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ECE 513:

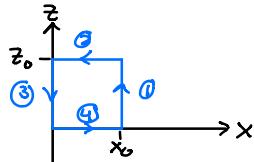
Faraday's Law states:

$$\oint \vec{E} \cdot d\vec{l} = - \frac{\partial}{\partial t} (\Phi_E)$$

Given an electric field:

$$\vec{E} = E_{ox} \cos(K_z z - \omega t) \hat{x}$$

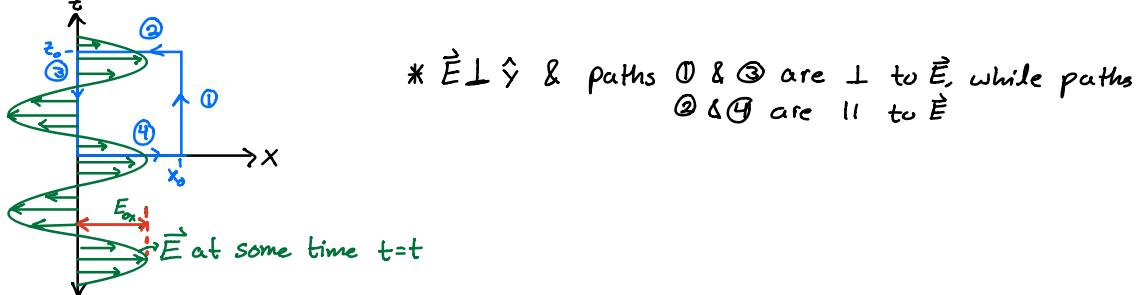
For rectangle 'e' given, let the path of integration in Faraday's Law be defined as:



$E$ -field is a wave, given by the function satisfying the wave equation:

\* Any function as  $f(\vec{K} \cdot \vec{r} - \omega t) \hat{n}$  satisfies the wave-equation. Where  $\vec{K}$  is direction of propagation of the wave, and  $\hat{n}$  is the polarization vector (plane of vibration).

$\therefore \vec{E}$  is a plane wave oscillating along  $x$ -axis & propagating along  $z$ -axis



as  $E$  propagates.  $|E| \in [-E_{ox}, E_{ox}]$

at  $z=0$  &  $z=z_0$

$$\vec{E}(0, t) = E_{ox} \cos(\omega t) \hat{x} \quad * \cos(-\omega t) = \cos(\omega t)$$

$$E(z_0, t) = E_{ox} \cos(z_0 K_z - \omega t) \hat{x}$$

Faraday's Law left hand side becomes:

$$\begin{aligned}
 \oint \vec{E} \cdot d\vec{l} &= \int_2 E(z_0, t) dx + \int_4 E(0, t) dx \quad * \text{ where } E(z, t) \text{ is independent of } x \\
 &= -X_0 E(z_0, t) + X_0 E(0, t) \\
 &= X_0 (E_{ox} \cos(\omega t) - E_{oz} \cos(\omega t - K_z z_0)) \quad * \text{ using Euler's formula:} \\
 &= X_0 E_{ox} \left( \frac{1}{2} [e^{j\omega t} + e^{-j\omega t} - e^{jK_z z_0} e^{-j\omega t} - e^{-jK_z z_0} e^{j\omega t}] \right) \quad \cos(\theta) = \frac{1}{2}(e^{j\theta} + e^{-j\theta}) \\
 &= X_0 E_{ox} \left( \frac{1}{2} [e^{j\omega t} (1 - e^{-jK_z z_0}) + e^{-j\omega t} (1 - e^{jK_z z_0})] \right) \\
 &= X_0 E_{ox} \left( \frac{1}{2} [e^{j\omega t} e^{-j\frac{K_z z_0}{2}} (e^{j\frac{K_z z_0}{2}} - e^{-j\frac{K_z z_0}{2}}) - e^{-j\omega t} e^{j\frac{K_z z_0}{2}} (e^{j\frac{K_z z_0}{2}} - e^{-j\frac{K_z z_0}{2}})] \right) \quad * \sin\theta = \frac{e^{j\theta} - e^{-j\theta}}{2j} \\
 &= j X_0 E_{ox} \left[ e^{j(\omega t - \frac{K_z z_0}{2})} \sin(\frac{K_z z_0}{2}) - e^{-j(\omega t - \frac{K_z z_0}{2})} \sin(\frac{K_z z_0}{2}) \right] \left( \frac{2j}{2j} \right) \\
 &= j X_0 E_{ox} \sin(\frac{K_z z_0}{2}) \left[ e^{j(\omega t - \frac{K_z z_0}{2})} - e^{-j(\omega t - \frac{K_z z_0}{2})} \right] \frac{2j}{2j} \\
 &= -2 X_0 E_{ox} \sin(\omega t - \frac{K_z z_0}{2}) \sin(\frac{K_z z_0}{2})
 \end{aligned}$$

