

Homework 1

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1 Preliminaries

1.1 2 Level System (with pulses)

1.1.1 Hamiltonian under unitary transformation

The matrix H written in the $\{|e\rangle, |g\rangle\}$ basis (naming $c_x = b_x \cos(\omega t)$ and $c_y = b_y \cos(\omega t)$, $b_x^2 + b_y^2 = 1$)

$$H = \begin{pmatrix} b_x \cos(\omega t) \\ b_y \cos(\omega t) \\ \gamma B \end{pmatrix} \cdot \vec{\sigma} = \begin{pmatrix} \gamma B & c_x - ic_y \\ c_x + ic_y & -\gamma B \end{pmatrix} \quad (1)$$

Is called **Hamiltonian** because is present in the Schrodinger equation:

$$\begin{pmatrix} |\dot{e}\rangle \\ |\dot{g}\rangle \end{pmatrix} = \begin{pmatrix} -i\gamma B & -c_y - ic_x \\ c_y - ic_x & i\gamma B \end{pmatrix} \begin{pmatrix} |e\rangle \\ |g\rangle \end{pmatrix} \quad (2)$$

Now, for the change of basis:

$$\begin{pmatrix} |\tilde{e}\rangle \\ |\tilde{g}\rangle \end{pmatrix} = \begin{pmatrix} e^{i\omega t} & 0 \\ 0 & e^{i\omega t} \end{pmatrix} \begin{pmatrix} |e\rangle \\ |g\rangle \end{pmatrix} \quad (3)$$

The time derivative is, by rule of chain:

$$\frac{d}{dt} \begin{pmatrix} |\tilde{e}\rangle \\ |\tilde{g}\rangle \end{pmatrix} = (i\omega) \begin{pmatrix} |\tilde{e}\rangle \\ |\tilde{g}\rangle \end{pmatrix} + e^{i\omega t} \begin{pmatrix} |\dot{e}\rangle \\ |\dot{g}\rangle \end{pmatrix} \quad (4)$$

And using the equation 2

$$\frac{d}{dt} \begin{pmatrix} |\tilde{e}\rangle \\ |\tilde{g}\rangle \end{pmatrix} = (i\omega \mathbb{I} + e^{i\omega t} \begin{pmatrix} -i\gamma B & -c_y - ic_x \\ c_y - ic_x & i\gamma B \end{pmatrix} e^{-i\omega t}) \begin{pmatrix} |\tilde{e}\rangle \\ |\tilde{g}\rangle \end{pmatrix} \quad (5)$$

By the definition of c_x and c_y it appears a product between $\cos(\omega t)$ and $e^{-i\omega t}$ that can be approximated if rapidly oscillating terms are neglected:

$$\cos(\omega t) e^{-i\omega t} = \frac{(e^{i\omega t} + e^{-i\omega t}) e^{-i\omega t}}{2} = \frac{1 + e^{-2i\omega t}}{2} \simeq \frac{1}{2} \quad (6)$$

Defining $c_{\pm} = c_x \pm ic_y$, the approximation of 6 makes it:

$$c_{\pm} e^{-i\omega t} = (c_x \pm ic_y) e^{-i\omega t} = (b_x \pm ib_y) \cos(\omega t) e^{-i\omega t} \simeq \frac{b_x \pm ib_y}{2} \quad (7)$$

Using 7, 5 can be simplified:

$$\frac{d}{dt} \begin{pmatrix} |\tilde{e}\rangle \\ |\tilde{g}\rangle \end{pmatrix} = \begin{pmatrix} -i(\gamma B - \omega) & -ie^{i\omega t} \frac{b_-}{2} \\ -ie^{i\omega t} \frac{b_+}{2} & i(\gamma B - \omega) \end{pmatrix} \begin{pmatrix} |\tilde{e}\rangle \\ |\tilde{g}\rangle \end{pmatrix} \quad (8)$$

And finally, defining $B_x = \frac{e^{i\omega t} b_-}{2}$ and $B_y = \frac{e^{i\omega t} b_+}{2}$, the hamiltonian for the $\{|\tilde{e}\rangle, |\tilde{g}\rangle\}$ basis is ($\Delta = \gamma B - \omega$, $B_{\pm} = B_x \pm iB_y$)

$$H = \begin{pmatrix} B_x \\ B_y \\ \Delta \end{pmatrix} \cdot \vec{\sigma} = \begin{pmatrix} \Delta & B_- \\ B_+ & -\Delta \end{pmatrix} \quad (9)$$

