

INTRODUCTION TO SYMPLECTIC GEOMETRY

NOTES FROM SUMMER SEMESTER 2025

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These notes are on the course V5D6 - Introduction to Symplectic Geometry¹, lectured jointly by Nathaniel Bottman and Laurent Côté. They are entirely my own, as are any potential mistakes in what follows. Proceed with caution. These notes are available on [GitHub](#) and [Overleaf](#).

This course covers the basic topics in symplectic geometry. We will not discuss Floer theory or J -holomorphic curves. The main reference for this course is [MS17] by McDuff and Salamon.

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¹https://www.math.uni-bonn.de/~lcote/V5D6_2025.html

1 Introduction to symplectic topology

LECTURE 1 (10.04.25).

In this lecture we covered linear symplectic geometry and the cotangent bundle as the first example of a symplectic manifold. References for this lecture are sections 2.1-2.3 and 3.1 of Mc-Duff and Salamon's *Introduction to Symplectic Topology*.

Definition. A *symplectic manifold* is a pair (M, ω) of an even-dimensional smooth manifold without boundary M , and a closed non-degenerate 2-form ω on M .

An equivalent condition to closedness and non-degeneracy would be that the highest wedge power of ω is a volume form. So in particular it follows that M is orientable.

Linear symplectic geometry

We now turn our attention to the simpler linear setting of symplectic geometry.

Definition. A symplectic vector space is a pair (V, ω) of a finite-dimensional vector space V over \mathbb{R} , and a non-degenerate skew-symmetric bilinear form ω on V .

Proposition 1.1. V is even-dimensional.

Proof. This follows immediately from the existence of a non-degenerate skew-symmetric bilinear form. \square

Example. Let $V = \mathbb{R}^{2n}$ and ω_0 be given by

$$\omega_0(v, w) = v^\top \begin{bmatrix} 0 & \text{Id}_{\mathbb{R}^n} \\ -\text{Id}_{\mathbb{R}^n} & 0 \end{bmatrix} w.$$

This shall be the canonical symplectic form on Euclidean space, and we call the specified matrix $-J_0$. If we denote by $x_1, \dots, x_n, y_1, \dots, y_n$ the standard basis on \mathbb{R}^{2n} , then we have

$$\omega_0(x_i, x_j) = \omega_0(y_i, y_j) = 0$$

for all i, j and

$$\omega_0(x_i, y_j) = \delta_{ij}.$$

Definition. A *linear symplectomorphism* is a linear isomorphism $\psi : (V, \omega) \rightarrow (V', \omega')$ such that $\psi^* \omega' = \omega$. Explicitly this means that

$$\omega(v, w) = \omega'(\psi v, \psi w).$$

Remark. We may also consider "symplectic linear maps" in general, but because the symplectic form is non-degenerate the pull-back condition requires these to be injective. Thus a symplectic linear map is just a symplectomorphism onto its image.

A linear symplectomorphism of $(\mathbb{R}^{2n}, \omega_0)$ is then represented by a matrix $A \in \text{GL}_{2n}(\mathbb{R})$ such that $A^\top(J_0)A = J_0$.

Definition. A matrix $A \in \text{GL}_{2n}(\mathbb{R})$ satisfying this equation is called symplectic. Denote the group of symplectic matrices by $\text{Sp}(2n)$ or $\text{Sp}(2n, \mathbb{R})$. A computation shows that $\det A = \pm 1$ for symplectic A . In fact, the Pfaffian may be used to show that $\det A = 1$.

compute this

Let (V, ω) be a symplectic vector space and $W \subset V$ a linear subspace.

Definition. The *symplectic complement* of W is

$$W^\omega = \{v \in V \mid \omega(v, w) = 0 \text{ for all } w \in W\}.$$

We then say that $W \subset V$ is

- *isotropic* if $W \subset W^\omega$,
- *coisotropic* if $W^\omega \subset W$,
- *symplectic* if $W \cap W^\omega = \emptyset$, and
- *Lagrangian* if $W = W^\omega$.

