

We have seen hints that our charts actually seem to define a basis for our manifold on \mathbb{R}^n . We can make this more specific:

Def. (Chart induced basis)

Let M be a manifold and $\gamma: M \rightarrow \mathbb{R}^n$ a chart. We define the following for $f \in C^\infty(M)$

$$\left(\frac{\partial}{\partial \gamma_i} \right)_p f \equiv \left(\frac{\partial f}{\partial \gamma_i} \right)_p = \left[\cancel{\frac{\partial}{\partial \gamma_i}} (f \circ \gamma^{-1}) \right] (\gamma(p)),$$

which is the i-th basis vector of the tangent space of the manifold with respect to chart γ .

Note, that $\left(\frac{\partial}{\partial \gamma_i} \right)_p$ is not ~~just~~ a partial derivative in the standard sense. Of course, as we write it this way, it is very similar to a partial derivative. However, as $\left(\frac{\partial}{\partial \gamma_i} \right)_p$ acts on f which is a function from ~~from~~ $M \rightarrow \mathbb{R}$, it makes no sense for $\left(\frac{\partial}{\partial \gamma_i} \right)_p$ to be ^a partial derivative.

Note however that $\frac{\partial}{\partial \gamma_i}$ is in fact a partial derivative. This makes sense, as

$$f \circ \gamma^{-1}: \mathbb{R}^n \rightarrow \mathbb{R}$$

is a standard function (scalar field) from \mathbb{R}^n to \mathbb{R} .

In later chapters, we will write $\frac{\partial}{\partial \gamma_i} (=) \frac{\partial}{\partial x_i}$, because it will be obvious what objects we are dealing with.

Again, the f -function lets us input the charts to extract the components. In the case of $\left(\frac{\partial}{\partial \gamma^i}\right)_p$ we find:

$$\begin{aligned} \text{f-th component in } \gamma\text{-basis} &= \left(\frac{\partial}{\partial \gamma^i}\right)_p \gamma^r = \left(\frac{\partial \gamma^r}{\partial \gamma^i}\right)_p := \underbrace{\left[\partial_{\gamma^i} (\gamma^r \circ \gamma^{-1})\right]}_{\vdots \quad \vdots = \delta^r_i} \gamma(p) \\ \text{of } \left(\frac{\partial}{\partial \gamma^i}\right)_p &\qquad \qquad \qquad \qquad \qquad \qquad \qquad \qquad \text{(this is not super obvious)} \end{aligned}$$

That property is very similar to standard partial derivatives.

Notice that the j -th component of $\left(\frac{\partial}{\partial \gamma^i}\right)_p$ is δ_j^i , meaning that $\left(\frac{\partial}{\partial \gamma^i}\right)_p$'s are the basis vectors themselves!

Note: By now, it wasn't

clear that $\left(\frac{\partial}{\partial \gamma^i}\right)_p$ is a

Proof. (that $\left(\frac{\partial}{\partial \gamma^i}\right)_p$ are basis of all $v_{\gamma(p)}$'s): ~~vector~~ vector in $T_p M$

$$\begin{aligned} v_{\gamma(p)}(f) &= (f \circ \gamma)'(0) \\ &= ((\underbrace{f \circ \gamma^{-1}}_{\mathbb{R}^n \rightarrow \mathbb{R}} \circ \underbrace{\gamma \circ \gamma}_{\mathbb{R} \rightarrow \mathbb{R}^n})')'(0) \end{aligned}$$

(see Def. on S.2 (5.3)).

$$\begin{aligned} &= \underbrace{\partial_{\gamma^i} (f \circ \gamma^{-1})}_{\mathbb{R}^n \rightarrow \mathbb{R}}|_{\gamma(\gamma^{-1}(0))} \underbrace{(\gamma^i \circ \gamma)'(0)}_{=: \dot{\gamma}^i} \\ &= \left(\frac{\partial f}{\partial \gamma^i}\right)_p \in \mathbb{R} \\ &= \dot{\gamma}^i \left(\frac{\partial}{\partial \gamma^i}\right)_p \end{aligned}$$

$$\Rightarrow v_{\gamma(p)} = \dot{\gamma}^i \left(\frac{\partial}{\partial \gamma^i}\right)_p$$

where $\dot{\gamma}^i$ are the components of $v_{\gamma(p)}$ and

$\left(\frac{\partial}{\partial \gamma^i}\right)_p$ are the basis vectors. Hence the basis for any vector on M are the "partial derivatives"!

