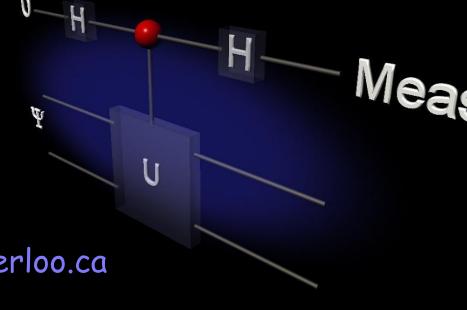
Introduction to Quantum Information Processing

CO481 CS467 PHYS467

Michele Mosca mmosca@iqc.uwaterloo.ca

Tuesdays and Thursdays 10am-11:15am





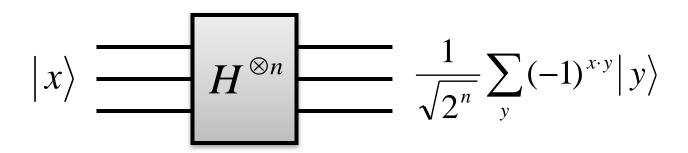




Reading

Chapter 6, sections 7.1.1, 7.1.3, 7.2, 7.3.1, 7.3.2, 7.3.3

Recall: Multi-qubit Hadamard



$$\frac{1}{\sqrt{2^n}} \sum_{y} (-1)^{x \cdot y} |y\rangle \qquad H^{\otimes n} \qquad |x\rangle$$

Quantum algorithms

- The algorithms we have seen have been computing "classical" functions on quantum superpositions
- This encoded information in the phases of the basis states: measuring basis states would provide little useful information
- But a simple quantum transformation translated the phase information into information that was measurable in the computational basis

Quantum factoring

- The security of many public key cryptosystems used in industry today relies on the difficulty of factoring large numbers into smaller factors.
- Factoring the integer N into smaller factors can be reduced to the following task:

Given integer a, find the smallest positive integer r so that $a^r \equiv 1 \mod N$

Complexity

•The best known rigorous classical algorithms use

$$e^{O(\sqrt{\log(N)\log\log(N)})}$$

operations

•The best known heuristic classical algorithms use

$$e^{O((\log(N)^{\frac{1}{3}}\log\log(N)^{\frac{2}{3}})}$$

operations

(Aside: how does factoring reduce to order-finding ?)

- The most common approach for factoring integers is the difference of squares technique:
 - "Randomly" find two integers x and y satisfying

$$x^2 = y^2 \mod N$$

- \triangleright So N divides $x^2 y^2 = (x y)(x + y)$
- \triangleright Hope that $\gcd(N, x y)$ is non-trivial
- If *r* is even, then let

$$x = a^{r/2} \mod N$$

$$x^2 = 1^2 \mod N$$

Quantum factoring

Since we know how to efficiently multiply by a mod N, we can efficiently implement

$$U_{a}|x\rangle = |ax\rangle$$

Note that

$$U_{a^r}|x\rangle = |a^r x\rangle = |x\rangle$$

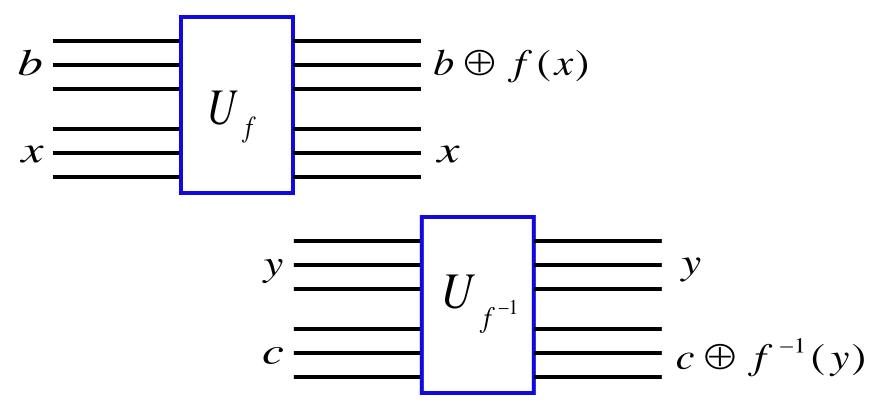
i.e.

