

Calculus Cheat Sheet

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Based on notes by Dr A.G. Walton

1 Basics

Formulae relating permutation and Kronecker delta:

$$\epsilon_{ijk}\epsilon_{klm} = \delta_{il}\delta_{jm} - \delta_{im}\delta_{jl}$$

$$\epsilon_{ijk}\epsilon_{ilm} = \delta_{jl}\delta_{km} - \delta_{jm}\delta_{kl}$$

Dot product

$$\mathbf{a} \cdot \mathbf{b} = a_i b_i$$

Cross product

$$[\mathbf{a} \times \mathbf{b}]_i = \epsilon_{ijk} a_j b_k$$

Triple scalar product

$$\mathbf{a} \cdot (\mathbf{b} \times \mathbf{c}) = \epsilon_{ijk} a_i b_j c_k$$

$$\mathbf{a} \cdot (\mathbf{b} \times \mathbf{c}) = 0 \Leftrightarrow a, b, c \text{ are coplanar.}$$

$$\mathbf{a} \cdot (\mathbf{b} \times \mathbf{c}) = (\mathbf{a} \times \mathbf{b}) \cdot \mathbf{c}$$

Triple vector product

$$\mathbf{a} \times (\mathbf{b} \times \mathbf{c}) = (\mathbf{a} \cdot \mathbf{c})\mathbf{b} - (\mathbf{a} \cdot \mathbf{b})\mathbf{c}$$

2 Gradient, divergence and curl

Gradient & directional derivative

$$\nabla \phi = \frac{\partial \phi}{\partial x} \hat{\mathbf{i}} + \frac{\partial \phi}{\partial y} \hat{\mathbf{j}} + \frac{\partial \phi}{\partial z} \hat{\mathbf{k}}, \frac{\partial \phi}{\partial \mathbf{s}} = \hat{\mathbf{s}} \cdot \nabla \phi$$

note that $\hat{\mathbf{s}}$ here must be unit vector.

Formal definition of gradient, S is a surface, τ is the volume of region enclosed by S :

$$\nabla \phi = \lim_{\tau \rightarrow 0} \frac{1}{\tau} \int_S \hat{\mathbf{n}} \phi dS$$

Laplace

$$\nabla^2 \phi := \nabla \cdot \nabla \phi = \frac{\partial^2}{\partial x_i^2} \phi$$

$$[\nabla^2 \mathbf{A}]_i = \frac{\partial^2}{\partial x_j^2} \mathbf{A}_i$$

Finding tangent plane at point P :

1. Change the equation of surface to $\phi(x, y, z) = c$ where c is a constant. Try to use addition instead of multiplication so that ϕ is more linear.
- 2.

$$\left. \frac{\partial \phi}{\partial x} \right|_P (x - x_P) + \left. \frac{\partial \phi}{\partial y} \right|_P (y - y_P) + \left. \frac{\partial \phi}{\partial z} \right|_P (z - z_P)$$

is the equation of tangent plane.

For a vector field \mathbf{A}

$$\text{div}(\mathbf{A}) = \nabla \cdot \mathbf{A}, \quad \text{curl}(\mathbf{A}) = \nabla \times \mathbf{A}$$

Note: 2D curl is defined assuming it is on $x - y$ plane of 3-D space, so

$$\text{curl}(\mathbf{A}) = \left(\frac{\partial A_2}{\partial x} - \frac{\partial A_1}{\partial y} \right) \hat{\mathbf{k}}$$

Strict definitions:

$$\text{div}(\mathbf{A}) = \lim_{\tau \rightarrow 0} \frac{1}{\tau} \int_S (\hat{\mathbf{n}} \cdot \mathbf{A}) dS \quad \text{curl}(\mathbf{A}) = \lim_{\tau \rightarrow 0} \frac{1}{\tau} \int_S (\hat{\mathbf{n}} \times \mathbf{A}) dS$$

