

Turbulence and Transport in Fusion Plasmas

Part II



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Monday Recap

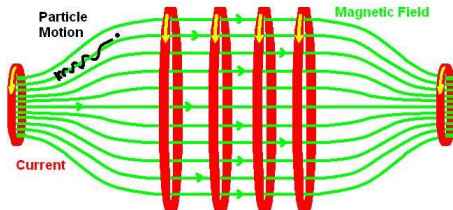
Yesterday, we covered

- how fusion energy will provide clean and safe energy
- the nature of Debye shielding and how to define plasmas
- what plasma waves and plasma drifts are

Next: how to design a confining magnetic field?

Magnetic Mirror

Early magnetic confinement: the **mirror machine**
(particles trapped between regions of large B)



Conserved: magnetic moment $\mu = \frac{mv_{\perp}^2}{2B}$, energy $E = \frac{1}{2}mv_{\parallel}^2 + \frac{1}{2}mv_{\perp}^2$

\Rightarrow particles moving from small B to larger B convert $v_{\parallel} \rightarrow v_{\perp}$
at $v_{\parallel} = 0$, particles are **reflected** (\leftrightarrow **magnetic mirror**)

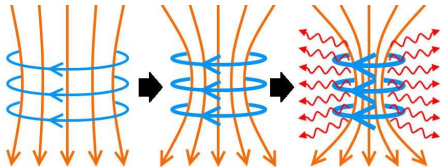
If $v_{\perp}/v < \sqrt{B_{\min}/B_{\max}}$, **no reflection, particle slips through**

Non-zero ν_{ei} : trapped particles can enter this *loss cone*

\Rightarrow *mirror machines have poor confinement due to end losses*

Z-Pinch

Another linear device (of Ocean's Eleven fame): the **Z-pinch**



current filaments
attract, compress field
 \Rightarrow increased n , but
only for short τ_E

Similar for lightning in weak rod; can use to shape metals, see

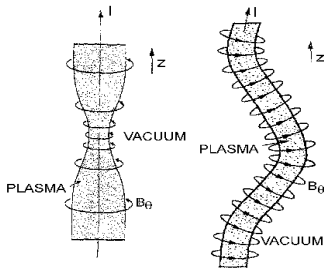
sciencedemonstrations.fas.harvard.edu/presentations/can-crusher-magnetic-implosion

However, no stable equilibrium!

Kink/sausage instabilities

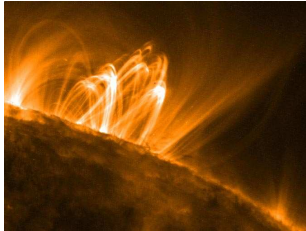
(different from transport-relevant instabilities in fusion reactors)

- **Sausage:** pinch effect
- **Kink:** prevents pinch from working



*A short excursion into a pinch-like phenomenon
in space plasmas . . .*

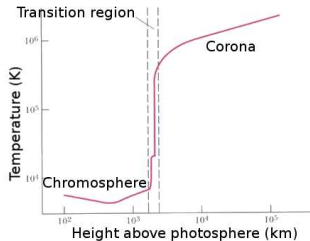
The Coronal Heating Problem



Decade-old question: *how can the solar corona be so hot?*

Heating mechanism:

- MHD waves?
- loop footpoint motion?
- kinetic Alfvén waves?
- turbulent reconnection?
- nanoflares?

