

Parity breaking correlation functions and fifth forces

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I. INTRODUCTION

Astrophysical and cosmological tests of gravity have been done before. Here we focus on a test of screening using clusters of galaxies. Previously various probes have been used such as

- X-ray measurements using hot intra-clusters gas dynamics
- Strong and Weak lensing

Screening can be accomplished in the following ways

- Weak coupling between the field and matter in regions of high density, thus inducing a weak fifth force as in symmetron theories.
- The field can acquire a large mass in high density environments, being short-ranged and undetectable and be light and long ranged in lower density regions, as for chameleon fields and $f(R)$ models.
- Finally one may change the kinetic contribution of the field to the Lagrangian with first and second order derivatives being important in a certain range, as happens in Vainshtein screening.

II. SYMMETRON

A. Intro

Here we summarise the results from [1, 2]. The scalar equation of motion for a symmetron is

$$\nabla^2\phi - \partial_\phi V + A^3(\phi)\partial_\phi\tilde{T} = 0 \quad (1)$$

where $A^2(\phi)$ is the conformal factor that relates the Jordan frame metric $\tilde{g}_{\mu\nu}$ to the Einstein frame metric

$$\tilde{g}_{\mu\nu} = A^2(\phi)g_{\mu\nu} \quad (2)$$

and the stress energy tensor is defined as the trace of the Jordan frame metric,

$$\begin{aligned} \tilde{T}_{\mu\nu} &= -\frac{2}{\sqrt{-\tilde{g}}}\frac{\delta L_m}{\delta g^{\mu\nu}} \\ \tilde{T} &= \tilde{T}_{\mu\nu}\tilde{g}^{\mu\nu} \end{aligned} \quad (3)$$

If we ignore backreaction of the scalar field on gravity, and consider spherically symmetric pressureless sources, the density $\rho = A^3\tilde{\rho}$ is conserved in the Einstein frame. The Laplacian operator also simplifies in this case to give

$$\nabla^2\phi = \frac{d^2\phi}{dr^2} + \frac{2}{r}\frac{d\phi}{dr} \quad (4)$$

For cases of roughly homogenous ρ , the field evolves according to an effective potential

$$V_{eff}(\phi) = V(\phi) + \rho A(\rho) \quad (5)$$

For the model which has a spontaneously broken \mathbb{Z}_2 .

B. Equations of motion

The general equation of motion

$$\partial_r^2 \phi + \frac{2}{r} \partial_r \phi = \partial_\phi V(\phi) + \partial_\phi A(\phi) \rho \quad (6)$$

For homogenous ρ such as the NFW profile, the field evolves according to

$$\partial_r^2 \phi + \frac{2}{r} \partial_r \phi = \partial_\phi V_{eff} \quad (7)$$

where

$$V_{eff} = \frac{1}{2} \left(\frac{\rho}{M^2} - \mu^2 \right) \phi^2 + \frac{1}{4} \lambda \phi^4 \quad (8)$$

We can separate the solution in terms of the solution inside and outside the object. We can start with separating the effective potential for top hat density profile

$$\rho = \begin{cases} \rho_0 & (r < R) \\ 0 & (r > R) \end{cases} \quad (9)$$

$$V_{eff} = \begin{cases} \frac{\rho \phi^2}{2M^2} & (r < R) \\ m_0^2 (\phi^2 - \phi_0^2)^2 / 2 & (r > R) \end{cases} \quad (10)$$

The solution for the scalar field can also be split in this way

$$\phi(r) = \begin{cases} C \frac{R}{r} \sinh \left(\frac{\sqrt{\rho} r}{M} \right) & (r < R) \\ D \frac{R}{r} e^{-m_0(r-R)} + \phi_0 & (r > R) \end{cases} \quad (11)$$

Coefficients C and D are fixed by matching the field and its radial derivative at the interface $r = R$

$$\begin{aligned} C &= \phi_0 \sqrt{\frac{\Delta R}{R}} \operatorname{sech} \left(\sqrt{\frac{R}{\Delta R}} \right) \\ D &= -\phi_0 \left(1 - \sqrt{\frac{\Delta R}{R}} \tanh \left(\sqrt{\frac{\Delta R}{R}} \right) \right) \end{aligned} \quad (12)$$

where $\frac{\Delta R}{R} \equiv \frac{M^2}{\rho R^2} = \frac{M^2}{6M_{pl}^2 \phi} = \frac{\phi_0}{6gM_{pl}\Phi}$ is the thin shell radius in analogy with chameleon screening. $\Phi \equiv \frac{\rho R^2}{6M_{pl}^2}$ is the gravitational potential of the source.