Properties:

- $\nabla(\phi_1 + \phi_2) = \nabla\phi_1 + \nabla\phi_2$
- $\text{div}(\mathbf{A} + \mathbf{B}) = \text{div}(\mathbf{A}) + \text{div}(\mathbf{B}), \text{curl}(\mathbf{A} + \mathbf{B}) = \text{curl}(\mathbf{A}) + \text{curl}(\mathbf{B})$
- $\nabla(\phi\psi) = \phi\nabla\psi + \psi\nabla\phi$
- $\text{div}(\phi\mathbf{A}) = \phi\text{div}(\mathbf{A}) + \nabla\phi \cdot \mathbf{A}, \text{curl}(\phi\mathbf{A}) = \phi\text{curl}(\mathbf{A}) + \nabla\phi \times \mathbf{A}$
- $\text{div}(\mathbf{A} \times \mathbf{B}) = \mathbf{B} \cdot \text{curl}(\mathbf{A}) - \mathbf{A} \cdot \text{curl}(\mathbf{B})$
- $\text{curl}(\mathbf{A} \times \mathbf{B}) = (\mathbf{B} \cdot \nabla)\mathbf{A} - (\mathbf{A} \cdot \nabla)\mathbf{B} + \text{div}(\mathbf{B})\mathbf{A} - \text{div}(\mathbf{A})\mathbf{B}$
- $\nabla(\mathbf{A} \cdot \mathbf{B}) = (\mathbf{B} \cdot \nabla)\mathbf{A} + (\mathbf{A} \cdot \nabla)\mathbf{B} + \mathbf{B} \times \text{curl}(\mathbf{A}) + \mathbf{A} \times \text{curl}(\mathbf{B})$
- $\text{curl}(\nabla\phi) = 0$
- $\text{div}(\text{curl}(\mathbf{A})) = 0$
- $\text{curl}(\text{curl}(\mathbf{A})) = \nabla(\text{div}(\mathbf{A})) - \nabla^2 \mathbf{A}$

Irrotational vector field: $\text{curl}(\mathbf{A}) = 0$

solenoid vector field: $\text{div}(\mathbf{A}) = 0$

3 Integrals

3.1 Line Integral

$$\int_{\gamma} f \, ds := \lim_{N \rightarrow \infty, \max(\delta s_n) \rightarrow 0} \sum_{n=1}^N f_n \delta s_n$$

where δs_n are lengths of segments on the path γ (must be smooth or piece-wise smooth). Function f is usually a scalar field, if it is vector field, integral is calculated element-wise.

Length of path γ can be calculated using integration $l = \int_{\gamma} ds$. If γ is $y = y(x)$, then the arc length between $(a, f(a)), (b, f(b))$ is

$$\int_a^b \sqrt{1 + \left(\frac{dy}{dx}\right)^2} \, dx$$

ds - Change of arc length, $\hat{\mathbf{t}} = \frac{d\mathbf{r}}{ds}$ - tangent vector, $d\mathbf{r} = \hat{\mathbf{t}} ds$ - path element. So we have another line integral defined for vector field \mathbf{F} :

$$\int_{\gamma} \mathbf{F} \cdot d\mathbf{r} := \int_{\gamma} (\mathbf{F} \cdot \hat{\mathbf{t}}) \, ds$$

Circulation: when the path is closed (same beginning and end point) denoted as $\oint_{\gamma} \mathbf{F} \cdot d\mathbf{r}$

Field \mathbf{F} is conservative (circulation around any closed path γ is 0) $\Leftrightarrow \mathbf{F} = \nabla \phi$ for some function ϕ . (ϕ is called potential) $\Leftrightarrow \text{curl}(\mathbf{F}) = 0$.

