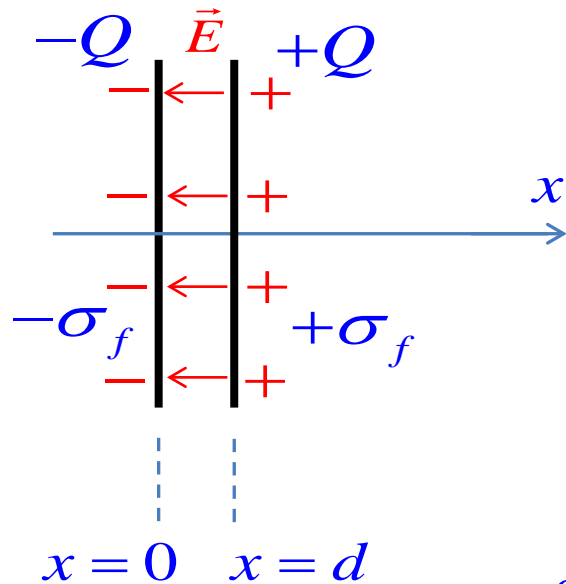


III Capacitors and Review of DC Circuits

1 Capacitors

*Parallel plate capacitor :
(plate dimensions $\ll d$)*



$$C = \frac{Q}{V} = \frac{\epsilon \oint \vec{E} \cdot d\vec{A}}{\int \vec{E} \cdot d\vec{l}}$$

- *assume Q (charge it up)*
- *find V in terms of Q using Gauss's Law*
- *obtain C from $C = \frac{Q}{V}$*

– for infinite sheet : $\vec{D} = \frac{\sigma_f}{2} \hat{n}$

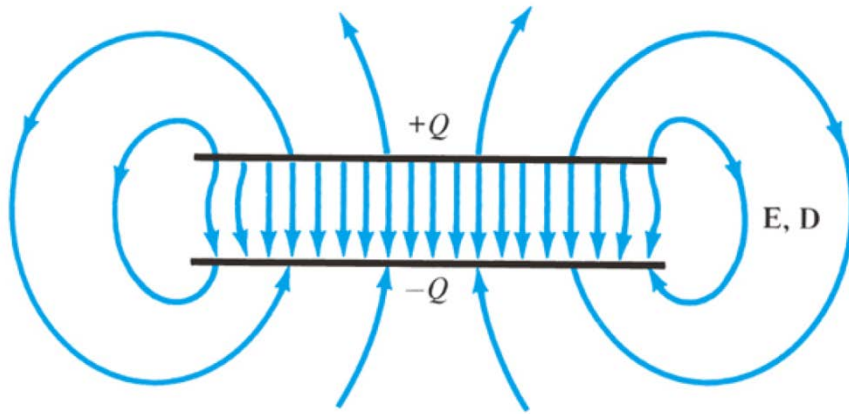
– then \vec{E} between plates : $\vec{E} = \frac{\sigma_f}{\epsilon_0} (-\hat{x}) = \frac{Q_f}{A\epsilon_0} (-\hat{x})$

$$\therefore \Delta V = V_d - V_0 = - \int_0^d \left(-\frac{Q_f}{A\epsilon_0} \right) \hat{x} = \frac{Q_f d}{A\epsilon_0}$$

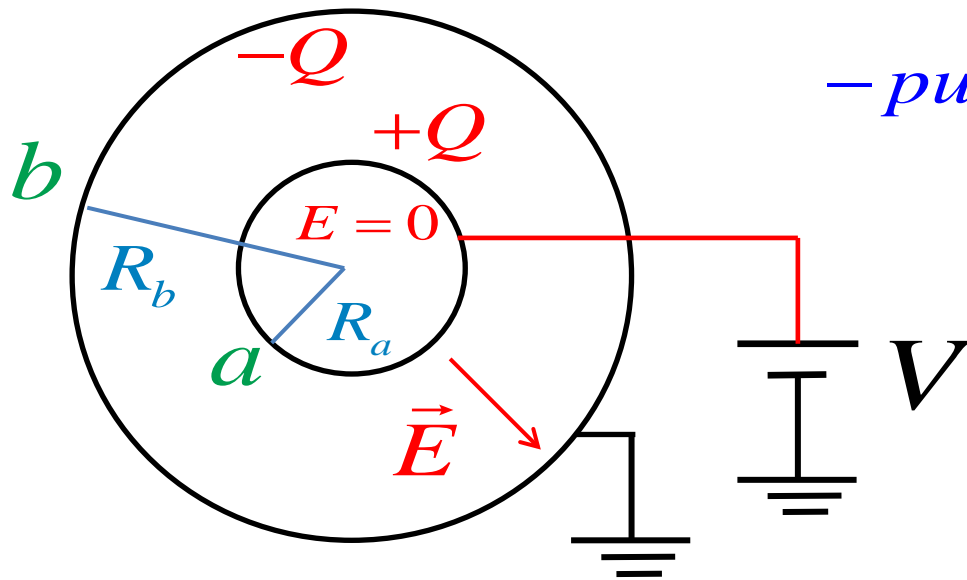
$$\Delta V = \frac{Q_f d}{A \epsilon_0}$$

$$C = \frac{Q}{V} \Rightarrow C = \frac{A \epsilon_0}{d}$$

-remember that it is assumed that any fringe field is negligible



Example : Capacitance of a pair of concentric, conducting spherical shells of radii R_a and R_b .