Note that so far we have not said anything about the screening environment. We are interested in the dark matter halos so we can use galaxy cluster measurements to test for screening. So we can assume $R = R_{vir}$ in this case and $M = M_{300}$ (or some definition of the halo mass).

C. Ratio of forces

The fifth force for the symmetron is

$$F_\phi = \frac{\phi \nabla \phi}{M_s^2} = -\frac{1}{M_s^2} \left(\frac{BR}{r} e^{-\sqrt{2}\mu r} + \phi_0 \right) \left(\frac{BR}{r} e^{-\sqrt{2}\mu r} \left(\frac{1}{r} + \sqrt{2}\mu \right) \right) \quad (13)$$

III. PARITY BREAKING CORRELATION FUNCTIONS

The overdensity of bright galaxies can be written as

$$\Delta_B(z, \hat{n}) = \Delta_B^{st}(z, \hat{n}) + \Delta_B^{rel}(z, \hat{n}) + \Delta_B^{lens}(z, \hat{n}) + \Delta_B^{AP}(z, \hat{n}) \quad (14)$$

where

$$\Delta_B^{st}(z, \hat{n}) = b_B \delta(z, \hat{n}) - \frac{1}{\mathcal{H}} \partial_r (\mathbf{v} \cdot \hat{n}) \quad (15)$$

$$\Delta_B^{rel}(z, \hat{n}) = \frac{1}{\mathcal{H}} (\partial_r \psi + \dot{\mathbf{v}} \cdot \hat{n}) - \left(\frac{\dot{\mathcal{H}}}{\mathcal{H}^2} + \frac{2}{r\mathcal{H}} - 1 + 5s_B \left(1 - \frac{1}{r\mathcal{H}} \right) \right) \mathbf{v} \cdot \hat{n} \quad (16)$$

$$\Delta_B^{lens}(z, \hat{n}) = (5s_B - 2) \int_0^r dr' \frac{(r - r')r'}{2r} \nabla_{\perp}^2(\phi + \psi) \quad (17)$$

$$\Delta_B^{AP}(z, \hat{n}) = (\partial_r - \partial_{\eta})(\Delta_B^{st} + \Delta_B^{rel} + \Delta_B^{lens}) \frac{\partial r(z, \Omega)}{\partial \Omega} \delta \Omega \quad (18)$$

The quantities with a B index denotes the value of a variable that will change for bright and faint galaxies. b_B is the bias of the bright galaxies and s_B is their effective number count slope. These expressions are valid in any conformal theory of gravity (where null geodesics are unchanged). Thus are perfectly suite for our purposes as all screened theories are conformally invariant. When timelike curves are also geodesics, we can use the Euler equation

$$\dot{\mathbf{v}} \cdot \hat{n} + \mathcal{H} \mathbf{v} \cdot \hat{n} + \partial_r \psi = 0 \quad (19)$$

Using this equation we can rewrite Eq (16) as

$$\Delta_B^{rel}(z, \hat{n}) = -\mathbf{v} \cdot \hat{n} \left(\frac{\dot{\mathcal{H}}}{\mathcal{H}^2} + \frac{2}{r\mathcal{H}} + 5s_B \left(1 - \frac{1}{r\mathcal{H}} \right) \right) \quad (20)$$

In presence of screened theories of gravity we include a new force which is sourced by the gravitational potential

$$\nabla \phi = -\alpha \frac{GM}{r} \quad (21)$$

Thus we can write a general parameterized form of the effect of screened fifth forces

$$\partial_r \psi \rightarrow (1 + \alpha) \partial_r \psi \quad (22)$$

Thus we can write the modified Euler equation as

$$\dot{\mathbf{v}} \cdot \hat{n} + \mathcal{H} \mathbf{v} \cdot \hat{n} + (1 + \alpha) \partial_r \psi = 0 \quad (23)$$

