Cluster Expansions of Thermal States using Tensor Networks

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Introduction

Introduction

- Overview condensed matter physics
  - Macroscopic and microscopic physical properties of matter
    - Metals
    - semiconductors
    - Liquids
    - Bose-Einstein Condensates
    - Magnets
  - Different disciplines
    - Experimental
    - Theoretical
    - Engineering

Introduction

- Overview condensed matter physics
- Strongly correlated materials [1]
  - Superconductors
  - Quantum spin liquids
  - Strange metals
  - Correlated topological matter

Introduction

- Overview condensed matter physics
- Strongly correlated materials
- How to proceed
  - Material synthesis and discovery
  - Analytical methods
  - Numerical methods

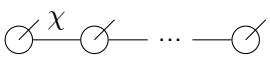
Simulating Quantum Many-body Systems

- Equations are known
- Curse of dimensionality
- Numerical methods
  - Exact diagonalisation
  - (post-) Hartree Fock methods, DFT methods
  - Monte Carlo methods
  - Tensor Networks

Tensor Networks

$$|\Psi\rangle = \sum_{i_1 i_2 \cdots i_n} C^{i_1 i_2 \cdots i_n} |i_1\rangle \otimes |i_2\rangle \otimes \cdots \otimes |i_n\rangle \,. \tag{1}$$

$$\begin{aligned} C^{i_1 i_2 \cdots i_n} &= w_l C^{i_1} C^{i_2} \cdots C^{i_n} w_r \\ &= \end{aligned} \tag{2}$$



- MPS
- Relevant corner Hilbert space

Operator Exponential

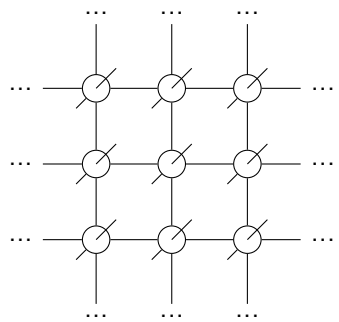
- (Real) Time evolution:

$$\hat{O} = e^{-i\hat{H}t} \qquad (3)$$

- Statistical ensembles:

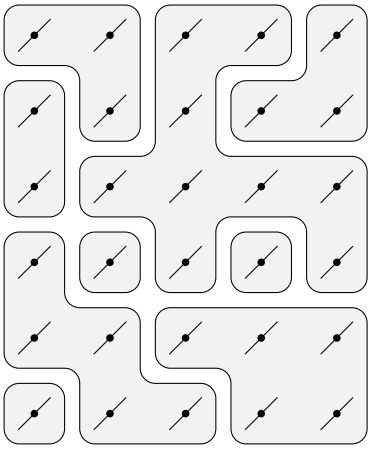
$$\hat{O} = \frac{e^{-\beta\hat{H}}}{\mathrm{Tr}\left(e^{-\beta\hat{H}}\right)} \qquad (4)$$

Imaginary time ($\beta = it$)

$\hat{O} =$  (5)

