

# IDENTICAL PARTICLES

## Exchange Forces

Suppose one particle is in state  $\psi_a(x)$ , and the other is in state  $\psi_b(x)$ , and these two states are orthogonal and normalised.

If the two particles are distinguishable, and number 1 is the one in state  $\psi_a$ , then the combined wave function is:

$$\psi(x_1, x_2) = \psi_a(x_1)\psi_b(x_2);$$

**Identical bosons:**

$$\psi_+(x_1, x_2) = \frac{1}{\sqrt{2}} [\psi_a(x_1)\psi_b(x_2) + \psi_b(x_1)\psi_a(x_2)];$$

**Identical fermions:**

$$\psi_-(x_1, x_2) = \frac{1}{\sqrt{2}} [\psi_a(x_1)\psi_b(x_2) - \psi_b(x_1)\psi_a(x_2)].$$

The expectation value of the square of the separation distance between the two particles:

$$\langle (x_1 - x_2)^2 \rangle = \langle x_1^2 \rangle + \langle x_2^2 \rangle - 2\langle x_1 x_2 \rangle.$$

# IDENTICAL PARTICLES

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**Case 1: Distinguishable particles.**

$$\langle x_1^2 \rangle = \int x_1^2 |\psi_a(x_1)|^2 dx_1 \int |\psi_b(x_2)|^2 dx_2 = \langle x^2 \rangle_a$$

(the expectation value of  $x^2$  in the one-particle state  $\psi_a$ ),

$$\langle x_2^2 \rangle = \int |\psi_a(x_1)|^2 dx_1 \int x_2^2 |\psi_b(x_2)|^2 dx_2 = \langle x^2 \rangle_b,$$

and

$$\langle x_1 x_2 \rangle = \int x_1 |\psi_a(x_1)|^2 dx_1 \int x_2 |\psi_b(x_2)|^2 dx_2 = \langle x \rangle_a \langle x \rangle_b.$$

In this case, then,

$$\left\langle (x_1 - x_2)^2 \right\rangle_d = \langle x^2 \rangle_a + \langle x^2 \rangle_b - 2 \langle x \rangle_a \langle x \rangle_b.$$

(Incidentally, the answer would—of course—be the same if particle 1 had been in state  $\psi_b$ , and particle 2 in state  $\psi_a$ .)

# IDENTICAL PARTICLES

## Exchange Forces:

**Case 2: Identical particles.**

$$\begin{aligned}\langle x_1^2 \rangle &= \frac{1}{2} \left[ \int x_1^2 |\psi_a(x_1)|^2 dx_1 \int |\psi_b(x_2)|^2 dx_2 \right. \\ &\quad + \int x_1^2 |\psi_b(x_1)|^2 dx_1 \int |\psi_a(x_2)|^2 dx_2 \\ &\quad \pm \int x_1^2 \psi_a(x_1)^* \psi_b(x_1) dx_1 \int \psi_b(x_2)^* \psi_a(x_2) dx_2 \\ &\quad \left. \pm \int x_1^2 \psi_b(x_1)^* \psi_a(x_1) dx_1 \int \psi_a(x_2)^* \psi_b(x_2) dx_2 \right] \\ &= \frac{1}{2} \left[ \langle x^2 \rangle_a + \langle x^2 \rangle_b \pm 0 \pm 0 \right] = \frac{1}{2} \left( \langle x^2 \rangle_a + \langle x^2 \rangle_b \right).\end{aligned}$$

Similarly,

$$\langle x_2^2 \rangle = \frac{1}{2} \left( \langle x^2 \rangle_b + \langle x^2 \rangle_a \right).$$

(Naturally,  $\langle x_2^2 \rangle = \langle x_1^2 \rangle$ , since you can't tell them apart.) But

# IDENTICAL PARTICLES

## Exchange Forces:

**Case 2: Identical particles.**

$$\begin{aligned}\langle x_1 x_2 \rangle &= \frac{1}{2} \left[ \int x_1 |\psi_a(x_1)|^2 dx_1 \int x_2 |\psi_b(x_2)|^2 dx_2 \right. \\ &\quad + \int x_1 |\psi_b(x_1)|^2 dx_1 \int x_2 |\psi_a(x_2)|^2 dx_2 \\ &\quad \pm \int x_1 \psi_a(x_1)^* \psi_b(x_1) dx_1 \int x_2 \psi_b(x_2)^* \psi_a(x_2) dx_2 \\ &\quad \left. \pm \int x_1 \psi_b(x_1)^* \psi_a(x_1) dx_1 \int x_2 \psi_a(x_2)^* \psi_b(x_2) dx_2 \right] \\ &= \frac{1}{2} (\langle x \rangle_a \langle x \rangle_b + \langle x \rangle_b \langle x \rangle_a \pm \langle x \rangle_{ab} \langle x \rangle_{ba} \pm \langle x \rangle_{ba} \langle x \rangle_{ab}) \\ &= \langle x \rangle_a \langle x \rangle_b \pm |\langle x \rangle_{ab}|^2,\end{aligned}$$

where

$$\langle x \rangle_{ab} \equiv \int x \psi_a(x)^* \psi_b(x) dx.$$

Thus

$$\left\langle (x_1 - x_2)^2 \right\rangle_{\pm} = \langle x^2 \rangle_a + \langle x^2 \rangle_b - 2\langle x \rangle_a \langle x \rangle_b \mp 2|\langle x \rangle_{ab}|^2.$$

# IDENTICAL PARTICLES

## Exchange Forces:

### Case 2: Identical particles.

Distinguishable particles:

$$\left\langle (x_1 - x_2)^2 \right\rangle_d = \langle x^2 \rangle_a + \langle x^2 \rangle_b - 2 \langle x \rangle_a \langle x \rangle_b.$$

Identical particles.:

$$\left\langle (x_1 - x_2)^2 \right\rangle_{\pm} = \langle x^2 \rangle_a + \langle x^2 \rangle_b - 2 \langle x \rangle_a \langle x \rangle_b \mp 2 |\langle x \rangle_{ab}|^2.$$

Difference:

$$\left\langle (\Delta x)^2 \right\rangle_{\pm} = \left\langle (\Delta x)^2 \right\rangle_d \mp 2 |\langle x \rangle_{ab}|^2;$$

**Identical bosons** (the upper signs) tend to be somewhat closer together, and **identical fermions** (the lower signs) somewhat farther apart, than distinguishable particles in the same two states.

$\langle x \rangle_{ab}$  *vanishes* unless the two wave functions actually *overlap*.

# IDENTICAL PARTICLES

## Exchange Forces:

### Case 2: Identical particles.

The *interesting* case is when the overlap integral is *not* zero.

The system behaves as though there were:

a “**force of attraction**” between identical bosons, **pulling them closer together**, and

a “**force of repulsion**” between identical fermions, **pushing them apart**.

We call it an **exchange force**, although it’s not really a force at all—no physical agency is pushing on the particles. **What is it then?**

It is a purely *geometrical* consequence of the symmetrisation requirement. It is also a strictly quantum mechanical phenomenon, with no classical counterpart.