

Aerial Manipulation using a Tilted-Motors Hexacopter with a Two DoF Robotic Arm

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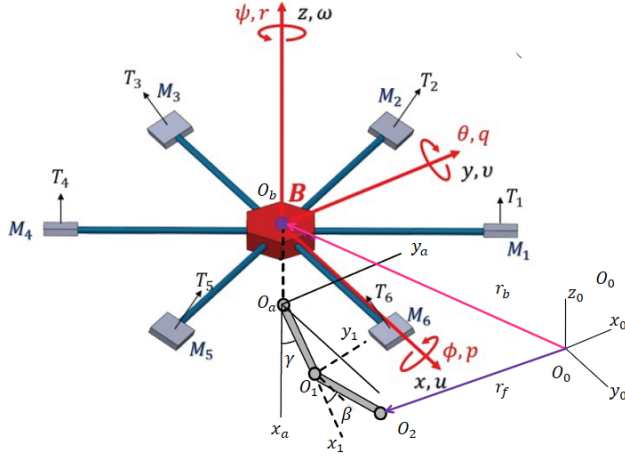


Figure 1: UAM/Arm system illustration. O_i represents the origin of a reference system with unit vectors x_i, y_i, z_i . The angle γ is taken from x_a to the first link and β from x_1 to the second link, both restricted to move in the plane (x_a, y_a) or (y_b, z_b) .

I. System Modeling

In this section, modeling of the UAM presented in Fig. 1 is formulated using Euler-Lagrange method. The following assumptions are considered:

- The 2 DoF robotic arm is modeled as a double compound pendulum.
- No payload is attached to the end-effector.
- Air drag is negligible.

A. Kinematic Model

Coordinate frames of the combined system are defined in Fig. 1. The frames O_0, O_b, O_a, O_1, O_2 are namely the inertial frame, the body frame of the hexarotor, the joint between the aerial vehicle and the manipulator, the joint between both links, and the end-effector attached frame.

Let the absolute position of the UAV be denoted by $r_b = [x, y, z]^T \in \mathbb{R}^3$, while the attitude is described by

the yaw-pitch-roll Euler Angles $\eta = [\phi, \theta, \psi] \in \mathbb{R}^3$. The position of the center of mass and angular velocity of link i with respect to O_0 can then be expressed by

$$r_{li} = r_b + \mathbf{R}_b r_{li}^b, \quad (1)$$

$$\omega_{li} = \omega_b + \mathbf{R}_b \omega_{li}^b, \quad (2)$$

where r_{li}^b and ω_{li}^b are the position of the center of mass and angular velocity of the link with respect to O_b , and

$$\mathbf{R}_b = \begin{bmatrix} C_\psi C_\theta & C_\psi S_\theta S_\phi - C_\phi S_\psi & S_\phi S_\psi + C_\phi C_\psi S_\theta \\ C_\theta S_\psi & C_\phi C_\psi + S_\psi S_\theta S_\phi & -C_\psi S_\phi + S_\psi S_\theta C_\phi \\ -S_\theta & C_\theta S_\phi & S_\theta C_\phi \end{bmatrix}, \quad (3)$$

is the rotation matrix denoting the orientation of O_b with respect to O_0 , with $C_* = \cos(*)$, $S_* = \sin(*)$.

If the angular velocity of the hexacopter with respect to O_0 is denoted by ω_b , while ω_b^b is the angular velocity of the UAV with respect to O_b , the following relations persist

$$\omega_b = \mathbf{A} \dot{\eta}, \quad (4)$$

$$\omega_b^b = \mathbf{R}_b^T \mathbf{A} \dot{\eta} = \mathbf{W} \dot{\eta}, \quad (5)$$

where

$$\mathbf{W} = \begin{bmatrix} 1 & 0 & -S_\theta \\ 0 & C_\phi & C_\theta S_\phi \\ 0 & -S_\phi & C_\theta C_\phi \end{bmatrix}, \quad (6)$$

maps from the time derivative of the Euler Angles to the angular velocity expressed in O_b .

Consider the following expression for the linear and angular velocities of the i -th link with respect to O_b

$$\dot{r}_{li}^b = \mathbf{J}_v^{li} \dot{q}, \quad (7)$$

$$\omega_{li}^b = \mathbf{J}_\omega^{li} \dot{q}, \quad (8)$$

where $q = [\gamma, \beta]^T \in \mathbb{R}^2$ is the vector of joint variables and \mathbf{J}_v^{li} and \mathbf{J}_ω^{li} are the contributions of the Jacobian columns relative to the joint velocities up to the current link i .

