

# Homework 4

David Miller

MAP5165: Methods in Applied Mathematics I

November 3, 2017

## Problem 1

Find an approximation solution for the IVP (here  $0 < \epsilon \ll 1$ ):

$$\frac{dy}{dt} + \epsilon y^2 - y = 0$$

subject to the initial condition  $y(0) = 1$ .

If we let  $y(t, \epsilon)$  be the solution to the problem we get that

$$y(t, \epsilon) \approx y(t, 0) + \epsilon \frac{\partial y(t, 0)}{\partial \epsilon} + \epsilon^2 \frac{\partial^2 y(t, 0)}{\partial \epsilon^2} + \mathcal{O}(\epsilon^3)$$

for some perturbation  $\epsilon$ . Now we take the partial with respect to the perturbation  $\epsilon$  of the original problem to get

$$\frac{\partial}{\partial \epsilon} \left( \frac{\partial}{\partial t} y + \epsilon y^2 - y \right) = 0, \quad \frac{\partial y(0, \epsilon)}{\partial \epsilon} = 1,$$

which can be rewritten as

$$\frac{\partial}{\partial t} \frac{\partial y(t, \epsilon)}{\partial \epsilon} = \frac{\partial y(t, \epsilon)}{\partial \epsilon} - y^2 - 2\epsilon y \frac{\partial y(t, \epsilon)}{\partial \epsilon}.$$

The unperturbed solution to this is  $y(t, 0) = e^t$ . Now if we define  $Y = \partial y(t, 0)/\partial \epsilon$  we get the following differential equation

$$\frac{dY}{dt} = Y - e^{2t}, \quad Y(0) = 0$$

where the general solution to this is  $Ae^t - e^{2t}$ . Using the initial condition sets  $A = 1$ , and therefore the solution is

$$Y = e^t - e^{2t}, \quad y \approx e^t + \epsilon(e^t - e^{2t}).$$

## Problem 2

I used an online source for guidance with this problem.

Reference: <https://www.iist.ac.in/sites/default/files/people/multiplescale.pdf>

Find an approximate solution for the IVP (here  $0 < \epsilon \ll 1$ ):

$$\frac{d^2 y}{dt^2} + \epsilon \frac{dy}{dt} + y = 0$$

subject to the initial condition  $y(0) = 0$  and  $\frac{dy}{dt}(0) = 1$ .

To apply the method of multiple scales with introduce the variables

$$T_0 = t, \quad T_1 = \epsilon t, \quad T_2 = \epsilon^2 t$$

where they each represent different time scales due to the dampening effecting the amplitude and phase shift of the oscillator. Using the chain rule we get

$$\begin{aligned} \frac{d}{dt} &= \frac{\partial}{\partial T_0} \frac{\partial T_0}{\partial t} + \frac{\partial}{\partial T_1} \frac{\partial T_1}{\partial t} + \frac{\partial}{\partial T_2} \frac{\partial T_2}{\partial t} + \mathcal{O}(\epsilon^3) \\ &= \frac{\partial}{\partial T_0} + \epsilon \frac{\partial}{\partial T_1} + \epsilon^2 \frac{\partial}{\partial T_2} \\ \frac{d^2}{dt^2} &= \frac{\partial^2}{\partial T_0^2} + 2\epsilon \frac{\partial^2}{\partial T_0 \partial T_1} + \epsilon^2 \left( \frac{\partial^2}{\partial T_0 \partial T_2} + \frac{\partial^2}{\partial T_1^2} \right) + \mathcal{O}(\epsilon^3) \end{aligned}$$

and thus the perturbed problem becomes

$$\frac{\partial^2 y}{\partial T_0^2} + 2\epsilon \frac{\partial^2 y}{\partial T_0 \partial T_1} + \epsilon^2 \left( \frac{\partial^2 y}{\partial T_0 \partial T_2} + \frac{\partial^2 y}{\partial T_1^2} \right) + \epsilon \left( \frac{\partial y}{\partial T_0} + \epsilon \frac{\partial y}{\partial T_1} + \epsilon^2 \frac{\partial y}{\partial T_2} \right) + y = 0$$

where we omit terms of  $\mathcal{O}(\epsilon^3)$ . To get an asymptotic approximation for  $y$  in the form

