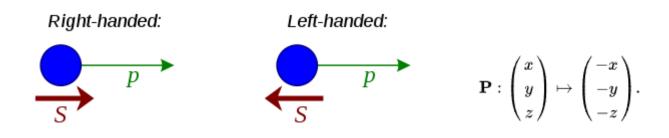
Measure polarization in potassium isotopes

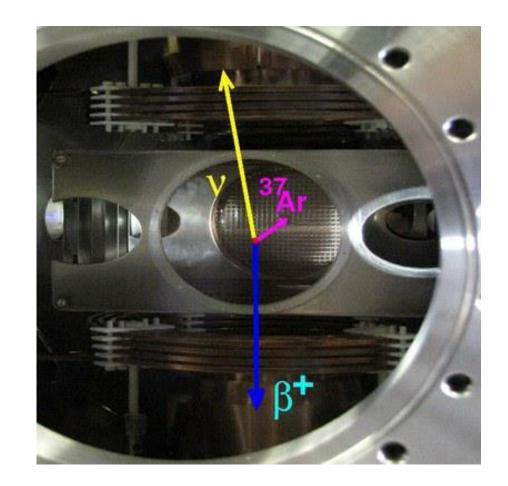
Dante Prins



Potassium 37 beta decay asymmetry experiment

Beta decay is known to violate parity symmetry, having a chirality dependance and a bias for left handedness. Our experiment polarized the parent atom and measures beta emission direction.