$$\therefore \oint \vec{E} \cdot d\vec{l} = 2 X_0 E_{ox} \sin\left(\frac{K_z z_0}{2} - \omega t\right) \sin\left(\frac{K_z z_0}{2}\right)$$

### 6.1.2) Continuing w/ Faraday's Law in Differential form

$$\vec{\nabla} \times \vec{E} = - \frac{\partial B}{\partial t}$$

where  $\vec{\nabla} \times \vec{E}$  can be computed in cartesian coordinates:

$$\text{where } \vec{E} = E_x \hat{x} + 0 \hat{y} + 0 \hat{z} \quad \& \quad \frac{\partial E}{\partial y} = \frac{\partial E}{\partial x} = 0 \quad * E(z, t) \text{ has no dependence on } x \& y$$

then :

$$\vec{\nabla} \times \vec{E} = \begin{pmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ 0 & 0 & 0 \end{pmatrix} = 0 \hat{x} + \frac{\partial E}{\partial z} \hat{y} + 0 \hat{z}$$

$$\frac{\partial E}{\partial z} = \frac{\partial}{\partial z} (E_{ox} \cos(z K_z - \omega t))$$

$$= -K_z E_{ox} \sin(z K_z - \omega t)$$

$$\therefore \vec{B} = \int (-K_z E_{ox} \sin(z K_z - \omega t)) \hat{y} dt$$

$$\vec{B} = K_z E_{ox} \frac{\cos(K_z z - \omega t)}{\omega}$$

$$\therefore \vec{B} = \frac{k_z}{\omega} E_{ox} \cos(K_z z - \omega t) \hat{y}$$

Faraday's Integral Law can then be proved using previous results:

$$\oint \vec{E} \cdot d\vec{l} = - \frac{d}{dt} \Phi_B$$

where  $\Phi_B = \int \vec{B} \cdot d\vec{a}$  of the same rectangular loop 'e':

$$\begin{aligned} & \int \vec{B} \cdot d\vec{a} \\ &= \int_0^{z_0} \int_0^{x_0} B dx dz (\hat{y} \cdot \hat{y}) \\ &= \int_0^{z_0} \int_0^{x_0} \frac{k_z}{\omega} E_{ox} \cos(K_z z - \omega t) dx dz \\ &= \frac{k_z}{\omega} E_{ox} x_0 \int_0^{z_0} \cos(K_z z - \omega t) dz \\ &= \frac{k_z}{\omega} E_{ox} x_0 \left[ \frac{\sin(K_z z - \omega t)}{K_z} \right]_0^{z_0} \\ & \Phi_B = \frac{E_{ox} x_0}{\omega} [\sin(K_z z_0 - \omega t) + \sin(\omega t)] \\ & \frac{\partial \Phi_B}{\partial t} = \frac{\partial}{\partial t} \left[ \frac{x_0 E_{ox}}{\omega} (\sin(K_z z_0 - \omega t) + \sin(\omega t)) \right] \\ & - \frac{\partial \Phi_B}{\partial t} = - \left( \frac{E_{ox} x_0}{\omega} \right) [\cos(K_z z_0 - \omega t) + \cos(\omega t)] \\ & \therefore - \frac{\partial \Phi_B}{\partial t} = x_0 E_{ox} [\cos(K_z z_0 - \omega t) + \cos(\omega t)] \end{aligned}$$

which is equivalent to results obtained from path integral.

