### Homework 11

# David Miller MAP5345: Partial Differential Equations I

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#### Problem 1

Consider polar coordinates  $(r, \theta)$ , whice are related to Cartesian coordinates via  $x = r\cos\theta$  and  $y = r\sin\theta$ .

a) Use the chain rule for partial derivatives to calculate  $\frac{\partial}{\partial r}$ ,  $\frac{\partial^2}{\partial r^2}$ ,  $\frac{\partial}{\partial \theta}$ ,  $\frac{\partial^2}{\partial \theta^2}$  in terms of x and y partial derivatives.

The coordinate transformation

$$x \mapsto r \cos(\theta), \quad y \mapsto r \sin(\theta)$$

yields the following

$$\begin{split} u_{\theta} &= u_x x_{\theta} + u_y y_{\theta} = -u_x r \sin(\theta) + u_y r \cos(\theta) \\ u_r &= u_x x_r + u_y y_r = -u_x \sin(\theta) + u_y \cos(\theta) \\ u_{\theta\theta} &= -r \cos(\theta) u_x - r \sin(\theta) \partial_{\theta} u_x - r \sin(\theta) u_y + r \cos(\theta) \partial_{\theta} u_y \\ &= -r \cos(\theta) u_x - r \sin(\theta) \left( -u_{xx} r \sin(\theta) + u_{xy} r \cos(\theta) \right) \\ &- r \sin(\theta) u_y + r \cos(\theta) \left( -r u_{xy} \sin(\theta) + u_{yy} r \cos(\theta) \right) \\ &= -r \left( \cos(\theta) u_x + \sin(\theta) u_y \right) + r^2 \left( \sin^2(\theta) u_{xx} - 2 \cos(\theta) \sin(\theta) u_{xy} + \cos^2(\theta) u_{yy} \right) \\ u_{rr} &= \cos(\theta) \partial_r u_x + \sin(\theta) \partial_r u_y \\ &= \cos^2(\theta) u_{xx} + 2 \cos(\theta) \sin(\theta) u_{xy} + \sin^2(\theta) u_{yy} \end{split}$$

b) Use your results to show that the Laplacian in polar coordinates is given by

$$\Delta u = u_{rr} + \frac{1}{r}u_r + \frac{1}{r^2}u_{\theta\theta}$$

Dividing the expression for  $u_{\theta\theta}$  yields

$$\frac{1}{r^2}u_{\theta\theta} = -\frac{1}{r}u_r + \sin^2(\theta)u_{xx} - 2\cos(\theta)\sin(\theta)u_{xy} + \cos^2(\theta)u_{yy}$$

Then adding the above to  $u_{rr}$  gets us

$$u_{rr} + \frac{1}{r^2}u_{\theta\theta} = -\frac{1}{r}u_r + u_{xx} + u_{yy}$$

$$\Rightarrow u_{xx} + u_{yy} = u_{rr} + \frac{1}{r}u_r + \frac{1}{r^2}u_{\theta\theta}$$

Using Julia, graph  $J_0(x)$ , i.e. the Bessel function of the first kind and of order n = 0. Choose a reasonable domain to graph the function, so that you can see at least a few roots of  $J_0$ . On the same graph, plot the function

$$\sqrt{\frac{2}{\pi x}}\cos(x - \pi/4 - n\pi/2)\tag{1}$$

and verify that the two are asymptotic for x >> 1. How many roots do you have to go out for the asymptotic formula to capture the roots pretty accurately? Do the same for  $J_1(x)$ , the first-kind Bessel function of order n = 1. Also, verify the asymptotic formula.

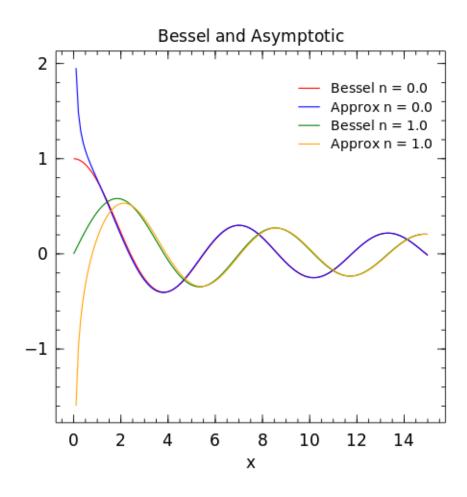


Figure 1: Bessel Function plotted against its asymptotic formula.

For a drum head with given wave-speed c and radius a, can you determine the fundamental frequency, and thus the 'pitch' of the sound that drum will produce? For a realistic drum with c = 400 m/s and a = 0.5, what is the fundamental frequency in Hz? Also find the first harmonic.

The fundamental frequency will be the first root of  $J_0$ , the 0-th order Bessel function, times the wave speed c. Similarly the n-th harmonic will be the n+1-th root of  $J_0$ . From this we get

Harmonic	Frequency (Hz)
0 (Fundamental)	1923.8608
1	4416.0624
2	6922.9816
3	9433.2272
4	11944.7336
:	i:

Consider the Bessel ODE derived in class. Put this ODE into the form of a Sturm-Liouville problem. What is the weighted inner product for which the functions  $R_m(r) = J_0(\sqrt{\lambda}r)$  are orthogonal?