– put $+Q$ on inner sphere

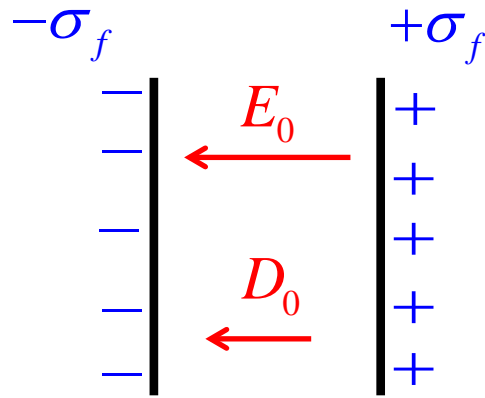
– \vec{E} is set up

– outer shell acquires $-Q$

- Start by finding E from Gauss's Law
- then find V_{BA} and finally C from $C = Q/V$

– get $C = \frac{ab}{4\pi\epsilon_0(b-a)}$

• *C without dielectric:*

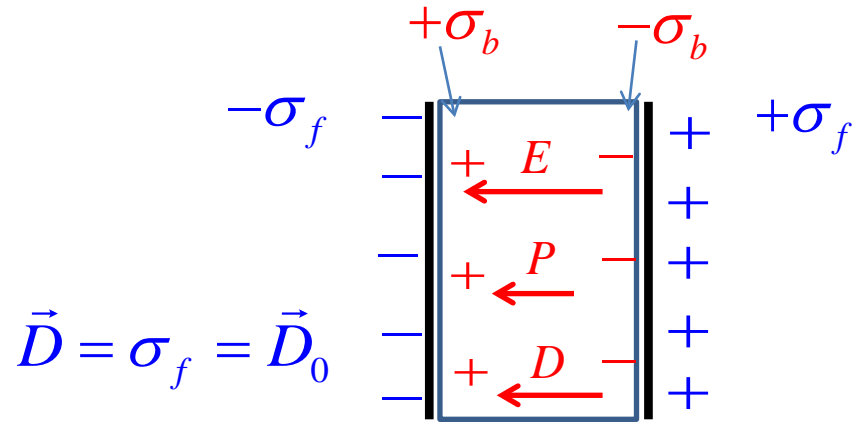


$$\vec{D}_0 = \epsilon_0 \vec{E}_0 = \sigma_f$$

$$\Delta V_0 = E_0 d$$

$$C_0 = \frac{Q}{\Delta V_0}$$

• *C with dielectric:*



$$\vec{D} = \sigma_f = \vec{D}_0$$

$$\vec{E} = \frac{\vec{D}}{\epsilon} = \frac{\vec{D}_0}{\epsilon_r \epsilon_0} = \frac{\vec{E}_0}{\epsilon_r} \Rightarrow E < E_0$$

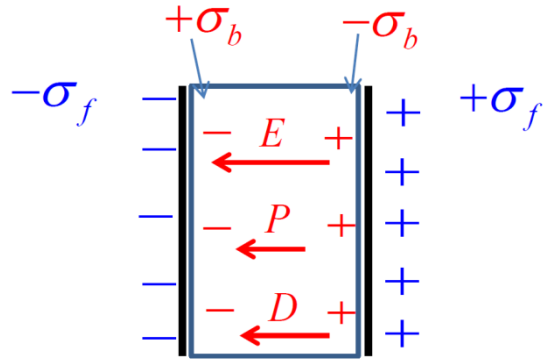
$$\vec{P} = \chi_e \epsilon_0 \vec{E} = (\epsilon_r - 1) \epsilon_0 \vec{E} = \frac{(\epsilon_r - 1)}{\epsilon_r} \epsilon_0 \vec{E}_0 = \epsilon_0 (\vec{E}_0 - \vec{E})$$

$$\Rightarrow E = E_0 - \frac{P}{\epsilon_0} \quad \text{or} \quad E = E_0 - E_b \quad \text{where} \quad E_b = \frac{P}{\epsilon_0} \equiv \text{local field}$$

$$\text{--also } |\sigma_b| = |P_n| = P = \frac{(\epsilon_r - 1)}{\epsilon_r} \epsilon_0 E_0 = \frac{(\epsilon_r - 1)}{\epsilon_r} D_0 = \frac{(\epsilon_r - 1)}{\epsilon_r} |\sigma_f| \quad \text{or} \quad \sigma_b = -\frac{(\epsilon_r - 1)}{\epsilon_r} \sigma_f$$

$$\text{and } E_b = \frac{|\sigma_b|}{2\epsilon_0} + \frac{|\sigma_b|}{2\epsilon_0} = \frac{|\sigma_b|}{\epsilon_0} \quad (\text{as for parallel plate capacitors with } \sigma_b)$$

– *C with dielectric continued* –



$$\Delta V = \int \vec{E} \cdot d\vec{l} = Ed = \frac{E_0}{\epsilon_r} d = \frac{\Delta V_0}{\epsilon_r}$$

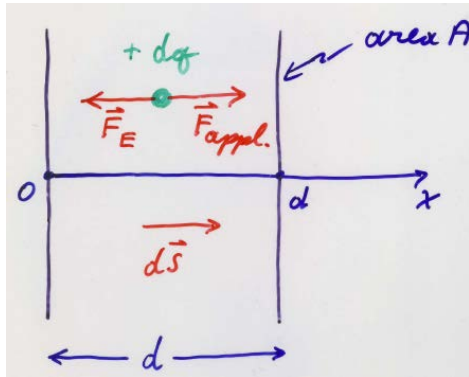
$\Rightarrow \Delta V$ is less than ΔV_0

$$C = \frac{Q}{\Delta V} = \frac{Q}{\left(\frac{\Delta V_0}{\epsilon_r} \right)} ; \quad \therefore \frac{C}{C_0} = \epsilon_r \quad \text{or} \quad C = \epsilon_r C_0$$

In addition, $W_E = U_C = \frac{1}{2} CV^2 = \frac{1}{2} QV = \frac{1}{2} \frac{Q^2}{C}$

(– see the following aside re development of W_E)

Aside:



$$\text{At } t = 0, \sigma = 0, \vec{E} = 0, \vec{F}_E = 0$$

\Rightarrow can move 1st dq across without doing work

Let $+q$ be on the RH plate at time t .

$$\text{Then } \sigma = \frac{q}{A}, \vec{E} = \frac{q}{A\epsilon_0}(-\hat{i})$$

$$\text{and } \vec{F}_E = |\vec{E}|dq(-\hat{i}) = \frac{q}{A\epsilon_0}dq(-\hat{i})$$

dW required to move $+dq$ across:

$$dW = \int_{x=0}^d \vec{F}_{ext} \cdot d\vec{s} = - \int_{x=0}^d \vec{F}_E \cdot d\vec{s} \quad (\text{note: } |F_E|(-\hat{i}) \cdot |dx| \hat{i} = -|F_E|dx)$$

$$= \frac{q dq}{A\epsilon_0} \int_0^d dx = \frac{q dq}{A\epsilon_0} d, \text{ but } C = \frac{A\epsilon_0}{d} \text{ so that } dW = \frac{q dq}{C}$$