$$\tilde{y}(t) = y_0 + \epsilon y_1 + \epsilon^2 y_2 \approx y(t)$$

we plug  $\tilde{y}(t)$  back into the perturbed equation and get

$$\begin{aligned} \frac{\partial^2 y_0}{\partial T_0^2} + \epsilon \frac{\partial^2 y_1}{\partial T_0^2} + \epsilon^2 \frac{\partial^2 y_2}{\partial T_0^2} + 2\epsilon \frac{\partial^2 y_0}{\partial T_0 \partial T_1} + 2\epsilon^2 \frac{\partial^2 y_0}{\partial T_0 \partial T_2} + 2\epsilon^2 \frac{\partial^2 y_1}{\partial T_0 \partial T_1} + \epsilon^2 \frac{\partial^2 y_0}{\partial T_1^2} + \\ 2\epsilon \left( \frac{\partial y_0}{\partial T_0} + \epsilon \frac{\partial y_0}{\partial T_1} + \epsilon \frac{\partial y_1}{\partial T_0} \right) + y_0 + \epsilon y_1 + \epsilon^2 y_2 = 0 \end{aligned}$$

where we omit terms of  $\mathcal{O}(\epsilon^3)$ . For the sake of brevity only first-order approximation will be done, but second order is done in a similar fashion. Gathering terms we get

$$\begin{aligned} \mathcal{O}(1) : \frac{\partial^2 y_0}{\partial T_0^2} + y_0 &= 0 \\ \mathcal{O}(\epsilon) : \frac{\partial^2 y_1}{\partial T_0^2} + y_1 + 2 \frac{\partial^2 y_0}{\partial T_0 \partial T_1} + 2 \frac{\partial y_0}{\partial T_0} &= 0 \end{aligned}$$

The initial conditions are  $y_0 = 1 \Rightarrow \frac{\partial y_0}{\partial T_0} = 0$  and  $y_1 = 0 \Rightarrow \frac{\partial y_1}{\partial T_0} = -\frac{\partial y_0}{\partial T_1}$  for  $T_0 = T_1 = 0$ . The general solution and its derivatives becomes

$$\begin{aligned} y_0 &= A(T_1)\cos(T_0) + B(T_0)\sin(T_0) \\ \frac{\partial y_0}{\partial T_0} &= -A(T_1)\sin(T_0) + B(T_0)\cos(T_0) \\ \frac{\partial^2}{\partial T_0 \partial T_1} &= -\sin(T_0)\frac{\partial A}{\partial T_1} + \cos(T_0)\frac{\partial B}{\partial T_1} \end{aligned}$$

Plugging this back in yields

$$\frac{\partial^2 y_1}{\partial T_0^2} + y_1 = 2\left(\frac{\partial A}{\partial T_1} + A\right)\sin(T_0) - 2\left(\frac{\partial B}{\partial T_1} + B\right)\cos(T_0)$$

The coefficients of  $\cos(T_0)$  and  $\sin(T_0)$  must vanish so we end up getting

$$\frac{\partial A}{\partial T_1} + A = 0, \quad \frac{\partial B}{\partial T_1} + B = 0 \quad \Rightarrow \quad A = ae^{-T_1}, \quad B = be^{-T_1}$$

Substituting this back in we get

$$y_0 = ae^{-T_1}\cos(T_0) + be^{-T_1}\sin(T_0)$$

where, from initial conditions, we get  $a = 1$  and  $b = 0$ . Therefore we obtain the perturbed solution

$$x = e^{-T_1}\cos(T_0) + \mathcal{O}(\epsilon) \quad \rightarrow \quad x = e^{-\epsilon t}\cos(t) + \mathcal{O}(\epsilon)$$

which is valid for time of order  $1/\epsilon$ .