If W is symplectic, then $(W, \omega|_W)$ is a symplectic vector space.

Lemma 1.2. We have the following two results,

$$\dim W + \dim W^\omega = \dim V, \quad \text{and} \quad (W^\omega)^\omega = W.$$

Proof. Consider the map

$$i_\omega : V \rightarrow V^*, \quad i_\omega(v)(w) = \omega(v, w).$$

Then this is an isomorphism, as ω is non-degenerate. Now see that

$$i_\omega(W^\omega) = \{\varphi \in V^* \mid \varphi(w) = 0 \text{ for all } w \in W\} = W^\perp.$$

Because i_ω is an isomorphism, this means that

$$\dim W + \dim W^\omega = \dim W + \dim W^\perp = \dim V.$$

See that by definition $W \subset (W^\omega)^\omega$ and as they have the same dimension they must then be equal. \square

Corollary. The following 3 results follow immediately,

- W is symplectic $\iff W^\omega$ is symplectic,
- W is isotropic $\iff W^\omega$ is coisotropic, and
- W is Lagrangian $\implies \dim W = \frac{1}{2} \dim V$.

Theorem 1.3. Any symplectic vector space (V, ω) of dimension $2n$ is isomorphic $(\mathbb{R}^{2n}, \omega_0)$.

Proof. It suffices to construct a *symplectic basis*, i.e. a basis $u_1, \dots, u_n, v_1, \dots, v_n$ of V such that prove this

$$\omega(u_i, u_j) = \omega(v_i, v_j) = 0,$$

for all i, j and

$$\omega(u_i, v_j) = \delta_{ij}.$$

To construct such a basis first choose any non-zero $u_1 \in V$, and because ω is non-degenerate, we can find some $v_1 \in V$ such that $\omega(u_1, v_1) = 1$. See that u_1, v_1 span a symplectic subspace. Repeating this process on the symplectic complement of their span then gets us a symplectic basis. \square

Symplectic manifolds

Definition. A symplectomorphism $(M, \omega_M) \rightarrow (N, \omega_N)$ is a diffeomorphism φ such that $\varphi^* \omega_N = \omega_M$. We denote the symplectomorphisms of (M, ω_M) by $\text{Symp}(M, \omega_M)$. This is a lie group. We say that a vector field $X \in \mathfrak{X}(M)$ is symplectic if $i_X \omega$ is closed. We denote these by $\chi(M, \omega)$.

Remark. By i_X we denote the interior derivative or contraction, i.e. for a k -form μ ,

$$i_X \mu(X_1, \dots, X_{k-1}) = \mu(X, X_1, \dots, X_{k-1}).$$

Sometimes this is also denoted by $X \lrcorner \mu$.

Proposition 1.4. Integrating symplectic vector fields results in a symplectomorphism. Specifically if we have a smooth families (φ_t) of smooth maps and (X_t) of smooth vector fields such that

$$\varphi_0 = \text{id} \quad \text{and} \quad \frac{d}{dt} \varphi_t = X_t \circ \varphi_t.$$

Then φ_t is symplectic for all t if and only if X_t is symplectic for all t .

Proof. prove this

A bit on the cotangent bundle

We consider the cotangent bundle T^*L of a closed manifold L .

Definition (Canonical 1-form). We denote by $\pi : T^*L \rightarrow L$ the projection of the fibers to the base space. Then we can take the pullback along the projection as follows. For $p \in L, \xi \in T_p^*(L)$ and a tangent vector $v \in T_{(p,\xi)}T^*L$, denote

$$\lambda_{\text{can}}(p, \xi)(v) := (\xi \circ d_p \pi)(v).$$

We call $\lambda_{\text{can}} \in \Omega^1(T^*L)$ the *canonical 1-form*.

