(k+1)/(2+(k-1)*Ma1^2)
R1Rs = ((2+(k-1)*Ma1^2)/(k+1))^0.5/Ma1
V1Vs=1/R1Rs
fLs1 = (1-Ma1^2)/(k*Ma1^2) + (k+1)/(2*k)*ln((k+1)*Ma1^2/(2+(k-1)*Ma1^2))
Ls1=fLs1*D/f
P02Ps = ((2+(k-1)*Ma2^2)/(k+1))^(0.5*(k+1)/(k-1))/Ma2
P2Ps=((k+1)/(2+(k-1)*Ma2^2))^0.5/Ma2
T2Ts=(k+1)/(2+(k-1)*Ma2^2)
R2Rs = ((2+(k-1)*Ma2^2)/(k+1))^0.5/Ma2
V2Vs=1/R2Rs
fLs2 = (1-Ma2^2)/(k*Ma2^2) + (k+1)/(2*k)*ln((k+1)*Ma2^2/(2+(k-1)*Ma2^2))
Ls2=fLs2*D/f
L=Ls1-Ls2
P02=P02Ps/P01Ps*P01
P2=P2Ps/P1Ps*P1
V2=V2Vs/V1Vs*V1
```

**Discussion** Note that the Mach number decreases nearly linearly along the duct.

 $Ma_2 = 1$ 

#### 12-143

**Solution** Air flowing at a subsonic velocity in a duct is accelerated by heating. The highest rate of heat transfer without affecting the inlet conditions is to be determined.

**Assumptions** 1 The assumptions associated with Rayleigh flow (i.e., steady one-dimensional flow of an ideal gas with constant properties through a constant cross-sectional area duct with negligible frictional effects) are valid. 2 Inlet conditions (and thus the mass flow rate) remain constant.

**Properties** We take the properties of air to be k = 1.4,  $c_p = 1.005$  kJ/kg·K, and R = 0.287 kJ/kg·K.

Analysis Heat transfer will stop when the flow is choked, and thus  $Ma_2 = V_2/c_2 = 1$ . The inlet density and stagnation temperature are

$$\rho_1 = \frac{P_1}{RT_1} = \frac{350 \text{ kPa}}{(0.287 \text{ kJ/kgK})(420 \text{ K})} = 2.904 \text{ kg/m}^3$$

$$T_{01} = T_1 \left( 1 + \frac{k-1}{2} \text{ Ma}_1^2 \right) = (420 \text{ K}) \left( 1 + \frac{1.4-1}{2} 0.6^2 \right) = 450.2 \text{ K}$$

Then the inlet velocity and the mass flow rate become

$$c_1 = \sqrt{kRT_1} = \sqrt{(1.4)(0.287 \text{ kJ/kg} \cdot \text{K})(420 \text{ K}) \left(\frac{1000 \text{ m}^2/\text{s}^2}{1 \text{ kJ/kg}}\right)} = 410.8 \text{ m/s}$$

$$V_1 = \text{Ma}_1 c_1 = 0.6(410.8 \text{ m/s}) = 246.5 \text{ m/s}$$

$$m_{air}^2 = \rho_1 A_{c1} V_1 = (2.904 \text{ kg/m}^3)(0.1 \times 0.1 \text{ m}^2)(246.5 \text{ m/s}) = 7.157 \text{ kg/s}$$

The Rayleigh flow functions corresponding to the inlet and exit Mach numbers are  $T_{02}/T_0^* = 1$  (since Ma<sub>2</sub> = 1).

$$\frac{T_{01}}{T_0^*} = \frac{(k+1)Ma_1^2[2 + (k-1)Ma_1^2]}{(1+kMa_1^2)^2} = \frac{(1.4+1)0.6^2[2 + (1.4-1)0.6^2]}{(1+1.4\times0.6^2)^2} = 0.8189$$

Therefore,

$$\frac{T_{02}}{T_{01}} = \frac{T_{02} / T_0^*}{T_{01} / T_0^*} = \frac{1}{0.8189} \rightarrow T_{02} = T_{01} / 0.8189 = (450.2 \text{ K}) / 0.8189 = 549.8 \text{ K}$$

Then the rate of heat transfer becomes

$$\mathcal{Q}^{k} = n_{kir}^{k} c_{p} (T_{02} - T_{01}) = (7.157 \text{ kg/s})(1.005 \text{ kJ/kg} \cdot \text{K})(549.8 - 450.2) \text{ K} = 716 \text{ kW}$$

**Discussion** It can also be shown that  $T_2 = 458$  K, which is the highest thermodynamic temperature that can be attained under stated conditions. If more heat is transferred, the additional temperature rise will cause the mass flow rate to decrease. We can also solve this problem using the Rayleigh function values listed in Table A-15.

 $Ma_2 = 1$ 

#### 12-144

**Solution** Helium flowing at a subsonic velocity in a duct is accelerated by heating. The highest rate of heat transfer without affecting the inlet conditions is to be determined.

**Assumptions** 1 The assumptions associated with Rayleigh flow (i.e., steady one-dimensional flow of an ideal gas with constant properties through a constant cross-sectional area duct with negligible frictional effects) are valid. 2 Inlet conditions (and thus the mass flow rate) remain constant.

**Properties** We take the properties of helium to be k = 1.667,  $c_p = 5.193$  kJ/kg·K, and R = 2.077 kJ/kg·K.

Analysis Heat transfer will stop when the flow is choked, and thus  $Ma_2 = V_2/c_2 = 1$ . The inlet density and stagnation temperature are

 $P_1 = 350 \text{ kPa}$  $T_1 = 420 \text{ K}$ 

$$\rho_1 = \frac{P_1}{RT_1} = \frac{350 \,\text{kPa}}{(2.077 \,\text{kJ/kgK})(420 \,\text{K})} = 0.4012 \,\text{kg/m}^3$$

$$T_{01} = T_1 \left( 1 + \frac{k-1}{2} \text{ Ma}_1^2 \right) = (420 \text{ K}) \left( 1 + \frac{1.667 - 1}{2} 0.6^2 \right) = 470.4 \text{ K}$$

Then the inlet velocity and the mass flow rate become

$$c_1 = \sqrt{kRT_1} = \sqrt{(1.667)(2.077 \text{ kJ/kg} \cdot \text{K})(420 \text{ K}) \left(\frac{1000 \text{ m}^2 / \text{s}^2}{1 \text{ kJ/kg}}\right)} = 1206 \text{ m/s}$$

$$V_1 = \text{Ma}_1 c_1 = 0.6(1206 \text{ m/s}) = 723.5 \text{ m/s}$$

$$m_{air}^{g} = \rho_1 A_{c1} V_1 = (0.4012 \text{ kg/m}^3)(0.1 \times 0.1 \text{ m}^2)(723.5 \text{ m/s}) = 2.903 \text{ kg/s}$$

The Rayleigh flow functions corresponding to the inlet and exit Mach numbers are  $T_{02}/T_0^* = 1$  (since Ma<sub>2</sub> = 1)

$$\frac{T_{01}}{T_0^*} = \frac{(k+1)\mathrm{Ma}_1^2[2 + (k-1)\mathrm{Ma}_1^2]}{(1+k\mathrm{Ma}_1^2)^2} = \frac{(1.667+1)0.6^2[2 + (1.667-1)0.6^2]}{(1+1.667\times0.6^2)^2} = 0.8400$$

Therefore,

$$\frac{T_{02}}{T_{01}} = \frac{T_{02} / T_0^*}{T_{01} / T_0^*} = \frac{1}{0.8400}$$
  $\rightarrow T_{02} = T_{01} / 0.8400 = (470.4 \text{ K}) / 0.8400 = 560.0 \text{ K}$ 

Then the rate of heat transfer becomes

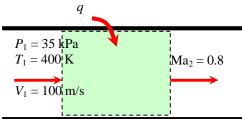
$$\mathcal{O}^{\mathbf{K}} = m_{\text{dir}}^{\mathbf{K}} c_p (T_{02} - T_{01}) = (2.903 \,\text{kg/s})(5.193 \,\text{kJ/kg} \cdot \text{K})(560.0 - 470.4) \,\text{K} = \mathbf{1350} \,\text{kW}$$

**Discussion** It can also be shown that  $T_2 = 420$  K, which is the highest thermodynamic temperature that can be attained under stated conditions. If more heat is transferred, the additional temperature rise will cause the mass flow rate to decrease. Also, in the solution of this problem, we cannot use the values of Table A-15 since they are based on k = 1.4.