Cluster Expansions

Cluser Exapansions

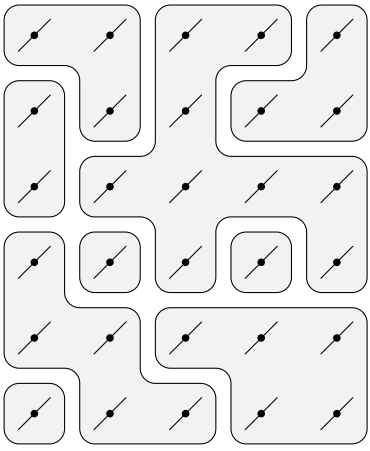


- $e^{\hat{H}} = \sum_{\{B\}} \bigotimes_i B_i$

$$e^{H(1)} = \quad (6)$$

$$e^{H(2)} = \quad + \quad (7)$$

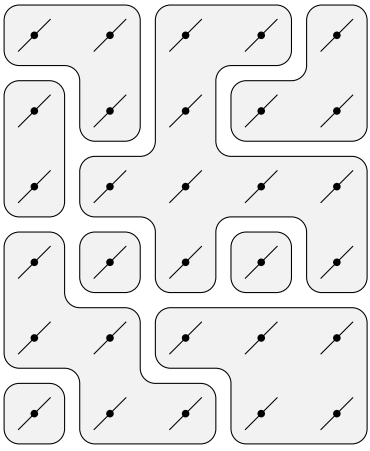
Cluser Exapansions

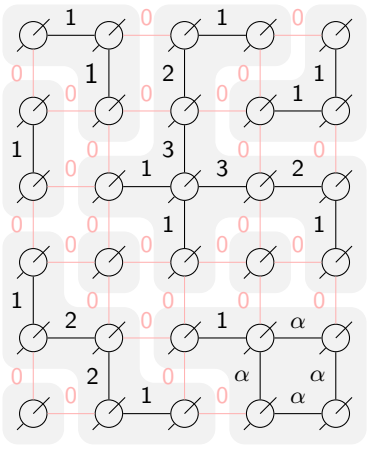


- $e^{\hat{H}} = \sum_{\{B\}} \bigotimes_i B_i$
- Finite number of blocks
- Encoded by 1 tensor

$$O^{abcd} = \quad \begin{array}{ccc} & b & \\ a & & i\ c \\ & j\ d & \end{array} \qquad (6)$$

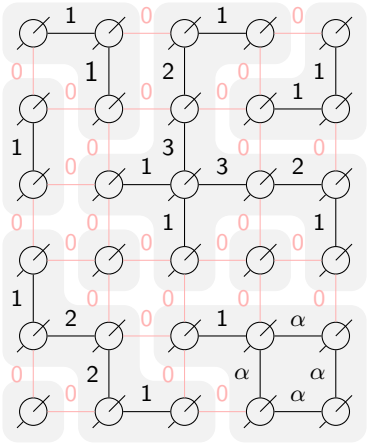
Cluser Exapansions





Cluser Exapansions

- Multiple choices for encoding
- Solvers
  - Linear
  - Nonlinear

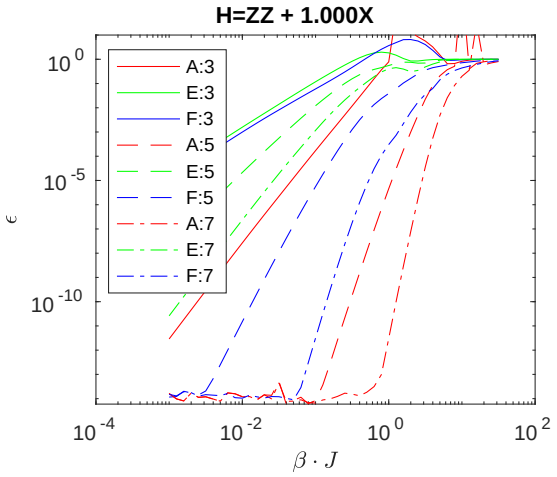


Advantages

- Doesn't break symmetry
- Thermodynamic limit
- Tensor Network toolbox

Results

1D: Transverse Field Ising (TFI)



- Relative error ϵ
- Different encodings:
  - A: Small
  - E: Strict
  - F: well-conditioned
- bond dimension

		Encoding	
		A	E/F
Order	3	5	10
	5	21	42
	7	85	170

Conclusion

- Large β-steps
- Real time evolution
- Encoding
- Truncation χ

2D TFI: Introduction

- Phase Transition
- Criticality
- Finite size scaling
  - Observables: m, S and ξ
  - Parameters: T_c, exponents
- $\Gamma = 2.5$
- VUMPS (χ,δ^{-1})