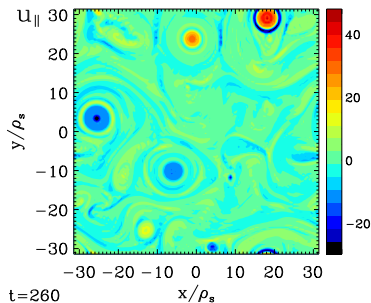
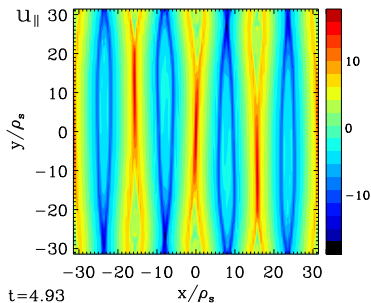


Driven Reconnection Turbulence

Using a fusion turbulence simulation code, study **magnetic reconnection** in a coronal loop (Pueschel ApJS 2014)

Driving turbulence through x - y -**periodic current sheets**,

$$\left. \frac{\partial g_j}{\partial t} \right|_{\text{drive}} = -\omega_{\text{dr}} \left(g_j(k_y = 0, t) - g_j(k_y = 0, t = 0) \right)$$

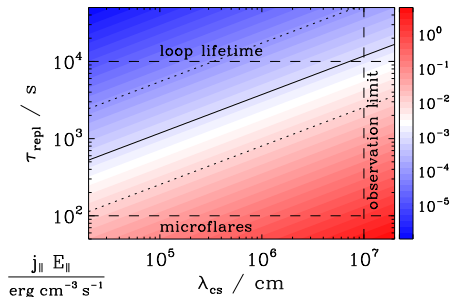


- **isotropization** despite driven current sheet
- formation of **plasmoids**, consistent with $S \gtrsim 10^4$ criterion

Application: Coronal Heating

Many turbulence simulations scanning over parameters:
extract **volumetric heating rate**

$$\frac{j_{\parallel} E_{\parallel}}{\text{erg cm}^{-3} \text{ s}^{-1}} = 1.5 \times 10^{-3} \left(\frac{n_{e0}}{10^9 \text{ cm}^{-3}} \right)^{0.375} \left(\frac{T_{e0}}{10^6 \text{ K}} \right)^{-0.1} \left(\frac{L_{\text{ref}}}{10^9 \text{ cm}} \right)^{0.2} \left(\frac{\beta}{3.5 \times 10^{-4}} \right)^{0.125} \\ \times \left(\frac{\nu_{ei}}{87 \text{ s}^{-1}} \right)^{-0.1} \left(\frac{m_e/m_i}{1/1836} \right)^{-0.25} \left(\frac{T_{i0}/T_{e0}}{1.0} \right)^{0.0} \left(\frac{B_{\text{rec}}}{5 \text{ G}} \right)^{1.8} \left(\frac{\tau_{\text{repl}}}{1100 \text{ s}} \right)^{-1.5} \left(\frac{\lambda_{\text{cs}}}{1.5 \times 10^5 \text{ cm}} \right)^{0.75}$$



Compare **observations**

Withbroe 77: $10^{-3} \frac{\text{erg}}{\text{cm}^3 \text{ s}}$
(or $L_{\text{ref}}^{-1} \times 10^6 \frac{\text{erg}}{\text{cm}^2 \text{ s}}$)

Ofman 98: $10^{-4} \frac{\text{erg}}{\text{cm}^3 \text{ s}}$

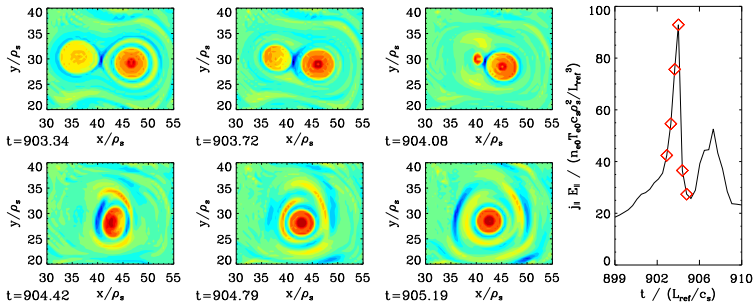
Guarrasi 14/Rosner 78:
 $2 \times 10^{-3} \frac{\text{erg}}{\text{cm}^3 \text{ s}}$

⇒ observations **well within predicted range** of heating!