**Solution** Air flowing at a subsonic velocity in a duct is accelerated by heating. For a specified exit Mach number, the heat transfer for a specified exit Mach number as well as the maximum heat transfer are to be determined.

Assumptions 1 The assumptions associated with Rayleigh flow (i.e., steady one-dimensional flow of an ideal gas with constant properties through a constant cross-sectional area duct with negligible frictional effects) are valid. 2 Inlet conditions (and thus the mass flow rate) remain constant.

t. We take the properties of air to be k = 1.4,  $c_p = 1.005$ 



Analysis

**Properties** 

 $kJ/kg \cdot K$ , and  $R = 0.287 kJ/kg \cdot K$ .

The inlet Mach number and stagnation temperature are

$$c_1 = \sqrt{kRT_1} = \sqrt{(1.4)(0.287 \text{ kJ/kg} \cdot \text{K})(400 \text{ K}) \left(\frac{1000 \text{ m}^2/\text{s}^2}{1 \text{ kJ/kg}}\right)} = 400.9 \text{ m/s}$$

$$Ma_1 = \frac{V_1}{c_1} = \frac{100 \text{ m/s}}{400.9 \text{ m/s}} = 0.2494$$

$$T_{01} = T_1 \left( 1 + \frac{k-1}{2} \text{ Ma}_1^2 \right) = (400 \text{ K}) \left( 1 + \frac{1.4 - 1}{2} 0.2494^2 \right) = 405.0 \text{ K}$$

The Rayleigh flow functions corresponding to the inlet and exit Mach numbers are (Table A-15):

$$Ma_1 = 0.2494$$
:  $T_{01}/T^* = 0.2559$   
 $Ma_2 = 0.8$ :  $T_{02}/T^* = 0.9639$ 

Then the exit stagnation temperature and the heat transfer are determined to be

$$\frac{T_{02}}{T_{01}} = \frac{T_{02} / T^*}{T_{01} / T^*} = \frac{0.9639}{0.2559} = 3.7667 \qquad \rightarrow \quad T_{02} = 3.7667 T_{01} = 3.7667 (405.0 \text{ K}) = 1526 \text{ K}$$

$$q = c_p \left( T_{02} - T_{01} \right) = \left( 1.005 \text{ kJ/kg} \cdot \text{K} \right) \left( 1526 - 405 \right) \text{ K} = 1126 \text{ kJ/kg} \cong \textbf{1130 kJ/kg}$$

Maximum heat transfer will occur when the flow is choked, and thus  $Ma_2 = 1$  and thus  $T_{02}/T^* = 1$ . Then,

$$\frac{T_{02}}{T_{01}} = \frac{T_{02} / T^*}{T_{01} / T^*} = \frac{1}{0.2559} \rightarrow T_{02} = T_{01} / 0.2559 = (405 \text{ K}) / 0.2559 = 1583 \text{ K}$$

$$q_{max} = c_p (T_{02} - T_{01}) = (1.005 \text{ kJ/kg} \cdot \text{K}) (1583 - 405) \text{ K} = 1184 \text{ kJ/kg} \cong \mathbf{1180} \text{ kJ/kg}$$

**Discussion** This is the maximum heat that can be transferred to the gas without affecting the mass flow rate. If more heat is transferred, the additional temperature rise will cause the mass flow rate to decrease.

**Solution** Air flowing at sonic conditions in a duct is accelerated by cooling. For a specified exit Mach number, the amount of heat transfer per unit mass is to be determined.

Assumptions The assumptions associated with Rayleigh flow (i.e., steady one-dimensional flow of an ideal gas with constant properties through a constant cross-sectional area duct with negligible frictional effects) are valid.

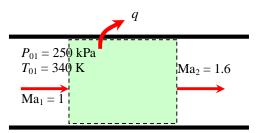
**Properties** We take the properties of air to be k = 1.4,  $c_p = 1.005$  kJ/kg·K, and R = 0.287 kJ/kg·K.

**Analysis** Noting that  $Ma_1 = 1$ , the inlet stagnation temperature is

$$T_{01} = T_1 \left( 1 + \frac{k-1}{2} \text{Ma}_1^2 \right) = (340 \text{ K}) \left( 1 + \frac{1.4 - 1}{2} 1^2 \right) = 408 \text{ K}$$

The Rayleigh flow functions  $T_0/T_0^*$  corresponding to the inlet and exit Mach numbers are (Table A-15):

$$Ma_1 = 1$$
:  $T_{01}/T_0^* = 1$   
 $Ma_2 = 1.6$ :  $T_{02}/T_0^* = 0.8842$ 



Then the exit stagnation temperature and heat transfer are determined to be

$$\frac{T_{02}}{T_{01}} = \frac{T_{02} / T_0^*}{T_{01} / T_0^*} = \frac{0.8842}{1} = 0.8842 \qquad \rightarrow \quad T_{02} = 0.8842 T_{01} = 0.8842 (408 \text{ K}) = 360.75 \text{ K}$$

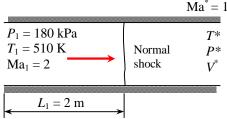
$$q = c_p (T_{02} - T_{01}) = (1.005 \,\mathrm{kJ/kg \cdot K})(360.75 - 408) \;\mathrm{K} = -47.49 \,\mathrm{kJ/kg} \cong -47.5 \,\mathrm{kJ/kg}$$

**Discussion** The negative sign confirms that the gas needs to be cooled in order to be accelerated.

**Solution** Combustion gases enter a constant-area adiabatic duct at a specified state, and undergo a normal shock at a specified location. The exit velocity, temperature, and pressure are to be determined.