In this case, given any path γ joining point A, B , we have

$$\int_{\gamma} f \, ds = \phi(B) - \phi(A)$$

Steps to find potential ϕ :

- $\frac{\partial \phi}{\partial x} = F_1$ so integrate F_1 w.r.t. x to find ϕ . The integration constant should be $C(y, z)$ (a function depending on y, z only)
- Then differentiation ϕ w.r.t y and compare with F_2 to solve for $C(y, z)$. It should be of the form $g(y, z) + C(z)$.
- Finally differentiation ϕ w.r.t. z . You may leave the constant c there as there is no way to get rid of it.

Another practical result is if there is a vector field \mathbf{B} s.t. $\text{curl}(\mathbf{B}) = \mathbf{A}$, then $\text{div}(\mathbf{A}) = 0$

Evaluation Line Integral

With parameterisation $x = x(t), y = y(t), z = z(t) t_0 \leq t \leq t_1$, we have

$$\int_{\gamma} \mathbf{F} \cdot d\mathbf{r} = \int_{t_0}^{t_1} \left(F_1 \frac{dx}{dt} + F_2 \frac{dy}{dt} + F_3 \frac{dz}{dt} \right) dt$$

3.2 Surface Integral

Some definitions: Convex surface: crossed by any straight line at most twice.

Closed surface: can divide the space into two non-connected regions. (interior & exterior)

Simply Connected Region: Any curve inside the region can be shrunk to a point (any point) without leaving the region. Otherwise the region is called multiply-connected.

Similar to line integral, we define surface integral on surface S as below

$$\int_S f dS := \lim_{N \rightarrow \infty, \max(\delta S_n) \rightarrow 0} \sum_{n=1}^N f_n \delta S_n$$

dS - area element (infinite small area of surface), $d\mathbf{S} = \hat{\mathbf{n}} dS$ - vector areal element ($\hat{\mathbf{n}}$ is normal to dS). Note area of S can be calculated using $\int_S 1 dS$.

Evaluating surface integral

$$I = \int_S f(P) dS$$

where P is a general point on surface S .

1. (From our notes) Projection

Choose a plane to project to (say the x, y plane) and find projection area Σ and normal vector $\hat{\mathbf{m}} = \hat{\mathbf{k}}$ to projected plane. Change variables in $f(x, y, z)$ if necessary to get rid of z (using the equation of surface S , NOT equation of Σ)

$$I = \int_{\Sigma} f(P) \frac{dx dy}{|\hat{\mathbf{n}} \cdot \hat{\mathbf{k}}|}$$

where $\hat{\mathbf{n}}$ is normal to S and it may depend on x, y, z . If the surface is given by $g(x, y, z) = c$ for some constant c , then $\hat{\mathbf{n}} = \nabla g / |\nabla g|$ (ALWAYS remember to check $\hat{\mathbf{n}}$ is pointing to exterior of region). Projection to other planes can be done similarly.

2. Parameterisation

Any surface can be parameterised by two parameters say $\mathbf{r} = \mathbf{r}(\theta, \phi)$.

For scalar function f :

$$\int_S f(\mathbf{r}) dS = \iint_A f(\mathbf{r}(\theta, \phi)) \left| \frac{\partial \mathbf{r}}{\partial \theta} \times \frac{\partial \mathbf{r}}{\partial \phi} \right| d\phi d\theta$$

where A is the corresponding area on $\theta - \phi$ plane.

For vector field \mathbf{F} :

$$\int_S \mathbf{F}(\mathbf{r}) \cdot d\mathbf{S} = \iint_A (\mathbf{F}(\mathbf{r}(\theta, \phi)) \cdot \hat{\mathbf{n}}) \left| \frac{\partial \mathbf{r}}{\partial \theta} \times \frac{\partial \mathbf{r}}{\partial \phi} \right| d\phi d\theta$$

where $\hat{\mathbf{n}}$ is unit normal to surface S .

3.3 Volume Integral

If τ is a region in 3D space, $\int_{\tau} f d\tau$ (defined similar to above) is volume integral where $d\tau = dx dy dz$ is volume element.