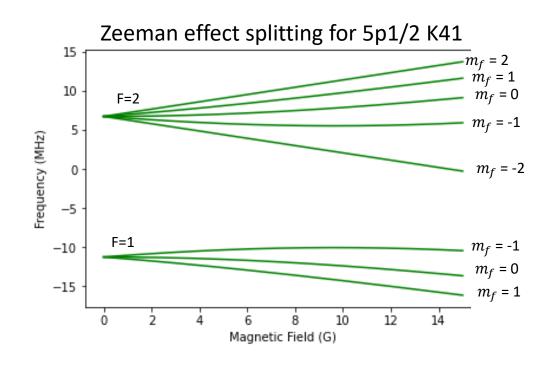


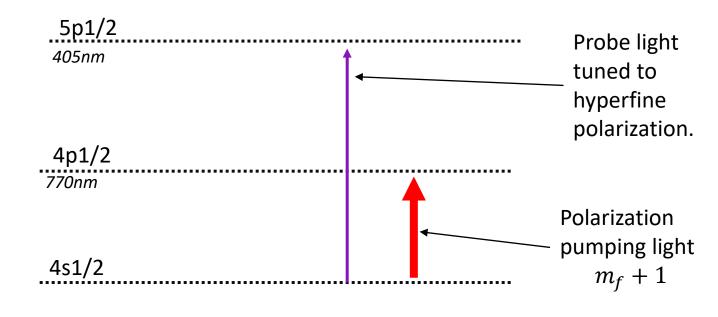
Zeeman effect at medium magnetic fields

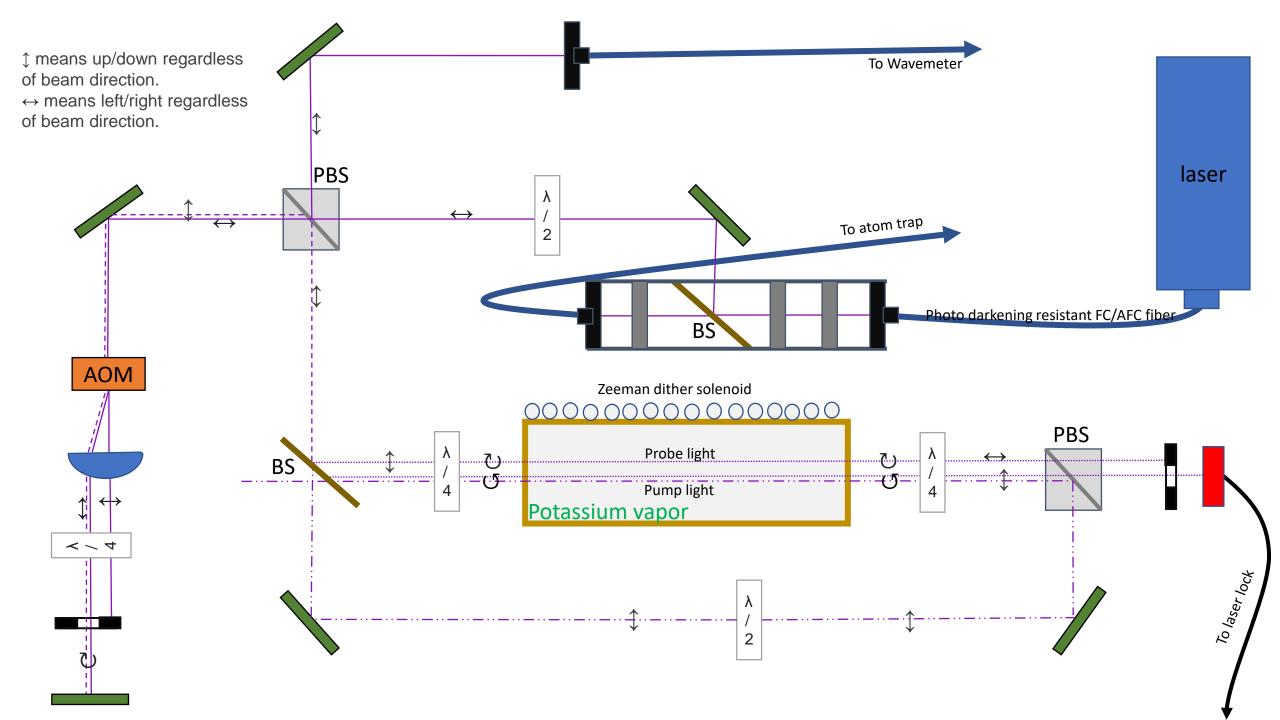
Polarization measurement

$$H=hAI_zJ_z+rac{hA}{2}(J_+I_-+J_-I_+)+\mu_{
m B}Bg_JJ_z+\mu_{
m N}Bg_II_z$$

In the presence of a magnetic field energies will shift allowing us to probe individual polarizations.