6.1.3) From Ampere's law in Integral form when  $\vec{J}=0$

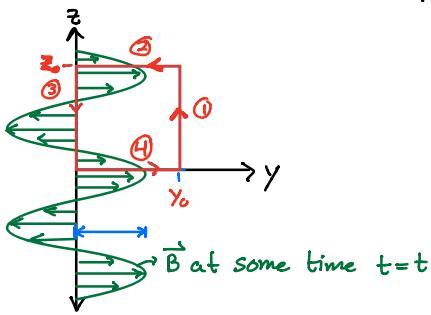
$$\oint \vec{B} \cdot d\vec{l} = \frac{1}{c^2} \frac{\partial \Phi_E}{\partial t}$$

From the  $\vec{B}$  found from 6.1.2 :

$$\vec{B} = \frac{k_z}{\omega} E_{ox} \cos(K_z z - \omega t) \hat{y}$$

We can see that  $\frac{k_z}{\omega} = v$ , where  $v$  is the velocity of the wave, and in free space represented by  $v=c$  (speed of light)

Similar to  $\vec{E}$ -field,  $\vec{B}$  is a plane wave (from satisfying wave equation by having the argument  $(\vec{k} \cdot \vec{r} - \omega t)$ ).  $\vec{B}$  then also propagates in  $z$ -direction &  $\vec{B} \perp \vec{E}$ , so it oscillates in  $y$ -direction. Then using similar arguments from 6.1.1 used for Faradays Law :



paths ① & ③ are  $\perp$  to  $\vec{B}$  (because  $\vec{B} \perp \vec{E}$ ) & ② & ④ are  $\parallel$  to  $\vec{B}$

$$\oint \vec{B} \cdot d\vec{l} = \int_2 B dy + \int_4 B dy$$

similar Faradays Law,  $\vec{B}(z,t)$ , at time  $t=t$ , during path ②:  $z=0$  & ④:  $z=z_0$   
 $\int_0^y B(y_0, t) dy + \int_{y_0}^0 B(0, t) dy$

$$= y_0 \left( \frac{1}{c} E_{ox} (\cos(k_z z_0 - \omega t) - \frac{1}{c} E_{ox} \cos(\omega t)) \right)$$

$$\oint \vec{B} \cdot d\vec{l} = \frac{y_0}{c} E_{ox} (\cos(k_z z_0 - \omega t) - \frac{1}{c} E_{ox} \cos(\omega t))$$

then Right-hand side :

$$\begin{aligned} \Phi_E &= \int_S \vec{E} \cdot d\vec{\alpha} \\ &= \int_0^{z_0} \int_0^y E_{ox} \cos(z k_z - \omega t) dy dz (\hat{x} \cdot \hat{x}) \end{aligned}$$

$$= y_0 E_{ox} \int_0^{z_0} \cos(z k_z - \omega t) dz$$

$$= \frac{y_0 E_{ox}}{k_z} \sin(z k_z - \omega t) \Big|_0^{z_0}$$

$$= \frac{y_0 E_{ox}}{k_z} (\sin(z_0 k_z - \omega t) - \sin(-\omega t))$$

$$\Phi_E = \frac{y_0 E_{ox}}{k_z} (\sin(z_0 k_z - \omega t) + \sin(\omega t))$$

$$\frac{\partial(\overline{\Phi}_e)}{\partial t} = \gamma_0 E_{ox} \frac{\omega}{k_z} (\cos(k_z - \omega t) + \cos(\omega t))$$

$$\therefore \frac{1}{c^2} \frac{\partial \overline{\Phi}_e}{\partial t} = \frac{\gamma_0}{c} E_{ox} (\cos(k_z - \omega t) + \cos(\omega t)) * \text{results are equivalent}$$

$$\therefore \oint \vec{B} \cdot d\vec{l} = \frac{1}{c^2} \frac{\partial \overline{\Phi}_e}{\partial t}$$