

Turbulence and Transport in Fusion Plasmas Part IV



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Wednesday Recap

Yesterday, we covered

- numerical treatment of the Horton-Holland dispersion relation
- different theory frameworks and their use
- Landau damping in kinetic theory
- what conditions have to be fulfilled so MHD and/or gyrokinetics can be used

Next: what coordinates and simulation domain should we use?

Group Work: Coordinates

45 minutes group work:

Find sources that explain

1 toroidal coordinates

2 the safety factor q as a measure of field-line pitch

and (roughly) familiarize yourself with those. Have a look at

3 all of

`www-fusion.ciemat.es/wiki/Toroidal_coordinates`

4 as much as you feel like of

`www-fusion.ciemat.es/wiki/Flux_coordinates`

5 bonus reading for those with a high pain threshold:

pages 1–3 of P. Xanthopoulos *et al.*,

Phys. Plasmas **13**, 092301 (2006)

Be prepared to present your findings.

Can you explain when/why/how field lines are (not) periodic?

Who can tell the group . . .

How are **toroidal coordinates** defined?

When are field lines periodic?
Why would we care?

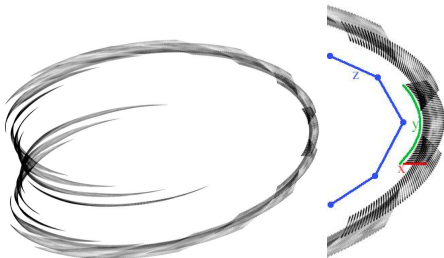
Is turbulence localized radially, toroidally, along the field?

What does all this mean for the stellarator?

Flux Tubes

Typical experiments/reactors: $k_{\perp}\rho_i \sim 0.1 - 1 \leftrightarrow n \sim 30 - 300$
Strong radial/toroidal localization \Rightarrow **flux tube** (Beer PoP 1995)

With radial domain $L_x \gtrsim \rho_{i,e} \ll R, a$, can do Taylor expansion of n, T, q , etc. profiles, e.g., $T(r \approx r_0) \approx T_0 + (r - r_0)dT/dr$
A little confusing: in flux tube, both T and dT/dr are constant!



Advantages

- cheaper (lower N_x)
- cleaner (Fourier)
- flexible (no fixed ρ^*)

Toroidal coordinates r, θ, ϕ transform to local x, y, z :

$$r = r_0 + x \quad \theta = z - \pi \quad (\text{circular flux surf.})$$

$$\phi = -\frac{q_0}{r_0}y + q_0 \left(1 + \frac{\hat{s}}{r_0}x\right) \theta \quad \hat{s} \equiv \frac{r_0}{q_0} \frac{dq}{dr}$$

Flux Tube vs. Flux Surface

Tokamak: a **single flux tube** represents entire **flux surface**

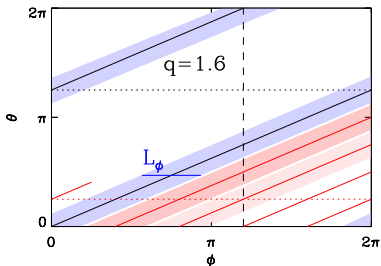
$\rho^* = \rho_i/R_0 \ll 1$ free to pick,

$\mathcal{M} = 2\pi/L_\phi \in \mathbb{N}$ arbitrary

\Rightarrow at $x = 0$, always periodic,
but not at $x \neq 0$ if $\hat{s} \neq 0$

Circular surface: parallel BC

$f(x, y, \pi) = f(x, y - 2\pi\hat{s}x, -\pi)$
(also called *twist-and-shift*)



Full system: entire flux surface (*but*: $n = 1$ means $k_y \sim k_{\parallel}$!)

\Rightarrow need to test convergence only for $k_y^{\min} \propto L_y^{-1}$ in $-\pi \leq z < \pi$

Stellarator: **flux tubes** starting at different ϕ **differ**

\Rightarrow for complete physics, need full-surface (or full-volume) code!