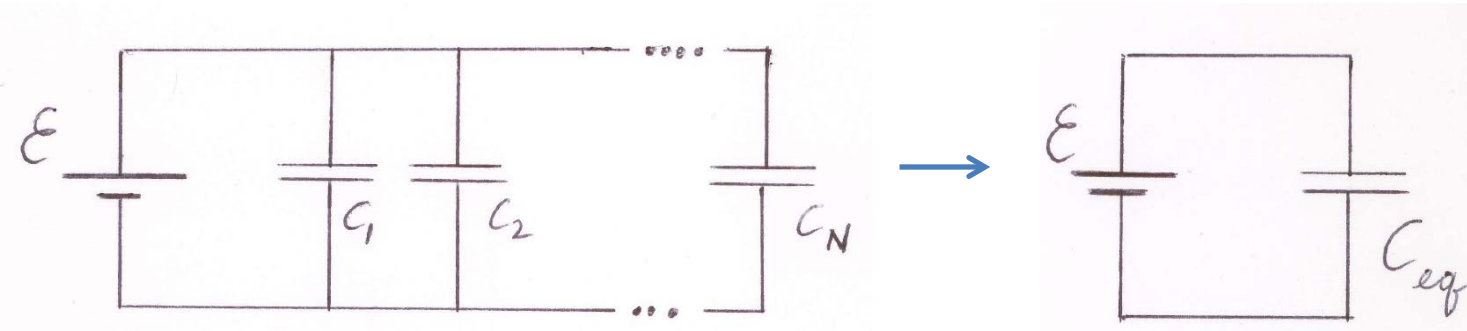
$$\therefore W = \int_{q=0}^Q \frac{q dq}{C} = \frac{Q^2}{2C} = U_C \quad (= \text{potential energy stored in } C)$$

(= work *you* or *battery* had to do)

End Aside

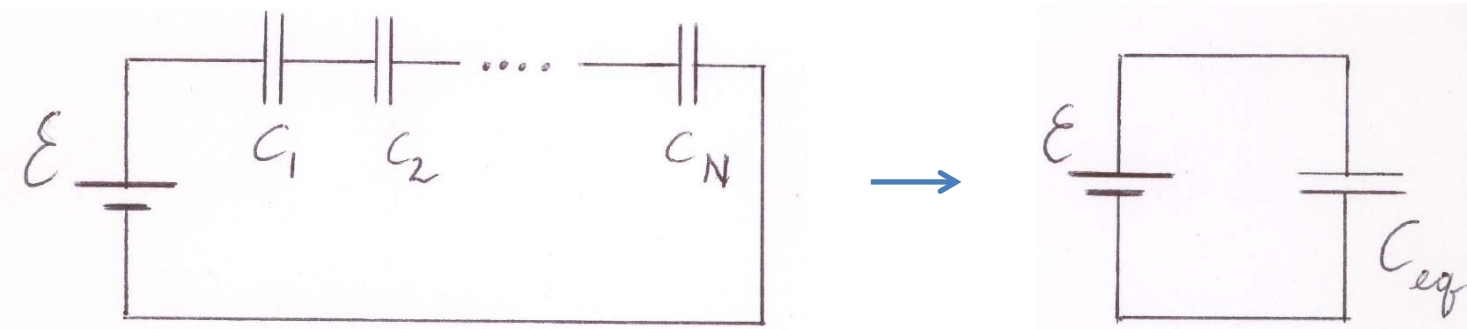
Capacitors in parallel and in series :

Parallel :



$$C_{eq} = \sum_{i=1}^N C_i$$

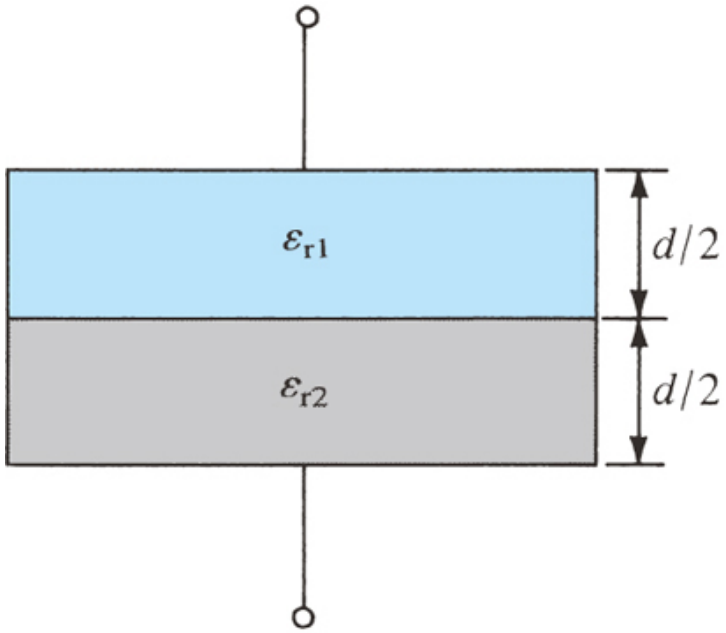
Series :