3.4 Some useful Integrals

Given curve $y = y(x), x \in [a, b]$, the arc length is

$$\int_a^b \sqrt{1 + \left(\frac{dy}{dx}\right)^2} dx$$

surface area of surface generated by revolving $y = y(x)$ between $x = a, b$ about x -axis is:

$$2\pi \int_a^b y \sqrt{1 + \left(\frac{dy}{dx}\right)^2} dx$$

for revolution about y -axis:

$$2\pi \int_a^b x \sqrt{1 + \left(\frac{dy}{dx}\right)^2} dx$$

volume of revolution about x -axis is

$$\int_a^b \pi y^2 dx$$

volume of revolution about y -axis is

$$\int_a^b \pi x^2 dy$$

4 Green's theorem, divergence theorem and Stokes theorem

Green's theorem gives an important connection between line integral and surface integral. If R is a closed plane region bounded by simple closed convex curve C (anti-clockwise) and $\mathbf{F} = \begin{pmatrix} P \\ Q \end{pmatrix}$:

$$\oint_C \mathbf{F} \cdot d\mathbf{r} = \oint_C (P dx + Q dy) = \iint_R \left(\frac{\partial Q}{\partial x} - \frac{\partial P}{\partial y} \right) dS$$

Green's theorem is proved directly for simply connected convex R but can be generalised to non-convex/multiply connected regions using suitable breakdown. For example, given a region R is bounded by an interior boundary C_0 (clockwise) and an exterior boundary C_1 (anti-clockwise):

$$\oint_{C_1} \mathbf{F} \cdot d\mathbf{r} - \oint_{C_0} \mathbf{F} \cdot d\mathbf{r} = \iint_R \left(\frac{\partial Q}{\partial x} - \frac{\partial P}{\partial y} \right) dS$$

If \mathbf{F} is undefined at some point P in region R , you should dig a circle with radius r around point P , calculate the integral and then let $r \rightarrow 0$.

Area of a region If a closed region has boundary curve C , then its area is given by

$$\frac{1}{2} \oint_C (x dy - y dx)$$

This integral can be calculated by parametrisation.

Flux

Flux of a surface S and vector field \mathbf{A} is

$$\int_S \mathbf{A} \cdot \hat{\mathbf{n}} dS$$

where $\hat{\mathbf{n}}$ is normal to surface S .

Divergence theorem

For **closed convex** surface S with normal $\hat{\mathbf{n}}$ and region τ enclosed by S that is simply connected. \mathbf{A} is a vector field with continuous derivative on τ :

$$\int_S \mathbf{A} \cdot d\mathbf{S} = \int_S \mathbf{A} \cdot \hat{\mathbf{n}} dS = \int_\tau \text{div}(\mathbf{A}) d\tau$$

This can be generalised to non-convex surfaces and multiply-connected regions just like Green's theorem.

You can use divergence theorem to evaluate $\int_\tau f(\mathbf{r}) d\tau$ but you have to first find function $\mathbf{A}(\mathbf{r})$ s.t. $\text{div} \mathbf{A} = f$. Try to pick \mathbf{A} as simple as possible (say, $\mathbf{A} = x\mathbf{i}$)

Green's identities: (ϕ, ψ are scalar functions, other symbols have the same meaning as in Divergence theorem)

$$\begin{aligned} \iint_S \phi \frac{\partial \psi}{\partial \mathbf{n}} dS &= \int_\tau \phi \nabla^2 \psi + \nabla \phi \cdot \nabla \psi d\tau \\ \iint_S \phi \frac{\partial \psi}{\partial \mathbf{n}} - \psi \frac{\partial \phi}{\partial \mathbf{n}} dS &= \int_\tau (\phi \nabla^2 \psi - \psi \nabla^2 \phi) d\tau \end{aligned}$$