Proposition 1.5. Let L be a smooth manifold, and define $\omega_{\text{can}} := -d\lambda_{\text{can}} \in \Omega^2(T^*L)$ the canonical symplectic form. Then $(T^*L, \omega_{\text{can}})$ is a symplectic manifold.

Proof. It is clear that ω_{can} is closed because it is exact. Verifying that ω_{can} is non-degenerate is left to the reader. \square

Remark. Let $x : U \rightarrow \mathbb{R}^n$ be a chart of L with coordinates x_1, \dots, x_n . At a point $p \in L$ see that the differentials $(dx_i)_{i=1}^n$ form a basis of the cotangent space T_p^*L and denoting $y_i = dx_i$ we get local coordinates $(x_1, \dots, x_n, y_1, \dots, y_n)$ for T^*U . Then one can define the canonical 1-form in local coordinates to be

$$\lambda_{\text{can}} = y \, dx.$$

Then $\omega_{\text{can}} = dx \wedge dy$ in local coordinates.

Proposition 1.6. Let $\sigma \in \Omega^1(L)$ be a 1-form. Consider it as a map $\sigma : L \rightarrow T^*L$. Then λ_{can} is described by the universal property that the pullback satisfies

$$\sigma^* \lambda_{\text{can}} = \sigma,$$

for any σ .

Proof. This is done in local coordinates on some chart $x : U \rightarrow \mathbb{R}^n$. We write

$$\sigma = \sum_{i=1}^n \sigma_i(x_1, \dots, x_n) dx_i.$$

Then this is a map $(x_1, \dots, x_n) \mapsto (x_1, \dots, x_n, \sigma_1, \dots, \sigma_n)$. Thus

$$\sigma^* \lambda_{\text{can}} = \sigma^* \left(\sum_{i=1}^n y_i dx_i \right) = \sum_{i=1}^n \sigma_i(x_1, \dots, x_n) dx_i = \sigma.$$

To verify that this uniquely determines λ it suffices to note that if $\sigma^* \mu = 0$ for all σ , then $\mu = 0$. \square

LECTURE 2 (17.04.25).

The references for this lecture are sections 3.2 and 3.3 of McDuff-Salamon.

Recap

write this

Theorem (Reynolds transport theorem). Let $\varphi_t : M \rightarrow M$ be a smooth family of diffeomorphism generated by vector fields X_t , and η_t a family of symplectic forms. Then:

$$\frac{d}{dt}(\varphi_t^* \eta_t) = \varphi_t^* \left(\frac{d\eta_t}{dt} + L_{X_t} \eta_t \right)$$

Moser's trick and its consequences

Theorem 2.1 (Moser's trick). Let M be a closed smooth even-dimensional manifold. Let $(\omega_t)_{t \in [0,1]}$ a family of symplectic forms on M and $(\sigma_t)_{t \in [0,1]}$ a family of smooth 1-forms such that

$$\frac{d}{dt} \omega_t = d\sigma_t.$$

Then there exists a family $(\psi_t)_{t \in [0,1]}$ of diffeomorphisms on M satisfying

$$\psi_t^* \omega_t = \omega_0.$$

To prove this we require the following lemma.

Lemma 2.2. Let M be a smooth $2n$ -manifold, $Q \subset M$ a closed submanifold, and $\omega_1, \omega_0 \in \Omega^2(M)$ closed forms such that for all points $q \in Q$, $\omega_{0,q}$ and $\omega_{1,q}$ are non-degenerate and agree on all of $T_q M$.

Proof.

□

prove this

We are now equipped to prove Moser's trick.

Proof.

□

prove this

Consequences of Moser's trick

Theorem (Darboux). Let (M, ω) be a symplectic $2n$ -manifold. Then for all points $p \in M$, there exists some neighborhood $U \ni p$ such that $(U, \omega|_U)$ is symplectomorphic an open $V \subset \mathbb{R}^{2n}$ equipped with the canonical symplectic form.

short bit on
Darboux's
theorem

Proof.