**Assumptions** 1 The assumptions associated with Fanno flow (i.e., steady, frictional flow of an ideal gas with constant properties through a constant cross-sectional area adiabatic duct) are valid. 2 The friction factor is constant along the duct.

**Properties** The specific heat ratio and gas constant of combustion gases are given to be k = 1.33 and R = 0.280 kJ/kg·K. The friction factor is given to be f = 0.010.



**Analysis** The Fanno flow functions corresponding to the inlet Mach number of 2 are calculated from the relations in Table A-16 for k = 1.33 to be

$$Ma_1 = 2$$
:  $(fL^*/D_h)_1 = 0.3402$   $T_1/T^* = 0.7018$ ,  $P_1/P^* = 0.4189$ 

First we check to make sure that the flow everywhere upstream the shock is supersonic. The required duct length from the inlet  $L_1^*$  for the flow to reach sonic conditions is  $L_1^* = 0.3402 \frac{D}{f} = 0.3402 \frac{0.10 \text{ m}}{0.010} = 3.40 \text{ m}$ , which is greater than the actual

length of 2 m. Therefore, the flow is indeed supersonic when the normal shock occurs at the indicated location. Also, using the actual duct length  $L_1$ , we have  $\frac{fL_1}{D_h} = \frac{(0.010)(2 \text{ m})}{0.10 \text{ m}} = 0.2000$ . Noting that  $L_1 = L_1^* - L_2^*$ , the function  $fL^*/D_h$  at the exit

state and the corresponding Mach number are  $\left(\frac{fL^*}{D_h}\right)_2 = \left(\frac{fL^*}{D_h}\right)_1 - \frac{fL_1}{D_h} = 0.3402 - 0.2000 = 0.1402 \rightarrow Ma_2 = 1.476.$ 

From the relations in Table A-16, at Ma<sub>2</sub> = 1.476:  $T_2/T^* = 0.8568$ ,  $P_2/P^* = 0.6270$ . Then the temperature, pressure, and velocity before the shock are determined to be

$$\frac{T_2}{T_1} = \frac{T_2 / T^*}{T_1 / T^*} = \frac{0.8568}{0.7018} = 1.2209$$
  $\rightarrow T_2 = 1.2209 T_1 = 1.2209 (510 \text{ K}) = 622.7 \text{ K}$ 

$$\frac{P_2}{P_1} = \frac{P_2 / P^*}{P_1 / P^*} = \frac{0.6270}{0.4189} = 1.4968$$
  $\rightarrow P_2 = 1.4968 P_1 = 1.4968 (180 \text{ kPa}) = 269.4 \text{ kPa}$ 

The normal shock functions corresponding to a Mach number of 1.476 are, from the relations in Table A-14,

$$Ma_2 = 1.476$$
:  $Ma_3 = 0.7052$ ,  $T_3/T_2 = 1.2565$ ,  $P_3/P_2 = 2.3466$ 

Then the temperature and pressure after the shock become

$$T_3 = 1.2565T_2 = 1.2565(622.7 \text{ K}) = 782.4 \text{ K}$$

$$P_3 = 2.3466P_2 = 2.3466(269.4 \text{ kPa}) = 632.3 \text{ kPa}$$

Sonic conditions exist at the duct exit, and the flow downstream of the shock is still Fanno flow. From the relations in Table A-16.

Ma<sub>3</sub> = 0.7052: 
$$T_3/T^* = 1.0767$$
,  $P_3/P^* = 1.4713$   
Ma<sub>4</sub> = 1:  $T_4/T^* = 1$ ,  $P_4/P^* = 1$ 

Then the temperature, pressure, and velocity at the duct exit are determined to be

$$\frac{T_4}{T_3} = \frac{T_4 / T^*}{T_3 / T^*} = \frac{1}{1.0767}$$
  $\rightarrow T_4 = T_3 / 1.0767 = (782.4 \text{ K}) / 1.0767 = 727 K$ 

$$\frac{P_4}{P_2} = \frac{P_4 / P^*}{P_2 / P^*} = \frac{1}{1.4713} \rightarrow P_4 = \frac{P_3 / 1.4713}{1.4713} = \frac{430 \text{ kPa}}{1.4713}$$

$$V_4 = \text{Ma}_4 c_4 = (1)\sqrt{kRT_4} = \sqrt{(1.33)(0.280 \text{ kJ/kg} \cdot \text{K})(727 \text{ K}) \left(\frac{1000 \text{ m}^2/\text{s}^2}{1 \text{ kJ/kg}}\right)} = \textbf{520 m/s}$$

**Discussion** It can be shown that  $L_3^* = 2.13$  m, and thus the total length of this duct is 4.13 m. If the duct is extended, the normal shock will move farther upstream, and eventually to the inlet of the duct.

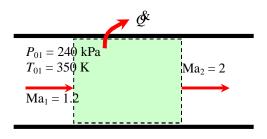
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**Solution** Air flowing at a supersonic velocity in a duct is accelerated by cooling. For a specified exit Mach number, the rate of heat transfer is to be determined.

**Assumptions** The assumptions associated with Rayleigh flow (i.e., steady one-dimensional flow of an ideal gas with constant properties through a constant cross-sectional area duct with negligible frictional effects) are valid.

**Properties** We take the properties of air to be k = 1.4,  $c_p = 1.005$  kJ/kg·K, and R = 0.287 kJ/kg·K.

**Analysis** Knowing stagnation properties, the static properties are determined to be



$$T_1 = T_{01} \left( 1 + \frac{k - 1}{2} \operatorname{Ma}_1^2 \right)^{-1} = (350 \,\mathrm{K}) \left( 1 + \frac{1.4 - 1}{2} 1.2^2 \right)^{-1} = 271.7 \,\mathrm{K}$$

$$P_1 = P_{01} \left( 1 + \frac{k - 1}{2} \operatorname{Ma}_1^2 \right)^{-k/(k - 1)} = (240 \,\mathrm{kPa}) \left( 1 + \frac{1.4 - 1}{2} 1.2^2 \right)^{-1.4/0.4} = 98.97 \,\mathrm{kPa}$$

$$\rho_1 = \frac{P_1}{RT_1} = \frac{98.97 \,\mathrm{kPa}}{(0.287 \,\mathrm{kJ/kgK})(271.7 \,\mathrm{K})} = 1.269 \,\mathrm{kg/m}^3$$

Then the inlet velocity and the mass flow rate become

$$c_1 = \sqrt{kRT_1} = \sqrt{(1.4)(0.287 \text{ kJ/kg} \cdot \text{K})(271.7 \text{ K}) \left(\frac{1000 \text{ m}^2/\text{s}^2}{1 \text{ kJ/kg}}\right)} = 330.4 \text{ m/s}$$

$$V_1 = \text{Ma}_1 c_1 = 1.2(330.4 \text{ m/s}) = 396.5 \text{ m/s}$$

$$m_{vir}^2 = \rho_1 A_{c1} V_1 = (1.269 \text{ kg/m}^3) [\pi (0.20 \text{ m})^2/4](396.5 \text{ m/s}) = 15.81 \text{ kg/s}$$

The Rayleigh flow functions  $T_0/T_0^*$  corresponding to the inlet and exit Mach numbers are (Table A-15):

$$Ma_1 = 1.8$$
:  $T_{01}/T_0^* = 0.9787$   
 $Ma_2 = 2$ :  $T_{02}/T_0^* = 0.7934$ 

Then the exit stagnation temperature is determined to be

$$\frac{T_{02}}{T_{01}} = \frac{T_{02} / T_0^*}{T_{01} / T_0^*} = \frac{0.7934}{0.9787} = 0.8107$$
  $\rightarrow T_{02} = 0.8107 T_{01} = 0.8107(350 \text{ K}) = 283.7 \text{ K}$ 

Finally, the rate of heat transfer is

$$\mathcal{G} = m_{\text{dir}}^{\text{K}} c_p \left( T_{02} - T_{01} \right) = \left( 15.81 \text{ kg/s} \right) \left( 1.005 \text{ kJ/kg} \cdot \text{K} \right) \left( 283.7 - 350 \right) \text{ K} = -1053 \text{ kW} \cong -1050 \text{ kW}$$

**Discussion** The negative sign confirms that the gas needs to be cooled in order to be accelerated. Also, it can be shown that the thermodynamic temperature drops to 158 K at the exit, which is extremely low. Therefore, the duct may need to be heavily insulated to maintain indicated flow conditions.