Taking the time derivative of 1 and substituting 7 yields

$$\dot{r}_{li} = \dot{r}_b - S(\mathbf{R}_b r_{li}^b) \omega_b + \mathbf{R}_b \mathbf{J}_v^{li} \dot{q}, \quad (9)$$

where $S(v)$ is the skew-symmetric matrix related to the vector v . By substituting 8 in 2 a similar expression for the angular velocity can be obtained

$$\omega_{li} = \omega_b + \mathbf{R}_b \mathbf{J}_\omega^{li} \dot{q}. \quad (10)$$

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The full kinematic model can then be rewritten in matrix form

$$\dot{r}_b = [\mathbf{I}_{3 \times 3}, \mathbf{0}_{3 \times 3}, \mathbf{0}_{3 \times 2}] \dot{\xi} = \mathbf{D}_r^b \dot{\xi} \quad (11)$$

$$\omega_b = [\mathbf{I}_{3 \times 3}, \mathbf{A}, \mathbf{0}_{3 \times 2}] \dot{\xi} = \mathbf{D}_\omega^b \dot{\xi} \quad (12)$$

$$\dot{r}_{li} = [\mathbf{I}_{3 \times 3}, -S(\mathbf{R}_b r_{li}^b \mathbf{A}), \mathbf{R}_b \mathbf{J}_v^{li}] \dot{\xi} = \mathbf{D}_r^{li} \dot{\xi} \quad (13)$$

$$\omega_{li} = [\mathbf{0}_{3 \times 3}, \mathbf{A}, \mathbf{R}_b \mathbf{J}_\omega^{li}] \dot{\xi} = \mathbf{D}_\omega^{li} \dot{\xi} \quad (14)$$

where

$$\xi = [x, y, z, \phi, \theta, \psi, \gamma, \beta]^T \in \mathbb{R}^8, \quad (15)$$

is the vector of generalized coordinates, $\mathbf{I}_{3 \times 3}$ is the (3×3) identity matrix, and $\mathbf{0}_{n \times m}$ is the $(n \times m)$ zero matrix.

Furthermore, the manipulator system is showed in 2. Hence, the rotation and translation operators up to the body frame are defined as follows

$$\mathbf{R}_2^1 = \mathbf{Z}(\beta), \quad (16) \quad d_2^1 = l_2 \begin{bmatrix} C_\beta \\ S_\beta \\ 0 \end{bmatrix}, \quad (17)$$

$$\mathbf{R}_1^a = \mathbf{Z}(\gamma), \quad (18) \quad d_1^a = l_1 \begin{bmatrix} C_\gamma \\ S_\gamma \\ 0 \end{bmatrix}, \quad (19)$$

$$\mathbf{R}_a^b = \mathbf{Y} \left(\frac{\pi}{2} \right), \quad (20) \quad d_a^b = a \begin{bmatrix} 0 \\ 0 \\ -1 \end{bmatrix}, \quad (21)$$

where \mathbf{R}_j^i and d_j^i represent the rotation and translation from frame O_j to O_i , respectively, $\mathbf{Y}(\cdot)$ and $\mathbf{Z}(\cdot)$ are the rotation matrix about the y and z axis, respectively, a is the distance from O_b to O_a , and l_i the total length of i -th link. Note that $d_b = r_b = [x, y, z]^T$ is the position of the hexacopter with respect to the inertial frame.

The double compound pendulum system is a really well studied problem, the position of the center of mass of each link with respect to O_a can derived by

$$r_{l1}^a = \left(c_1 \begin{bmatrix} 0 \\ C_\gamma \\ S_\gamma \end{bmatrix} \right), \quad (22)$$

$$r_{l2}^a = \mathbf{R}_1^a \left(c_2 \begin{bmatrix} 0 \\ C_\beta \\ S_\beta \end{bmatrix} \right) + d_1^a, \quad (23)$$

where c_i is the distance from the beginning of link i to its center of mass. Then, adding the rotation about the y axis \mathbf{R}_a^b and translation d_a^b yields

$$\begin{aligned} r_{l1}^b &= \mathbf{R}_a^b r_{l1}^a + d_a^b \\ &= \begin{bmatrix} 0 \\ c_1 S_\gamma \\ -a - c_1 C_\gamma \end{bmatrix}, \end{aligned} \quad (24)$$

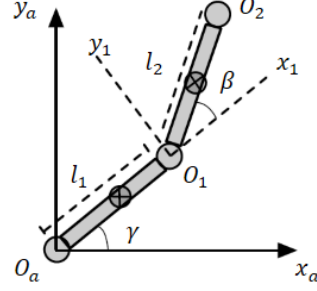


Figure 2: Double compound pendulum system.

$$\begin{aligned} r_{l2}^b &= \mathbf{R}_a^b r_{l2}^a + d_a^b \\ &= \begin{bmatrix} 0 \\ l_1 S_\gamma + c_2 S_{\gamma\beta} \\ -a - l_1 C_\gamma - c_2 C_{\gamma\beta} \end{bmatrix}, \end{aligned} \quad (25)$$

where $C_{\gamma\beta} = \cos(\gamma + \beta)$ and $S_{\gamma\beta} = \sin(\gamma + \beta)$.