□

prove this

Theorem (Moser's stability theorem). Let M be a closed manifold, and $(\omega_t)_{t \in [0,1]}$ a smooth family of symplectic forms such that $[\omega_t] \in H_{\text{dR}}^2(M, \mathbb{R})$ is fixed independently of t . Then there exists a family of diffeomorphism $(\psi_t)_{t \in [0,1]}$ such that

$$\psi_t^* \omega_t = \omega_0.$$

Proof.

□

prove this

Definition. We say that a smooth submanifold $Q \subset (M, \omega)$ is symplectic, (co)isotropic, or Lagrangian if for all $q \in Q$ the tangent space $T_q Q \subset T_q M$ satisfies the corresponding property with respect to ω .

Recall specifically that $T_q Q \subset T_q M$ is Lagrangian if $(T_q Q)^\omega = T_q Q$ or equivalently if $\dim Q = \dim M/2$ and $\omega_q|_{T_q Q} = 0$.

Example $(T^*L, \omega_{\text{can}})$. As previously discussed, for a manifold L the cotangent bundle and canonical form give a symplectic manifold $(T^*L, \omega_{\text{can}})$. Now the zero section $L \subset T^*L$ and the cotangent space $T_q^*L \subset T_q T^*L$ are Lagrangian submanifolds.

tikz this

In fact we will next show that any closed Lagrangian submanifold is locally like T^*L .

Theorem (Weinstein neighborhood theorem). Let $L^n \subset (M^{2n}, \omega)$ be a Lagrangian submanifold. Then there exist open sets

$$L \subset U \subset M \quad \text{and} \quad L \subset V \subset T^*L,$$

and a symplectomorphism $\varphi : (V, \omega_{\text{can}}) \rightarrow (U, \omega|_U)$ such that $\varphi|_L = \text{id}_L$.

Proof.

□

prove this

2 Dynamics and integrable systems

LECTURE 3 (24.04.25).

In this lecture we covered Hamiltonian mechanics and vector fields, as well as the Poisson bracket. The goal of this lecture was to introduce some the historical and physical motivation for symplectic geometry. In the notes for this lecture we omit some of the physical calculations for brevity as they will not be relevant going further. We attempt to retain the physical and historical motivation behind symplectic geometry.

References for this lecture are sections 1.1 and 3.1 of [MS17]. (*Note: From this point onward the references are specifically for the third edition of the book, I am unsure which edition the previous ones refer to.*)

Hamiltonian mechanics

Consider a physical system, the configurations of which may be described by a point $x \in \mathbb{R}^n$. For example, the positions of three celestial objects may be described by a point in \mathbb{R}^9 . We are interested in trajectories in this space, denoted by $t \mapsto x(t)$. Now suppose that we have a function

$$L : \mathbb{R} \times \mathbb{R}^n \times \mathbb{R}^n \rightarrow \mathbb{R}, \quad L = L(t, x, v)$$

such that trajectories are the critical points of the "action functional":

$$I(x) := \int_{t_0}^{t_1} L(t, x(t), \dot{x}(t)) dt.$$

In fact, such a function models the difference between the kinetic and potential energies of the system. (*Note: In the tradition of the physicists we will abuse notation and denote $L(t, x(t), v(t))$ by $L(t, x, v)$.*)

Lemma 3.1. A minimal path $x : [t_0, t_1] \rightarrow \mathbb{R}^n$ satisfies the Euler-Lagrange equations:

$$\frac{d}{dt} \frac{\partial L}{\partial v} = \frac{\partial L}{\partial x},$$

where by $\partial L / \partial v$ we refer to $(\partial L / \partial v_1, \dots, \partial L / \partial v_n)$, and similarly for x .