$$U_a^r = I$$

Remember that $|\chi\rangle$ represents the state corresponding to the binary representation of x (e.g. for four qubits, $|2\rangle$ represents $|0010\rangle$)

(Aside: more on reversible computing)

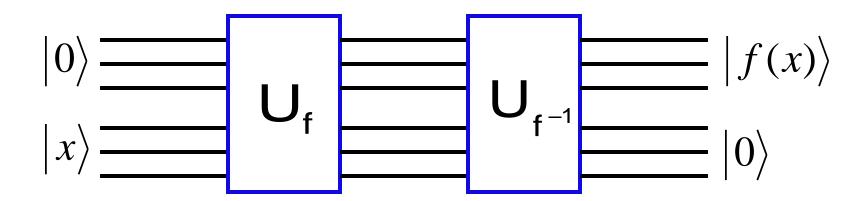
If we know how to efficiently compute f and f^{-1} then we can efficiently and reversibly map



(Aside: more on reversible computing)

And therefore, for such invertible f, we can efficiently map

$$|x\rangle \rightarrow |f(x)\rangle$$



Finding "

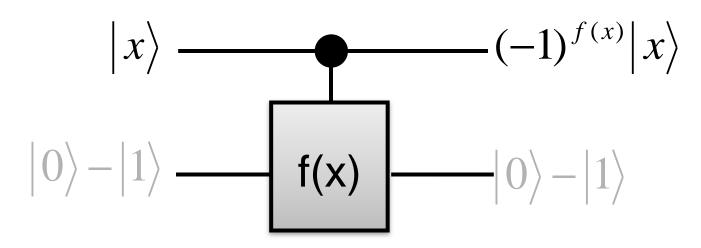
For most integers k, a good estimate of $\frac{k}{r}$ (with error at most $\frac{1}{2r^2}$) allows us to determine r (even if we don't know k).

(using continued fractions)

So what?

How would we convert the eigenvalues into something measurable?

Recall the "trick":



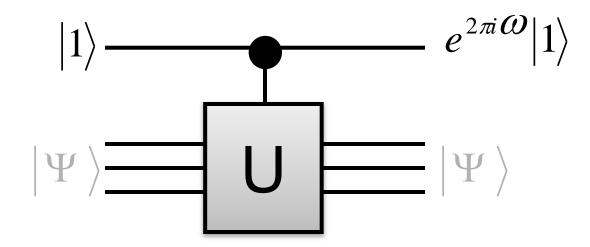
$$|x\rangle(|0\rangle - |1\rangle) \rightarrow |x\rangle(|f(x)\rangle - |f(x) \oplus 1\rangle)$$

$$= |x\rangle(-1)^{f(x)}(|0\rangle - |1\rangle)$$

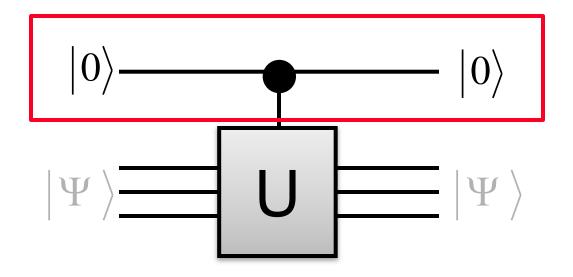
$$= (-1)^{f(x)}|x\rangle(|0\rangle - |1\rangle)$$

Eigenvalue kick-back (Kitaev)

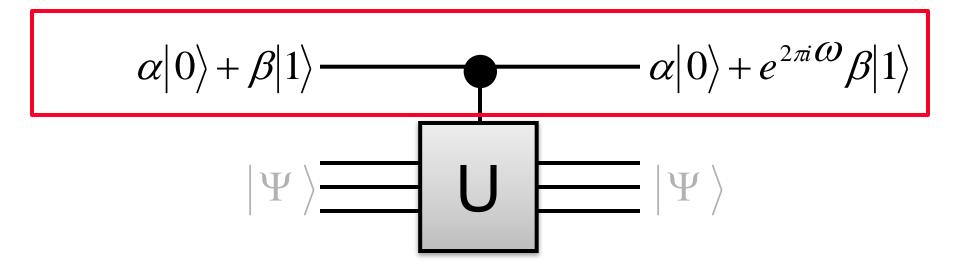
· Consider a unitary operation **U** with eigenvalue $e^{2\pi i \omega}$ and eigenvector $|\Psi
angle$