Nanoflares

Nanoflares: short ($t \gtrsim 20$ s) spikes in observed heating

Look at details of turbulence: **merging current filaments**



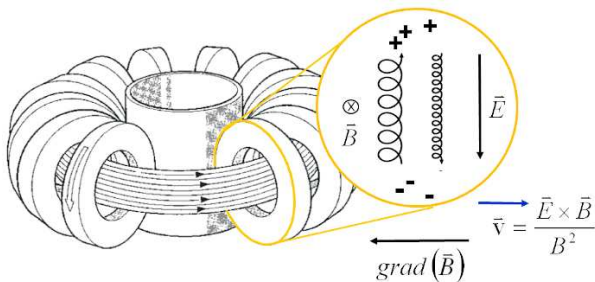
\Rightarrow **mergers** are the **cause** of heating rate **spikes**

Apply to Active Regions: $7 \text{ s} \lesssim t_{\text{merge}} \lesssim 50 \text{ s} \Rightarrow$ **match!**

Similar, but kinetic, process to the compression in a Z-pinch

Toroidal Confinement I

Bending a linear B -field into a torus. . . *what could go wrong?*



⇒ **near-instantaneous loss of all particles to reactor wall**

- two new plasma drifts (more later):

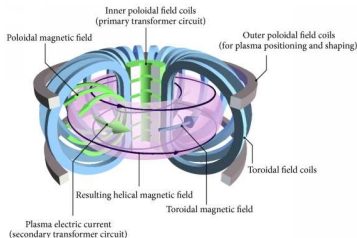
$$v_{\nabla B} \propto q_{i,e}^{-1} \text{ and } v_{E \times B} \propto q_{i,e}^0$$

- $v_{\nabla B}$ separates charges, $v_{E \times B}$ lets them escape

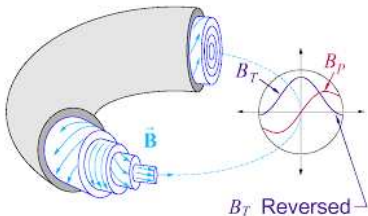
Since 1950s, many different concepts to deal with this problem

Toroidal Confinement II

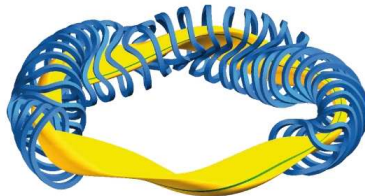
Tokamak: Soviet invention, transformer induces current



Reversed-field pinch (RFP): cheap tokamak, poor stability



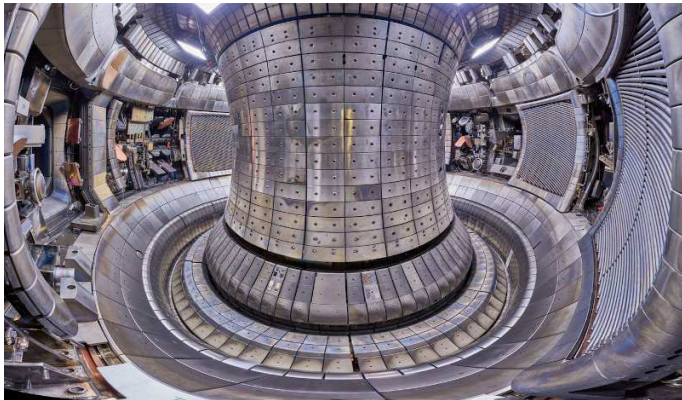
Stellarator: U.S. invention, complex coils, optimizable



Spherical Tokamak (ST): high efficiency, poor stability



Tokamak: ASDEX Upgrade



- fits 2–8 workers
- wall tiles
- divertor
- diagnostics
- ICRH antennae
- NBI shaft
- ECRH system
- pellet injector

Questions & Discussion

What's unclear? Who needs help? Who needs fresh tea?