$$\frac{1}{C_{eq}} = \sum_{i=1}^N \frac{1}{C_i}$$

Example

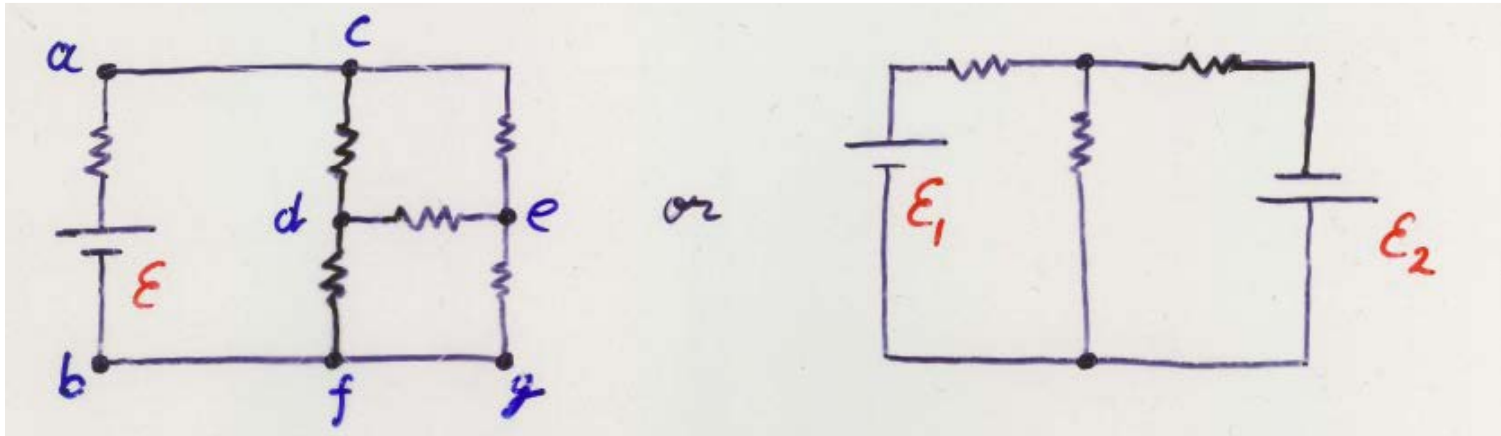
Determine the capacitance of the following :



2 Review of DC Circuits

OHM's Law: $V = I R$

and KIRCHHOFF's Rules (*for more complicated circuits*):

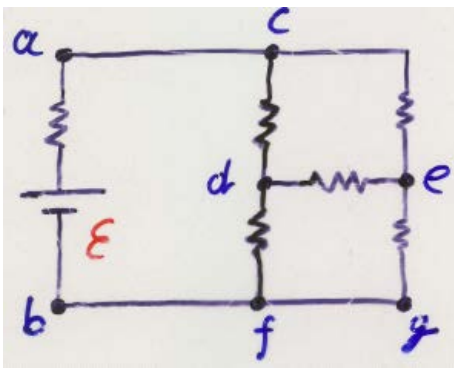


Branch point or node point:

a point where 3 or more conductors are joined; eg., c, d, f, e

Closed Loop:

closed conducting path; eg., acdfba, acegfba



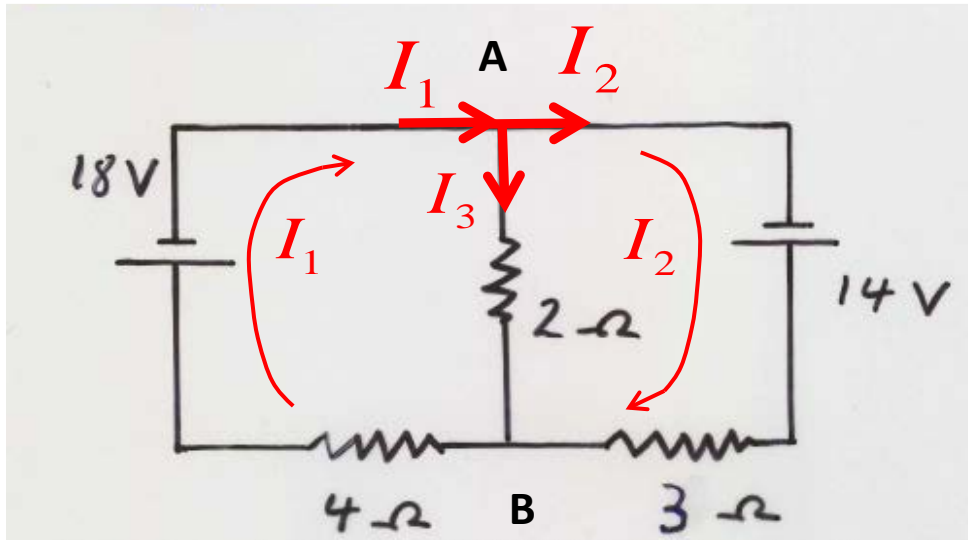
Point or Current *rule* : Algebraic sum of currents meeting at a branch point is zero.

$$\sum_i I_i = 0 \quad (\text{currents entering } +ve, \text{ currents leaving } -ve)$$

Voltage or Loop *rule* : Algebraic sum of potential differences in any loop must equal zero.
(independent (non – crossing) loops)

$$\sum_i \mathcal{E}_i - \sum_k (IR)_k = 0$$

Example Solve for all currents. Find V_{AB} .



Kirchhoff's current, point, or node rule:

- *point A:* $\sum I_{in} = \sum I_{out} \Rightarrow I_1 - I_2 - I_3 = 0$ or $I_1 = I_2 + I_3$ (i)
- *point B:* $I_2 + I_3 = I_1$; nothing new

Kirchhoff's loop rule:

- *loop 1:* $-18 - 2I_3 - 4I_1 = 0$ (ii)
 - *loop 2:* $14 - 3I_2 + 2I_3 = 0$ (iii)
- $\left\{ \begin{array}{l} \bullet 3 \text{ equations} \\ \text{(linear, nonhomogeneous)} \\ \bullet 3 \text{ unknowns} \\ \text{--use Cramer's rule} \end{array} \right.$

Solve to get $I_1 = -2.39$ A, $I_2 = 1.84$ A, $I_3 = -4.23$ A

To solve a set of linear, nonhomogeneous equations:

$$a_{11}x_1 + a_{12}x_2 + \dots + a_{1n}x_n = b_1$$

$$a_{21}x_1 + a_{22}x_2 + \dots + a_{2n}x_n = b_2$$

.

.