Under this modification to the Euler equation we can rewrite the Eq (16) as

$$\begin{aligned} \Delta_B^{rel(MG)}(z, \hat{n}) &\equiv \Delta_B^{MG}(z, \hat{n}) + \Delta_B^{rel}(z, \hat{n}) \\ \Delta_B^{MG}(z, \hat{n}) &\equiv \left(1 - \frac{1}{1 + \alpha^2} \right) \left(\frac{\dot{\mathbf{v}} \cdot \hat{n}}{\mathcal{H}} + \mathbf{v} \cdot \hat{n} \right) \\ \Delta_B^{rel}(z, \hat{n}) &\equiv - \left(\frac{\mathcal{H}}{\mathcal{H}^2} + \frac{2}{r\mathcal{H}} + 5s_B \left(1 - \frac{1}{r\mathcal{H}} \right) \right) \vec{v} \cdot \hat{n} \end{aligned} \quad (24)$$

Using these we can write down correlation functions as normal. The only terms that will change will be terms that previously dependent on the relativistic part.

Lets start with writing down the correlation function in normal GR.

$$\begin{aligned}\xi^{rel}(r, r', \theta) &= \langle \Delta_B^{st}(z, \hat{n}) \Delta_F^{rel}(z', \hat{n}') \rangle \\ &= A \int \frac{d^3 k}{(2\pi)^2} e^{i\vec{k}(\vec{x}' - \vec{x})} \frac{(k\eta_0)^{n_s-1}}{k^3} \left[\left[\frac{\dot{\mathcal{H}}(r')}{\mathcal{H}^2(r')} + \frac{2}{r'\mathcal{H}(r')} + 5s_B(r') \left(1 - \frac{1}{r'\mathcal{H}(r')} \right) \right] i(\hat{\mathbf{k}} \cdot \hat{\mathbf{n}} T_V(k, r')) \right. \\ &\quad \left. \left(b_B T_D(k, r) - \frac{k}{\mathcal{H}(r)} (\hat{\mathbf{k}} \cdot \hat{\mathbf{n}})^2 T_V(k, r) \right) - (r \leftrightarrow r') \right]\end{aligned}\quad (25)$$

Now we can look at the correlation function for modified gravity.

$$\begin{aligned}\langle \Delta_B^{st}(z, \hat{n}) \Delta_F^{rel(MG)}(z', \hat{n}') \rangle &= A \int \frac{dk}{2\pi} \frac{d\Omega_k}{2\pi} e^{i\vec{k}(\vec{x}' - \vec{x})} \frac{(k\eta_0)^{n_s-1}}{k^2} \left[\left(1 - \frac{1}{1+\alpha^2} \right) \left((i\vec{k} \cdot \hat{n}) \left(\frac{\dot{T}_v(k, r)}{\mathcal{H}} + T_v(k, r) \right) \right) \right] \\ &\quad \times \left[\left(b_B T_D(k', r') - \frac{k'}{\mathcal{H}} ((\vec{k}' \cdot \hat{n}')^2 T_v(k', r')) \right) \right] \\ &= \frac{A}{2\pi} \left(1 - \frac{1}{1+\alpha^2} \right) \int d\Omega_k e^{i\vec{k}(\vec{x}' - \vec{x})} (i\vec{k} \cdot \hat{n}) \frac{(k\eta_0)^{n_s-1}}{k} \frac{b_B}{\mathcal{H}} \dot{T}_v(k, r) T_D(k', r') \\ &\quad - \frac{A}{2\pi} \left(1 - \frac{1}{1+\alpha^2} \right) \int d\Omega_k e^{i\vec{k}(\vec{x}' - \vec{x})} (i\vec{k} \cdot \hat{n}) \frac{(k\eta_0)^{n_s-1}}{k} (\vec{k}' \cdot \hat{n}')^2 \frac{k'}{\mathcal{H}^2} \dot{T}_v(k, r) T_v(k', r') \\ &\quad + \frac{A}{2\pi} \left(1 - \frac{1}{1+\alpha^2} \right) \int d\Omega_k e^{i\vec{k}(\vec{x}' - \vec{x})} (i\vec{k} \cdot \hat{n}) \frac{(k\eta_0)^{n_s-1}}{k} b_B(r') T_v(k, r) T_D(k', r') \\ &\quad - \frac{A}{2\pi} \left(1 - \frac{1}{1+\alpha^2} \right) \int d\Omega_k e^{i\vec{k}(\vec{x}' - \vec{x})} (i\vec{k} \cdot \hat{n}) \frac{(k\eta_0)^{n_s-1}}{k} (\vec{k}' \cdot \hat{n}')^2 \frac{k'}{\mathcal{H}^2} T_v(k, r) T_v(k', r') \\ &= \frac{A}{2\pi} \left(1 - \frac{1}{1+\alpha^2} \right) \int d\Omega_k e^{i\vec{k}(\vec{x}' - \vec{x})} (i\vec{k} \cdot \hat{n}) \frac{(k\eta_0)^{n_s-1}}{k} \underbrace{\left[\frac{b_B}{\mathcal{H}} \dot{T}_v(k, r) T_D(k', r') \right]}_{T1} \\ &\quad - \underbrace{(\vec{k}' \cdot \hat{n}')^2 \frac{k'}{\mathcal{H}^2} \dot{T}_v(k, r) T_v(k', r')}_{T2} + \underbrace{b_B(r') T_v(k, r) T_D(k', r')}_{T3} - \underbrace{(\vec{k}' \cdot \vec{n}')^2 \frac{k'}{\mathcal{H}^2} T_v(k, r) T_v(k', r')}_{T4}\end{aligned}\quad (26)$$