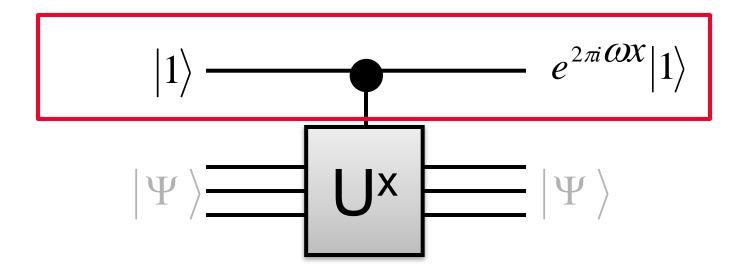
$$|1\rangle|\Psi\rangle \to |1\rangle U|\Psi\rangle = |1\rangle e^{2\pi i \omega}|\Psi\rangle$$
$$= e^{2\pi i \omega}|1\rangle|\Psi\rangle$$

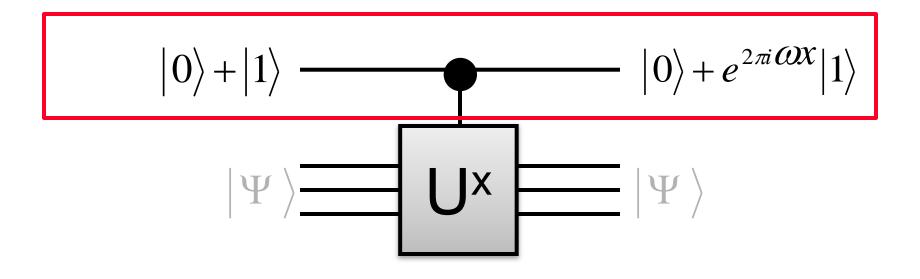


• As a relative phase, $e^{2\pi i \omega}$ becomes measurable



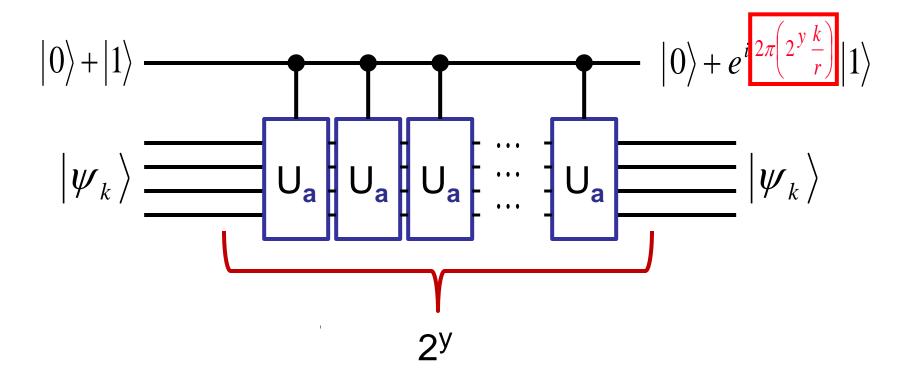
• If we exponentiate \mathbf{U} , we get multiples of ω





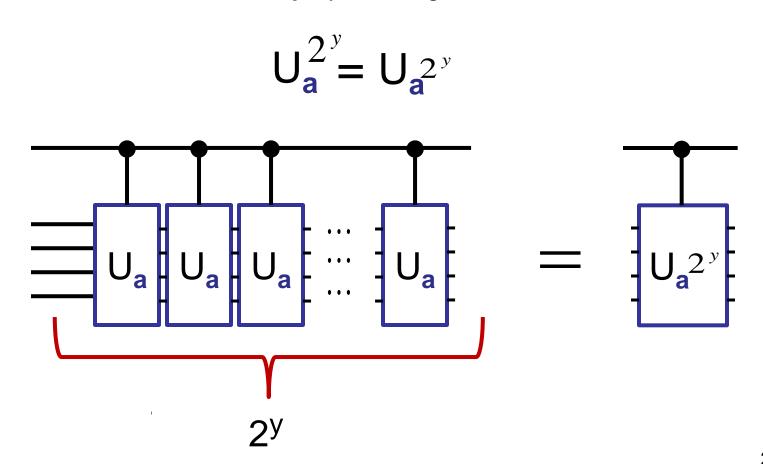
Inefficient exponentiation

We can effect a relative phase shift of $e^{i 2\pi \left(2^{y} \frac{k}{r}\right)}$