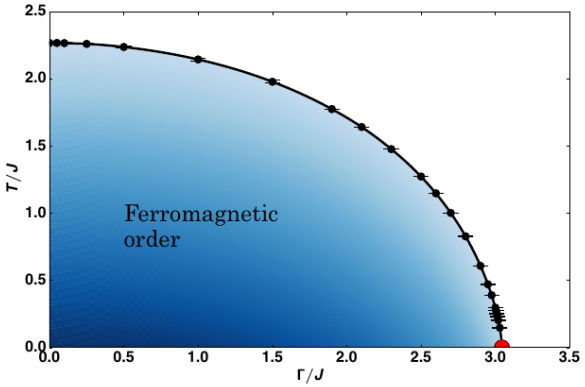
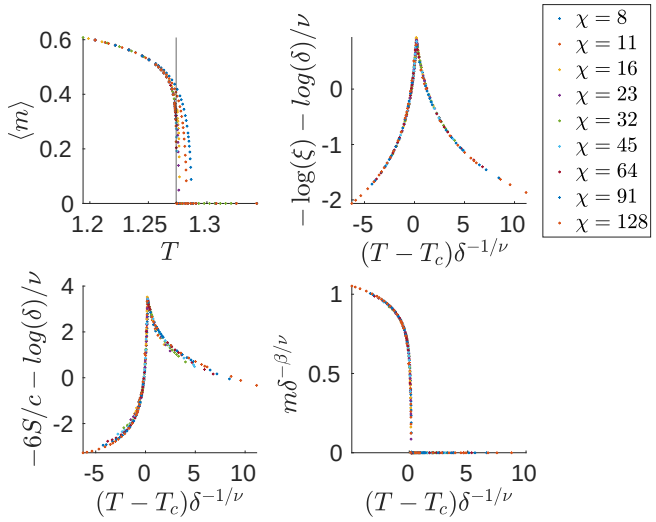


Figure taken from [2]

TFI Phase Diagram: $\Gamma = 2.5$



	T_c
Fit	1.2736(6)
QMC	1.2737(6)
TN	1.2737(2)

Data from [3]

Conclusion and Outlook

Conclusion

- Construction fast and stable
- Cluster expansions work well in 1D and 2D
- Real time evolution

Outlook

- 3D
- Incorperating internal symmetries
- Lattices

References I

A. Alexandradinata, N. P. Armitage, A. Baydin, W. Bi, Y. Cao, H. J. Changlani, E. Chertkov, E. H. d. S. Neto, L. Delacretaz, I. E. Baggari, G. M. Ferguson, W. J. Gannon, S. A. A. Ghorashi, B. H. Goodge, O. Goulko, G. Grissonnanche, A. Hallas, I. M. Hayes, Y. He, E. W. Huang, A. Kogar, D. Kumah, J. Y. Lee, A. Legros, F. Mahmood, Y. Maximenko, N. Pellatz, H. Polshyn, T. Sarkar, A. Scheie, K. L. Seyler, Z. Shi, B. Skinner, L. Steinke, K. Thirunavukkuarasu, T. V. Trevisan, M. Vogl, P. A. Volkov, Y. Wang, Y. Wang, D. Wei, K. Wei, S. Yang, X. Zhang, Y.-H. Zhang, L. Zhao, A. Zong, The Future of the Correlated Electron Problem (oct 2020).
arXiv:2010.00584.
URL http://arxiv.org/abs/2010.00584

References II


S. Hesselmann, S. Wessel, Thermal Ising transitions in the vicinity of two-dimensional quantum critical points, PHYSICAL REVIEW B 93 (2016) 155157.
doi:10.1103/PhysRevB.93.155157.


P. Czarnik, P. Corboz, Finite correlation length scaling with infinite projected entangled pair states at finite temperature, Physical Review B 99 (2019) 245107.
doi:10.1103/PhysRevB.99.245107.

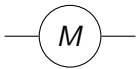
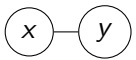
Tensor Networks

Tensor Networks: Introduction

$$|\Psi\rangle = \sum_{i_1 i_2 \cdots i_n} C^{i_1 i_2 \cdots i_n} |i_1\rangle \otimes |i_2\rangle \otimes \cdots \otimes |i_n\rangle \,. \tag{6}$$

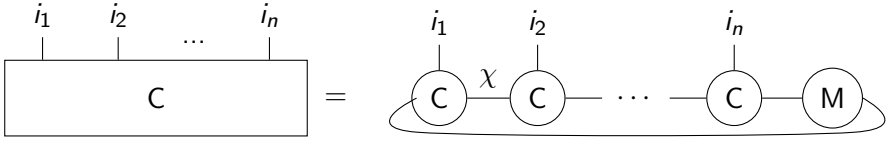
$$C^{i_1 i_2 \cdots i_n} = Tr(C^{i_1} C^{i_2} \cdots C^{i_n} M). \tag{7}$$