The angular velocity of each link with respect to the body frame can be expressed as

$$\omega_{l1}^b = \mathbf{R}_a^b \begin{bmatrix} 0 \\ 0 \\ \dot{\gamma} \end{bmatrix} = \begin{bmatrix} \dot{\gamma} \\ 0 \\ 0 \end{bmatrix}, \quad (26)$$

$$\omega_{l2}^b = \mathbf{R}_a^b \begin{bmatrix} 0 \\ 0 \\ \dot{\gamma} + \dot{\beta} \end{bmatrix} = \begin{bmatrix} \dot{\gamma} + \dot{\beta} \\ 0 \\ 0 \end{bmatrix}. \quad (27)$$

By equating the time derivative of 24 and 25 with 7, the Jacobian contributions can be defined

$$\mathbf{J}_v^{l1} = \begin{bmatrix} 0 & 0 \\ c_1 C_\gamma & 0 \\ c_1 S_\gamma & 0 \end{bmatrix}, \quad (28)$$

$$\mathbf{J}_v^{l2} = \begin{bmatrix} 0 & 0 \\ l_1 C_\gamma + c_2 C_{\gamma\beta} & c_2 C_{\gamma\beta} \\ l_1 S_\gamma + c_2 S_{\gamma\beta} & c_2 S_{\gamma\beta} \end{bmatrix}. \quad (29)$$

Also 26 and 26 with 8

$$\mathbf{J}_\omega^{l1} = \begin{bmatrix} 1 & 0 \\ 0 & 0 \\ 0 & 0 \end{bmatrix}, \quad (30) \quad \mathbf{J}_\omega^{l2} = \begin{bmatrix} 1 & 1 \\ 0 & 0 \\ 0 & 0 \end{bmatrix}. \quad (31)$$

B. Dynamic Model

The system is described by the general Matrix Form of Euler-Lagrange Equations

$$\mathbf{D}(\xi) \ddot{\xi} + \mathbf{C}(\xi, \dot{\xi}) \dot{\xi} + \mathbf{G}(\xi) = \mathbf{u}, \quad (32)$$

For the Euler-Lagrange method, it is required to establish the kinetic and potential energies of the system, thus obtaining the Lagrangian defined by

$$\mathcal{L} = \mathcal{K} - \mathcal{U} \quad (33)$$

where \mathcal{K} is the total kinetic energy and \mathcal{U} is the total potential energy. Let the contribution of the vehicle to the kinetic energy be denoted by \mathcal{K}_H and the contribution from each link \mathcal{K}_{li} , such that

$$\mathcal{K} = \mathcal{K}_H + \sum_{i=1}^2 \mathcal{K}_{li}. \quad (34)$$

The contribution of the UAV is given by

$$\mathcal{K}_H = \frac{1}{2} m_b \dot{r}_b^T \dot{r}_b + \frac{1}{2} \omega_b^T \mathbf{R}_b \mathbf{H}_b \mathbf{R}_b^T \omega_b, \quad (35)$$

in which m_b is the mass of the vehicle and \mathbf{H}_b is the inertia tensor expressed in the body frame, where, by symmetry, \mathbf{H}_b is diagonal and constant.

The contribution of each link is given by

$$\mathcal{K}_{li} = \frac{1}{2} m_{li} \dot{r}_{li}^T \dot{r}_{li} + \frac{1}{2} \omega_{li}^T \mathbf{R}_b \mathbf{R}_{li}^b \mathbf{H}_{li} (\mathbf{R}_{li}^b)^T \mathbf{R}_b^T \omega_{li}, \quad (36)$$

where \mathbf{H}_{li} is the constant inertia tensor of the link i expressed in a frame O_{li} attached to its center of mass and with the axes being coincident with the inertia central axes, while \mathbf{R}_{li}^b is the rotation matrix that transforms from O_{li} to O_b . From Fig. 2, the frames that are parallel to O_{l1} and O_{l2} are O_1 and O_2 , respectively. Thus, the rotation matrices can be defined as

$$\mathbf{R}_{l1}^b = \mathbf{R}_a^b \mathbf{R}_1^a, \quad (37) \quad \mathbf{R}_{l2}^b = \mathbf{R}_a^b \mathbf{R}_1^a \mathbf{R}_2^1. \quad (38)$$