1.1.2 Time evolution for the excited state

The matrix of 9 can be splitted into 2 operators ($r^2 = \Delta^2 + B_x^2 + B_y^2$):

$$H = \frac{1}{2} \begin{pmatrix} \Delta + r & B_- \\ B_+ & r - \Delta \end{pmatrix} - \frac{1}{2} \begin{pmatrix} r - \Delta & -B_- \\ -B_+ & \Delta + r \end{pmatrix} \quad (10)$$

The both operators **are projectors** for orthogonal states:

$$\begin{aligned} \frac{1}{2} \begin{pmatrix} \Delta + r & B_- \\ B_+ & r - \Delta \end{pmatrix} &\Rightarrow v_+ = \frac{1}{\sqrt{2}} \begin{pmatrix} \sqrt{\Delta + r} \\ \frac{B_+}{\sqrt{\Delta + r}} \end{pmatrix} \\ \frac{1}{2} \begin{pmatrix} r - \Delta & -B_- \\ -B_+ & \Delta + r \end{pmatrix} &\Rightarrow v_- = \frac{1}{\sqrt{2}} \begin{pmatrix} \frac{-B_-}{\sqrt{\Delta + r}} \\ \sqrt{\Delta + r} \end{pmatrix} \end{aligned} \quad (11)$$

Therefore, it was obtained a **Spectral Decomposition** for H

$$H = v_+ v_+^\dagger - v_- v_-^\dagger \quad (12)$$

So, the evolution operator e^{iHt} for the system is ($\hbar = 1$)

$$e^{-it} v_+ v_+^\dagger + e^{it} v_- v_-^\dagger = \begin{pmatrix} r \cos(t) - i\Delta \sin(t) & -iB_- \sin(t) \\ iB_+ \sin(t) & r \cos(t) + i\Delta \sin(t) \end{pmatrix} \quad (13)$$

If the initial condition is:

$$|\phi(0)\rangle = |\tilde{e}\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix} \quad (14)$$

The state after time evolution is:

$$|\phi(t)\rangle = e^{-iHt} |\phi(0)\rangle = \begin{pmatrix} r\cos(t) - i\Delta\sin(t) \\ iB_+\sin(t) \end{pmatrix} \quad (15)$$

1.1.3 Can the evolution to reach the ground state?

With the definition given in 15, it can be answered this question.

If in a given time $|\phi(t)\rangle = |g\rangle$ this implies that $\langle g | \phi(t) \rangle = 1$:

$$B_+B_-\sin^2(t) = 1 \Rightarrow \sin^2(t) = \frac{1}{B_+B_-} \quad (16)$$

Using that $\cos^2(t) = 1 - \sin^2(t)$ and 16, it can be write down the another condition: $\langle e | \phi(t) \rangle = 0$

$$r^2\cos^2(t) + \Delta^2\sin^2(t) = 0 \Rightarrow r^2(1 - \frac{1}{B_+B_-}) + \frac{\Delta^2}{B_+B_-} = 0 \quad (17)$$

If $r^2 = \Delta^2 + B_+B_-$, the numerator 17 gives the condition:

$$r^2B_+B_- - r^2 + \Delta^2 = 0 \Rightarrow (r^2 - 1)B_+B_- = 0 \Rightarrow r^2 = 1 \quad (18)$$

So, for the ground state can be reached, it must needed

$$r^2 = \Delta^2 + B_+B_- = (\gamma B - \omega) + \frac{b_x^2 + b_y^2}{4} = 1 \quad (19)$$

1.2 Time Evolution in Two Level System

For the next Hamiltonian ($|n|^2 = n_x^2 + n_y^2 + n_z^2 = 1$):

$$H = \gamma B \hat{n} \cdot \vec{\sigma} = \begin{pmatrix} \gamma B n_z & \gamma B(n_x - in_y) \\ \gamma B(n_x + in_y) & -\gamma B n_z \end{pmatrix} \quad (20)$$