Tensor Networks: Graphical Notation

conventional	Einstein	tensor notation
$\vec{x}$	x_α	
M	$M_{\alpha\beta}$	
$\vec{x} \cdot \vec{y}$	$x_\alpha y_\alpha$	

Tensor Networks: MPS

$$C^{i_1 i_2 \cdots i_n} = Tr(C^{i_1} C^{i_2} \cdots C^{i_n} M) \tag{8}$$



(9)

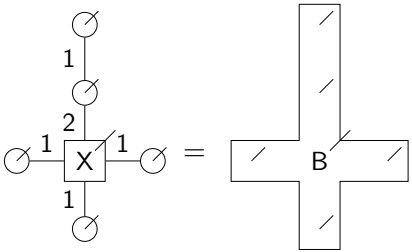
Tensor Networks: Operators

$$\hat{O} = \cdots \cdots \tag{10}$$

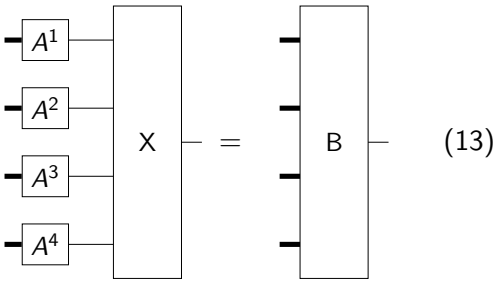
$$\hat{O}\,|\Psi\rangle = \cdots \chi \cdots \cdots \chi \cdots = \cdots \chi^2 \cdots \tag{11}$$

Linear Solver

Linear Solver: Inversion Scheme



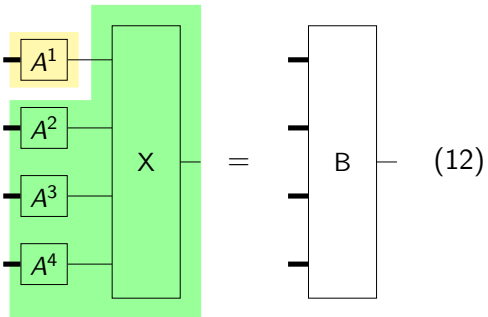
(12)



(13)

Linear Solver: Inversion Scheme

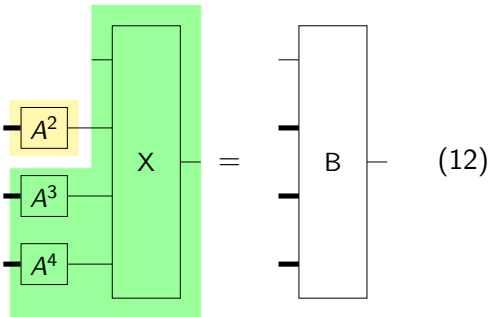
- Invert A^i separately
  - Fast
  - Numerically unstable



(12)

Linear Solver: Inversion Scheme

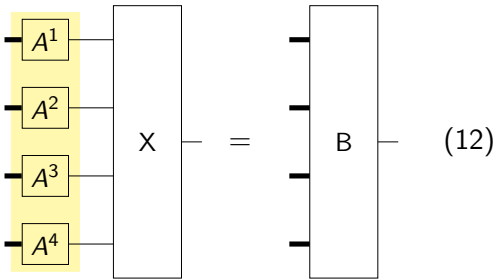
- Invert A^i separately
  - Fast
  - Numerically unstable



(12)