Adding each kinetic energy contribution, the total kinetic energy 34 can be expressed as

$$\mathcal{K} = \frac{1}{2} \dot{\xi}^T \mathbf{D}(\xi) \dot{\xi}, \quad (39)$$

where $\mathbf{D}(\xi) \in \mathbb{R}^{8 \times 8}$ is the symmetric and positive definite inertia matrix from 32, which can be written using 11-14

$$\begin{aligned} \mathbf{D}(\xi) = & m_b (\mathbf{D}_r^b)^T \mathbf{D}_r^b + (\mathbf{D}_\omega^b)^T \mathbf{R}_b \mathbf{H}_b \mathbf{R}_b^T \mathbf{D}_\omega^b + \\ & \sum_{i=1}^2 (m_{li} (\mathbf{D}_r^{li})^T \mathbf{D}_r^{li} + (\mathbf{D}_\omega^{li})^T \mathbf{R}_b \mathbf{H}_{li}^b (\mathbf{R}_b^b)^T \mathbf{D}_\omega^{li}) \end{aligned} \quad (40)$$

in which $\mathbf{H}_{li}^b = \mathbf{R}_{li}^b \mathbf{H}_{li} (\mathbf{R}_{li}^b)^T$.

The Coriolis Matrix $\mathbf{C}(\xi, \dot{\xi})$ from 32 can be obtained from the Christoffel Symbols of the First Kind using the inertia matrix, hence, the k, j -th element of $\mathbf{C}(\xi, \dot{\xi}) \in \mathbb{R}^{8 \times 8}$ is defined by

$$C_{kj} = \sum_{i=1}^8 \frac{1}{2} \left[\frac{\partial d_{kj}}{\partial \xi_i} + \frac{\partial d_{ki}}{\partial \xi_j} - \frac{\partial d_{ij}}{\partial \xi_k} \right] \dot{\xi}_i, \quad (42)$$

where d_{kj} is the k, j -th element of $\mathbf{D}(\xi)$ and ξ_i is the i -th element of ξ .

On the other hand, the potential energy of the system can be defined as

$$\mathcal{U} = m_b g (\hat{z}^T r_b) + \sum_{i=1}^2 [m_{li} g (\hat{z}^T r_{li})], \quad (43)$$

where \hat{z} is the unit vector for the z axis of O_0 . Then the gravity vector $\mathbf{G}(\xi)$ k -th element is

$$G_k = \frac{\partial \mathcal{U}}{\partial \xi_k}. \quad (44)$$

According to [1, 2] the generalized torque related to the joint variable q_i can be actuated directly through joint actuators. However, the generalized torques in x, y, z, ϕ, θ and ψ need to be transformed. Hence, the input vector \mathbf{u} is

$$\mathbf{u} = \mathbf{B}(\xi) \begin{bmatrix} F_M \\ \tau_q \end{bmatrix}, \quad (45)$$

where $\tau_q = [\tau_\gamma, \tau_\beta]^T$ is the torque acting on each joint variable,

$$\mathbf{B}(\xi) = \begin{bmatrix} \mathbf{R}_b & 0 & 0 \\ 0 & \mathbf{W}^{-1} & 0 \\ 0 & 0 & \mathbf{I}_{2 \times 2} \end{bmatrix} \in \mathbb{R}^{8 \times 8}, \quad (46)$$

is the matrix presented in [3] but extended to implement the torques related to q , and F_M is the vector of distributed forces produced by the motors

$$F_M = [F_x \ F_y \ F_z \ \tau_\phi \ \tau_\theta \ \tau_\psi]^T. \quad (47)$$

Each motor provides a thrust T_i and torque Q_i related to its angular rate, thus, a simplified motor model is considered

$$T_i \approx k_T \Omega_i^2, \quad (48)$$

$$Q_i \approx k_Q \Omega_i^2, \quad (49)$$

where k_T and k_Q are the thrust and torque factor respectively, while Ω_i denotes the angular rate of the i -th motor.

In the hexacopter with tilted-rotors, the forces and torques provided by each motor have components over all of the axis. The total mapping of the forces for the six motors is arranged in the matrix

$$\mathbf{M}_\varphi = \begin{bmatrix} M_{\varphi 11} & \dots & M_{\varphi 16} \\ \vdots & \ddots & \vdots \\ M_{\varphi 61} & \dots & M_{\varphi 66} \end{bmatrix} \in \mathbb{R}^{6 \times 6}, \quad (50)$$

which depends on the inclination φ and the angle α_i that each motor has with respect to the y_b axis. The

derivation and whole components are shown in [4]. Finally, the relation between the angular rates and the distributed forces is given by

$$F_M = \mathbf{M}_\varphi \boldsymbol{\Omega}^2, \quad (51)$$

where $\boldsymbol{\Omega}^2$ is the vector containing the square of the angular rates.

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