The procedure done in 10 can be done two, finding H as a sum of 2 projectors, and by the way, its spectral decomposition:

$$H = \frac{\gamma B}{2} \left[\begin{pmatrix} 1 + n_z & n_x - in_y \\ n_x + in_y & 1 - n_z \end{pmatrix} - \begin{pmatrix} 1 - n_z & -n_x + in_y \\ -n_x - in_y & 1 + n_z \end{pmatrix} \right] \quad (21)$$

Using it, it can be found the evolution operator, that has the same form of the found in 13 ($n_{\pm} = n_x \pm in_y$):

$$e^{-iH\tau} = \begin{pmatrix} \cos(\frac{\gamma B\tau}{2}) - in_z \sin(\frac{\gamma B\tau}{2}) & -in_- \sin(\frac{\gamma B\tau}{2}) \\ in_+ \sin(\frac{\gamma B\tau}{2}) & \cos(\frac{\gamma B\tau}{2}) + in_z \sin(\frac{\gamma B\tau}{2}) \end{pmatrix} \quad (22)$$

Defining $\theta = \gamma b\tau$, the evolution operator can be understood as a rotating operator:

$$e^{-iH\tau} = \begin{pmatrix} \cos(\frac{\theta}{2}) - in_z \sin(\frac{\theta}{2}) & -in_- \sin(\frac{\theta}{2}) \\ in_+ \sin(\frac{\theta}{2}) & \cos(\frac{\theta}{2}) + in_z \sin(\frac{\theta}{2}) \end{pmatrix} \quad (23)$$

The operator of 23 will be analysed the next special cases:

- If $\hat{n} = \hat{x}$, $n_x = 1$, $n_y = n_z = 0$, 23 becomes

$$e^{-iH\tau} = \begin{pmatrix} \cos(\frac{\theta}{2}) & -i \sin(\frac{\theta}{2}) \\ i \sin(\frac{\theta}{2}) & \cos(\frac{\theta}{2}) \end{pmatrix} \quad (24)$$

- If $\hat{n} = \hat{y}$, $n_y = 1$, $n_z = n_x = 0$, 23 becomes

$$e^{-iH\tau} = \begin{pmatrix} \cos(\frac{\theta}{2}) & -\sin(\frac{\theta}{2}) \\ \sin(\frac{\theta}{2}) & \cos(\frac{\theta}{2}) \end{pmatrix} \quad (25)$$

- If $\hat{n} = \hat{z}$, $n_z = 1$, $n_x = n_y = 0$, 23 becomes

$$e^{-iH\tau} = \begin{pmatrix} \cos(\frac{\theta}{2}) - i \sin(\frac{\theta}{2}) & 0 \\ 0 & \cos(\frac{\theta}{2}) + i \sin(\frac{\theta}{2}) \end{pmatrix} \quad (26)$$

| | | | |
|-------------------------------|---|--|--|
| $ \phi(0)\rangle = e\rangle$ | $\hat{n} = \hat{x}$ | $\hat{n} = \hat{y}$ | $\hat{n} = \hat{z}$ |
| $\theta = \frac{\pi}{2}$ | $ \phi(t)\rangle = \frac{ e\rangle + i g\rangle}{\sqrt{2}}$ | $ \phi(t)\rangle = \frac{ e\rangle + g\rangle}{\sqrt{2}}$ | $ \phi(t)\rangle = \frac{1-i}{\sqrt{2}} e\rangle$ |
| $\theta = \pi$ | $ \phi(t)\rangle = i g\rangle$ | $ \phi(t)\rangle = g\rangle$ | $ \phi(t)\rangle = -i e\rangle$ |

| | | | |
|-------------------------------|---|--|--|
| $ \phi(0)\rangle = g\rangle$ | $\hat{n} = \hat{x}$ | $\hat{n} = \hat{y}$ | $\hat{n} = \hat{z}$ |
| $\theta = \frac{\pi}{2}$ | $ \phi(t)\rangle = \frac{ g\rangle - i e\rangle}{\sqrt{2}}$ | $ \phi(t)\rangle = \frac{ g\rangle - e\rangle}{\sqrt{2}}$ | $ \phi(t)\rangle = \frac{1+i}{\sqrt{2}} g\rangle$ |
| $\theta = \pi$ | $ \phi(t)\rangle = -i e\rangle$ | $ \phi(t)\rangle = - e\rangle$ | $ \phi(t)\rangle = i g\rangle$ |