.

$$a_{n1}x_1 + a_{n2}x_2 + \dots + a_{nn}x_n = b_n$$

By Cramer's Rule :

$$x_1 = \frac{D_1}{D}, \quad x_2 = \frac{D_2}{D}, \quad \dots, \quad x_n = \frac{D_n}{D}$$

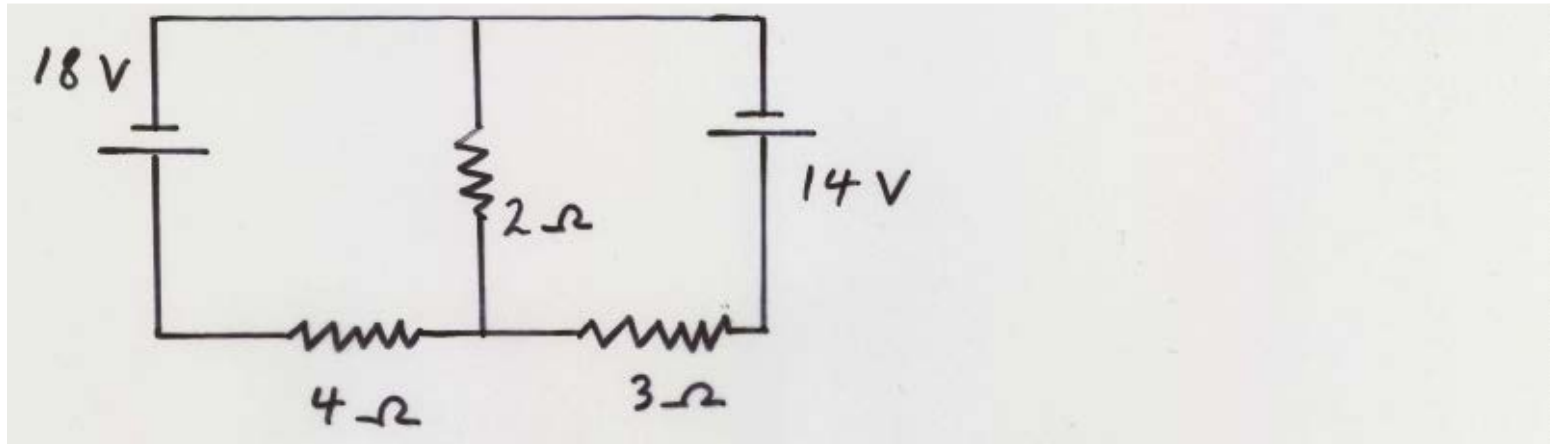
where

$$D = \begin{vmatrix} a_{11} & \dots & a_{1n} \\ a_{21} & \dots & a_{2n} \\ \cdot & & \\ \cdot & & \\ \cdot & & \\ a_{n1} & \dots & a_{nn} \end{vmatrix}$$

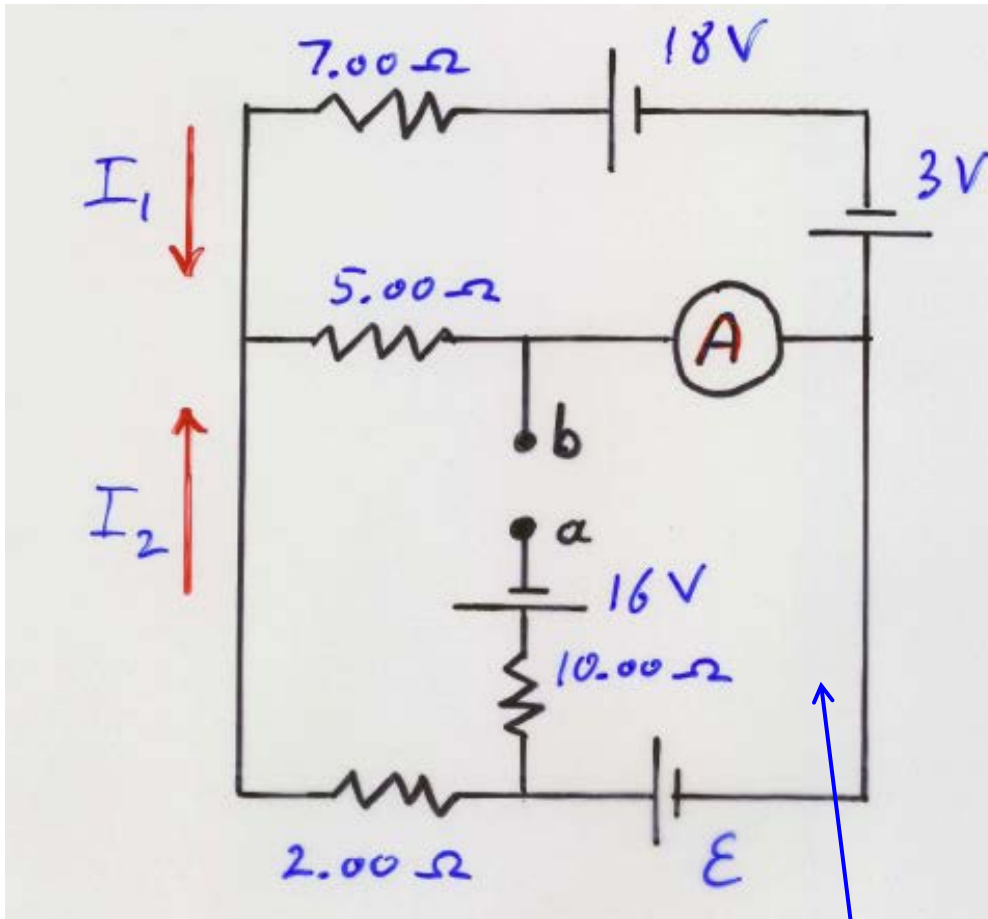
and where D_1, \dots, D_n are determinants obtained by replacing in D the k^{th} column by the column with elements b_1, b_2, b_n .

Example (same as previous example)

Solve for all currents using only the loop rules. Find $I_{2\Omega}$.



Example



The ammeter reading is 2.00 A.

Find I_1 , I_2 , \mathcal{E} and V_{ab} .