 $Ma_2 = 1$ 

D = 6 cm

 $L_1$ 

Duct lenght L, m

 $T^*$ 

 $P^*$ 

26

**Solution** Choked subsonic airflow in a constant cross-sectional area adiabatic duct is considered. The effect of duct length on the mass flow rate and the inlet conditions is to be investigated as the duct length is doubled.

**Assumptions** 1 The assumptions associated with Fanno flow (i.e., steady, frictional flow of an ideal gas with constant properties through a constant cross-sectional area adiabatic duct) are valid. 2 The friction factor remains constant along the duct.

**Properties** We take the properties of air to be k = 1.4,  $c_p = 1.005$  kJ/kg·K, and R = 0.287 kJ/kg·K. The average friction factor is given to be f = 0.02.

 $P_1 = 100 \text{ kPa}$ 

 $V_1 = 120 \text{ m/s}$ 

 $T_1 = 400 \text{ K}$ 

Analysis We use EES to solve the problem. The flow is choked, and thus Ma<sub>2</sub> =1. The inlet Mach number is

$$c_1 = \sqrt{kRT_1} = \sqrt{(1.4)(0.287 \text{ kJ/kg} \cdot \text{K})(400 \text{ K}) \left(\frac{1000 \text{ m}^2/\text{s}^2}{1 \text{ kJ/kg}}\right)} = 400.9 \text{ m/s}$$

$$Ma_1 = \frac{V_1}{c_1} = \frac{120 \text{ m/s}}{400.9 \text{ m/s}} = 0.2993$$

Corresponding to this Mach number we have, from Table A-16,  $fL^*/D_h = 5.3312$ , Therefore, the original duct length is

$$L = L_1^* = 5.3312 \frac{D}{f} = 5.3312 \frac{0.06 \text{ m}}{0.02} = 16.0 \text{ m}$$

Then the initial mass flow rate becomes

15

18

19

20

21

22

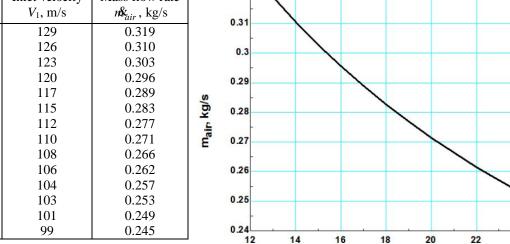
23

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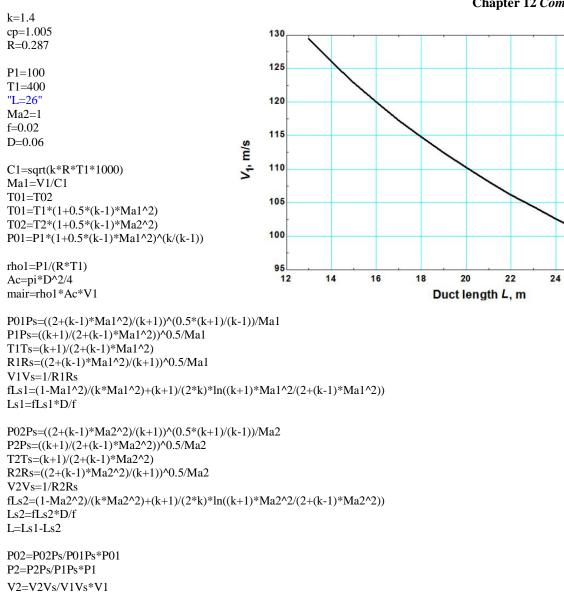
26

$\rho_1 = \frac{P_1}{RT_1} = \frac{100 \text{ kPa}}{(0.287 \text{ kJ/kgK})(400 \text{ K})} = 0.8711 \text{ kg/m}^3$							
$n_{air} = \rho$	$A_{c1}V_1 = (0.8711)$	$kg/m^3)[\pi(0.06 m)]$	$(2)^2 / 4$ (120 m/s) = 0.296 kg/s				
Duct length	Inlet velocity	Mass flow rate	0.32				
<i>L</i> , m	$V_1$ , m/s	<i>n</i> & <sub>air</sub> , kg∕s	0.31				
13	129	0.319	5.5.1				
14	126	0.310					



The EES program is listed below, along with a plot of inlet velocity vs. duct length:

26



**Discussion** Note that once the flow is choked, any increase in duct length results in a decrease in the mass flow rate and the inlet velocity.



**Solution** Using EES (or other) software, the shape of a converging-diverging nozzle is to be determined for specified flow rate and stagnation conditions. The nozzle and the Mach number are to be plotted.

**Assumptions** 1 Air is an ideal gas with constant specific heats. 2 Flow through the nozzle is steady, one-dimensional, and isentropic. 3 The nozzle is adiabatic.

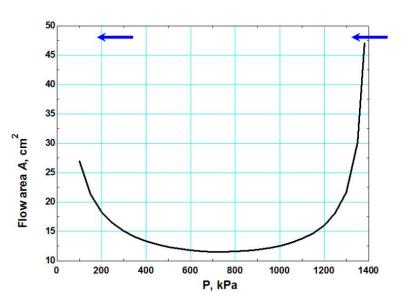
**Properties** The specific heat ratio of air at room temperature is 1.4.

Analysis The problem is solved using EES, and the results are tabulated and plotted below.