Efficient exponentiation

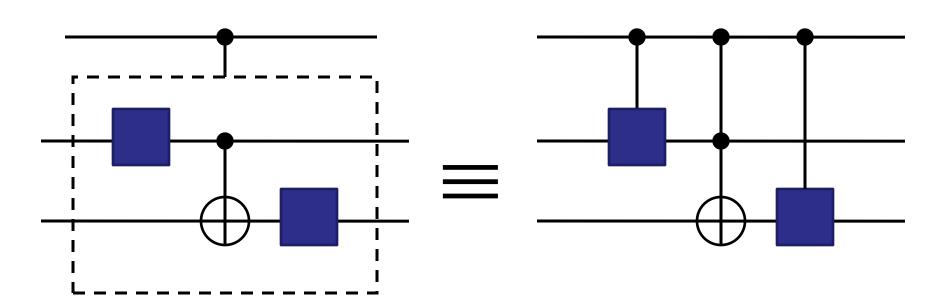
But we can also do it **efficiently** by noticing that



How do we implement c-U ?

Replace every gate G in the circuit with a c-G.

For example,



Next step?

• We thus know how, given an eigenvector with eigenvalue $e^{2\pi i\left(2^y\frac{k}{r}\right)}$, to construct

$$= \left(\left| 0 \right\rangle + e^{2\pi i (2^{n-1}\frac{k}{r})} \left| 1 \right\rangle \right) \otimes \left(\left| 0 \right\rangle + e^{2\pi i (2^{n-2}\frac{k}{r})} \left| 1 \right\rangle \right) \otimes \cdots \otimes \left(\left| 0 \right\rangle + e^{2\pi i (\frac{k}{r})} \left| 1 \right\rangle \right)$$

Useful identity

We can show that

$$\left(|0\rangle + e^{2\pi i(2^{n-1}\omega)} |1\rangle \right) \otimes \left(|0\rangle + e^{2\pi i(2^{n-2}\omega)} |1\rangle \right) \otimes \cdots \otimes \left(|0\rangle + e^{2\pi i(\omega)} |1\rangle \right)$$

$$=\sum_{y=0}^{2^{n}-1}e^{2\pi i\omega y}|y\rangle$$

 Suppose we wish to estimate a number ω ∈ [0,1) given the quantum state

$$\sum_{y=0}^{2^{n}-1} e^{2\pi i \omega y} |y\rangle$$

Note that in binary we can express

$$\omega = 0.x_1x_2x_3...$$

$$2\omega = x_1.x_2x_3...$$

$$2^{n-1}\omega = x_1 x_2 x_3 \dots x_{n-1} x_n x_{n+1} \dots$$

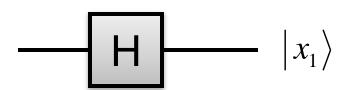
• Since $e^{2\pi ik} = 1$ for any integer k, we have

$$e^{2\pi i(2\omega)} = e^{2\pi i(x_1.x_2x_3...)} = e^{2\pi ix_1}e^{2\pi i(0.x_2x_3...)} = e^{2\pi i(0.x_2x_3...)}$$

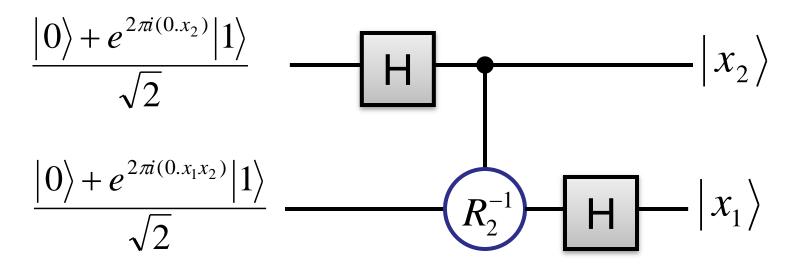
$$e^{2\pi i(2^k\omega)} = e^{2\pi i(0.x_{k+1}x_{k+2}...)}$$