Ballooning Representation

Exercise: derive **parallel BC** in $k_{x,y}$ **Fourier space**

$$f(k_x, k_y, \pi) = (-1)^{\mathcal{N}} f(k_x + \mathcal{N} k_x^{\text{shift}}, k_y, -\pi)$$

$$\mathcal{N} = 2\pi \hat{s} k_y / k_x^{\text{shift}} \stackrel{\text{commonly}}{\pm 1}$$

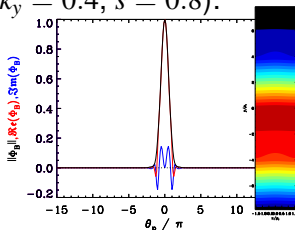
- real space: y shift
- k -space: k_x shift
- can be used to stitch together k_x

Ballooning space

Extended parallel coordinate:
ballooning angle θ_p , usually
 $= 0$ at $k_x = 0$ (Candy PoP 2004)

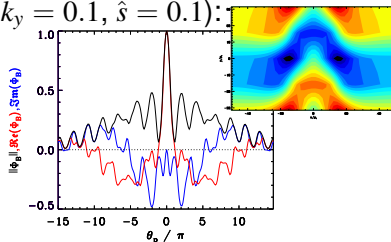
“**ballooned**” mode

$(k_y = 0.4, \hat{s} = 0.8)$:



“**slab-like**” mode

$(k_y = 0.1, \hat{s} = 0.1)$:



Group Work: Locality

20 minutes group work

- 1 Determine whether the flux tube is likely to be valid for the same machines/radii as in the Orderings group work:
Is the Taylor expansion a good approximation throughout $L_x \sim 100\rho_i$?

Note: this is a simple estimate! For real applications, more thorough studies (e.g., comparing local, global) may be needed.

Questions & Discussion

Anything unclear so far?

Group Work: Equilibria

MHD: **magnetic equilibria** — no MHD instability, fluxes from neoclassical collisions or microturbulence

1.5 hours group work:

- 1 Download J.W. Haverkort's write-up on equilibria:

<http://homepage.tudelft.nl/20x40/documents/Equilibria.pdf>

- 2 Work through

- a Sec. 1.1
- b Sec. 1.2
- c Appendix A

for good understanding of the **Grad-Shafranov** equation

- 3 get R.L. Miller *et al.*, Phys. Plasmas **5**, 973 (1998), read Secs. 1–3, make notes about what is unclear, distill key findings

Then reconvene in the plenum to discuss everyone's findings

The Zoo of Instabilities

Microinstabilities: drift waves driven by pressure gradients

Note: all of them have **critical gradients**

| | ITG | ETG | TEM | KBM | MT |
|----------------------------|-------------------|----------------------|------------------------|-----------------------------|-------------------|
| drive | ∇T_i | ∇T_e | $\nabla T_e, \nabla n$ | $\nabla T_{i,e} + \nabla n$ | ∇T_e |
| ρ_j scale | i | e | i | i | i |
| ω sign ¹ | + | - | (-) | + | - |
| $\beta \nearrow$ | $\gamma \searrow$ | $\gamma \rightarrow$ | $\gamma \rightarrow$ | $\gamma \nearrow$ | $\gamma \nearrow$ |
| Φ vs. A_{\parallel} | \gg | \gg | \gg | \gg | \lesssim |
| parity ² | +(-) | +(-) | +(-) | +(-) | - |
| slab branch ³ | ✓ | ✓ | × | × | ✓ |
| zonal flows ⁴ | ✓ | (✓) | ×, ✓ | × | (✓) |

Cause of turbulence & transport in fusion experiments

Each of the above is relevant to all of tokamak, stellarator, RFP

¹+(-) drifts in ion(electron)-direction; some use opposite nomenclature!

²+(-) means even (odd) $\Phi(z)$ and odd (even) A_{\parallel}

³slab mode: parallel motion important, $|\theta_p| \gg \pi$

⁴nonlinear saturation mechanism, discussed later

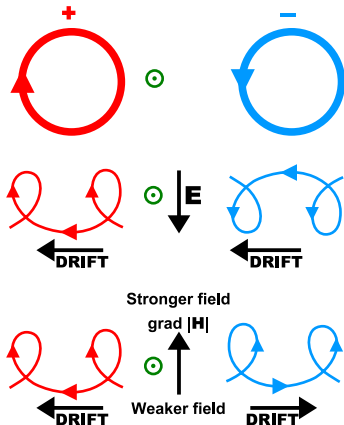
Plasma Drifts

To understand **drift-wave instabilities**, recall drifts:
gradients cause perpendicular **drifts at constant v**