k=1.4 Cp=1.005 "kJ/kg.K" R=0.287 "kJ/kg.K" P0=1400 "kPa" T0=200+273 "K" m=3 "kg/s" rho\_0=P0/(R\*T0) rho=P/(R\*T) T=T0\*(P/P0)^((k-1)/k) V=SQRT(2\*Cp\*(T0-T)\*1000) A=m/(rho\*V)\*10000 "cm2" C=SQRT(k\*R\*T\*1000) Ma=V/C

2	/							
1.5								
1								
0.5					_			
0								
0	- 2	200	400	600	800 , kPa	1000	1:	200

Pressure	Flow area	Mach number	
P, kPa	$A, cm^2$	Ma	
1400	8	0	
1350	30.1	0.229	
1300	21.7	0.327	
1250	18.1	0.406	
1200	16.0	0.475	
1150	14.7	0.538	
1100	13.7	0.597	
1050	13.0	0.655	
1000	12.5	0.710	
950	12.2	0.766	
900	11.9	0.820	
850	11.7	0.876	
800	11.6	0.931	
750	11.5	0.988	
700	11.5	1.047	
650	11.6	1.107	
600	11.8	1.171	
550	12.0	1.237	
500	12.3	1.308	
450	12.8	1.384	
400	13.3	1.467	
350	14.0	1.559	
300	15.0	1.663	
250	16.4	1.784	
200	18.3	1.929	
150	21.4	2.114	
100	27.0	2.373	



**Discussion** The shape is not actually to scale since the horizontal axis is pressure rather than distance. If the pressure decreases linearly with distance, then the shape *would* be to scale.



**Solution** Steam enters a converging nozzle. The exit pressure, the exit velocity, and the mass flow rate versus the back pressure for a specified back pressure range are to be plotted.

Assumptions 1 Steam is to be treated as an ideal gas with constant specific heats. 2 Flow through the nozzle is steady, one-dimensional, and isentropic. 3 The nozzle is adiabatic.

**Properties** The ideal gas properties of steam are R = 0.462 kJ/kg.K,  $c_p = 1.872 \text{ kJ/kg.K}$ , and k = 1.3.

**Analysis** We use EES to solve the problem. The stagnation properties in this case are identical to the inlet properties since the inlet velocity is negligible. Since the flow is isentropic, they remain constant throughout the nozzle,

$$P_0 = P_i = 6 \text{ MPa}$$

$$T_0 = T_i = 700 \text{ K}$$

The critical pressure is determined to be

$$P^* = P_0 \left(\frac{2}{k+1}\right)^{k/(k-1)} = (6 \text{ MPa}) \left(\frac{2}{1.3+1}\right)^{1.3/0.3} = 3.274 \text{ MPa}$$

Then the pressure at the exit plane (throat) is

$$P_e = P_b$$

for

$$P_b$$
 ≥ 3.274 MPa

$$P_e = P^* = 3.274 \text{ MPa}$$
 for

$$P_b < 3.274 \text{ MPa (choked flow)}$$

Thus the back pressure does not affect the flow when  $3 < P_b < 3.274$  MPa. For a specified exit pressure  $P_e$ , the temperature, velocity, and mass flow rate are

Temperature

$$T_e = T_0 \left(\frac{P_e}{P_0}\right)^{(k-1)/k} = (700 \text{ K}) \left(\frac{P_e}{6}\right)^{0.3/1}$$

Velocity 
$$V = \sqrt{2c_p (T_0 - T_e)} = \sqrt{2(1.872 \text{ kJ/kg} \cdot \text{K})(700 - T_e) \left(\frac{1000 \text{ m}^2/\text{s}^2}{1 \text{ kJ/kg}}\right)}$$

Density

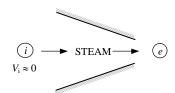
$$\rho_e = \frac{P_e}{RT_e} = \frac{P_e}{(0.462 \text{ kPa} \cdot \text{m}^3 / \text{kg} \cdot \text{K})T_e}$$

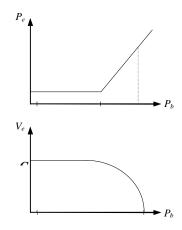
Mass flow rate

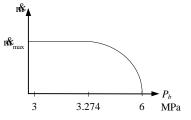
$$n = \rho_e V_e A_e = \rho_e V_e (0.0008 \,\mathrm{m}^2)$$

The results of the calculations are tabulated as follows:

$P_b$ , MPa	$P_e$ , MPa	$T_e$ , K	$V_e$ , m/s	$\rho_e$ , kg/m <sup>3</sup>	m%, kg/s
6.0	6.0	700	0	18.55	0
5.5	5.5	686.1	228.1	17.35	3.166
5.0	5.0	671.2	328.4	16.12	4.235
4.5	4.5	655.0	410.5	14.87	4.883
4.0	4.0	637.5	483.7	13.58	5.255
3.5	3.5	618.1	553.7	12.26	5.431
3.274	3.274	608.7	584.7	11.64	5.445
3.0	3.274	608.7	584.7	11.64	5.445







Discussion Once the back pressure drops below 3.274 MPa, the flow is choked, and not remain constant from then on.

**Solution** An expression for the ratio of the stagnation pressure after a shock wave to the static pressure before the shock wave as a function of k and the Mach number upstream of the shock wave is to be found.

**Analysis** The relation between  $P_1$  and  $P_2$  is

$$\frac{P_2}{P_1} = \frac{1 + k \text{Ma}_2^2}{1 + k \text{Ma}_1^2} \longrightarrow P_2 = P_1 \left( \frac{1 + k \text{Ma}_1^2}{1 + k \text{Ma}_2^2} \right)$$

We substitute this into the isentropic relation

$$\frac{P_{02}}{P_2} = (1 + (k-1) \operatorname{Ma}_2^2 / 2)^{k/(k-1)}$$

which yields

$$\frac{P_{02}}{P_1} = \left(\frac{1 + k \text{Ma}_1^2}{1 + k \text{Ma}_2^2}\right) \left(1 + (k - 1) \text{Ma}_2^2 / 2\right)^{k/(k-1)}$$

where

$$\operatorname{Ma}_{2}^{2} = \frac{\operatorname{Ma}_{1}^{2} + 2/(k-1)}{2k\operatorname{Ma}_{2}^{2}/(k-1) - 1}$$

Substituting,

$$\boxed{\frac{P_{02}}{P_{1}} = \left(\frac{(1+k\mathrm{Ma}_{1}^{2})(2k\mathrm{Ma}_{1}^{2}-k+1)}{k\mathrm{Ma}_{1}^{2}(k+1)-k+3}\right)\left(1+\frac{(k-1)\mathrm{Ma}_{1}^{2}/2+1}{2k\mathrm{Ma}_{1}^{2}/(k-1)-1}\right)^{k/(k-1)}}$$

**Discussion** Similar manipulations of the equations can be performed to get the ratio of other parameters across a shock.



**Solution** Using the compressible flow relations, the one-dimensional compressible flow functions are to be evaluated and tabulated as in Table A-13 for air.

**Properties** The specific heat ratio is given to be k = 1.4 for air.

Analysis The compressible flow functions listed below are expressed in EES and the results are tabulated.