Generalised integration by part: (R is region enclosed by curve C)

$$\iint_R \phi \nabla^2 \psi dx dy = \oint_C \phi \frac{\partial \psi}{\partial \mathbf{n}} dS - \iint_R \nabla \phi \cdot \nabla \psi dx dy$$

Gauss's flux theorem: S is closed surface with normal $\hat{\mathbf{n}}$ and O means origin.

$$\int_S \frac{\hat{\mathbf{n}} \cdot \mathbf{r}}{|\mathbf{r}|^3} dS = \begin{cases} 0 & \text{if } O \text{ is exterior} \\ 4\pi & \text{if } O \text{ is interior} \end{cases}$$

Stokes theorem S is **open** surface with boundary curve γ . \mathbf{A} is continuously differentiable on S ,

$$\oint_{\gamma} \mathbf{A} \cdot d\mathbf{r} = \int_S \text{curl}(\mathbf{A}) \cdot \hat{\mathbf{n}} \, dS$$

Note LHS is independent of the surface S chosen.

5 Curvilinear system

For coordinate system $\mathbf{x} = (x_1, x_2, x_3)$, transformation to new coordinate system $\mathbf{u} = (u_1, u_2, u_3)$ is possible if

$$\det(J(u_x)) = \det \begin{pmatrix} \frac{\partial u_1}{\partial x_1} & \frac{\partial u_1}{\partial x_2} & \frac{\partial u_1}{\partial x_3} \\ \frac{\partial u_2}{\partial x_1} & \frac{\partial u_2}{\partial x_2} & \frac{\partial u_2}{\partial x_3} \\ \frac{\partial u_3}{\partial x_1} & \frac{\partial u_3}{\partial x_2} & \frac{\partial u_3}{\partial x_3} \end{pmatrix} \neq 0$$

Also remember that $J(x_u)J(u_x) = I$, so $\det(J(x_u))\det(J(u_x)) = 1$. This means if $\det(J(u_x)) \neq 0$, the above condition is satisfied.

For a given point $P = (x_1, x_2, x_3)$, by finding $\nabla u_i(P)$ and normalising (change to a unit vector). We have $\hat{\mathbf{u}}_i$, the unit normal vector to the surface $u_i = u_i(P)$ (where $u_i(P)$ is a constant). So

$$\hat{\mathbf{u}}_i = \frac{\nabla u_i(P)}{|\nabla u_i(P)|}$$

system is orthogonal of $\hat{\mathbf{u}}_i$ are mutually orthogonal.

In general, tangential vectors(to the line where only u_i varies) $\hat{\mathbf{e}}_i$ are used more. In Cartesian coordinates, $\hat{\mathbf{e}}_i$ are $\hat{i}, \hat{j}, \hat{k}$. To find $\hat{\mathbf{e}}_i$, use this formulae:

$$\frac{\partial \mathbf{r}}{\partial u_i} = h_i \hat{\mathbf{e}}_i, \text{ where } h_i = \left| \frac{\partial \mathbf{r}}{\partial u_i} \right|, \mathbf{r} = x\hat{i} + y\hat{j} + z\hat{k}$$

h_i are called length scale. Showing $\hat{\mathbf{e}}_i$ are mutually orthogonal also proves orthogonality of the coordinate system.

In orthogonal system:

$$\hat{\mathbf{e}}_i = \hat{\mathbf{u}}_i$$

Path element

$$d\mathbf{r} = \sum h_i du_i \hat{\mathbf{e}}_i$$

(for orthogonal system)

$$(ds)^2 = (d\mathbf{r} \cdot d\mathbf{r}) = \sum h_i^2 du_i^2$$

Volume element

$$dV = h_1 h_2 h_3 du_1 du_2 du_3$$

Area element (on the surface where u_1 is constant)

$$dS = h_2 h_3 du_2 du_3$$

Gradient

$$\nabla = \sum \frac{1}{h_i} \hat{e}_i \frac{\partial}{\partial u_i}$$

from this we have $\hat{e}_i = \nabla \mathbf{u}_i h_i$

Divergence

$$\text{div}(\mathbf{A}) = \frac{1}{h_1 h_2 h_3} \left\{ \sum \frac{\partial}{\partial u_i} (A_i h_j h_k) \right\}$$