Lets focus on one term first, $T1$ and do the angular integral first

$$\begin{aligned}T1 &= \int d\Omega_k e^{i\vec{k}(\vec{x}' - \vec{x})} (i\vec{k} \cdot \hat{n}) \frac{(k\eta_0)^{n_s-1}}{k} \frac{b_B}{\mathcal{H}} \dot{T}_v(k, r) T_D(k', r') \\ &= \frac{(4\pi)^2}{3} \int d\Omega_k \sum_{LM} \sum_{m=-1}^1 i^{L+1} j_L(kd) Y_{LM}^*(\vec{k}) Y_{lm}(\hat{N}) Y_{1m}^*(\hat{k}) Y_{1m}(\hat{n}) \frac{(k\eta_0)^{n_s-1}}{k} \frac{b_B}{\mathcal{H}} \dot{T}_v(k, r) T_D(k', r') \\ &= \frac{(4\pi)^2}{3} \sum_{M,m=-1}^1 (-1) \delta_{Mm} Y_{1M}(\hat{N}) Y_{1m}(\hat{n}) \frac{(k\eta_0)^{n_s-1}}{k} \frac{b_B}{\mathcal{H}} j_i(kd) \dot{T}_v(k, r) T_d(k', r')\end{aligned}\quad (27)$$

Similarly $T3$ is

$$T3 = \frac{(4\pi)^2}{3} \sum_{M,m=-1}^1 (-1) \delta_{Mm} Y_{1M}(\hat{N}) Y_{1m}(\hat{n}) \frac{(k\eta_0)^{n_s-1}}{k} b_B j_i(kd) T_v(k, r) T_d(k', r')\quad (28)$$

Next lets compute $T2$.

$$\begin{aligned}
T3 &= \int d\Omega_k e^{i\vec{k}(\vec{x}' - \vec{x})} (i\vec{k} \cdot \hat{n}) (\vec{k} \cdot \hat{n})^2 \frac{(k\eta_0)^{n_s-1}}{k^2} \frac{k}{\mathcal{H}^2} \dot{T}_v(k, r) T_v(k, r') \\
&= \int d\Omega_k \left(4\pi \sum_{LM} i^{L+1} j_L(kd) Y_{LM}^*(\hat{k}) Y_{LM}(\hat{N}) \right) \left(\frac{4\pi}{3} \sum_{m=-1}^1 Y_{1m}^*(\hat{k}) Y_{1m}(\hat{n}) \right) \\
&\quad \times \left(\frac{8\pi}{15} \sum_{m'=-2}^2 Y_{2m'}^*(\hat{n}') Y_{2m'}(\hat{k}) + \frac{1}{3} \right) \left(\frac{(k\eta_0)^{n_s-1}}{k^2} \frac{k}{\mathcal{H}^2} \dot{T}_v(k, r) T_v(k, r) \right)
\end{aligned} \tag{29}$$