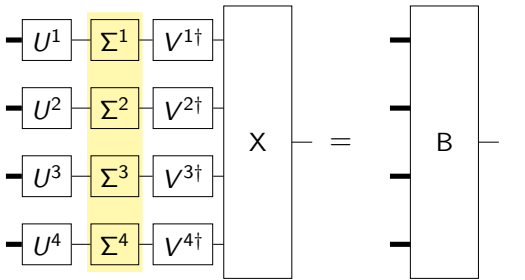
Linear Solver: Inversion Scheme

- Invert A^i separately
- Full inversion
  - Slow
  - Stable for pseudoinverse



Linear Solver: Inversion Scheme

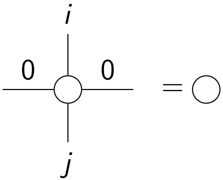
- Invert A^i separately
- Full inversion
- Sparse full inversion
  - $A^i = U^i \Sigma^i V^{i\dagger}$

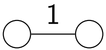


(12)

Construction

Notation

$O^{00} =$  (13)

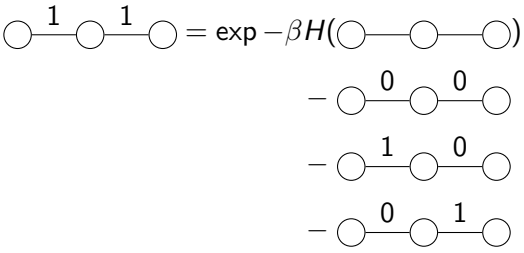
$O^{01} O^{10} =$  (14)

General idea

$$\bigcirc = \exp\left(-\beta H(\bigcirc)\right) \tag{15}$$

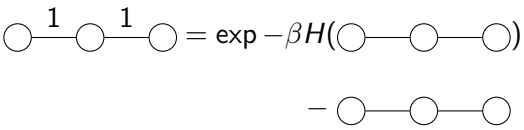
$$\bigcirc\!\!\overset{1}{-}\!\!\bigcirc = \exp -\beta H(\bigcirc\!\!-\!\!\bigcirc - \bigcirc\!\!\overset{0}{-}\!\!\bigcirc) \tag{16}$$

General idea



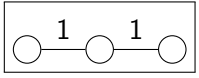
(17)

General idea



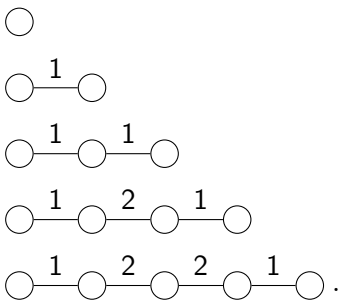
(17)

General idea



(17)

1D: Variant A



(18a)

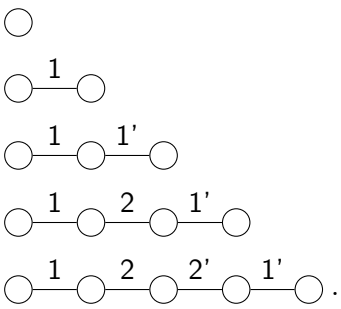
(18b)

(18c)

(18d)

(18e)

1D: Variant E



(19a)

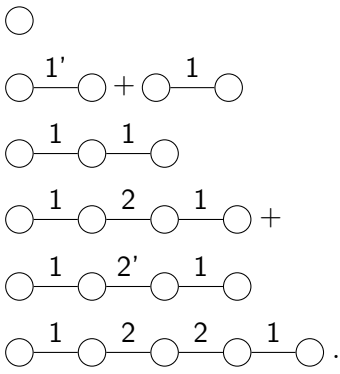
(19b)

(19c)

(19d)

(19e)

1D: Variant F



(20a)

(20b)

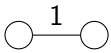
(20c)

(20d)

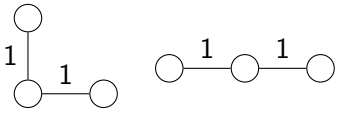
(20e)

$$O^{0000} = \quad = \bigcirc \tag{21}$$

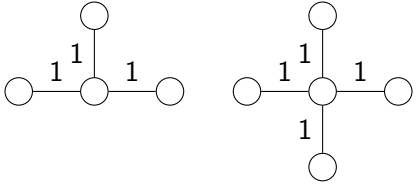
2D: Linear Blocks



(22a)

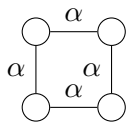


(22b)



(22c)