—can get $I_1 = 0.714 \text{ A}$

$$I_2 = 1.29 \text{ A}$$

$$\mathcal{E} = 12.6 \text{ V}$$

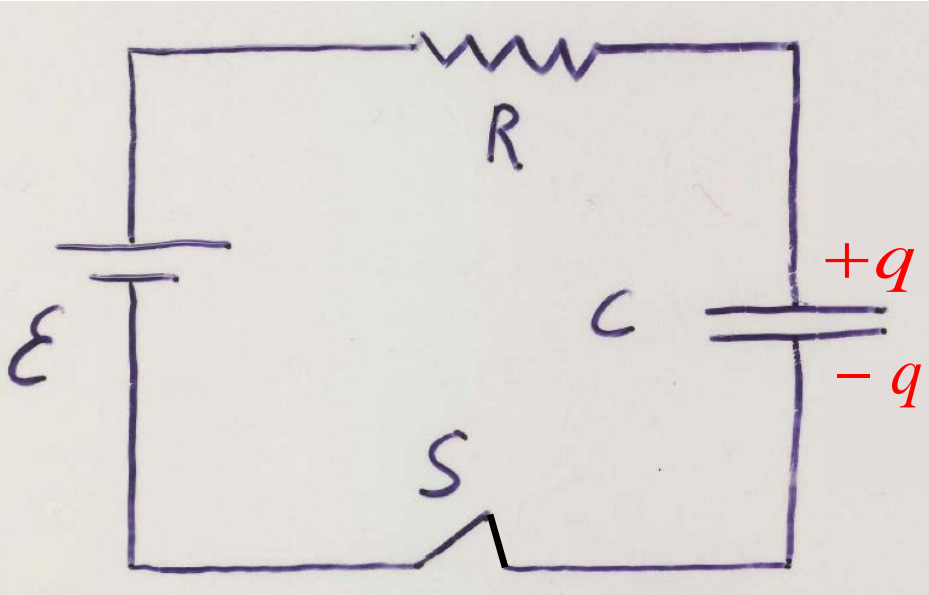
Although I is 0 can still write down a loop equation:

$$12.6 - 10.00 \times 0 - 16 - V_{ab} = 0$$

$$(V_{ab} = -3.4 \text{ V})$$

RC circuits (series, DC)

Charging:



$t < 0$:

$$Q_C = 0; V_C = 0; V_R = 0$$

– at $t = 0$ the switch is closed

– at $t = 0^+$:

$$Q_C = 0; V_C = 0; V_R = \mathcal{E}$$

$$I = \frac{V_R}{R} = \frac{\mathcal{E}}{R}$$

– as $t \rightarrow \infty$ C, is fully charged: $I = 0$; $V_R = 0$; $V_C = \mathcal{E}$ and $Q = C\mathcal{E}$

– at some time t :

q on C is $q(t)$; current is $i(t)$

K's loop rule:

$$\mathcal{E} - iR - \frac{q}{C} = 0$$

$$\mathcal{E} - iR - \frac{q}{C} = 0$$

but $i = \frac{dq}{dt}$ so that $\mathcal{E} - \frac{dq}{dt}R - \frac{q}{C} = 0$

or $(\mathcal{E} - \frac{q}{C})dt - Rdq = 0$

Separate variables $\frac{dq}{C\mathcal{E} - q} = \frac{1}{RC}dt$

or $\int_{q_i}^{q_{\text{final}}} \frac{dq}{C\mathcal{E} - q} = \frac{1}{RC} \int_{t_i}^{t_f} dt$

Here $q_i = 0$, $q_{\text{final}} = q$, $t_i = 0$, $t_f = t$

$$\therefore \int_0^q \frac{dq'}{C\mathcal{E} - q'} = \frac{1}{RC} \int_0^t dt'$$