And replacing for $\theta = \frac{\pi}{2}$ and $\theta = \pi$, it will be obtained the next tables for $|\phi_0\rangle = |e\rangle$ and $|\phi_0\rangle = |g\rangle$:

1.3 Final State after Sequences

For an initial state ($|\alpha|^2 + |\beta|^2 = 1$):

$$|\varphi_0\rangle = \alpha |e\rangle + \beta |g\rangle = \begin{pmatrix} \alpha \\ \beta \end{pmatrix} \quad (27)$$

It can be done a Ramsey Sequence ($\varphi = \frac{\gamma B \tau}{2}$)

$$\begin{aligned}
 & \left(\frac{\pi}{2}_x\right) : \begin{pmatrix} \alpha \\ \beta \end{pmatrix} \rightarrow \begin{pmatrix} \frac{\alpha+i\beta}{\sqrt{2}} \\ \frac{\beta-i\alpha}{\sqrt{2}} \end{pmatrix} \\
 & (\tau) : \begin{pmatrix} \frac{\alpha+i\beta}{\sqrt{2}} \\ \frac{\beta-i\alpha}{\sqrt{2}} \end{pmatrix} \rightarrow \begin{pmatrix} \frac{(\cos\varphi - in_z \sin\varphi)(\alpha+i\beta) - in_- \sin\varphi(\beta-i\alpha)}{\sqrt{2}} \\ \frac{in_+ \sin\varphi(\alpha+i\beta) + (\cos\varphi + in_z \sin\varphi)(\beta-i\alpha)}{\sqrt{2}} \end{pmatrix} \\
 & \left(\frac{\pi}{2}_x\right) : \begin{pmatrix} \frac{(\cos\varphi - isin\varphi)(\alpha+i\beta)}{\sqrt{2}} \\ \frac{(\cos\varphi + isin\varphi)(\beta-i\alpha)}{\sqrt{2}} \end{pmatrix} \rightarrow \begin{pmatrix} \frac{e^{-i\varphi}(\alpha+i\beta) + ie^{i\varphi}(\beta-i\alpha)}{2} \\ \frac{e^{i\varphi}(\beta-i\alpha) - ie^{-i\varphi}(\alpha+i\beta)}{2} \end{pmatrix}
 \end{aligned} \quad (28)$$

It can be done a kind of Echo Sequence ($\varphi = \frac{\gamma B \tau}{2}$)

$$\begin{aligned}
 & \left(\frac{\pi}{2}_x\right) : \begin{pmatrix} \alpha \\ \beta \end{pmatrix} \rightarrow \begin{pmatrix} \frac{\alpha+i\beta}{\sqrt{2}} \\ \frac{\beta-i\alpha}{\sqrt{2}} \end{pmatrix} \\
 (\mathcal{T}) : & \begin{pmatrix} \frac{\alpha+i\beta}{\sqrt{2}} \\ \frac{\beta-i\alpha}{\sqrt{2}} \end{pmatrix} \rightarrow \begin{pmatrix} \frac{(\cos\varphi - in_z \sin\varphi)(\alpha+i\beta) - in_- \sin\varphi(\beta-i\alpha)}{\sqrt{2}} \\ \frac{in_+ \sin\varphi(\alpha+i\beta) + (\cos\varphi + in_z \sin\varphi)(\beta-i\alpha)}{\sqrt{2}} \end{pmatrix} \\
 (\pi_x) : & \begin{pmatrix} \frac{(\cos\varphi - isin\varphi)(\alpha+i\beta)}{\sqrt{2}} \\ \frac{(\cos\varphi + isin\varphi)(\beta-i\alpha)}{\sqrt{2}} \end{pmatrix} \rightarrow \begin{pmatrix} -ie^{i\varphi} \left(\frac{\beta-i\alpha}{\sqrt{2}}\right) \\ ie^{-i\varphi} \left(\frac{\alpha+i\beta}{\sqrt{2}}\right) \end{pmatrix} \\
 (\mathcal{T}) : & \begin{pmatrix} -e^{i\varphi} \left(\frac{\alpha+i\beta}{\sqrt{2}}\right) \\ -e^{-i\varphi} \left(\frac{\beta-i\alpha}{\sqrt{2}}\right) \end{pmatrix} \rightarrow \begin{pmatrix} -e^{-i\varphi} e^{i\varphi} \frac{\alpha+i\beta}{\sqrt{2}} \\ -e^{i\varphi} e^{-i\varphi} \frac{\beta-i\alpha}{\sqrt{2}} \end{pmatrix} \\
 & \left(\frac{\pi}{2}_x\right) : \begin{pmatrix} -\frac{\alpha+i\beta}{\sqrt{2}} \\ -\frac{\beta-i\alpha}{\sqrt{2}} \end{pmatrix} \rightarrow \begin{pmatrix} \frac{-(\alpha+i\beta+i(\beta-i\alpha))}{2} \\ \frac{-(\beta-i\alpha-i(\alpha+i\beta))}{2} \end{pmatrix}
 \end{aligned} \tag{29}$$