- electric field: “E cross B”,
 $\mathbf{v}_E = c(\mathbf{E} \times \mathbf{B})/B^2$
- inhomogeneous guide field:
“grad B” & curvature,
 $\mathbf{v}_{\nabla B} = v_{\perp}^2 (\mathbf{b} \times \nabla B)/(2B\Omega_j)$
 $\mathbf{v}_c = v_{\parallel}^2 (\nabla \times \mathbf{b})_{\perp}/\Omega_j$

Key properties

\mathbf{v}_E in same direction for i, e
 $\mathbf{v}_{\nabla B, c}$ opposite for i, e

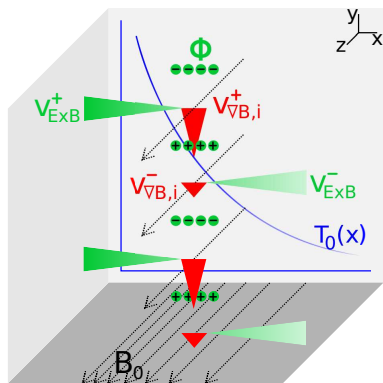


(adapted from: Wikipedia)

ITG & ETG Modes

ITG: Coppi PoF 1967 (linear), Dimits PoP 2000 (nonlinear)

ETG: Liu PRL 1971 (linear), Jenko PoP 2000 (nonlinear)



Toroidal ITG mode:

$$\delta\Phi \rightarrow v_E \rightarrow \delta T_i \rightarrow v_{\nabla B} \rightarrow \delta\Phi$$

- slab: k_{\parallel} instead of ∇B_0
- ETG linearly isomorphic
 $\Rightarrow \gamma_{\text{ITG}}/v_{\text{ti}} = \gamma_{\text{ETG}}/v_{\text{te}}$
- also called η_i (η_e) mode:
 ∇n can stabilize
- nonlinear toroidal ITG:
zonal flows (ETG, slab
ITG: somewhat less)

Characteristic scales

$$k_y \rho_{s,e} \sim 0.1 - 1, k_x \sim 0 - k_y, \gamma \sim 0.1 - 1 v_{\text{ti,te}}/L_{\text{Ti,e}}, \omega \sim \pm\gamma$$

Trapped-Electron Modes

Linear: Coppi PRL 1974, nonlinear: Ernst PoP 2004

∇T vs. ∇n drive: Ernst PoP 2009

∇n -driven TEM:

$$\delta\Phi \rightarrow v_E \rightarrow \delta n \rightarrow v_{\nabla B}^{e,i} \rightarrow \delta\Phi$$

- no slab equivalent
- $\nu_{ei} \gtrsim \gamma$: “dissipative” DTEM
- $\nu_{ei} < \gamma$: “collisionless” CTEM
- “ion” iTEM: Plunk JPP 2017
- “ubiquitous” UTEM ($\omega > 0$):
Coppi PoFB 1990

Electrons trapped on outboard ($\nabla B_0 \parallel \nabla n, T$)

∇T_e -driven TEM
works correspondingly

- ∇T TEM and ETG driven by ∇T_e , can be indistinguishable
- trapped-ion TIM of limited relevance

Characteristic scales

$$k_y \rho_s \sim 0.2 - 2, k_x \sim 0 - k_y, \gamma \sim 0.1 - 1 c_s / L_{n,Te}, \omega \sim -\gamma$$

Kinetic Ballooning Modes

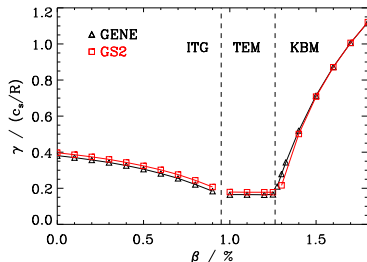
Linear: Tsai PoFB 1993, Hirose PRL 1994

Nonlinear: Pueschel PoP 2008 & 2010, Ishizawa NF 2013

KBM (also: “Alfvénic” AITG):

kinetic version of MHD ideal ballooning at high $n > 10$

Destabilized at high β



Driven by total gradient

$$\nabla p \sim \nabla n + \nabla T_i + \nabla T_e$$

$$\beta_{\text{crit}}^{\text{KBM}}(k_y \rightarrow 0) \rightarrow \beta_{\text{crit}}^{\text{MHD}} \text{ at}$$

$$\alpha_{\text{MHD}} = \beta q_0^2 R_0 (dp/dr)/p \approx 0.6\hat{s}$$

McKinney JPP 2021: saturation requires $\beta < \beta_{\text{crit}}^{\text{KBM}}(k_y^{\text{min}})$