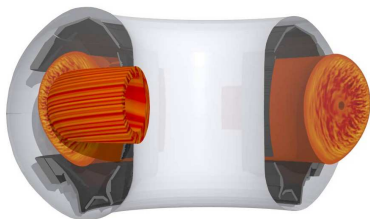
10-minute self-evaluated quiz:

- 1 What elements are to be used as fuel in fusion power plants, and why?
- 2 What are the safety considerations for fusion reactors, and how restrictive are they?
- 3 Why are modern fusion experiments toroidal and not linear?
- 4 What is the Debye length, and how does it affect whether an ionized gas is a plasma?

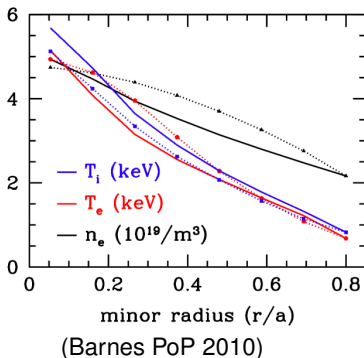
Quiz: Answers

- 1 Highest reaction cross-section: deuterium-tritium
(*note*: alternate concepts rely on less efficient aneutronic reactions)
- 2 No runaway reactions, only true concern is tritium inventory; no substantial worry (but need smart choices for materials for reactor walls, steel)
- 3 Linear machines are subject to large end losses, poor confinement
- 4 Inside the Debye length, charge is shielded, no collective behavior, but gas-like

Transport



Fusion plasmas: radial distribution of n , T_i , T_e , ...



Transport: heat/particle/momentum losses on time scale of heat/particle/momentum confinement time τ
(\longleftrightarrow triple product $nT\tau_E$)

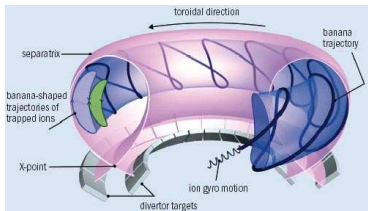
- What are the causes for these losses?
- How can we describe/predict profile evolution?

(Neo-)classical vs. Anomalous Flux

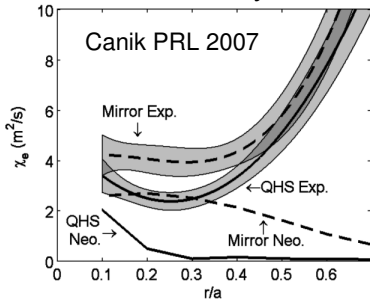
Historically, main heat loss mechanism thought to be
classical \rightarrow *neoclassical* \rightarrow *anomalous*

Classical: ignorably small (in 1950s, vision of tabletop reactors)

Neoclassical = collisions on banana orbits: small-ish



Anomalous: usually dominates



Historically, *anomalous* = *no idea where it comes from*,
but now used as synonym for **turbulence**

Plasma Instability

Fluctuation measurements: plasma “bubbles” like boiling water

⇒ *How does a physicist describe a bubbling system?*

Plasma Instability

Fluctuation measurements: plasma “bubbles” like boiling water

⇒ *How does a physicist describe a bubbling system?*

... we look at a **small perturbation from equilibrium**

E.g., density

$$n \rightarrow n_0 + n_1 \quad \epsilon n_0 \sim n_1$$

⇒ can solve with ansatz

$$n_1 = n_1(t=0) \exp(i\omega_c t)$$

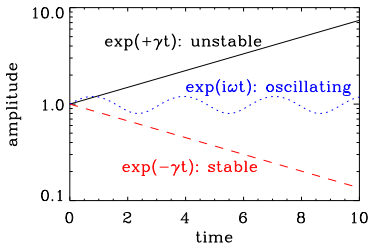
In equilibrium (fusion: MHD),

(ω_c : complex frequency)

$$\frac{\partial n_0}{\partial t} = 0$$

$$\frac{\partial n_1}{\partial t} = An_1 + Bn_1^2 + \dots$$

Linearization: drop
terms $\propto \epsilon^2$ and higher



Describes eigenmode growing exponentially at rate γ forever
→ *what stops this growth?*