This might seem overwhelming, so let us
make
~~take~~ an example, namely polar coordinates.

Consider $S^1 = \{x^2 + y^2 = 1 \mid x, y \in \mathbb{R}\}$ and

the chart:

$$\varphi(x) = (r, \varphi) = (r^1, r^2)(\begin{pmatrix} x \\ y \end{pmatrix})$$

with $r = \sqrt{x^2 + y^2}$ and $\varphi = \arctan(y/x)$.

Let us compute $(\frac{\partial}{\partial r^1})_p = (\frac{\partial}{\partial r})_p$ and $(\frac{\partial}{\partial r^2})_p = (\frac{\partial}{\partial \varphi})_p$!

$$(\frac{\partial}{\partial r})_p(f) = \frac{\partial f}{\partial r}|_p = \left(\partial_r f(\varphi^{-1}) \right)|_{\varphi(p)}$$

Notice, that $\varphi^{-1} : (r, \varphi) \mapsto (x, y) = (r \cos \varphi, r \sin \varphi)$.

$$\begin{aligned} &= \left(\partial_r f(x(r, \varphi), y(r, \varphi)) \right)|_{\varphi(p)} \\ \text{Chain rule} \quad &\downarrow \quad \left[\partial_x f(x, y) \frac{\partial x(r, \varphi)}{\partial r} + \partial_y f(x, y) \frac{\partial y(r, \varphi)}{\partial r} \right] |_{\varphi(p)} \\ \text{for standard} \quad &= (r \cos \varphi, r \sin \varphi) \\ \text{functions} \quad &= (\sqrt{x^2 + y^2}, \arctan(\frac{y}{x})) \\ \text{as } M = S^1 \subseteq \mathbb{R}^2 \quad &= (\sqrt{x^2 + y^2}, \arctan(\frac{y}{x})) \end{aligned}$$

$$= [\partial_x f(x, y) \cos \varphi + \partial_y f(x, y) \sin \varphi]|_{\varphi(p)}$$

$$\begin{aligned} &= \partial_x f(x, y) \cos(\arctan(\frac{y}{x})) \\ &\quad + \partial_y f(x, y) \sin(\arctan(\frac{y}{x})) \end{aligned}$$

$$\begin{aligned} x^2 + y^2 = 1 \quad &= \partial_x f \frac{x}{\sqrt{x^2 + y^2}} + \partial_y f \frac{y}{\sqrt{x^2 + y^2}} \\ &= (x \partial_x + y \partial_y) f \end{aligned}$$

Similarly: $(\frac{\partial}{\partial \varphi})_r f = (x \partial_y - y \partial_x) f$

Notice that, as $M \subseteq \mathbb{R}^2$, we ~~wrote~~ wrote standard partial derivatives ∂_x, ∂_y instead of $(\frac{\partial}{\partial x})_p$ and $(\frac{\partial}{\partial y})_p$.

Proof. (that $\partial_x = (\frac{\partial}{\partial x})_p$ on \mathbb{R}^n)

$$(\frac{\partial}{\partial x})_p f = (\partial_x f(\gamma^{-1}))|_{\gamma(p)}$$

with ^{the} identity chart $\gamma: (\mathbb{X}) \mapsto (\mathbb{Y})$, so:

$$(\frac{\partial}{\partial x})_p f = (\partial_x f(x, y)). \quad \square$$

So, to conclude this, we can write:

$$(\frac{\partial}{\partial r})_p = x(\frac{\partial}{\partial x})_p + y(\frac{\partial}{\partial y})_p \quad ; \quad p = (x, y)$$

and $(\frac{\partial}{\partial \theta})_p = x(\frac{\partial}{\partial y})_p - y(\frac{\partial}{\partial x})_p$.