2D: Nonlinear Blocks



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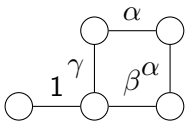
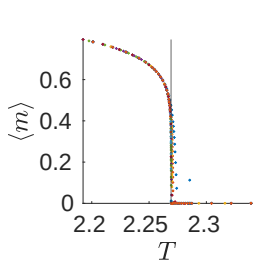
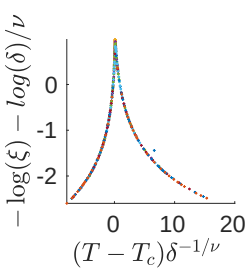


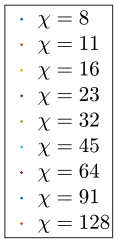
(24)

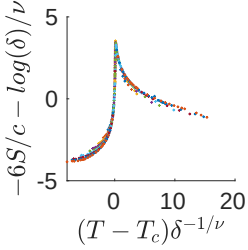
TFI Collapses

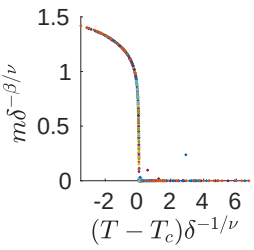
TFI Phase Diagram: Classical Ising



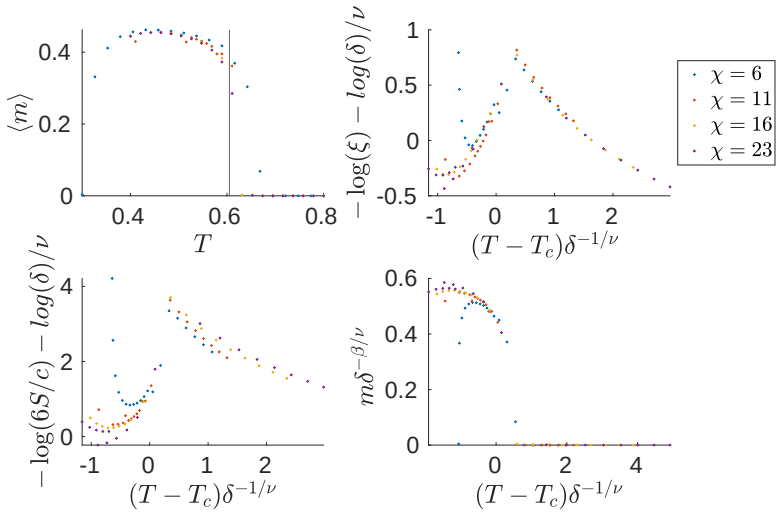






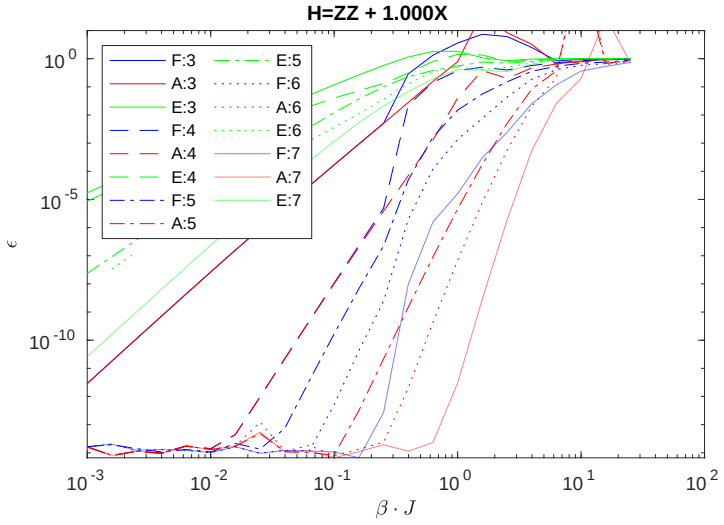


	T_c
Fit	2.691(9)
Exact	2.691853

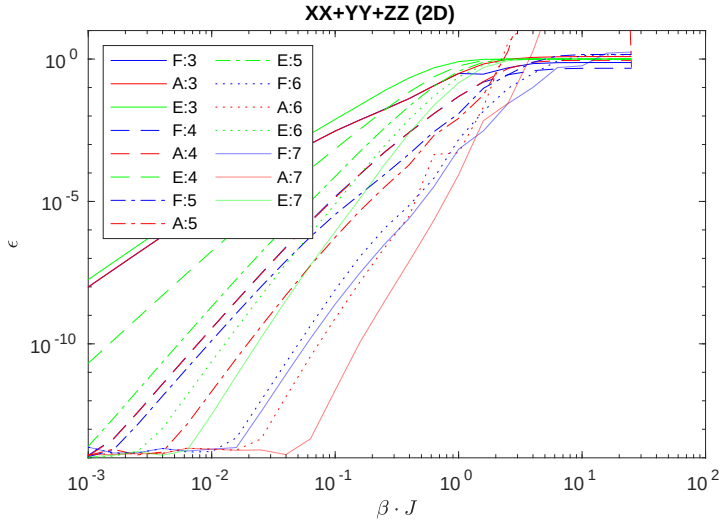


Direct Results

1D: Transverse Field Ising (TFI): full

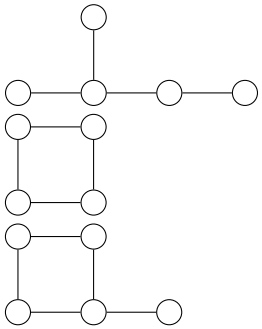


1D: Heisenberg XXX



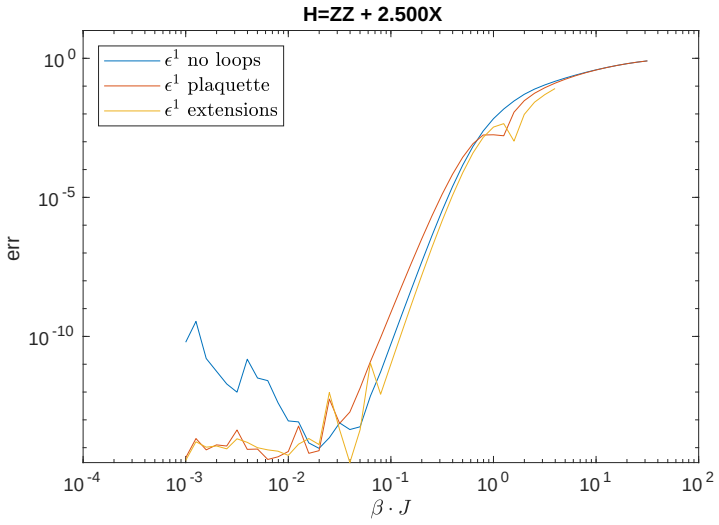
2D: Encodings + Error Measure