Curl

$$\text{curl}(\mathbf{A}) = \frac{1}{h_1 h_2 h_3} \begin{vmatrix} h_1 \hat{e}_1 & h_2 \hat{e}_2 & h_3 \hat{e}_3 \\ \frac{\partial}{\partial u_1} & \frac{\partial}{\partial u_2} & \frac{\partial}{\partial u_3} \\ h_1 A_1 & h_2 A_2 & h_3 A_3 \end{vmatrix}$$

Laplacian

$$\nabla^2 \Phi = \frac{1}{h_1 h_2 h_3} \left\{ \sum \frac{\partial}{\partial u_i} \left(\frac{h_j h_k}{h_i} \frac{\partial \Phi}{\partial u_i} \right) \right\}$$

5.1 Cartesian, Cylindrical and Spherical

Three systems are all orthogonal.

Definition of Cylindrical coordinates:

$$x = r \cos \theta, y = r \sin \theta, z = z$$

Definition of Spherical coordinates:

$$x = r \sin \theta \cos \phi, y = r \sin \theta \sin \phi, z = r \cos \theta$$

where $r \geq 0$ is radius, $\theta \in [0, \pi]$ is rotation above x-y plane and $\phi \in [0, 2\pi]$ is rotation with in x-y plane.

Tables below are ordered by (x, y, z) , (r, θ, z) , (r, θ, ϕ)

Coordinate system	h_i	\hat{e}_i
Cartesian	1,1,1	$\hat{i}, \hat{j}, \hat{k}$
Cylindrical	1	$\cos \theta \hat{i} + \sin \theta \hat{j}$
	r	$-\sin \theta \hat{i} + \cos \theta \hat{j}$
	1	\hat{k}
Spherical	1	$\sin \theta \cos \phi \hat{i} + \sin \theta \sin \phi \hat{j} + \cos \theta \hat{k}$
	r	$\cos \theta \cos \phi \hat{i} + \cos \theta \sin \phi \hat{j} - \sin \theta \hat{k}$
	$r \sin \theta$	$-\sin \theta \sin \phi \hat{i} + \sin \theta \cos \phi \hat{j}$

Table 1: Table of tangential vectors and length scales

Coordinate system	gradient	div	Laplacian
Cartesian	$\hat{\mathbf{i}} \frac{\partial}{\partial x} + \hat{\mathbf{j}} \frac{\partial}{\partial y} + \hat{\mathbf{k}} \frac{\partial}{\partial z}$	$\frac{\partial A_1}{\partial x} + \frac{\partial A_2}{\partial y} + \frac{\partial A_3}{\partial z}$	$\frac{\partial^2 \Phi}{\partial x^2} + \frac{\partial^2 \Phi}{\partial y^2} + \frac{\partial^2 \Phi}{\partial z^2}$
Cylindrical	$\hat{\mathbf{r}} \frac{\partial}{\partial r} + \frac{\hat{\boldsymbol{\theta}}}{r} \frac{\partial}{\partial \theta} + \hat{\mathbf{k}} \frac{\partial}{\partial z}$	$\frac{\partial A_1}{\partial r} + \frac{A_1}{r} + \frac{1}{r} \frac{\partial A_2}{\partial \theta} + \frac{\partial A_3}{\partial z}$	$\frac{\partial^2 \Phi}{\partial r^2} + \frac{1}{r} \frac{\partial \Phi}{\partial r} + \frac{1}{r^2} \frac{\partial^2 \Phi}{\partial \theta^2} + \frac{\partial^2 \Phi}{\partial z^2}$
Spherical	$\hat{\mathbf{r}} \frac{\partial}{\partial r} + \frac{\hat{\boldsymbol{\theta}}}{r} \frac{\partial}{\partial \theta} + \frac{\hat{\boldsymbol{\phi}}}{r \sin \theta} \frac{\partial}{\partial \phi}$	$\frac{1}{r^2 \sin \theta} \left\{ \frac{\partial}{\partial r} (r^2 \sin \theta A_1) + \frac{\partial}{\partial \theta} (r \sin \theta A_2) + \frac{\partial}{\partial \phi} (r A_3) \right\}$	$\frac{\partial^2 \Phi}{\partial r^2} + \frac{2}{r} \frac{\partial \Phi}{\partial r} + \frac{\cot \theta}{r^2} \frac{\partial \Phi}{\partial \theta} + \frac{1}{r^2} \frac{\partial^2 \Phi}{\partial \theta^2} + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 \Phi}{\partial \phi^2}$