• If $\omega = 0.x_1$ then we can do the following

$$\frac{\left|0\right\rangle + e^{2\pi i(0.x_1)}\left|1\right\rangle}{\sqrt{2}} = \frac{\left|0\right\rangle + \left(-1\right)^{x_1}\left|1\right\rangle}{\sqrt{2}}$$

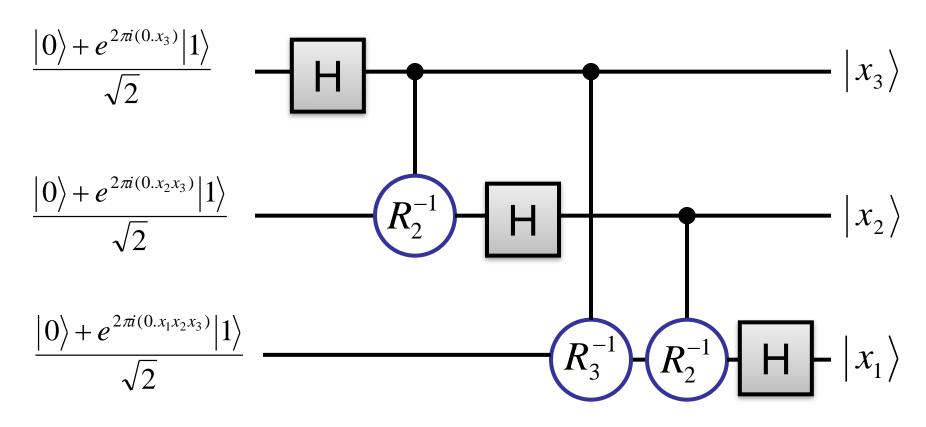


• So if $\omega = 0.x_1x_2$ then we can do the following



$$R_k = \begin{bmatrix} 1 & 0 \\ 0 & e^{2\pi i/2^k} \end{bmatrix}$$

• So if $\omega = 0.x_1x_2x_3$ then we can do the following



 Generalizing this network (and reversing the order of the qubits at the end) gives us a network with O(n²) gates that implements

$$\sum_{y=0}^{2^{n}-1} e^{2\pi i \frac{x}{2^{n}}} y | y \rangle \mapsto | x \rangle$$

Discrete Fourier transform

 The discrete Fourier transform maps vectors of dimension N by transforming the Xth elementary vector according to

$$(0,0,...,0,1,0,...0) \mapsto (1,e^{2\pi i \frac{x}{N}},e^{2\pi i \frac{2x}{N}},...,e^{2\pi i \frac{(N-1)x}{N}})$$

 The quantum Fourier transform maps vectors in a Hilbert space of dimension N according to

$$|x\rangle \mapsto \sum_{y=0}^{N-1} e^{2\pi i \frac{x}{N}y} |y\rangle$$

Discrete Fourier transform

 Thus we have illustrated how to implement (the inverse of) the quantum Fourier transform in a Hilbert space of dimension 2ⁿ

Estimating arbitrary ω ε [0,1)

- What if ω is not necessarily of the form $\frac{x}{2^n}$ for some integer x?
- The QFT will map

$$\sum_{x=0}^{2^n-1} e^{2\pi i \omega z} |z\rangle$$

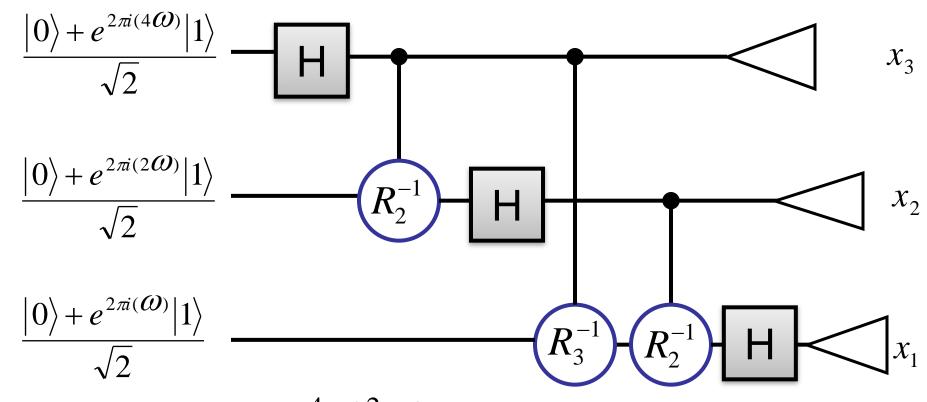
to a superposition

$$\left|\widetilde{\omega}\right\rangle = \sum_{y} \alpha_{y} \left| y \right\rangle$$