$$\int_0^q \frac{dq'}{C\mathcal{E} - q'} = \frac{1}{RC} \int_0^t dt'$$

But $\int \frac{dx}{a+bx} = \frac{1}{b} \ln(a+bx)$

- here $a = C\mathcal{E}$ and $b = -1$

$$\therefore -\ln(C\mathcal{E} - q') \Big|_0^q = \frac{1}{RC} t \Big|_0^t$$

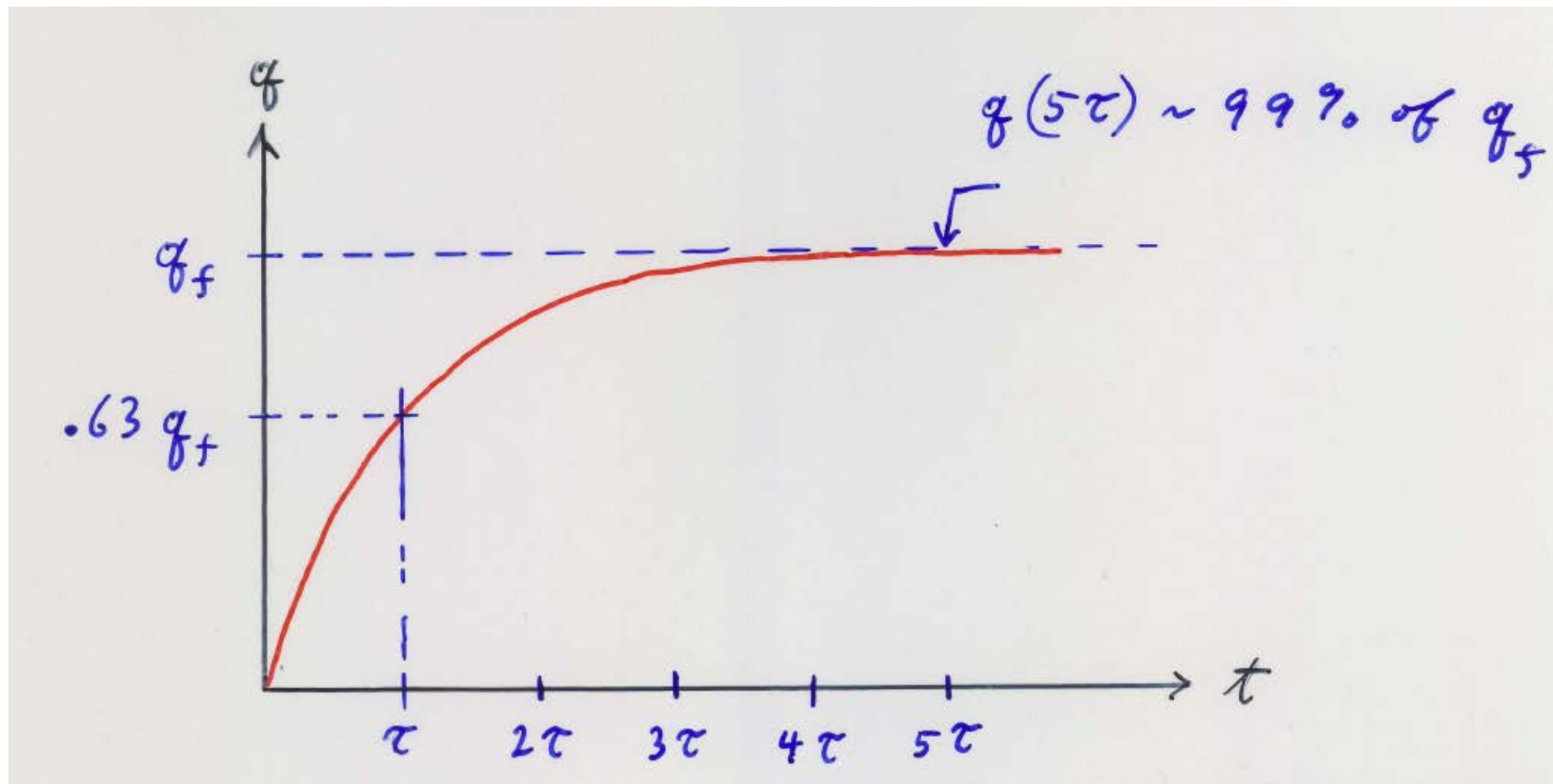
$$\text{or } \ln\left(\frac{C\mathcal{E} - q}{C\mathcal{E}}\right) = -\frac{t}{RC}$$

$$\therefore q(t) = C\mathcal{E} \left(1 - e^{-\frac{t}{RC}}\right)$$

$$\text{or } q(t) = q_{\text{final}} \left(1 - e^{-\frac{t}{RC}}\right)$$

$RC \equiv \text{time constant of the circuit} = \tau$

$$q(t) = q_{final} \left(1 - e^{-\frac{t}{RC}} \right)$$



$$q(t) = q_{final} \left(1 - e^{-\frac{t}{RC}} \right) \quad \therefore v_C(t) = \frac{q(t)}{C} = \mathcal{E} \left(1 - e^{-\frac{t}{RC}} \right)$$

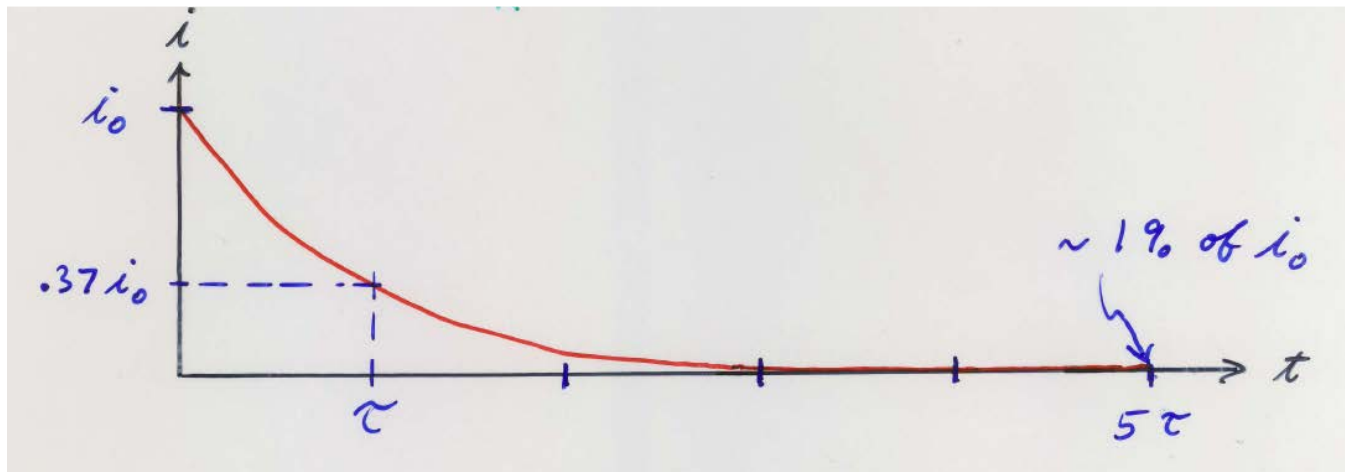
Note that when $t = 0$, $v_C = 0$

$t \rightarrow \infty$, $v_C = \mathcal{E}$

$v_C(t)$ follows the same behaviour with t as $q(t)$.

$$i(t) = \frac{dq(t)}{dt} = -C\mathcal{E}e^{-\frac{t}{RC}} \left(\frac{-1}{RC} \right)$$

$$i(t) = \frac{\mathcal{E}}{R} e^{-\frac{t}{RC}} = i_0 e^{-\frac{t}{RC}}$$



Example

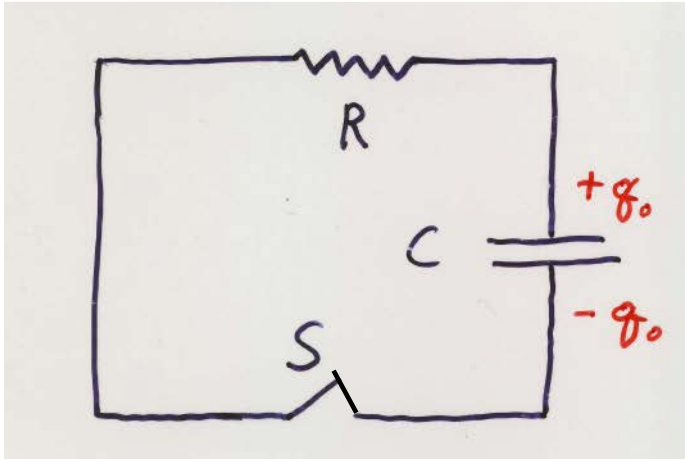
A $15\text{ k}\Omega$ resistor and capacitor are connected in series and a 12 V potential is suddenly applied. The potential across the capacitor rises to 5.0 V in $1.3\text{ }\mu\text{s}$.