Proof. See Lemma 1.1.1 in [MS17]. □

This is the Lagrangian formulation of mechanics. To transform this into the Hamiltonian formulation of mechanics we apply the Legendre transformation to our coordinates. We replace v_i by y_i where

$$y_i = \frac{\partial L}{\partial v_i}(x, v).$$

This is a valid coordinate transformation as long as the Legendre condition,

$$\det \left(\frac{\partial^2 L}{\partial v_i \partial v_j} \right)_{ij} \neq 0,$$

holds. For clarity of notation we denote $G_i(t, x, y) = v_i$, and avoid reference to v so that we may take y as a given.

Definition (Hamiltonian). We define the Hamiltonian to be

$$H(t, x, y) = \sum_{i=1}^n y_i \cdot G_i(t, x, y) - L(t, x, G(t, x, y)).$$

This represents the total energy within the system.

We omit the computation but one may see that

$$\frac{\partial H}{\partial x_i} = -\dot{y}_i \quad \text{and} \quad \frac{\partial H}{\partial y_i} = \dot{x}_i.$$

These are called Hamilton's equations. Note that while the Lagrangian formalism is concerned with vectors in space and the tangent bundle, in Hamilton's reformulation we focus on the cotangent bundle.

Write $z(t) = (x(t), y(t))$. Then we may rewrite Hamilton's equations as

$$\dot{z} = -J_0 \nabla H(z), \quad J_0 = \begin{bmatrix} 0 & -\text{id}_n \\ \text{id}_n & 0 \end{bmatrix}.$$

The Poisson bracket

In what follows we assume that H is independent of t , so $H(t, x(t), y(t)) = H(x(t), y(t))$.

finish this
lecture

LECTURE 4 (08.05.25).

The references for this lecture are sections 1.2-1.6 of [Eva23]. In this lecture we went over integrable systems and Hamiltonian torus actions. We will denote by $M = (M^{2n}, \omega)$ a symplectic manifold.

Integrable Hamiltonian system

Definition (Poisson commutativity). We say that smooth functions $F, G \in C^\infty(M)$ *Poisson commute* if $\{F, G\} = 0$, i.e. if $\omega(X_F, X_G)$ vanishes everywhere. We will refer to this simply as commuting whenever the meaning is clear from the context.

Lemma 4.1 (Invariance of commuting functions). If $F, G \in C^\infty(M)$ commute then F is invariant under the flow of X_G .

Proof. We consider the behavior of f on integral curves of X_G .

more detail

$$\frac{d}{dt} F(\varphi_{X_G}^t(x)) = dF(X_{G, \varphi_G^t(x)}) = -\omega(X_{F, x}, X_{G, x}) = -\{F, G\}(x) = 0.$$

□

Lemma 4.2 (Commutativity of flows of Hamiltonian vector fields). The flows φ_F^t and φ_G^t commute if and only if $\{F, G\}$ is locally constant.

Proof. Recall that the flows of two vector fields commute if and only if their Lie bracket vanishes identically. We further recall that $[X_F, X_G] = X_{\{F, G\}}$, and by the non-degeneracy of ω the vector field $X_{\{F, G\}}$ vanishes if and only if $\omega(X_{\{F, G\}}, -) = d\{F, G\}$ vanishes. But $d\{F, G\}$ vanishes identically exactly when $\{F, G\}$ is locally constant. □

Definition (Induced Hamiltonian group action). Given a map

$$\mathbf{H} = (H_1, \dots, H_k) : M \rightarrow \mathbb{R}^k$$

such that $\{H_i, H_j\} = 0$ for all $1 \leq i, j \leq k$ we can define an induced Hamiltonian group action $\Psi : \mathbb{R}^k \curvearrowright M$ by:

$$\Psi : \mathbb{R}^k \times M \rightarrow M, \quad \Psi^t(x) = \varphi_{H_1}^{t_1} \circ \dots \circ \varphi_{H_k}^{t_k}(x).$$

Definition (Integrable system). We say that $\mathbf{H} = (H_1, \dots, H_n)$ is a complete commuting Hamiltonian system if $\{H_i, H_j\} = 0$. If \mathfrak{H} is proper and has a dense set of regular values then we say that is an *integrable system*. (A map is proper if the preimages of compact sets are compact.)

give an ex-
ample from
the book
and Lie
group ver-
sion

Lemma 4.3 (Orbits of Hamiltonian actions). If $\mathbf{H} : M \rightarrow \mathbb{R}^k$ generates a Hamiltonian \mathbb{R}^k -action, then the orbits of this action are isotropic submanifolds. Furthermore if M contains a regular point of \mathbf{H} then $k \leq n$.