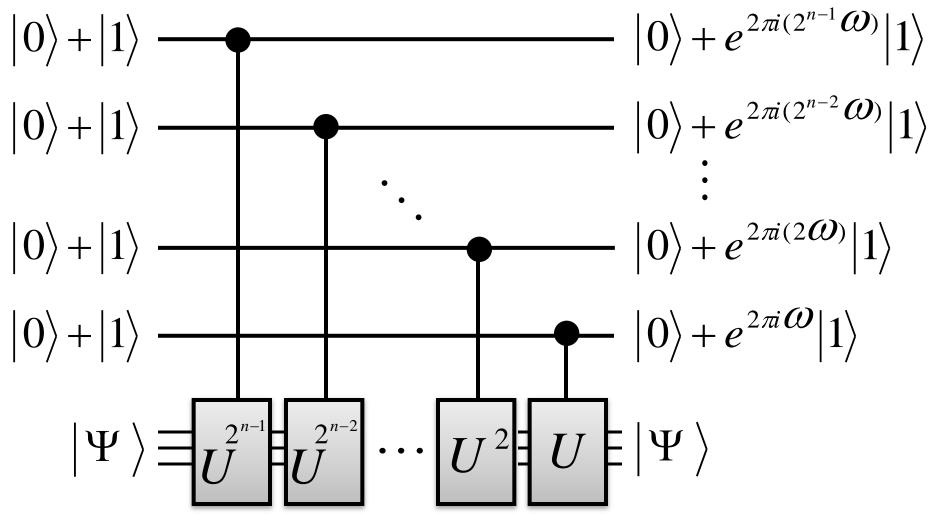
where

here
$$\Pr{ob}\left(\left|\frac{y}{N} - \omega\right| \le \frac{1}{N}\right) \ge \frac{8}{\pi^2} \qquad \left|\alpha_y\right| \in O\left(\frac{1}{\left|\frac{y}{N} - \omega\right|}\right)$$