The Sturm-Liouville equation

$$(p(x)u'(x))' - q(x)u(x) = -\lambda m(x)u(x)$$

has inner product (orthogonal functions)

$$\langle f, g \rangle = \int m(x) f(x) \overline{g(x)} dx$$

The ODE we have in class can be multiplied by r to obtain

$$R''(r) + \frac{1}{r}R'(r) + R(r)\left(\lambda - \frac{n^2}{r^2}\right) = 0$$

$$rR''(r) + R'(r) + rR(r)\left(\lambda - \frac{n^2}{r^2}\right) = 0$$
(multiply by  $r$ )
$$\Rightarrow u = R(r), \quad p = r, \quad q = \frac{n^2}{r}, \quad m = r$$

Therefore r is the weight that makes the Bessel functions orthogonal. The resulting inner product for the Bessel functions is

$$\langle R_i, R_j \rangle = 2\pi \int r R_i R_j \, dr$$

where the  $2\pi$  comes from the integral about the  $\theta$  domain  $[0, 2\pi]$ .

Consider radial vibrations of a circular drum of radius a=0.5 m. The medium's speed of sound given by c=400 m/s.

a) Now, suppose that the initial displacement is  $u_0 = a^2 - r^2$  and there is no initial velocity. Write don the exact solution to the IBVP.

The general solution of the PDE under rotational invariance is

$$u(r,t) = \sum_{n=1}^{\infty} (A_n \cos(\sqrt{\lambda_n}ct) + B_n \sin(\sqrt{\lambda}ct)) J_0(\sqrt{\lambda_n}r)$$

but under zero initial velocity we get  $B_n = 0$  and are left with

$$u(r,t) = \sum_{n=1}^{\infty} A_n \cos(\sqrt{\lambda_n} ct) J_0(\sqrt{\lambda_n} r)$$

Using projection to determine  $A_n$  we arrive at

$$A_{n} = \frac{\langle a^{2} - r^{2}, J_{0}(\sqrt{\lambda_{n}}r) \rangle}{\langle J_{0}(\sqrt{\lambda_{n}}r), J_{0}(\sqrt{\lambda_{n}}r) \rangle} = \frac{\int_{0}^{0.5} (a^{2} - r^{2}) J_{0}(\sqrt{\lambda_{n}}r) r \, dr}{\int_{0}^{0.5} J_{0}(\sqrt{\lambda_{n}}r) J_{0}(\sqrt{\lambda_{n}}r) r \, dr}$$

Therefore the exact solution is

$$\sum_{n=1}^{\infty} A_n \cos(\sqrt{\lambda_n} ct) J_0(\sqrt{\lambda_n} r), \quad A_n = \frac{\int_0^{0.5} (a^2 - r^2) J_0(\sqrt{\lambda_n} r) r \, dr}{\int_0^{0.5} J_0(\sqrt{\lambda_n} r) J_0(\sqrt{\lambda_n} r) r \, dr}$$

b) Perform the projection numerically, using the trapezoidal rule, in order to find a numerical approximation to the first three coefficients in your solution.

Using the trapezoidal rule to perform quadrature we get the following coefficients

$$A_1 \approx 0.0375909$$
,  $A_2 \approx 0.0194997$ ,  $A_3 \approx 0.0047067$ 

c) Visualize a few different time-values of this approximate solution in Julia.

## **Decay Coefficients**

Find the lower bounds of the following functions

**Theorem 1.** If  $f_{ext} \in C^n$  then there exists a M such that  $c_k \leq \frac{M}{|k|^n}$ 

**Theorem 2.** If  $f_{ext} \in L^2$  and if  $c_k = \frac{\langle f, e^{ikx} \rangle}{\langle e^{ikx}, e^{ikx} \rangle}$  and there exists  $M, \epsilon > 0$  such that  $|c_k| \le \frac{M}{|k|^{1+n+\epsilon}}$  then  $f_{ext} \in C^n$ 

a) |x|

We know that  $f_{ext} \in C^0$  and therefore  $f_{ext} \not\in C^1$  and the decay rate is bounded below by 0. From Theorem 2 we then know that  $\not\exists M$  or  $\epsilon$  such that  $|c_k| \leq \frac{M}{|k|^{2+\epsilon}}$  and therefore the decay rate is bounded above by 2.

*b*) *x* 

We know that  $f_{ext}$  is piecewise polynomial and therefor  $f_{ext} \not\in C^0$  and the decay rate is bounded below by some constant M. From Theorem 2 we then know that  $\not\exists M$  or  $\epsilon$  such that  $|c_k| \leq \frac{M}{|k|^{1+\epsilon}}$  and therefore is bounded above by 1.

c) 
$$\pi^2 - x^2$$

Since  $f_{ext} \in C^0$  the proof follows similarly from part (a). Therefore the decay rate is bounded below by 0 and above by 2.

d) 
$$\sqrt{\pi^2 - x^2}$$

Since  $f_{ext} \in C^0$  the proof follows similarly from part (a). Therefore the decay rate is bounded below by 0 and above by 2.

e) 
$$x(\pi^2 - x^2)$$

We know that  $f_{ext} \in C^1$  and therefore  $f_{ext} \notin C^2$  and the decay rate is bounded below by 1. From Theorem 2 we then know that  $\not\exists M$  or  $\epsilon$  such that  $|c_k| \leq \frac{M}{|k|^{3+\epsilon}}$  and therefore the decay rate is bounded above by 3.

$$f) (\pi^2 - x^2)^2$$

We know that  $f_{ext} \in C^2$  and therefore  $f_{ext} \notin C^3$  and the decay rate is bounded below by 2. From Theorem 2 we then know that  $\not\exists M$  or  $\epsilon$  such that  $|c_k| \leq \frac{M}{|k|^{4+\epsilon}}$  and therefore the decay rate is bounded above by 4.