Proof. _____ □ prove this

Corollary. If $k = n$ then the orbits of regular points are Lagrangian.

Proof. _____ □ prove this

3 Lagrangian submanifolds and toric symplectic manifolds

LECTURE 5 (15.05.25).

No reference have been given for this lecture. In this lecture we covered toric symplectic manifolds.

Definition (Symplectic toric manifold). We say $M = (M^{2n}, \omega)$ is *toric* if it admits an integrable system such that the \mathbb{R}^n action factors through $\mathbb{R}^n/\mathbb{Z}^n$.

So there exists a map $\mathbb{H}(= (H_1, \dots, H_n)) : M \rightarrow \mathbb{R}^n$ which reduces to $\mathbb{R}^n/\mathbb{Z}^n$.

Example. Let $M = \mathbb{C}^n \cong \mathbb{R}^{2n}$ and \mathbb{T}^n act on \mathbb{C}^n by rotation. So

$$(e^{i\theta_1}, \dots, e^{i\theta_n})(z_1, \dots, z_n) \mapsto (e^{i\theta_1}z_1, \dots, e^{i\theta_n}z_n).$$

Then this action is induced by the moment map

$$\mu = \mathbb{H} = (H_1, \dots, H_n), \quad H_i(x_1, \dots, x_n, y_1, \dots, y_n) = \frac{1}{2}(x_i^2 + y_i^2).$$

The image of μ is just $(\mathbb{R}_{\geq 0})^n$ and the preimage of a point in the open locus is just \mathbb{T}^n .

We now move our definitions to abstract torii.

Definition (Updated symplectic toric manifold). A (k) -torus is a k -dimensional connected compact Lie group. Such a group is non-canonically isomorphic to $(S^1)^k$.

Let (M^{2n}, ω) be a symplectic manifold, and T be a k -torus for some $1 \leq k \leq n$. Then an action $T \curvearrowright M$ is *Hamiltonian* if there exists a moment map $\mu : M \rightarrow \mathfrak{t}^\vee$ inducing $T \curvearrowright M$. Here \mathfrak{t}^\vee denotes the dual of the Lie algebra \mathfrak{t} . Then M is toric.

Given a Hamiltonian action $\Psi : T \times M \rightarrow M$ and $v \in \mathfrak{t}$, we define the vector field $v^\# \in \mathcal{X}(M)$ to be

$$v_p^\# = \left. \frac{d}{dt} \right|_{t=0} \Psi(e^{tv} p).$$

Lemma 5.1. Given $\xi \in \mathfrak{t}$ the vector field $v^\#$ is symplectically dual to $-\mathrm{d}\mu(v)$.

Proof. We work on a basis and identify $T = \mathbb{R}^k / \mathbb{Z}^k$ with $\mu = (H_1, \dots, H_k)$. We want to show that

$$w(v^\#, -) = -\mathrm{d}\mu(v).$$

It suffices to check this on the basis of ∂_{x_i} , and we get

$$w(\partial_{x_i}^\#, -) = -\mathrm{d}H_i$$

which is how $\partial_{x_i}^\#$ is defined. \square

Lemma 5.2. Let $M = (M^{2n}, \omega)$ be symplectic and equipped with a Hamiltonian torus action $T \curvearrowright M$, with moment map μ . Given a morphism of $\varphi : T' \rightarrow T$ torii the moment map of the induced action $\mu' : M \rightarrow (\mathfrak{t}')^\vee$ satisfies

$$\varphi\mu' = \varphi^* \circ \mu.$$

Proof. For $v \in \mathfrak{t}'$ we take the pushforward $(\mathrm{d}\varphi(v))^\#$. Now $(\mathrm{d}\varphi(v))^\#$ is dual to $-\mathrm{d}\mu(\mathrm{d}\varphi(v)) = -\mathrm{d}(\varphi^*\mu)(v)$. \square