• For any real $\omega \in [0,1)$

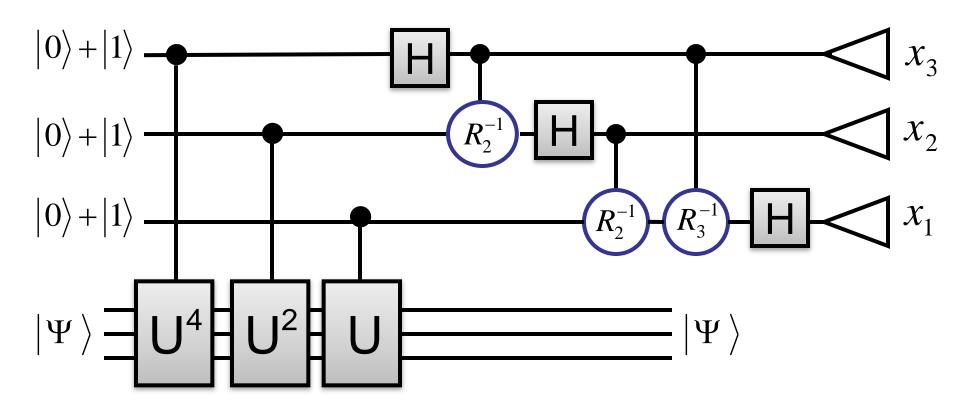


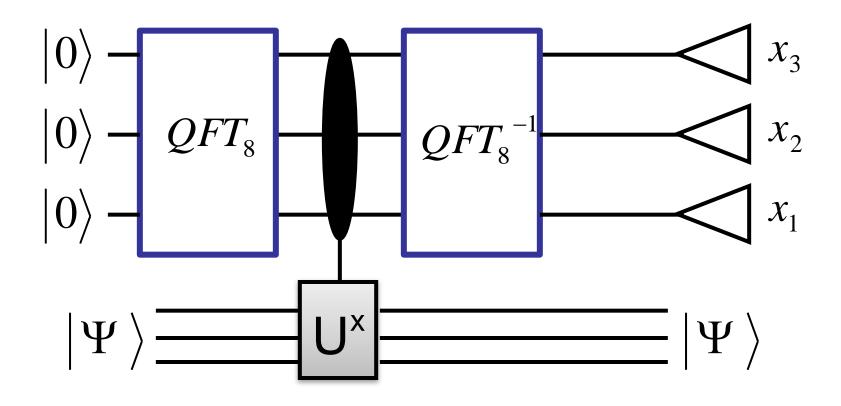
• With high probability $\frac{4x_1+2x_2+x_3}{8} \approx \alpha$



Phase estimation

$$\frac{2^{n-1}x_1+2^{n-2}x_2+\cdots x_n}{2^n} \approx \omega$$





- Given **U** with eigenvector $|\Psi\>\rangle$ and eigenvalue ${\bf e}^{2\pi{\rm i}\omega}$, we thus have an algorithm that maps

$$|0\rangle |\Psi\rangle \rightarrow |\tilde{\omega}\rangle |\Psi\rangle$$

• Given **U** with eigenvectors $|\Psi_{\bf k}\rangle$ and respective eigenvalues ${\bf e}^{2\pi i \Theta_{\bf k}}$ we thus have an algorithm that maps

$$|0\rangle\!|\Psi_{k}\rangle\!\rightarrow\!|\widetilde{\omega}_{k}\rangle\!|\Psi_{k}\rangle$$

and therefore

$$\left|0\right\rangle\!\sum_{k}\alpha_{k}\!\left|\Psi_{k}\right.\right\rangle\!=\!\sum_{k}\alpha_{k}\!\left|0\right\rangle\!\!\left|\Psi_{k}\right.\right\rangle\!\to\!\sum_{k}\alpha_{k}\!\left|\widetilde{\omega}_{k}\right\rangle\!\!\left|\Psi_{k}\right.\right\rangle$$

Measuring the first register of

$$\sum_{\mathbf{k}} \alpha_{\mathbf{k}} \big| \widetilde{\omega}_{\mathbf{k}} \big\rangle \big| \Psi_{\mathbf{k}} \big\rangle$$

is equivalent to measuring $|\widetilde{\omega}_{\mathbf{k}}\rangle$ with probability $|lpha_{k}|^{2}$.

Eigenvectors

We know the eigenvalues of U_a are of the form

$$e^{2\pi i \frac{k}{r}}$$

$$U_{a} | \psi_{k} \rangle = e^{i2\pi \frac{k}{r}} | \psi_{k} \rangle$$

$$| \psi_{k} \rangle = \sum_{j=0}^{r-1} e^{-i2\pi j \frac{k}{r}} | a^{j} \rangle$$

Checking the eigenvalues

$$U_a | \psi_k \rangle = \sum_{j=0}^{r-1} e^{-i2\pi j \frac{k}{r}} U_a | a^j \rangle$$

$$= \sum_{j=0}^{r-1} e^{-i2\pi j \frac{k}{r}} |a^{j+1}\rangle = e^{i2\pi \frac{k}{r}} \left(\sum_{j=1}^{r} e^{-i2\pi j \frac{k}{r}} |a^{j}\rangle \right)$$

$$= e^{i2\pi \frac{k}{r}} \left(\sum_{j=0}^{r-1} e^{-i2\pi j \frac{k}{r}} |a^j\rangle \right) = e^{i2\pi \frac{k}{r}} |\psi_k\rangle$$

Note that

$$|1\rangle = \sum_{k=0}^{r-1} \frac{1}{\sqrt{r}} |\psi_k\rangle$$

Eigenvalue Estimation Factoring Algorithm

$$|0\rangle|1\rangle \mapsto \sum_{k=0}^{r-1} \sum_{x} |x\rangle|\psi_k\rangle$$

$$\mapsto \sum_{k=0}^{r-1} \sum_{x} e^{2\pi i k x/r} |x\rangle |\psi_k\rangle$$

$$\sum_{k} \left(\int_{\frac{k}{r}} |\psi_{k}\rangle \right)$$

Circuit for Eigenvalue Estimation Factoring Algorithm