Hyperfine structure splitting varies by isotope

Isotope	Abundance
K37	1.23s half life
K39	93%
K41	7%

$$H_{HFS} = A \mathbf{I} \cdot \mathbf{J}$$

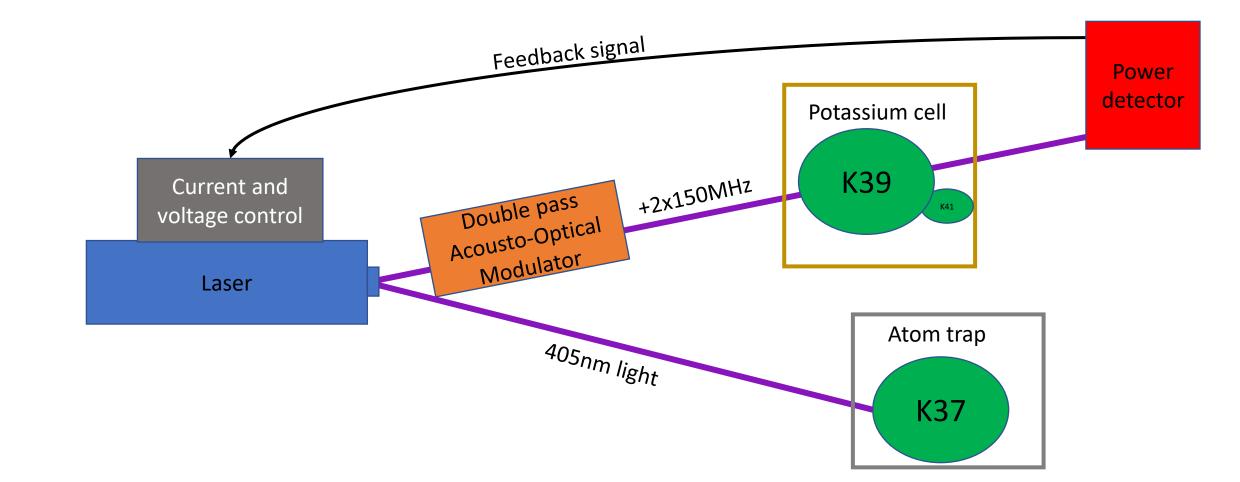
$$K37 A_{4s_{1/2}} = 240.3 \text{MHz}$$

 $K39 A_{4s_{1/2}} = 461.7 \text{MHz}$
 $K41 A_{4s_{1/2}} = 254.0 \text{MHz}$

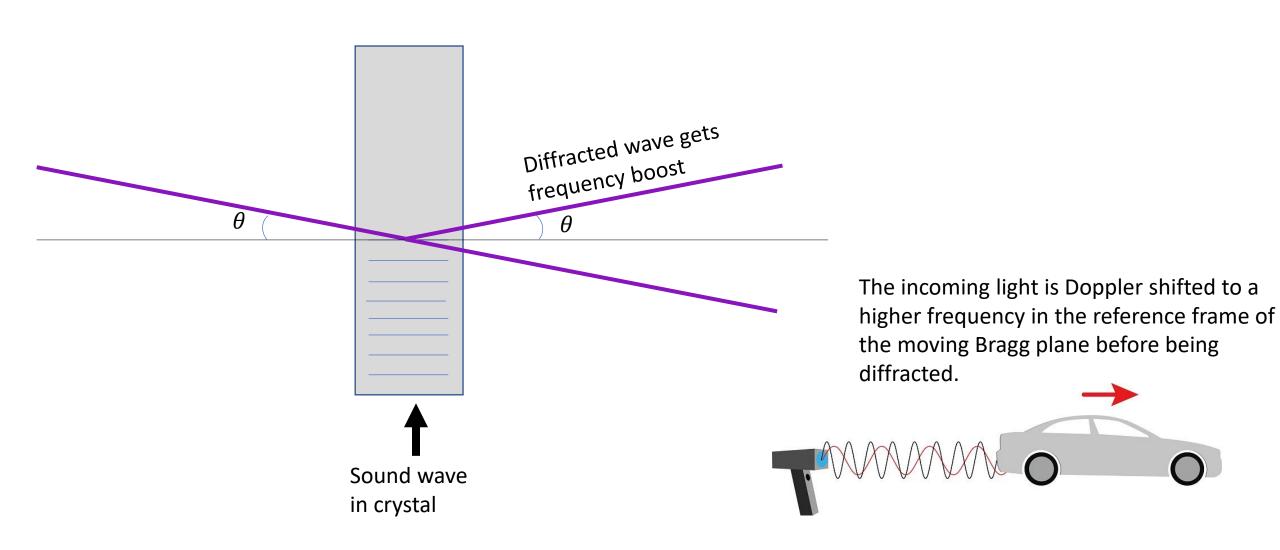
How do we lock the trap to K39 spectra but scan K37?