$$Ma^* = Ma \sqrt{\frac{k+1}{2+(k-1)Ma^2}} \qquad \frac{A}{A^*} = \frac{1}{Ma} \left[ \left( \frac{2}{k+1} \right) \left( 1 + \frac{k-1}{2} Ma^2 \right) \right]^{0.5(k+1)/(k-1)}$$

$$\frac{P}{P_0} = \left( 1 + \frac{k-1}{2} Ma^2 \right)^{-k/(k-1)} \qquad \frac{\rho}{\rho_0} = \left( 1 + \frac{k-1}{2} Ma^2 \right)^{-1/(k-1)}$$

$$\frac{T}{T_0} = \left( 1 + \frac{k-1}{2} Ma^2 \right)^{-1}$$

Air:

k=1.4

 $PP0=(1+(k-1)*M^2/2)^{-(-k/(k-1))}$ 

 $TT0=1/(1+(k-1)*M^2/2)$ 

 $DD0=(1+(k-1)*M^2/2)^{-1/(k-1)}$ 

 $Mcr=M*SQRT((k+1)/(2+(k-1)*M^2))$ 

 $AAcr=((2/(k+1))*(1+0.5*(k-1)*M^2))^(0.5*(k+1)/(k-1))/M$ 

Ma	Ma <sup>*</sup>	$A/A^*$	$P/P_0$	$ ho/ ho_0$	$T/T_0$
1.0	1.0000	1.0000	0.5283	0.6339	0.8333
1.5	1.3646	1.1762	0.2724	0.3950	0.6897
2.0	1.6330	1.6875	0.1278	0.2300	0.5556
2.5	1.8257	2.6367	0.0585	0.1317	0.4444
3.0	1.9640	4.2346	0.0272	0.0762	0.3571
3.5	2.0642	6.7896	0.0131	0.0452	0.2899
4.0	2.1381	10.7188	0.0066	0.0277	0.2381
4.5	2.1936	16.5622	0.0035	0.0174	0.1980
5.0	2.2361	25.0000	0.0019	0.0113	0.1667
5.5	2.2691	36.8690	0.0011	0.0076	0.1418
6.0	2.2953	53.1798	0.0006	0.0052	0.1220
6.5	2.3163	75.1343	0.0004	0.0036	0.1058
7.0	2.3333	104.1429	0.0002	0.0026	0.0926
7.5	2.3474	141.8415	0.0002	0.0019	0.0816
8.0	2.3591	190.1094	0.0001	0.0014	0.0725
8.5	2.3689	251.0862	0.0001	0.0011	0.0647
9.0	2.3772	327.1893	0.0000	0.0008	0.0581
9.5	2.3843	421.1314	0.0000	0.0006	0.0525
10.0	2.3905	535.9375	0.0000	0.0005	0.0476

**Discussion** The tabulated values are useful for quick calculations, but be careful – they apply only to one specific value of k, in this case k = 1.4.



**Solution** Using the compressible flow relations, the one-dimensional compressible flow functions are to be evaluated and tabulated as in Table A-13 for methane.

**Properties** The specific heat ratio is given to be k = 1.3 for methane.

Analysis The compressible flow functions listed below are expressed in EES and the results are tabulated.

$$Ma^* = Ma \sqrt{\frac{k+1}{2+(k-1)Ma^2}} \qquad \frac{A}{A^*} = \frac{1}{Ma} \left[ \left( \frac{2}{k+1} \right) \left( 1 + \frac{k-1}{2} Ma^2 \right) \right]^{0.5(k+1)/(k-1)}$$

$$\frac{P}{P_0} = \left( 1 + \frac{k-1}{2} Ma^2 \right)^{-k/(k-1)} \qquad \frac{\rho}{\rho_0} = \left( 1 + \frac{k-1}{2} Ma^2 \right)^{-1/(k-1)}$$

$$\frac{T}{T_0} = \left( 1 + \frac{k-1}{2} Ma^2 \right)^{-1}$$

# Methane:

k=1.3

PP0=(1+(k-1)\*M^2/2)^(-k/(k-1))

 $TT0=1/(1+(k-1)*M^2/2)$ 

 $DD0=(1+(k-1)*M^2/2)^{-1/(k-1)}$ 

 $Mcr=M*SQRT((k+1)/(2+(k-1)*M^2))$ 

 $AAcr=((2/(k+1))*(1+0.5*(k-1)*M^2))^(0.5*(k+1)/(k-1))/M$ 

Ma	Ma <sup>*</sup>	$A/A^*$	$P/P_0$	$\rho/\rho_0$	$T/T_0$
1.0	1.0000	1.0000	0.5457	0.6276	0.8696
1.5	1.3909	1.1895	0.2836	0.3793	0.7477
2.0	1.6956	1.7732	0.1305	0.2087	0.6250
2.5	1.9261	2.9545	0.0569	0.1103	0.5161
3.0	2.0986	5.1598	0.0247	0.0580	0.4255
3.5	2.2282	9.1098	0.0109	0.0309	0.3524
4.0	2.3263	15.9441	0.0050	0.0169	0.2941
4.5	2.4016	27.3870	0.0024	0.0095	0.2477
5.0	2.4602	45.9565	0.0012	0.0056	0.2105
5.5	2.5064	75.2197	0.0006	0.0033	0.1806
6.0	2.5434	120.0965	0.0003	0.0021	0.1563
6.5	2.5733	187.2173	0.0002	0.0013	0.1363
7.0	2.5978	285.3372	0.0001	0.0008	0.1198
7.5	2.6181	425.8095	0.0001	0.0006	0.1060
8.0	2.6350	623.1235	0.0000	0.0004	0.0943
8.5	2.6493	895.5077	0.0000	0.0003	0.0845
9.0	2.6615	1265.6040	0.0000	0.0002	0.0760
9.5	2.6719	1761.2133	0.0000	0.0001	0.0688
10.0	2.6810	2416.1184	0.0000	0.0001	0.0625

**Discussion** The tabulated values are useful for quick calculations, but be careful – they apply only to one specific value of k, in this case k = 1.3.

# Fundamentals of Engineering (FE) Exam Problems

### 12-155

An aircraft is cruising in still air at 5°C at a velocity of 400 m/s. The air temperature at the nose of the aircraft where stagnation occurs is

- (a) 5°C
- (b) 25°C
- (c) 55°C
- (d) 80°C
- (e) 85°C

Answer (e) 85°C

**Solution** Solved by EES Software. Solutions can be verified by copying-and-pasting the following lines on a blank EES screen. (Similar problems and their solutions can be obtained easily by modifying numerical values).