It can be done another kind of Echo Sequence ($\varphi = \frac{\gamma B \tau}{2}$)

$$\begin{aligned}
 & \left(\frac{\pi}{2}_x\right) : \begin{pmatrix} \alpha \\ \beta \end{pmatrix} \rightarrow \begin{pmatrix} \frac{\alpha+i\beta}{\sqrt{2}} \\ \frac{\beta-i\alpha}{\sqrt{2}} \end{pmatrix} \\
 (\mathcal{T}) : & \begin{pmatrix} \frac{\alpha+i\beta}{\sqrt{2}} \\ \frac{\beta-i\alpha}{\sqrt{2}} \end{pmatrix} \rightarrow \begin{pmatrix} \frac{(\cos\varphi - in_z \sin\varphi)(\alpha+i\beta) - in_- \sin\varphi(\beta-i\alpha)}{\sqrt{2}} \\ \frac{in_+ \sin\varphi(\alpha+i\beta) + (\cos\varphi + in_z \sin\varphi)(\beta-i\alpha)}{\sqrt{2}} \end{pmatrix} \\
 (\pi_x) : & \begin{pmatrix} \frac{(\cos\varphi - isin\varphi)(\alpha+i\beta)}{\sqrt{2}} \\ \frac{(\cos\varphi + isin\varphi)(\beta-i\alpha)}{\sqrt{2}} \end{pmatrix} \rightarrow \begin{pmatrix} -ie^{i\varphi} \left(\frac{\beta-i\alpha}{\sqrt{2}}\right) \\ ie^{-i\varphi} \left(\frac{\alpha+i\beta}{\sqrt{2}}\right) \end{pmatrix} \\
 (\mathcal{T}) : & \begin{pmatrix} -e^{i\varphi} \left(\frac{\alpha+i\beta}{\sqrt{2}}\right) \\ -e^{-i\varphi} \left(\frac{\beta-i\alpha}{\sqrt{2}}\right) \end{pmatrix} \rightarrow \begin{pmatrix} -e^{-i\varphi} e^{i\varphi} \frac{\alpha+i\beta}{\sqrt{2}} \\ -e^{i\varphi} e^{-i\varphi} \frac{\beta-i\alpha}{\sqrt{2}} \end{pmatrix} \\
 & \left(\frac{\pi}{2}_y\right) : \begin{pmatrix} -\frac{\alpha+i\beta}{\sqrt{2}} \\ -\frac{\beta-i\alpha}{\sqrt{2}} \end{pmatrix} \rightarrow \begin{pmatrix} \frac{-(\alpha+i\beta+\beta-i\alpha)}{2} \\ \frac{-(\beta-i\alpha-\alpha-i\beta)}{2} \end{pmatrix}
 \end{aligned} \tag{30}$$