Isotope frequency shift is corrected by Acousto-Optical Modulator

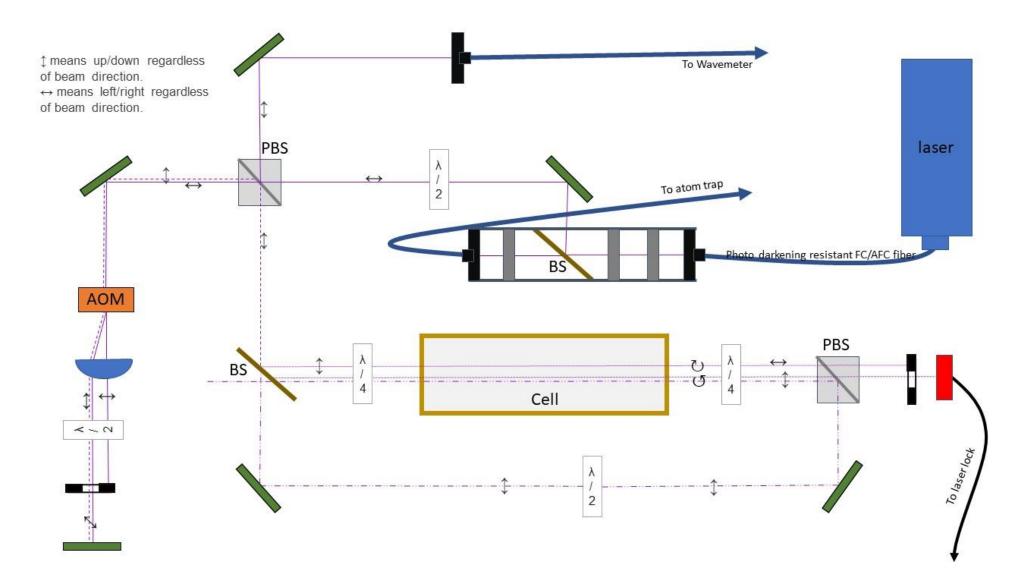
Picture of aom module



Acousto-Optical Modulator physics



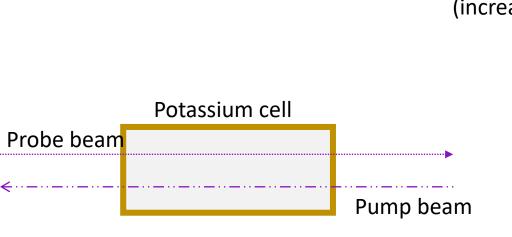
Beam polarization

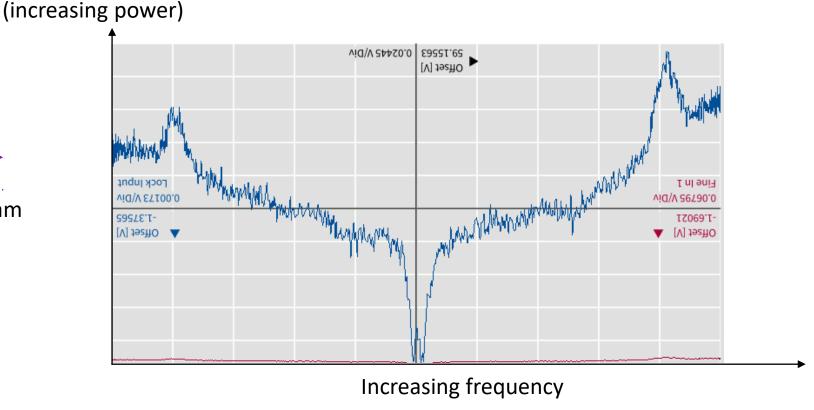


K cell sat spectroscopy

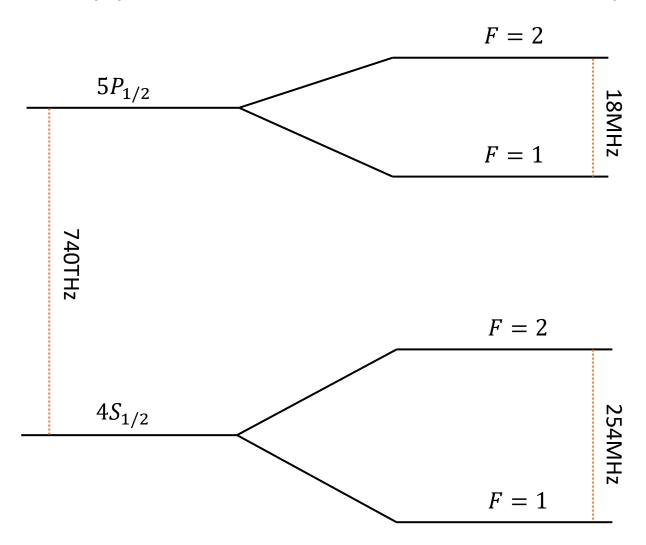
We know the cell will absorb light at transitions frequencies, how do we arrive at the actual absorption signal?