Before proceeding we note that there are two types of terms in this expression. First there are terms that are quadratic in Y_{lm} 's and second, terms that are cubic in Y_{lm} 's.

$$\begin{aligned}
T3 &= \frac{(4\pi)^2}{3} \int d\Omega_k \sum_{LM} \sum_{m=-1}^1 \sum_{m'=-2}^2 i^{L+1} j_L(kd) \left(\frac{8\pi}{15} Y_{LM}^*(\hat{k}) Y_{LM}(\hat{N}) Y_{1m}^*(\hat{k}) Y_{1m}(\hat{n}) Y_{2m'}^*(\hat{n}') Y_{2m'}(\hat{k}) \right. \\
&\quad \left. + \frac{1}{3} Y_{LM}^*(\hat{k}) Y_{LM}(\hat{N}) Y_{1m}^*(\hat{k}) Y_{1m}(\hat{n}) \right) \left(\frac{(k\eta_0)^{n_s-1}}{k^2} \frac{k}{\mathcal{H}^2} \dot{T}_v(k, r) T_v(k, r) \right) \\
&= \left(\frac{4\pi}{3} \right)^2 \left(\sum_{LM} \sum_{m=-1}^1 \sum_{m'=-2}^2 (\delta_{iL} \delta_{Mm}) \right) \left(i^{L+1} j_L(kd) \frac{(k\eta_0)^{n_s-1}}{k^2} \frac{k}{\mathcal{H}^2} \dot{T}_v(k, r) T_v(k, r) \right) \\
&\quad + \frac{(4\pi)^2 8\pi}{45} \left(\sum_{LM} \sum_{m=-1}^1 \sum_{m'=-2}^2 \sqrt{\frac{(2L+1)15}{4\pi}} \begin{pmatrix} L & 2 & 1 \\ -M & -m & -m' \\ 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} L & 2 & 1 \\ 0 & 0 & 0 \end{pmatrix} \right. \\
&\quad \left. \times \left(i^{L+1} j_L(kd) \frac{(k\eta_0)^{n_s-1}}{k^2} \frac{k}{\mathcal{H}^2} \dot{T}_v(k, r) T_v(k, r) \right) \right)
\end{aligned} \tag{30}$$

Lets evaluate the Wigner 3j's explicitly. First we list some of their properties that we will need. For Wigner symbol of the form $\begin{pmatrix} j_1 & j_2 & j_3 \\ m_1 & m_2 & m_3 \end{pmatrix}$ we have the following properties. The Wigner symbols are zero unless

$$\begin{aligned}
P1 : m_i &\in \{-j_i, -j_i + 1, \dots, j_i\} \\
P2 : m_1 + m_2 + m_3 &= 0 \\
P3 : |j_1 - j_2| &\leq j_3 \leq j_1 + j_2 \\
P4 : (j_1 + j_2 + j_3) &\text{ is an integer}
\end{aligned} \tag{31}$$

Using $P3$ we know that $L = 2, 3$.

$$\begin{aligned}
T_\Omega &\equiv \left(\sum_{LM} \sum_{m=-1}^1 \sum_{m'=-2}^2 \sqrt{\frac{(2L+1)15}{4\pi}} \begin{pmatrix} L & 2 & 1 \\ -M & -m & -m' \\ 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} L & 2 & 1 \\ 0 & 0 & 0 \end{pmatrix} \right) \\
&= \sum_{m=-1}^1 \sum_{m'=-2}^2 \left(\sqrt{\frac{75}{4\pi}} \sum_{M=-2}^2 \begin{pmatrix} 2 & 2 & 1 \\ -M & -m & -m' \\ 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} 2 & 2 & 1 \\ 0 & 0 & 0 \end{pmatrix} + \sum_{M=-3}^3 \sqrt{\frac{105}{4\pi}} \begin{pmatrix} 3 & 2 & 1 \\ -M & -m & -m' \\ 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} 3 & 2 & 1 \\ 0 & 0 & 0 \end{pmatrix} \right)
\end{aligned} \tag{32}$$