Saturation and Turbulence

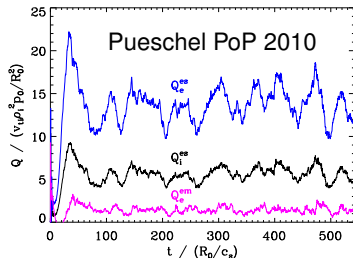
As n_1 grows, eventually $n_1^2 \sim n_1$, and **nonlinearity** becomes important

In plasmas, typically

$$\frac{\partial n_1}{\partial t} = A n_1 + \frac{\partial n_1}{\partial x} \frac{\partial n_1}{\partial y}$$

- here, quadratic nonlinearity
- saturation: **balance** of linear, nonlinear terms
- generally, **cannot solve analytically**

Solve equations numerically



quasi-stationary state
after linear growth, saturation
 \Rightarrow **statistical analysis**

Turbulence: very difficult, decades of research in many fields

Group Work: Turbulence Basics

45 minutes group work:

Find online sources that explain

- chaos
- random walk
- the Kolmogorov spectrum

and be prepared to give a short summary.

Did your sources explain what the Lyapunov exponent is?

Volunteers ... Who wants to tell us:

What does **chaos**
mean to a physicist?

What is and when do we
observe **random walk**?

What is the
Lyapunov exponent?

What is the
Kolmogorov spectrum?

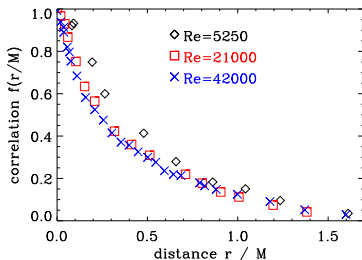
Kolmogorov Turbulence

Kolmogorov 1941

turbulence theory based
on **self-similarity**

Stewart 1951

wind tunnel experiments:
no viscosity dependence
of low- ν (high-Re) turbulence



dissipation rate $\epsilon = 2\nu \int k^2 E dk$, with $[E] = \frac{m^3}{s^2}$, $[\epsilon] = \frac{m^2}{s^3}$, $[\nu] = \frac{m^2}{s}$

dimensional analysis: $E(k) \propto u(k)^2 \lambda(k) [k\lambda]^n$ (self-similarity/scale-free)

to get rid of ν , use $[\lambda] = m = [\frac{\nu^{3/4}}{\epsilon^{1/4}}]$ and $[u] = \frac{m}{s} = [\nu^{1/4} \epsilon^{1/4}]$

\Rightarrow to make energy ν -independent, choose $n = -5/3$:

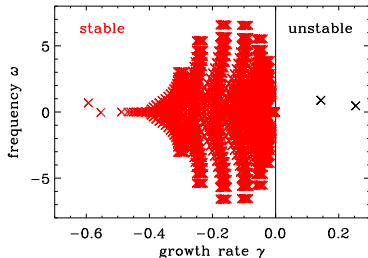
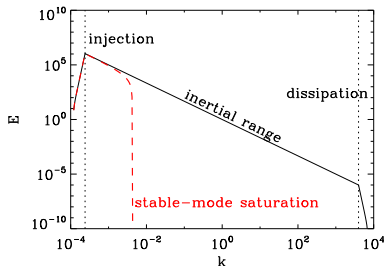
$$E(k) \propto k^n \nu^{\frac{3n+5}{4}} \epsilon^{\frac{1-n}{4}} = k^{-5/3} \epsilon^{2/3}$$

(Kolmogorov inertial range, confirmed in many experiments)

Beyond Kolmogorov

Kolmogorov 1941: transfer energy conservatively across scales

Full solution to linear operator A



However, instability-driven plasma systems can have energy **injection/dissipation/removal at the same scales!**