Table 2: Table of div, gradient and Laplacian for vector valued function \mathbf{A} or scalar valued function Φ

5.2 Change of variables

If we parameterise surface S by (u_1, u_2) ,

$$dS = |J| du_1 du_2, \text{ where } J = \frac{\partial \mathbf{r}}{\partial u_1} \times \frac{\partial \mathbf{r}}{\partial u_2}$$

If S is on x-y plane

$$J = J(\mathbf{x}_u), \text{ where } \mathbf{x} = (x, y)$$

If S is $z = f(x, y)$ and we parameterise using x, y :

$$|J| = \sqrt{1 + |\nabla f|^2}$$

So surface area of any surface S is

$$\int_{\Sigma} \sqrt{1 + |\nabla f|^2} dx dy$$

where Σ is projection of S onto x-y plane.

6 Calculus of Variation

Vanishing lemma If g is a continuous function s.t. for every smooth function $\eta(x)$ with $\eta(x_1) = \eta(x_2) = 0$,

$$\int_{x_1}^{x_2} g(x) \eta(x) dx = 0$$

then $g \equiv 0$

Target of this section: find function $y = y(x)$ that minimises the integral

$$I := \int_{x_1}^{x_2} L(x, y, y') dx$$

L is a functional.

1D E-L equation and special cases

Full equation

$$\frac{\partial L}{\partial y} = \frac{d}{dx} \left\{ \frac{\partial L}{\partial y'} \right\}$$

L independent of y

$$\frac{\partial L}{\partial y'} = \text{constant}$$

L independent of y'

$$\frac{\partial L}{\partial y} = 0$$

L independent of x :

$$L - y' \frac{\partial L}{\partial y'} = \text{constant}$$

Finding extrema of $I = \int_{x_1}^{x_2} L(x, y, y') dx$ where L is a functional

1. Use E-L equation to obtain a differential equation.
2. Solve it, remember to include the integration constants.
3. Check that $I''(0) > 0$ or $I''(0) < 0$
4. use boundary conditions to determine integration constants

Multivariate E-L If we are finding extrema of

$$I = \int_{t_1}^{t_2} L(t, x_1(t), x_1'(t), \dots, x_n(t), x_n'(t)) dt$$

E-L equation becomes a set of equations

$$\frac{\partial L}{\partial x_i} = \frac{d}{dt} \left\{ \frac{\partial L}{\partial x_i'} \right\}$$

With constraint

$$J = \int_{t_1}^{t_2} g(t, x_1(t), x_1'(t), \dots, x_n(t), x_n'(t)) dt = \text{constant}$$

E-L becomes

$$\frac{\partial}{\partial x_i} (L + \lambda g) - \frac{d}{dt} \frac{\partial}{\partial x_i'} (L + \lambda g) = 0$$

Leave λ there and after finishing regular steps, plug in integral J to determine λ . However, the order you determine the constants is not restricted, you may as well find λ in terms of an integration constant c first, then solve c using the restriction.