Now we note that

$$\begin{aligned}
\begin{pmatrix} 2 & 2 & 1 \\ 0 & 0 & 0 \end{pmatrix} &= 0 \\
\begin{pmatrix} 3 & 2 & 1 \\ 0 & 0 & 0 \end{pmatrix} &= -\sqrt{\frac{3}{35}}
\end{aligned} \tag{33}$$

Further using $P2$ we can reduce the sum over M, m, m' to a sum over m', m as $M = -(m + m')$. i.e we can write M in the Wigner symbol as $\delta_{M, -(m+m')}$. Thus we are left with

$$\begin{aligned}
T_\Omega &= -\sqrt{\frac{9}{4\pi}} \sum_{m=-1}^1 \sum_{m'=-2}^2 \left(\delta_{M,-(m+m')} \begin{pmatrix} 3 & 2 & 1 \\ -m & -m' \end{pmatrix} \right) \\
&= -\sqrt{\frac{9}{4\pi}} \sum_{m'=-2}^2 \left(\underbrace{\begin{pmatrix} 3 & 2 & 1 \\ \delta_{M,-(m+m')} & 1 & -m \end{pmatrix}}_{W1} + \underbrace{\begin{pmatrix} 3 & 2 & 1 \\ \delta_{M,-(m+m')} & 0 & -m \end{pmatrix}}_{W2} + \underbrace{\begin{pmatrix} 3 & 2 & 1 \\ \delta_{M,-(m+m')} & -1 & -m \end{pmatrix}}_{W3} \right)
\end{aligned} \tag{34}$$

We can compute each of these terms individually,

$$\begin{aligned}
W1 &= -\sqrt{\frac{9}{4\pi}} \left(\underbrace{\begin{pmatrix} 3 & 2 & 1 \\ 1 & 1 & -2 \end{pmatrix}}_0 + \underbrace{\begin{pmatrix} 3 & 2 & 1 \\ 0 & 1 & -1 \end{pmatrix}}_{-\sqrt{\frac{1}{35}}} + \underbrace{\begin{pmatrix} 3 & 2 & 1 \\ -1 & 1 & 0 \end{pmatrix}}_{2\sqrt{\frac{2}{105}}} + \underbrace{\begin{pmatrix} 3 & 2 & 1 \\ -2 & 1 & 1 \end{pmatrix}}_{-\sqrt{\frac{2}{21}}} + \underbrace{\begin{pmatrix} 3 & 2 & 1 \\ -3 & 1 & 2 \end{pmatrix}}_0 \right) \\
&= \left(\sqrt{\frac{9}{4\pi}} \right) \left(\frac{-1 + 2\sqrt{\frac{2}{3}} - \sqrt{\frac{10}{3}}}{\sqrt{35}} \right)
\end{aligned} \tag{35}$$

$$\begin{aligned}
W2 &= -\sqrt{\frac{9}{4\pi}} \underbrace{\begin{pmatrix} 3 & 2 & 1 \\ 2 & 0 & -2 \end{pmatrix}}_0 + \underbrace{\begin{pmatrix} 3 & 2 & 1 \\ 1 & 0 & -1 \end{pmatrix}}_{-\sqrt{\frac{2}{21}}} + \underbrace{\begin{pmatrix} 3 & 2 & 1 \\ 0 & 0 & 0 \end{pmatrix}}_{2\sqrt{\frac{2}{105}}} + \underbrace{\begin{pmatrix} 3 & 2 & 1 \\ -1 & 0 & 1 \end{pmatrix}}_{-\sqrt{\frac{1}{35}}} + \underbrace{\begin{pmatrix} 3 & 2 & 1 \\ -2 & 0 & 2 \end{pmatrix}}_0 \\
&= -\sqrt{\frac{9}{4\pi}} \sqrt{\frac{1}{35}}
\end{aligned} \tag{36}$$