Note: Does this seem familiar?

This relation can be shown in a more general sense.

Proof.

$$\begin{aligned} (\frac{\partial}{\partial \gamma_i})_p f &= (\partial_{\gamma_i} (f \circ \gamma_i^{-1}))(\gamma_i(p)) \\ &= (\partial_{\gamma_i} ((f \circ \gamma^{-1}) \circ (\gamma \circ \gamma_i^{-1}))) \gamma_i(p) \\ &= \underbrace{\partial_{\gamma_i} (f \circ \gamma^{-1})}_{=\gamma_i(p)}|_{\gamma(\gamma_i^{-1}(p))} \underbrace{\partial_{\gamma_i} (\gamma \circ \gamma_i^{-1})}_{=\gamma(p)}|_{\gamma(p)} \\ &=: (\frac{\partial}{\partial \gamma_i})_p f \\ &=: (\frac{\partial \gamma^i}{\partial \gamma_i})_p \end{aligned}$$

$$= (\frac{\partial \gamma^i}{\partial \gamma_i})_p$$

$$\Rightarrow (\frac{\partial}{\partial \gamma_i})_p = (\frac{\partial \gamma^i}{\partial \gamma_i})_p \quad ||(*)$$

Let us now revisit the example on 5.5 and 5.6, where we obtained:

$$1) \nabla_{\gamma_1 p} (\gamma_0) = 1$$

$$2) \nabla_{\gamma_1 p} (\gamma_{IR}) = \begin{pmatrix} 0 \\ 1 \end{pmatrix}$$

and notice that for the case 1), we have one basis vector $\left(\frac{\partial}{\partial \gamma_0}\right)_p$ and for the case 2), we have two basis vectors $\left(\frac{\partial}{\partial \gamma_{IR}^1}\right)_p$ and $\left(\frac{\partial}{\partial \gamma_{IR}^2}\right)_p$, @ $p = (1, 0)$

We can use the relation (*):

$$\nabla_{\gamma_1 p} = 1 \cdot \left(\frac{\partial}{\partial \gamma_0}\right)_p$$

$$= \left(\frac{\partial \gamma_{IR}^1}{\partial \gamma_0}\right)_p \left(\frac{\partial}{\partial \gamma_{IR}^1}\right)_p + \left(\frac{\partial \gamma_{IR}^2}{\partial \gamma_0}\right)_p \left(\frac{\partial}{\partial \gamma_{IR}^2}\right)_p$$

where we note that

$$\left(\frac{\partial \gamma_{IR}^1}{\partial \gamma_0}\right)_p = \partial_{\gamma_0} (\gamma_{IR}^1 \circ \gamma_0^{-1})|_{\gamma_0(p)}$$

with $\gamma_0(p) = \gamma_0(x, y) = \arctan(y/x)$

and $(\gamma_{IR}^1 \circ \gamma_0^{-1})(\lambda)$

$$= \gamma_{IR}^1 (\cos \lambda, \sin \lambda) = \cos \lambda$$

$$\Rightarrow \partial_{\gamma_0} (\gamma_{IR}^1 \circ \gamma_0^{-1})|_{\gamma_0(p)}$$

$$= \partial_\lambda \cos \lambda = -\sin \lambda \mid_{\arctan(\text{***})} = 0$$

Similarly :

$$\left(\frac{\partial^2 \psi}{\partial q_0^2} \right)_P = 1.$$

$$\Rightarrow V_{8,P} = \left(\frac{\partial}{\partial q_0} \right)_P = \left(\frac{\partial}{\partial q^2} \right)_P.$$

□

Hence, we were just expressing $V_{8,P}$ in different basis!