```
k=1.4
Cp=1.005 "kJ/kg.K"
T1=5 "C"
Vel1= 400 "m/s"
T1_stag=T1+Vel1^2/(2*Cp*1000)
```

"Some Wrong Solutions with Common Mistakes:"

W1\_Tstag=T1 "Assuming temperature rise"

W2\_Tstag=Vel1^2/(2\*Cp\*1000) "Using just the dynamic temperature"

W3 Tstag=T1+Vel1^2/(Cp\*1000) "Not using the factor 2"

## 12-156

Air is flowing in a wind tunnel at 25°C, 80 kPa, and 250 m/s. The stagnation pressure at a probe inserted into the flow stream is

- (a) 87 kPa
- (b) 93 kPa
- (c) 113 kPa
- (d) 119 kPa
- (e) 125 kPa

Answer (c) 113 kPa

**Solution** Solved by EES Software. Solutions can be verified by copying-and-pasting the following lines on a blank EES screen. (Similar problems and their solutions can be obtained easily by modifying numerical values).

```
k=1.4

Cp=1.005 "kJ/kg.K"

T1=25 "K"

P1=80 "kPa"

Vel1= 250 "m/s"

T1_stag=(T1+273)+Vel1^2/(2*Cp*1000) "C"

T1_stag/(T1+273)=(P1_stag/P1)^((k-1)/k)
```

# "Some Wrong Solutions with Common Mistakes:"

 $T11\_stag/T1=(W1\_P1stag/P1)^{((k-1)/k)}; \quad T11\_stag=T1+Vel1^{2}/(2*Cp*1000) \text{ "Using deg. C for temperatures"} \\ T12\_stag/(T1+273)=(W2\_P1stag/P1)^{((k-1)/k)}; \quad T12\_stag=(T1+273)+Vel1^{2}/(Cp*1000) \text{ "Not using the factor 2"} \\ T13\_stag/(T1+273)=(W3\_P1stag/P1)^{(k-1)}; \quad T13\_stag=(T1+273)+Vel1^{2}/(2*Cp*1000) \text{ "Using wrong isentropic relation"} \\ T13\_stag/(T1+273)=(W3\_P1stag/P1)^{(k-1)}; \quad T13\_stag=(T1+273)+Vel1^{2}/(2*Cp*1000) \text{ "Using wrong isentropic relation"} \\ T13\_stag/(T1+273)=(W3\_P1stag/P1)^{(k-1)}; \quad T13\_stag=(T1+273)+Vel1^{2}/(2*Cp*1000) \text{ "Using wrong isentropic relation"} \\ T13\_stag/(T1+273)=(W3\_P1stag/P1)^{(k-1)/k}; \quad T13\_stag=(T1+273)+Vel1^{2}/(2*Cp*1000) \text{ "Using wrong isentropic relation"} \\ T13\_stag/(T1+273)=(W3\_P1stag/P1)^{(k-1)/k}; \quad T13\_stag=(T1+273)+Vel1^{2}/(2*Cp*1000) \text{ "Using wrong isentropic relation"} \\ T13\_stag/(T1+273)=(W3\_P1stag/P1)^{(k-1)/k}; \quad T13\_stag=(T1+273)+Vel1^{2}/(2*Cp*1000) \text{ "Using wrong isentropic relation"} \\ T13\_stag/(T1+273)=(W3\_P1stag/P1)^{(k-1)/k}; \quad T13\_stag=(T1+273)+Vel1^{2}/(2*Cp*1000) \text{ "Using wrong isentropic relation"} \\ T13\_stag/(T1+273)=(W3\_P1stag/P1)^{(k-1)/k}; \quad T13\_stag/(T1+273)+Vel1^{2}/(2*Cp*1000) \text{ "Using wrong isentropic relation"} \\ T13\_stag/(T1+273)=(W3\_P1stag/P1)^{(k-1)/k}; \quad T13\_stag/(T1+273)+Vel1^{2}/(2*Cp*1000) \text{ "Using wrong isentropic relation"} \\ T13\_stag/(T1+273)=(W3\_P1stag/P1)^{(k-1)/k}; \quad T13\_stag/(T1+273)+Vel1^{2}/(2*Cp*1000) \text{ "Using wrong isentropic relation"} \\ T13\_stag/(T1+273)=(W3\_P1stag/P1)^{(k-1)/k}; \quad T13\_stag/(T1+273)+Vel1^{2}/(2*Cp*1000) \text{ "Using wrong isentropic relation"} \\ T13\_stag/(T1+273)=(W3\_P1stag/P1)^{(k-1)/k}; \quad T13\_stag/(T1+273)+Vel1^{2}/(2*Cp*1000) \text{ "Using wrong isentropic relation"} \\ T13\_stag/(T1+273)=(W3\_P1stag/P1)^{(k-1)/k}; \quad T13\_stag/(T1+273)+Vel1^{2}/(2*Cp*1000) \text{ "Using wrong isentropic relation"} \\ T13\_stag/(T1+273)=(W3\_P1stag/P1)^{(k-1)/k}; \quad T13\_stag/(T1+273)+Vel1^{2}/(2*Cp*1000) \text{ "Using wrong isentropic relation"} \\ T13\_stag/(T1+273)+Vel1^{2}/(T1+273)+Vel1^{2}/(T1+273)+Vel1^{2}/(T1+273)$ 

An aircraft is reported to be cruising in still air at -20°C and 40 kPa at a Mach number of 0.86. The velocity of the aircraft is

- (a) 91 m/s
- (b) 220 m/s
- (c) 186 m/s
- (d) 280 m/s
- (e) 378 m/s

Answer (d) 280 m/s

**Solution** Solved by EES Software. Solutions can be verified by copying-and-pasting the following lines on a blank EES screen. (Similar problems and their solutions can be obtained easily by modifying numerical values).

k=1.4

Cp=1.005 "kJ/kg.K"

R=0.287 "kJ/kg.K"

T1=-20+273 "K"

P1=40 "kPa"

Mach=0.86

VS1=SQRT(k\*R\*T1\*1000)

Mach=Vel1/VS1

"Some Wrong Solutions with Common Mistakes:"

W1\_vel=Mach\*VS2; VS2=SQRT(k\*R\*T1) "Not using the factor 1000"

W2\_vel=VS1/Mach "Using Mach number relation backwards"

W3\_vel=Mach\*VS3; VS3=k\*R\*T1 "Using wrong relation"

### 12-158

Air is flowing in a wind tunnel at 12°C and 66 kPa at a velocity of 230 m/s. The Mach number of the flow is

- (a) 0.54
- (b) 0.87
- (c) 3.3
- (d) 0.36
- (e) 0.68

Answer (e) 0.68

**Solution** Solved by EES Software. Solutions can be verified by copying-and-pasting the following lines on a blank EES screen. (Similar problems and their solutions can be obtained easily by modifying numerical values).