- large-scale unstable modes inject energy ($\nabla T_{e,i}$, ∇n)
- $10^4 +$ large-scale **stable modes** remove energy

\Rightarrow **no true inertial range**, no $5/3$ law

Fluid Model for Plasmas I

For very simple plasma description (Horton PoF 1981, 1988), can use (normalized) **nonlinear two-field fluid model** with

- electric potential $\Phi(x, y)$
- (ion or electron) pressure $p(x, y)$

$$\begin{aligned}\frac{\partial \Phi}{\partial t} &= (\delta - \nabla_{\perp}^2)^{-1} \left(\overbrace{-\frac{\partial \Phi}{\partial y} + \epsilon \frac{\partial p}{\partial y}}^{\text{drifts}} + \nu \nabla_{\perp}^2 \Phi - \hat{z} \cdot (\nabla_{\perp} \Phi \times \nabla_{\perp}^3 \Phi) \right) \\ \frac{\partial p}{\partial t} &= \underbrace{(-1 - \eta) \frac{\partial \Phi}{\partial y}}_{\text{drive}} - \chi \nabla_{\perp}^4 p - \hat{z} \cdot (\nabla_{\perp} \Phi \times \nabla_{\perp} p)\end{aligned}$$

With: $\epsilon \sim 1/\nabla n$, “collisionalities” ν, χ , coordinate $\hat{z} \parallel \mathbf{B}$, radial coordinate x , toroidal coordinate y

Physics: toroidal ion-/electron-temperature-gradient modes, accessing free energy in background ∇T via $\eta \equiv d \ln T / d \ln n$

Going to Fourier space easy for linear terms, $\partial/\partial x, y \rightarrow ik_{x,y}$, *but what happens to the nonlinearities?*

Group Work: Fourier Convolution

30 minutes group work:

What does the expression

$$\{f, \Phi\} \equiv \frac{\partial f}{\partial x} \frac{\partial \Phi}{\partial y} - \frac{\partial f}{\partial y} \frac{\partial \Phi}{\partial x}$$

look like when in Fourier space,

$f(x, y) \rightarrow f(k_x, k_y)$ and $\Phi(x, y) \rightarrow \Phi(k_x, k_y)$, with e.g.

$$\mathcal{F}(\partial_y \Phi) = \int d\mathbf{x} e^{-i\mathbf{k}\mathbf{x}} \sum_{\mathbf{k}'} e^{i\mathbf{k}'\mathbf{x}} i k'_y \Phi(\mathbf{k}')$$

Hint 1: use separate transforms (with k' , k'') for $\partial_x f$, $\partial_y \Phi$ etc.

Hint 2: $\int \exp iax \rightarrow \delta_a$

Hint 3: the final form should have a single sum over $k'_{x,y}$

Solution: Fourier Convolution

$$\begin{aligned}\mathcal{F}(\{f, \Phi\}) &= \int d\mathbf{x} e^{-i\mathbf{k}\mathbf{x}} \left[\frac{\partial f}{\partial x} \frac{\partial \Phi}{\partial y} - \frac{\partial f}{\partial y} \frac{\partial \Phi}{\partial x} \right] \\ \mathcal{F}(\{f, \Phi\}) &= \int d\mathbf{x} e^{-i\mathbf{k}\mathbf{x}} \left[\left(\sum_{\mathbf{k}'} e^{i\mathbf{k}'\mathbf{x}} i k'_y \Phi(\mathbf{k}') \right) \left(\sum_{\mathbf{k}''} e^{i\mathbf{k}''\mathbf{x}} i k''_x f(\mathbf{k}'') \right) - \dots \right] \\ \mathcal{F}(\{f, \Phi\}) &= - \int d\mathbf{x} \sum_{\mathbf{k}'} \sum_{\mathbf{k}''} e^{-i(\mathbf{k}-\mathbf{k}'-\mathbf{k}'')\mathbf{x}} \left[k'_y \Phi(\mathbf{k}') k''_x f(\mathbf{k}'') - \dots \right] \\ \mathcal{F}(\{f, \Phi\}) &= - \sum_{\mathbf{k}'} \sum_{\mathbf{k}''} \delta_{\mathbf{k}-\mathbf{k}'-\mathbf{k}''} (k'_y k''_x - k'_x k''_y) \Phi(\mathbf{k}') f(\mathbf{k}'') \\ \mathcal{F}(\{f, \Phi\}) &= \sum_{\mathbf{k}'} (k'_x k_y - k_x k'_y) \Phi(\mathbf{k}') f(\mathbf{k} - \mathbf{k}')\end{aligned}$$