Note: $\beta/\beta_{\text{crit}}$ common figure of merit for electromagnetic effects

Characteristic scales

$$k_y \rho_s \sim 0 - 0.5, k_x \sim 0, \gamma \sim 0.1 - 1 c_s/L_p, \omega \gg \gamma \text{ near } \beta_{\text{crit}}$$

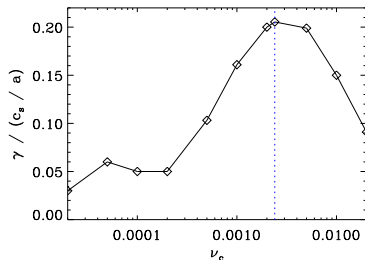
Microtearing Modes

Linear: Hazeltine PoF 1975, Drake PoF 1977

Nonlinear: Doerk PRL 2011, Guttentfelder PRL 2011

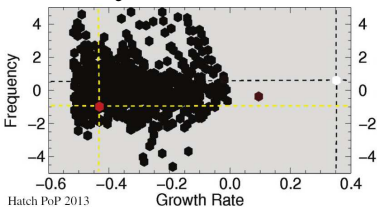
Global tearing: driven by ∇j , while **MT:** driven by β , ∇T_e

Energy access via ν_c
or curvature (collisionless)



Nonlinearly, pure Q_e^{em}
(no $Q_{e,i}^{\text{es}}$ or particle flux)

Hatch PRL 2012 & PoP 2013:
Subdominant MT responsible for Q_e^{em} in ITG turbulence



\Rightarrow first-ever example of
important stable mode

Characteristic scales

$k_y \rho_s \sim 0.01 - 0.5$ (slab), $0.1 - 1$ (tor'l), $\gamma \sim 0.1 - 1 c_s / L_{Te}$, $\omega < 0$

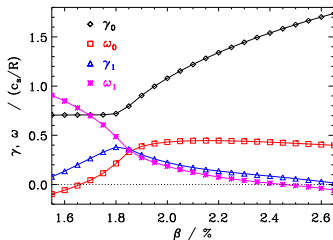
Hybrid Modes

Already mentioned: ∇T TEM & ETG can join forces

Kammerer PoP 2008, Pueschel PoP 2008:

hybrid modes combining properties of **two instabilities**

E.g., can continuously
transform KBM into TEM



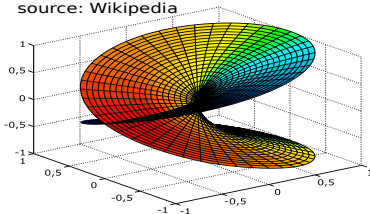
Consequences:

- turbulence regime boundaries can have odd behavior
- **subcritical** linear excitation (e.g., KBM below β_{crit})

Mathematically, related to
exceptional points

Riemann surface

source: Wikipedia



Walking circle in parameter
space can give different mode

Group Work: Characteristic Scales

1 hour group work:

- 1 for same machines as earlier, calculate characteristic
 - a diamagnetic frequency
 - b parallel transit frequency
 - c wavelengths corresponding to $k_y \rho_s = 0.3$, $k_y \rho_e = 0.3$
 - d β and ballooning threshold $\beta_{\text{crit}}^{\text{MHD}}$

in SI or cgs units (can look up diamagnetic frequency in A.J. Brizard, Rev. Mod. Phys. **79**, 421 (2007))

Who can say what their physical relevance is?

- 2 where feasible, estimate importance of instabilities for the above machines/radii: ITG, ETG, TEM, KBM

Questions & Discussion

Anything unclear that we talked about?

Any feedback for the instructor?

Anyone still awake?