Symplectic reduction

Lemma 5.3 (Linear symplectic reduction). Let (V, ω) be a symplectic vector space and $F \hookrightarrow V$ be a linear subspace. We denote by $\pi : F \rightarrow F/(F \cap F^\omega) = \bar{F}$ the projection map. Then there exists a linear symplectic form $\bar{\omega}$ on the quotient \bar{F} satisfying the formula $\pi^*\bar{\omega} = \omega|_F$.

Proof. First we check that this defines a form. See that for $v \in \ker \pi = F \cap F^\omega$ and $u \in F$ we have $\omega(v, u) = 0$, so taking

$$\bar{\omega}(-, -) = \omega(\pi^{-1}(-), (-))$$

does not depend on the choice of preimage. A form defined this way is clearly also anti-symmetric.

To show that $\bar{\omega}$ is non-degenerate consider a non-zero vector $v \in \bar{F}$ and consider a vector $v' \in \pi^{-1}(v)$. As v is non-zero, $v' \in F \setminus F^\omega$ so there exists some $u \in F$ such that $\bar{\omega}(v, \pi(u)) = \omega(v', u) \neq 0$. \square

Lemma 5.4 (Symplectic reduction for the torus). Let $T \curvearrowright (M^{2n}, \omega)$ be a Hamiltonian action with moment map $\mu : M \rightarrow \mathfrak{t}^\vee$. If $c \in \mathfrak{t}^\vee$ is a regular value then

$$M//_c T := \mu^{-1}(c)/T \quad \text{and} \quad \pi^*\bar{\omega} = \omega|_{\mu^{-1}(c)}$$

define a symplectic manifold and form.

maybe something on why these are the same?

something's not right here

bit more detail

Proof. We again work with a basis and set $T = \mathbb{R}^k / \mathbb{Z}^k$. We identify $\mathfrak{t}^\vee = \mathbb{R}^k$. Because c is regular $\mu^{-1}(c)$ is also a manifold. Given $\xi \in \mu^{-1}(c)$ we denote \mathcal{O}_ξ by its orbit.

why?

By linear symplectic reduction it suffices to show that

$$T_x \mathcal{O}_\xi = T_x \mu^{-1}(c) \cap (T_x \mu^{-1}(c))^\omega.$$

□

finish this proof

Example. Consider $S^1 \curvearrowright \mathbb{C}^n$ by $(e^{i\theta}, z) \mapsto (e^{i\theta} z_1, \dots, e^{i\theta} z_n)$. This is generated by $\mu = \frac{1}{2} \sum_i (x_i^2 + y_i^2)$. Consider a point z such that each $z_i \neq 0$. Then z is regular and

$$\mathbb{R}^{2n} //_z S^1 := \mu^{-1}(z) / S^1 \cong \mathbb{CP}^{n-1}.$$

Constructing new symplectic manifolds

Consider an exact sequence of torii

$$0 \rightarrow D^k \rightarrow E = \mathbb{T}^n \rightarrow F^{n-k} \rightarrow 0,$$

and let $\mathbb{T}^n = E \curvearrowright \mathbb{C}^n$ be the standard action of rotating every coordinate individually. This is generated by $\mu_i = \frac{1}{2}(x_i^2 + y_i^2)$. We denote $\mu_D : \mathbb{C}^n \rightarrow \mathfrak{d}^\vee \cong \mathbb{R}^k$ be the induced map.

LECTURE 6

4 Symplectic capacities and embeddings

LECTURE 7

LECTURE 8

LECTURE 9

LECTURE 10

5 Hamiltonian dynamics

LECTURE 11

References

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