a) Calculate the time constant. $(2.4 \times 10^{-6}\text{ s})$

b) Find the capacitance of the capacitor. (160 pF)

c) When will v_c be $\frac{1}{2}\mathcal{E}$? $(1.66\text{ }\mu\text{s})$

Discharging:



*Let the charge initially be
 q_0 or $q_i = \mathcal{E}C$*

With the switch open:

$$I = 0, \quad V_R = 0, \quad V_C = \mathcal{E} = \frac{q_0}{C}$$

At $t = 0$ close the switch: $i_0 = \frac{V_C}{R} = \frac{q_0}{CR}$

$V_R + V_C = 0$, so that at time t

$$-i(t)R + \frac{q(t)}{C} = 0$$

or
$$\frac{dq(t)}{dt} R + \frac{q(t)}{C} = 0$$

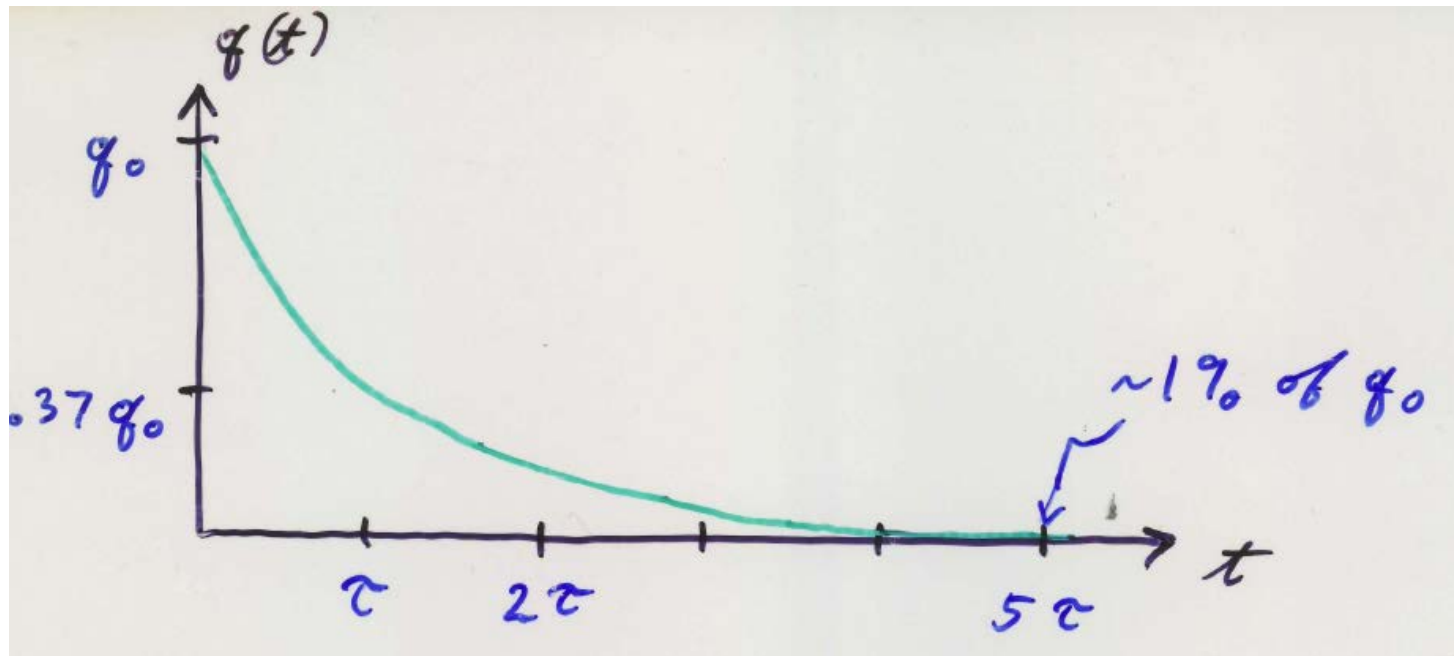
where we have set
 $i(t) = -dq/dt$

$$\frac{dq(t)}{dt} R + \frac{q(t)}{C} = 0$$

$$\int_{q_0}^q \frac{dq'}{q'} = -\frac{1}{RC} \int_0^t dt'$$

$$\ln \frac{q}{q_0} = -\frac{t}{RC}$$

$$q(t) = q_0 e^{-\frac{t}{RC}}$$



$$q(t) = q_0 e^{-\frac{t}{RC}}$$

$$\begin{aligned} \text{then } v_C(t) &= \frac{q(t)}{C} = \frac{q_0}{C} e^{-\frac{t}{RC}} \\ &= \frac{\mathcal{E}C}{C} e^{-\frac{t}{RC}} = \mathcal{E} e^{-\frac{t}{RC}} \end{aligned}$$

$$\begin{aligned} \text{here } i(t) &= -\frac{dq}{dt} = \frac{q_0}{RC} e^{-\frac{t}{RC}} \\ &= \frac{\mathcal{E}C}{RC} e^{-\frac{t}{RC}} = i_0 e^{-\frac{t}{RC}} \end{aligned}$$

$$\text{here } v_R(t) = -i(t)R = -\frac{\mathcal{E}}{R} e^{-\frac{t}{RC}} R = -\mathcal{E} e^{-\frac{t}{RC}}$$

Note that at $t = 0$

$$q = q_0, \quad i = \frac{\mathcal{E}}{R}, \quad v_R = -\mathcal{E}, \quad v_C = +\mathcal{E}$$