k=1.4

Cp=1.005 "kJ/kg.K"

R=0.287 "kJ/kg.K"

T1=12+273 "K"

P1=66 "kPa"

Vel1=230 "m/s"

VS1=SORT(k\*R\*T1\*1000)

Mach=Vel1/VS1

"Some Wrong Solutions with Common Mistakes:"

W1\_Mach=Vel1/VS2; VS2=SQRT(k\*R\*(T1-273)\*1000) "Using C for temperature"

W2\_Mach=VS1/Vel1 "Using Mach number relation backwards"

W3\_Mach=Vel1/VS3; VS3=k\*R\*T1 "Using wrong relation"

Consider a converging nozzle with a low velocity at the inlet and sonic velocity at the exit plane. Now the nozzle exit diameter is reduced by half while the nozzle inlet temperature and pressure are maintained the same. The nozzle exit velocity will

(a) remain the same

(b) double

(c) quadruple

(d) go down by half

(e) go down to one-fourth

Answer (a) remain the same

### 12-160

Air is approaching a converging-diverging nozzle with a low velocity at 12°C and 200 kPa, and it leaves the nozzle at a supersonic velocity. The velocity of air at the throat of the nozzle is

(a) 338 m/s

(b) 309 m/s

(c) 280 m/s

(d) 256 m/s

(e) 95 m/s

Answer (b) 309 m/s

**Solution** Solved by EES Software. Solutions can be verified by copying-and-pasting the following lines on a blank EES screen. (Similar problems and their solutions can be obtained easily by modifying numerical values).

k=1.4
Cp=1.005 "kJ/kg.K"
R=0.287 "kJ/kg.K"
"Properties at the inlet"
T1=12+273 "K"
P1=200 "kPa"
Vel1=0 "m/s"
To=T1 "since velocity is zero"
Po=P1
"Throat properties"
T\_throat=2\*To/(k+1)
P\_throat=Po\*(2/(k+1))^(k/(k-1))
"The velocity at the throat is the velocity of sound,"
V\_throat=SQRT(k\*R\*T\_throat\*1000)

"Some Wrong Solutions with Common Mistakes:"

W1\_Vthroat=SQRT(k\*R\*T1\*1000) "Using T1 for temperature"

W2\_Vthroat=SQRT(k\*R\*T2\_throat\*1000); T2\_throat=2\*(To-273)/(k+1) "Using C for temperature"

W3 Vthroat=k\*R\*T throat "Using wrong relation"

Argon gas is approaching a converging-diverging nozzle with a low velocity at 20°C and 120 kPa, and it leaves the nozzle at a supersonic velocity. If the cross-sectional area of the throat is 0.015 m<sup>2</sup>, the mass flow rate of argon through the nozzle is

(a) 0.41 kg/s

(b) 3.4 kg/s

(c) 5.3 kg/s

(d) 17 kg/s

(e) 22 kg/s

Answer (c) 5.3 kg/s

**Solution** Solved by EES Software. Solutions can be verified by copying-and-pasting the following lines on a blank EES screen. (Similar problems and their solutions can be obtained easily by modifying numerical values).

k=1.667

Cp=0.5203 "kJ/kg.K"

R=0.2081 "kJ/kg.K"

A=0.015 "m^2"

"Properties at the inlet"

T1=20+273 "K"

P1=120 "kPa"

Vel1=0 "m/s"

To=T1 "since velocity is zero"

Po=P1

"Throat properties"

 $T_{throat}=2*To/(k+1)$ 

 $P_{throat}=Po*(2/(k+1))^{(k/(k-1))}$ 

rho throat=P throat/(R\*T throat)

"The velocity at the throat is the velocity of sound,"

V\_throat=SQRT(k\*R\*T\_throat\*1000)

m=rho\_throat\*A\*V\_throat

"Some Wrong Solutions with Common Mistakes:"

 $W1\_mass = rho\_throat*A*V1\_throat; V1\_throat = SQRT(k*R*T1\_throat*1000); T1\_throat = 2*(To-273)/(k+1) "Using C for temp" = 2*(To-273)/(k+1) "Usin$ 

W2\_mass=rho2\_throat\*A\*V\_throat; rho2\_throat=P1/(R\*T1) "Using density at inlet"

Carbon dioxide enters a converging-diverging nozzle at 60 m/s, 310°C, and 300 kPa, and it leaves the nozzle at a supersonic velocity. The velocity of carbon dioxide at the throat of the nozzle is

(a) 125 m/s

(b) 225 m/s

(c) 312 m/s

(d) 353 m/s

(e) 377 m/s

Answer (d) 353 m/s

**Solution** Solved by EES Software. Solutions can be verified by copying-and-pasting the following lines on a blank EES screen. (Similar problems and their solutions can be obtained easily by modifying numerical values).

```
k=1.289
Cp=0.846 "kJ/kg.K"
R=0.1889 "kJ/kg.K"
"Properties at the inlet"
T1=310+273 "K"
P1=300 "kPa"
Vel1=60 "m/s"
To=T1+Vel1^2/(2*Cp*1000)
To/T1 = (Po/P1)^{(k-1)/k}
"Throat properties"
T_{throat}=2*To/(k+1)
P_{throat}=Po*(2/(k+1))^{(k/(k-1))}
"The velocity at the throat is the velocity of sound,"
V throat=SQRT(k*R*T throat*1000)
"Some Wrong Solutions with Common Mistakes:"
W1 Vthroat=SORT(k*R*T1*1000) "Using T1 for temperature"
W2 Vthroat=SQRT(k*R*T2 throat*1000); T2 throat=2*(T throat-273)/(k+1) "Using C for temperature"
W3_Vthroat=k*R*T_throat "Using wrong relation"
```

# 12-163

Consider gas flow through a converging-diverging nozzle. Of the five statements below, select the one that is incorrect:

- (a) The fluid velocity at the throat can never exceed the speed of sound.
- (b) If the fluid velocity at the throat is below the speed of sound, the diversion section will act like a diffuser.
- (c) If the fluid enters the diverging section with a Mach number greater than one, the flow at the nozzle exit will be supersonic.
- (d) There will be no flow through the nozzle if the back pressure equals the stagnation pressure.
- (e) The fluid velocity decreases, the entropy increases, and stagnation enthalpy remains constant during flow through a normal shock.

Answer (c) If the fluid enters the diverging section with a Mach number greater than one, the flow at the nozzle exit will be supersonic.

Combustion gases with k = 1.33 enter a converging nozzle at stagnation temperature and pressure of 350°C and 400 kPa, and are discharged into the atmospheric air at 20°C and 100 kPa. The lowest pressure that will occur within the nozzle is

(a) 13 kPa

(b) 100 kPa

(c) 216 kPa

(d) 290 kPa

(e) 315 kPa

Answer (c) 216 kPa

**Solution** Solved by EES Software. Solutions can be verified by copying-and-pasting the following lines on a blank EES screen. (Similar problems and their solutions can be obtained easily by modifying numerical values).

k=1.33

Po=400 "kPa"

"The critical pressure is"

P throat= $Po*(2/(k+1))^{k}(k/(k-1))$ 

"The lowest pressure that will occur in the nozzle is the higher of the critical or atmospheric pressure."

"Some Wrong Solutions with Common Mistakes:"

W2\_Pthroat=Po\*(1/(k+1))^(k/(k-1)) "Using wrong relation"

W3\_Pthroat=100 "Assuming atmospheric pressure"

# **Design and Essay Problems**

# 12-165 to 12-167

**Solution** Students' essays and designs should be unique and will differ from each other.