decreasing absorption





Hyperfine structure frequencies for K39

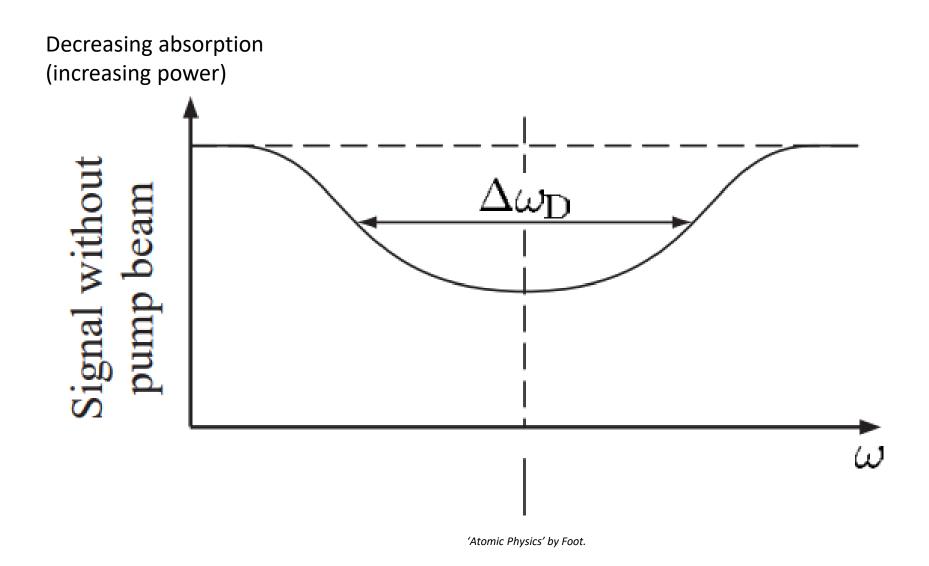


Hyperfine transition $ 4S_{\frac{1}{2}}F_I\rangle$ to $ 5P_{\frac{1}{2}}F_{II}\rangle$	Shift from 24701.382 <i>cm</i> ⁻¹
$F_I = 1$ $F_{II} = 2$	295MHz
$F_I = 1$ $F_{II} = 1$	277MHz
$F_I = 2$ $F_{II} = 2$	-166MHz
$F_I = 2$ $F_{II} = 1$	-184MHz

$$H_{HFS} = A \, \mathbf{I \cdot J}$$

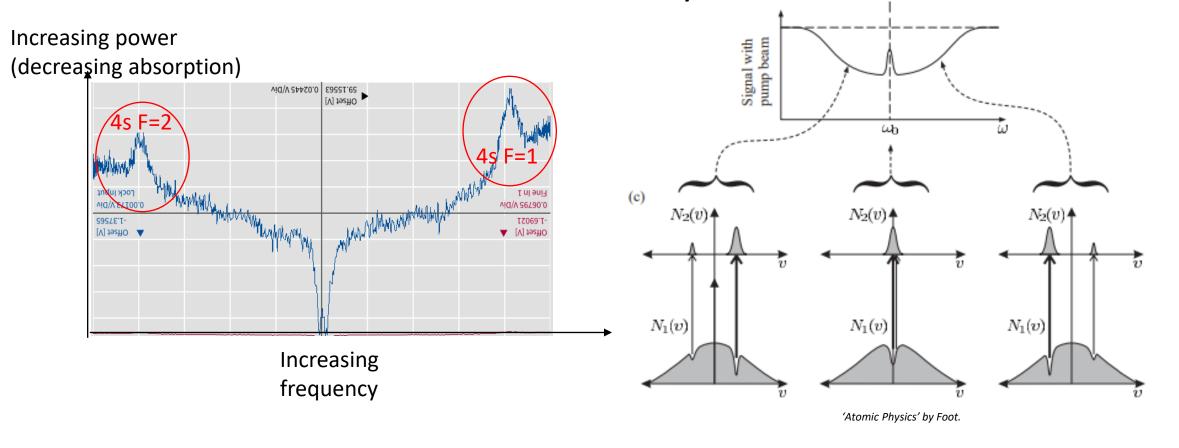
= 126.9MHz
= 8.99MHz

Doppler shift broadened absorption



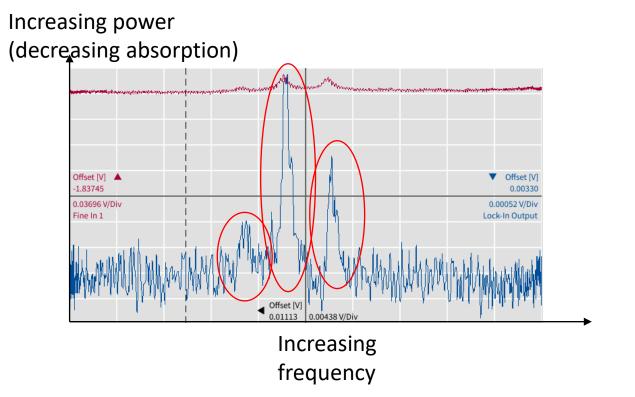
Pump beam hole burning

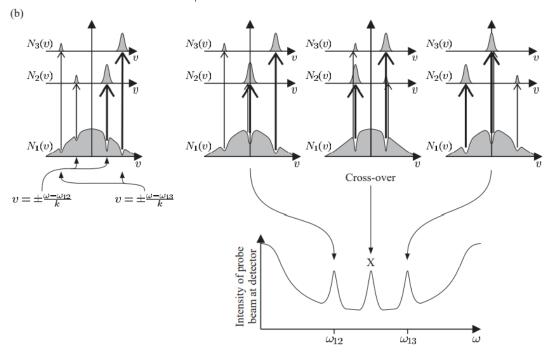
Now that our absorptions are Doppler smeared, we use the pump beam to select atoms with zero velocity.