In the processes 28, 29 and 30, the operators were applicated and in the next step, is simplified the result state. Therefore, The

final and normalised vectors after the processes are:

$$\begin{aligned}
|\phi\rangle_{Ramsey} &= \frac{1}{\sqrt{2}} \begin{pmatrix} \alpha + i\beta \\ \beta - i\alpha \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} \alpha \\ \beta \end{pmatrix} + \frac{i}{\sqrt{2}} \begin{pmatrix} \beta \\ -\alpha \end{pmatrix} \\
|\phi\rangle_{Echo1} &= \begin{pmatrix} \frac{-\alpha-i\beta}{\sqrt{2}} \\ \frac{-\beta+i\alpha}{\sqrt{2}} \end{pmatrix} = \frac{-1}{\sqrt{2}} \begin{pmatrix} \alpha \\ \beta \end{pmatrix} + \frac{-i}{\sqrt{2}} \begin{pmatrix} \beta \\ -\alpha \end{pmatrix} \\
|\phi\rangle_{Echo2} &= \begin{pmatrix} \frac{(\alpha+i\beta)(i-1)}{2} \\ \frac{(\beta-i\alpha)(i-1)}{2} \end{pmatrix} = \frac{i-1}{2} \begin{pmatrix} \alpha \\ \beta \end{pmatrix} + \frac{-1-i}{2} \begin{pmatrix} \beta \\ -\alpha \end{pmatrix}
\end{aligned} \tag{31}$$

For the 3 vectors, the probability to obtain again the initial state of 27 can be obtained by the splitting made in 31 for each one of them. The probabilities will be:

$$\begin{aligned}
P_{0,Ramsey} &= \left(\frac{1}{\sqrt{2}}\right)^2 = \frac{1}{2} \\
P_{0,Echo1} &= \left(\frac{-1}{\sqrt{2}}\right)^2 = \frac{1}{2} \\
P_{0,Echo2} &= \left|\frac{i-1}{2}\right|^2 = \left(\frac{\sqrt{2}}{2}\right)^2 = \frac{1}{2}
\end{aligned} \tag{32}$$

2 Central Spin and Random Fluctuating Field

In the suggested problem (spin $\frac{1}{2}$ under a magnetic field B_0 , is valid a Hamiltonian similar to the defined in 20. If the rotating frame is along the \hat{z} axis, $n_z = 1$ and $n_x = n_y = 0$:

$$H = \gamma B \hat{\sigma}_z = \begin{pmatrix} \gamma B & 0 \\ 0 & -\gamma B \end{pmatrix} \tag{33}$$

The evolution operator derived by 33 is

$$e^{-iH\tau} = \begin{pmatrix} \cos(\frac{\gamma B\tau}{2}) - i\sin(\frac{\gamma B\tau}{2}) & 0 \\ 0 & \cos(\frac{\gamma B\tau}{2}) + i\sin(\frac{\gamma B\tau}{2}) \end{pmatrix} \tag{34}$$

2.1 Static Magnetic Field

The both analysis are made an initial state given by 27 and the tables developed in a previous section:

2.1.1 State after Ramsey sequence

As it has done in 28, it is made a Ramsey Sequence, now with the Hamiltonian of 34:

$$\begin{aligned}
 & \left(\frac{\pi}{2}_x\right) : \begin{pmatrix} \alpha \\ \beta \end{pmatrix} \rightarrow \begin{pmatrix} \frac{\alpha+i\beta}{\sqrt{2}} \\ \frac{\beta-i\alpha}{\sqrt{2}} \end{pmatrix} \\
 & (\tau) : \begin{pmatrix} \frac{\alpha+i\beta}{\sqrt{2}} \\ \frac{\beta-i\alpha}{\sqrt{2}} \end{pmatrix} \rightarrow \begin{pmatrix} \frac{(\cos\varphi-i\sin\varphi)(\alpha+i\beta)}{\sqrt{2}} \\ \frac{(\cos\varphi+i\sin\varphi)(\beta-i\alpha)}{\sqrt{2}} \end{pmatrix} \\
 & \left(\frac{\pi}{2}_x\right) : \begin{pmatrix} \frac{e^{-i\varphi}(\alpha+i\beta)}{\sqrt{2}} \\ \frac{e^{i\varphi}(\beta-i\alpha)}{\sqrt{2}} \end{pmatrix} \rightarrow \begin{pmatrix} \frac{e^{-i\varphi}(\alpha+i\beta)+ie^{i\varphi}(\beta-i\alpha)}{2} \\ \frac{e^{i\varphi}(\beta-i\alpha)-ie^{-i\varphi}(\alpha+i\beta)}{2} \end{pmatrix}
 \end{aligned} \tag{35}$$