Why do we care, or: what are the advantages of Fourier space?

Fluid Model for Plasmas II

In Fourier space, the model becomes

$$\frac{\partial \Phi(k_{\perp})}{\partial t} = [\delta(k_y) + k_{\perp}^2]^{-1} \left(-ik_y \Phi(k_{\perp}) + ik_y \epsilon p(k_{\perp}) - \nu k_{\perp}^2 \Phi(k_{\perp}) - \frac{1}{2} \sum_{k'_{\perp}} (k'_{\perp} \times \hat{z} \cdot k_{\perp}) [(k_{\perp} - k'_{\perp})^2 - k'_{\perp}{}^2] \Phi(k'_{\perp}) \Phi(k_{\perp} - k'_{\perp}) \right)$$

$$\frac{\partial p(k_{\perp})}{\partial t} = -ik_y (1 + \eta) \Phi(k_{\perp}) - \chi k_{\perp}^4 p(k_{\perp}) - \sum_{k'_{\perp}} (k'_{\perp} \times \hat{z} \cdot k_{\perp}) \Phi(k'_{\perp}) p(k_{\perp} - k'_{\perp})$$

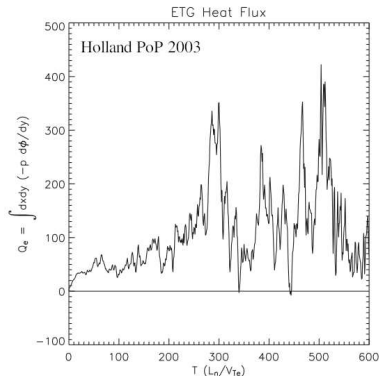
δ : adiabatic response; $\delta = \frac{1}{0}$ at $\frac{k_y}{k_y=0} > 0$ for ITG, $\delta = 1$ for ETG
(all just a trick to get more or less correct nonlinear behavior)

Later, more physics background, proper derivation; for now,
how do we use a model like this and what can we learn?

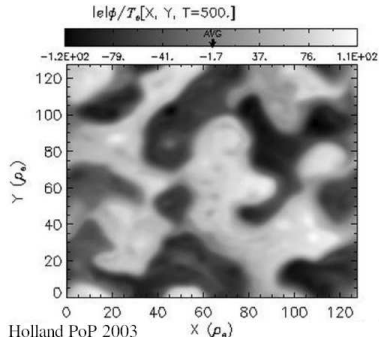
Instability and Saturation

Commonly, pick arbitrary initial condition (Φ, p) , evolve in time

Below, ETG (we'll talk more about drift-wave physics later)



- short linear phase
- linear: mostly $k_x = 0$ to access ∇T



- nonlinearity rises early
- nonlinear coupling creates $k_x > 0$ eddies

Questions & Discussion

Anything unclear so far?

Group Work: Code Prep

30 minutes group work:

- In a programming language of your choice, write a code that plots contours of the function

$$f(z) = (z + 0.2i)(z - 0.2i)(z + 0.3 - 0.5i)$$

in the complex plane ($z \in \mathbb{C}$)

- Be prepared to read data from an ASCII text file and create line and contour plots based on that data; example data:

1.33 -1.8337E+01

1.34 -1.8591E+01

1.35 811.429

1.36 40238.2