$W3$ is the same as $W1$. Thus T_Ω is

$$T_\Omega = -\sqrt{\frac{9}{4\pi}} \left(\frac{-2 + 2\sqrt{2/3} - \sqrt{10/3}}{\sqrt{35}} \right) \tag{37}$$

A. Fourier space transformations

Here we write down all the ingredients we need and their Fourier space transformation

$$\begin{aligned}
\partial_r(\vec{v} \cdot \hat{n}) &= -k(\hat{k} \cdot \hat{n})^2 V(k, \eta) \sim -k(\hat{k} \cdot \hat{n})^2 T_V(k, r) \\
\vec{v} \cdot \hat{n} &= i(\hat{k} \cdot \hat{n}) V(k, \eta) \sim i(\hat{k} \cdot \hat{n}) T_V(k, r) \\
\dot{\vec{v}} \cdot \hat{n} &= i(\hat{k} \cdot \hat{n}) \dot{V}(k, \eta) \sim i(\hat{k} \cdot \hat{n}) \dot{T}_V(k, r)
\end{aligned} \tag{38}$$

IV. SHOOTING METHOD

Our differential equation is

$$\partial_r^2 \phi + \frac{2}{r} \partial_r \phi = \left(\frac{\rho}{M_s^2} - \mu^2 \right) \phi + \lambda \phi^3 \tag{39}$$

The BC's are

$$\begin{aligned}\partial_r \phi(0) &= 0 \\ \phi(\infty) &= 1\end{aligned}\tag{40}$$

We can convert this second order DE into two first order DE's

$$\begin{aligned}y_1 &= \phi \\ y_2 &= \partial_r \phi\end{aligned}\tag{41}$$

Thus the DE is

$$y'_2 + \frac{2y_2}{r} = \left(\frac{\rho}{M_s^2} - \mu^2 \right) y_1 + \lambda y_1^3\tag{42}$$

V. CHECKS IN CLUSTERS

If all galaxies in a cluster are screened we have nothing to measure. So we must check that out to the screening radius we see galaxies around clusters.