Crossover effect

Every pair of peaks has a composite peak halfway in between from the doppler shifts of the two side peaks.

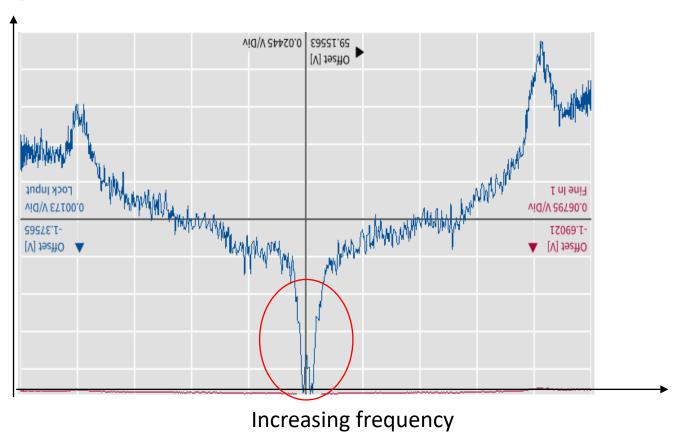




'Atomic Physics' by Foot.

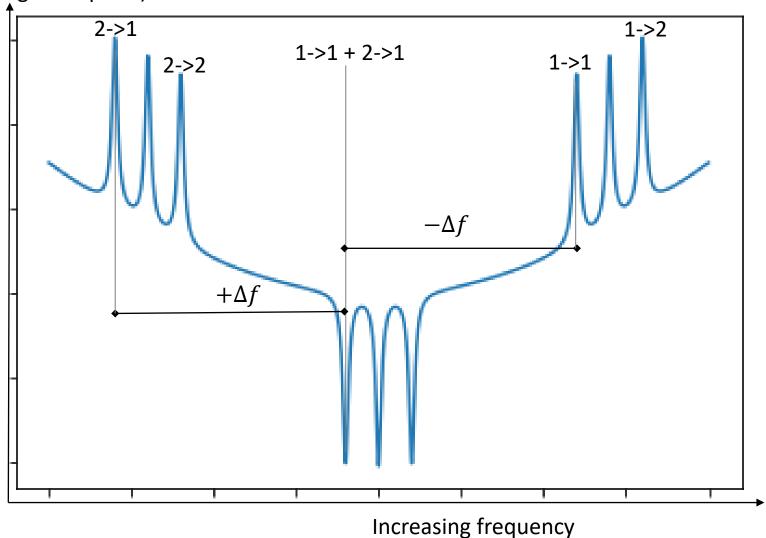
Enhanced absorption

Increasing power (decreasing absorption)



Enhanced absorption

Increasing power (decreasing absorption)

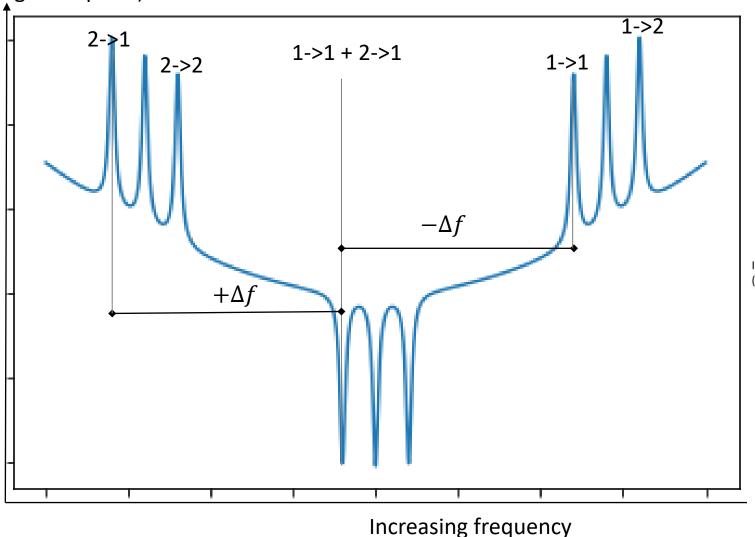


Because the two Doppler crossed frequencies start from different hyperfine ground states the pump does not steal from the probe.