- Relative error ϵ more challenging
- Encodings based on A (order 5)
  - No loops
  - +Plaquette
  - +Extensions



	χ
no loops	21
plaquette	27
extensions	43

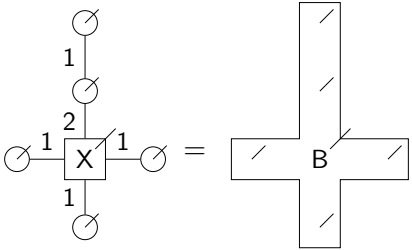
2D: Transverse Field Ising



Solvers

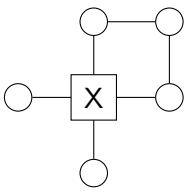
Linear solver

- $AX = B$
- Invert leg per leg
- Pseudoinverse

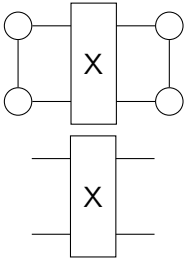


(25)

Linear Solver: Applicability



(26)



(27a)

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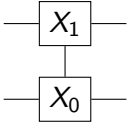
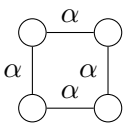


(27b)

Nonlinear Solver

- Nonlinear least squares
- Jacobian
- Permutations



(28)

Sequential Linear Solver

- Based on linear solver
- Sweep over unknown tensors
- Permutations