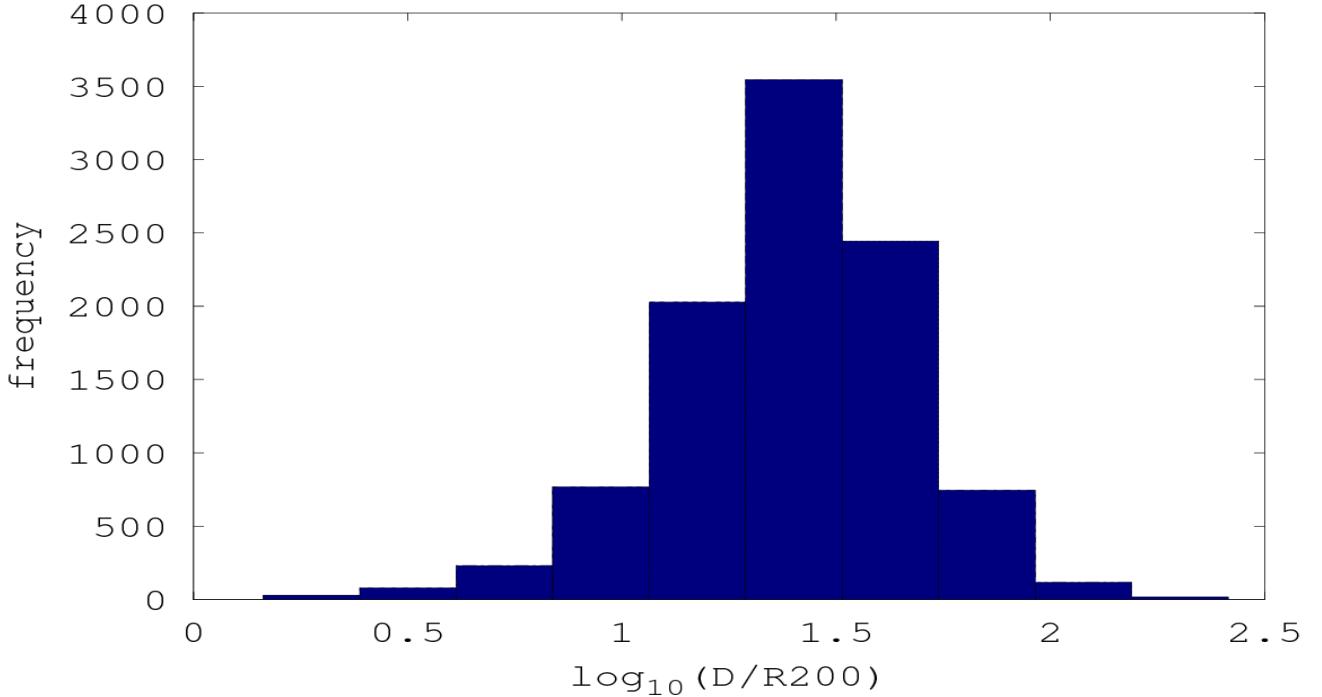


FIG. 1. SDSS

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- [1] K. Hinterbichler and J. Khoury, Phys. Rev. Lett. **104**, 231301 (2010), arXiv:1001.4525 [hep-th].
[2] J. Clampitt, B. Jain, and J. Khoury, JCAP **1201**, 030 (2012), arXiv:1110.2177 [astro-ph.CO].

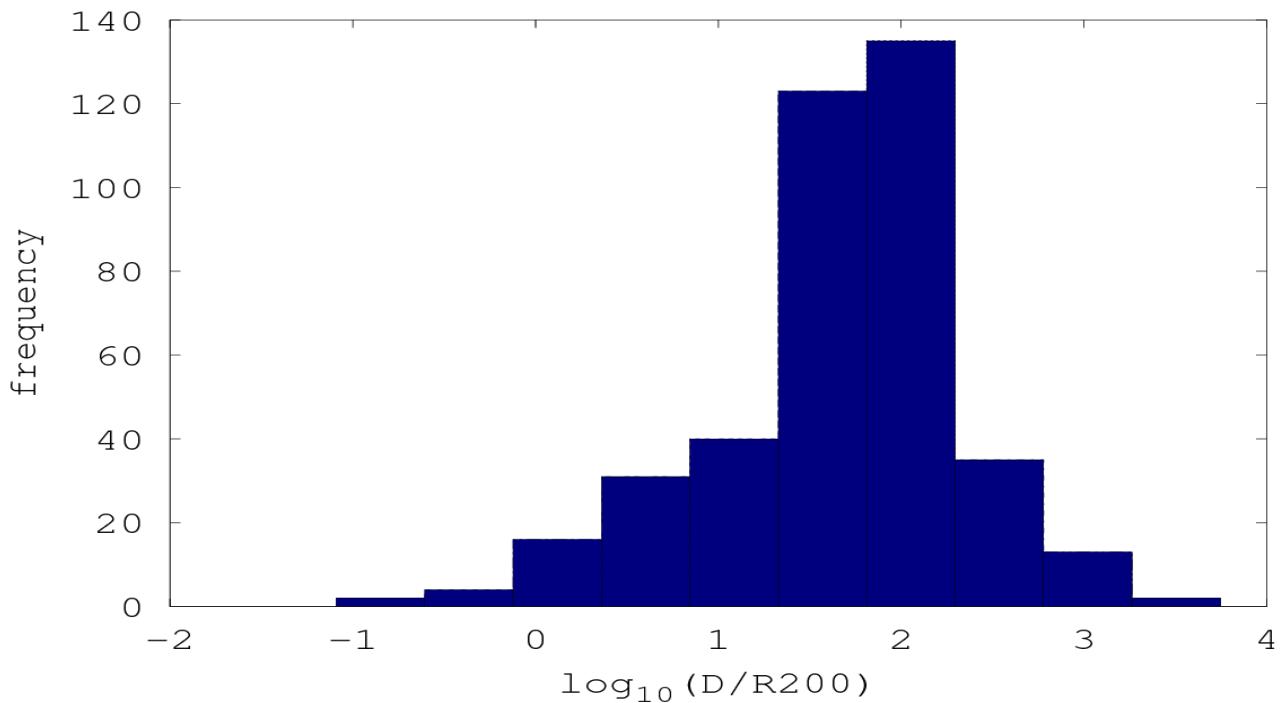


FIG. 2. XMM