E-L for surface integral The function $f(\mathbf{r})$ where $\mathbf{r} = x\mathbf{i} + y\mathbf{j}$ that maximises

$$I = \int_R L(\mathbf{r}, f(\mathbf{r}), \nabla f(\mathbf{r})) dx dy$$

can be found by solving the following equation

$$\frac{\partial L}{\partial f} = \text{div}(\nabla_{\nabla f} L) = \text{div}(\hat{\mathbf{i}} \frac{\partial L}{\partial f_x} + \hat{\mathbf{j}} \frac{\partial L}{\partial f_y})$$

where f_x, f_y are partial derivatives of f .

Isoperimetric inequality For any simple curve with area A and perimeter l ,

$$4\pi A \leq l^2$$

equality holds iff curve is circle.

7 Integration Techniques

Ways of showing $f(x) \equiv 0$ using integration:

1. Use vanishing lemma
2. Prove $\int_{x_1}^{x_2} f^2 dx = 0$. Or if you know $f(x) \geq 0$, prove $\int_{x_1}^{x_2} f dx = 0$

Some inverse trig integrals

$$\begin{aligned}\int \frac{dx}{\sqrt{a^2 - x^2}} &= \sin^{-1}\left(\frac{x}{a}\right) + c \\ \int \frac{dx}{a^2 + x^2} &= \frac{1}{a} \tan^{-1}\left(\frac{x}{a}\right) + c \\ \int \frac{dx}{x\sqrt{x^2 - a^2}} &= \frac{1}{a} \sec^{-1}\left(\frac{|x|}{a}\right) + c\end{aligned}$$

Some inverse hyperbolic integrals

$$\begin{aligned}\int \frac{dx}{\sqrt{a^2 + x^2}} &= \sinh^{-1}\left(\frac{x}{a}\right) + c \\ \int \frac{dx}{\sqrt{x^2 - a^2}} &= \cosh^{-1}\left(\frac{x}{a}\right) + c \\ \int \frac{dx}{a^2 - x^2} &= \frac{1}{a} \tanh^{-1}\left(\frac{x}{a}\right) + c \quad (x^2 < a^2) \\ \int \frac{dx}{a^2 - x^2} &= \frac{1}{a} \coth^{-1}\left(\frac{x}{a}\right) + c \quad (x^2 > a^2)\end{aligned}$$

Derivatives of trigs:

$$\begin{aligned}(\sin(x))' &= \cos(x) & (\cos(x))' &= -\sin(x) & (\tan(x))' &= \sec^2(x) \\ (\csc(x))' &= -\cot(x) \csc(x) & (\sec(x))' &= \tan(x) \sec(x) & (\cot(x))' &= -\csc^2(x)\end{aligned}$$

Reduction formulas:

$$\begin{aligned}I_n &= \int \cos^n(x) dx = \frac{1}{n} \cos^{n-1}(x) \sin(x) + \frac{n-1}{n} I_{n-2} \\ I_n &= \int \sin^n(x) dx = -\frac{1}{n} \sin^{n-1}(x) \cos(x) + \frac{n-1}{n} I_{n-2} \\ I_n &= \int x^n e^{ax} dx = \frac{x^n e^{ax}}{a} - \frac{1}{n} I_{n-1}\end{aligned}$$

Product to sum rules for trigs:

$$\begin{aligned}\sin(x) \cos(y) &= \frac{1}{2} [\sin(x+y) + \sin(x-y)] \\ \cos(x) \sin(y) &= \frac{1}{2} [\sin(x+y) - \sin(x-y)] \\ \cos(x) \cos(y) &= \frac{1}{2} [\cos(x+y) + \cos(x-y)] \\ \sin(x) \sin(y) &= -\frac{1}{2} [\cos(x+y) - \cos(x-y)]\end{aligned}$$