Instead, the pump beam states randomly decay towards the opposing hyperfine, increasing the population the probe sees.

Enhanced absorption

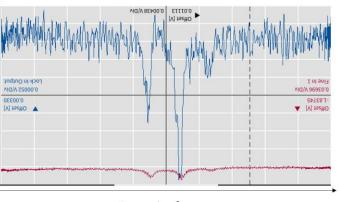
Increasing power (decreasing absorption)



Atoms moving towards the probe beam absorb the pump beam in the 1->1 transition, which randomly decays, populating the F=2 ground state for the probe beam which is absorbed in the 2->1 transition.

Atoms moving away from the probe beam absorb the pump beam in the 2->1 transition, which randomly decays, populating the F=1 ground state for the probe beam which is absorbed in the 1->2 transition.

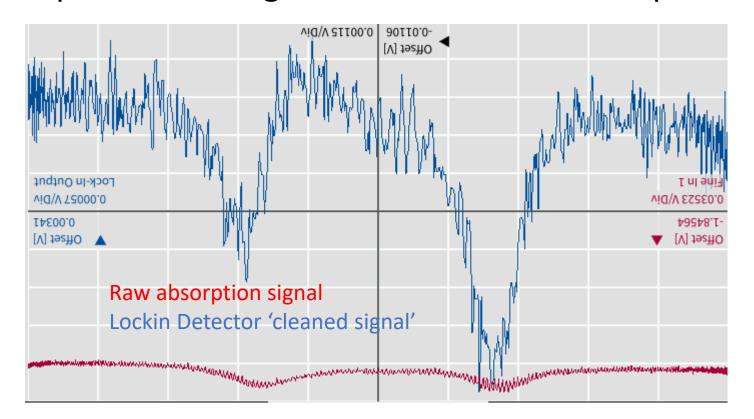
Increasing power (decreasing absorption)

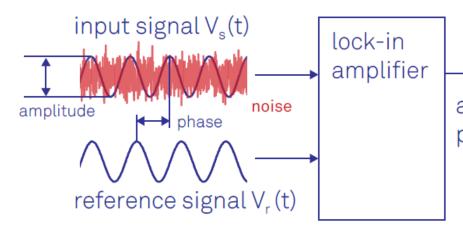


Increasing frequency

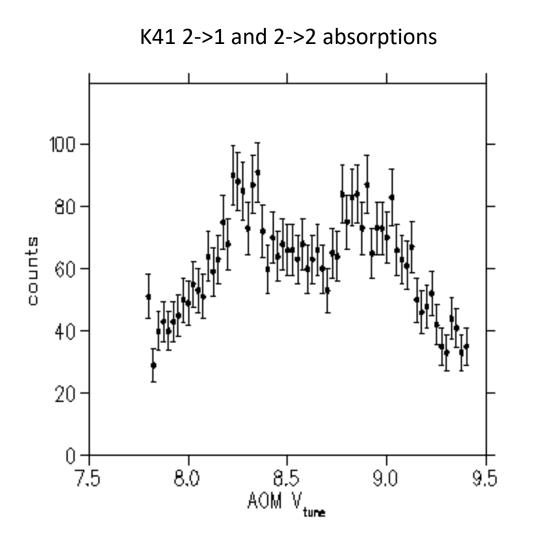
Locking the laser frequency using absorption peaks

Using a solenoid and the Zeeman effect to produce slight high-frequency energy oscillations in the cell, correlating absorption peaks to our magnetic scan we can produce a cleaner signal and lock to the peak with a Proportional Integral Derivative control loop.

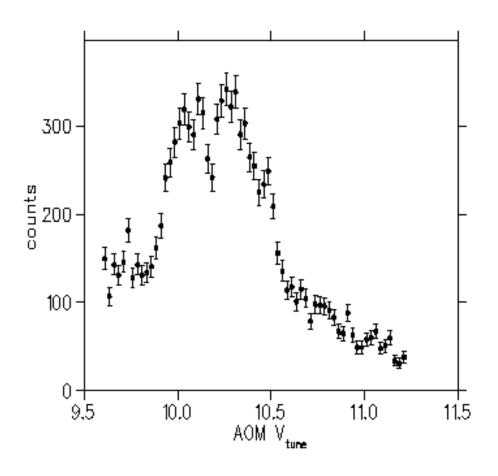




Very early data



K41 1->1 and 1->2 absorptions



Thank you!