Such that the final state is:

$$|\phi(\tau)\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} \alpha + i\beta \\ \beta - i\alpha \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} \alpha \\ \beta \end{pmatrix} + \frac{i}{\sqrt{2}} \begin{pmatrix} \beta \\ -\alpha \end{pmatrix} \tag{36}$$

And the probability to be in the initial state is:

$$P_0 = \left(\frac{1}{\sqrt{2}}\right)^2 = \frac{1}{2} \tag{37}$$

2.1.2 State after Echo sequence

As it has done in 29, it is made a Echo Sequence, now with the Hamiltonian of 34:

$$\begin{aligned}
 \left(\frac{\pi}{2}_x\right) : \begin{pmatrix} \alpha \\ \beta \end{pmatrix} &\rightarrow \begin{pmatrix} \frac{\alpha+i\beta}{\sqrt{2}} \\ \frac{\beta-i\alpha}{\sqrt{2}} \end{pmatrix} \\
 (\tau) : \begin{pmatrix} \frac{\alpha+i\beta}{\sqrt{2}} \\ \frac{\beta-i\alpha}{\sqrt{2}} \end{pmatrix} &\rightarrow \begin{pmatrix} \frac{(\cos\varphi-i\sin\varphi)(\alpha+i\beta)}{\sqrt{2}} \\ \frac{(\cos\varphi+i\sin\varphi)(\beta-i\alpha)}{\sqrt{2}} \end{pmatrix} \\
 (\pi_x) : \begin{pmatrix} \frac{e^{-i\varphi}(\alpha+i\beta)}{\sqrt{2}} \\ \frac{e^{i\varphi}(\beta-i\alpha)}{\sqrt{2}} \end{pmatrix} &\rightarrow \begin{pmatrix} -ie^{i\varphi}\left(\frac{\beta-i\alpha}{\sqrt{2}}\right) \\ ie^{-i\varphi}\left(\frac{\alpha+i\beta}{\sqrt{2}}\right) \end{pmatrix} \\
 (\tau) : \begin{pmatrix} -e^{i\varphi}\left(\frac{\alpha+i\beta}{\sqrt{2}}\right) \\ -e^{-i\varphi}\left(\frac{\beta-i\alpha}{\sqrt{2}}\right) \end{pmatrix} &\rightarrow \begin{pmatrix} -e^{-i\varphi}e^{i\varphi}\frac{\alpha+i\beta}{\sqrt{2}} \\ -e^{i\varphi}e^{-i\varphi}\frac{\beta-i\alpha}{\sqrt{2}} \end{pmatrix} \\
 \left(\frac{\pi}{2}_x\right) : \begin{pmatrix} -\frac{\alpha+i\beta}{\sqrt{2}} \\ -\frac{\beta-i\alpha}{\sqrt{2}} \end{pmatrix} &\rightarrow \begin{pmatrix} \frac{-(\alpha+i\beta+i(\beta-i\alpha))}{2} \\ \frac{-(\beta-i\alpha-i(\alpha+i\beta))}{2} \end{pmatrix}
 \end{aligned} \tag{38}$$

Such that the final state is:

$$|\phi(\tau)\rangle = \frac{1}{\sqrt{2}} \begin{pmatrix} -\alpha - i\beta \\ -\beta + i\alpha \end{pmatrix} = \frac{-1}{\sqrt{2}} \begin{pmatrix} \alpha \\ \beta \end{pmatrix} + \frac{-i}{\sqrt{2}} \begin{pmatrix} \beta \\ -\alpha \end{pmatrix} \tag{39}$$

And the probability to be in the initial state is the same that the obtained in 37, that is $P_0 = \frac{1}{2}$.

2.2 Static and Random Magnetic Fields

If it is added a random element in the Magnetic Field, it will be observed more effects. Making a Taylor expansion for the exponential function:

$$e^{i\phi} \simeq 1 + i\phi - \phi^2 + \dots \tag{40}$$

It can be obtained by this the expectation value, if as usual $\langle \phi \rangle = 0$

$$\langle e^{i\phi} \rangle = 1 - \langle \phi^2 \rangle \tag{41}$$

- 2.2.1 Expectation Value of Random Function
- 2.2.2 Spectral Densities for Different Noises
- 2.2.3 Ramsey sequence for noisy system
- 2.2.4 Echo sequence for noisy system