### Problem 3

**NOTE:** For the sake of brevity, some algebra is omitted in the writeup of Problem 3.

Find an approximate formula for each root of the following algebraic equations (here  $0 < \epsilon \ll 1$ ):

(a)  $xe^{-x} = \epsilon$

For the regularly perturbed root  $x = 0$  we let

$$x \approx x_0 + \epsilon x_1 + \epsilon^2 x_2$$

Plugging back into the original system we get

$$(x_0 + \epsilon x_1 + \epsilon^2 x_2) - \epsilon(1 + x_0 + \epsilon x_1 + \epsilon^2 x_2)$$

where we used the Taylor series expansion for  $e^x$ . Collecting terms we get

$$\begin{aligned} \mathcal{O}(1) : x - 0 &= 0 \\ \mathcal{O}(\epsilon) : x_1 - 1 &= 0 \Rightarrow x_1 = 1 \\ \mathcal{O}(\epsilon^2) : x_2 - x_1 &= 0 \Rightarrow x_2 = 1 \end{aligned}$$

where we get

$$x = \epsilon + \epsilon^2$$

for the root  $x = 0$ . Now considering the other perturbed roots, we take the log of both sides to get

$$x = \log(x) + \log\left(\frac{1}{\epsilon}\right)$$

where  $x$  is the perturbed root. We can treat the perturbed root as a fixed point and apply fixed point iteration to determine  $x$

$$\begin{aligned} x_1 &= \log\left(\frac{1}{\epsilon}\right) \\ x_2 &= \log\left(\frac{1}{\epsilon}\right) + \log\left(\log\left(\frac{1}{\epsilon}\right)\right) \\ x_3 &= \log\left(\frac{1}{\epsilon}\right) + \log\left(\log\left(\frac{1}{\epsilon}\right) + \log\left(\log\left(\frac{1}{\epsilon}\right)\right)\right) \\ &= \log\left(\frac{1}{\epsilon}\right) + \log\left(\log\left(\frac{1}{\epsilon}\right)\right) + \frac{\log\left(\log\left(\frac{1}{\epsilon}\right)\right)}{\log\left(\frac{1}{\epsilon}\right)} + \mathcal{O}\left(\left(\frac{\log\left(\log\left(\frac{1}{\epsilon}\right)\right)}{\log\left(\frac{1}{\epsilon}\right)}\right)^2\right) \end{aligned}$$

As  $\epsilon \rightarrow 0$  we get that the expression converges.

$$(b) \ x^3 - x + \epsilon = 0$$

Unperturbed Roots:  $x^3 - x = 0 \Rightarrow x = 0, \pm 1$

Taylor Expanding these roots we get

$$\begin{aligned} x_{+1} &= 1 + \epsilon x_1 + \epsilon^2 x_2 + \mathcal{O}(\epsilon^3) \\ x_{-1} &= -1 + \epsilon x_1 + \epsilon^2 x_2 + \mathcal{O}(\epsilon^3) \\ x_0 &= \epsilon x_1 + \epsilon^2 x_2 + \mathcal{O}(\epsilon^3) \end{aligned}$$

Plugging back into the perturbed problem we get the following

$$\begin{aligned} x_{+1}^3 - x_{+1} - \epsilon &= (1 + \epsilon x_1 + \epsilon^2 x_2 + \mathcal{O}(\epsilon^3))^3 - (1 + \epsilon x_1 + \epsilon^2 x_2 + \mathcal{O}(\epsilon^3)) + \epsilon \\ &= (1 + 3\epsilon x_1 + \epsilon^2(3x_1^2 + 3x_2 + \mathcal{O}(\epsilon^3))) - (1 + \epsilon x_1 + \epsilon^2 x_2 + \mathcal{O}(\epsilon^3)) + \epsilon \\ \mathcal{O}(1) : 1 - 1 &= 0 \\ \mathcal{O}(\epsilon) : 3x_1 - x_1 + 1 &= 0 \Rightarrow x_1 = -\frac{1}{2} \\ \mathcal{O}(\epsilon^2) : 3x_1^2 + 3x_2 - x_2 &= 0 \Rightarrow x_2 = -\frac{3}{8} \end{aligned}$$

$$\begin{aligned} x_{-1}^3 - x_{-1} - \epsilon &= (-1 + \epsilon x_1 + \epsilon^2 x_2 + \mathcal{O}(\epsilon^3))^3 - (-1 + \epsilon x_1 + \epsilon^2 x_2 + \mathcal{O}(\epsilon^3)) + \epsilon \\ &= (-1 + 3\epsilon x_1 + \epsilon^2(3x_2 - 3x_1^2 + \mathcal{O}(\epsilon^3))) - (-1 + \epsilon x_1 + \epsilon^2 x_2 + \mathcal{O}(\epsilon^3)) + \epsilon \\ \mathcal{O}(1) : -1 + 1 &= 0 \\ \mathcal{O}(\epsilon) : 3x_1 - x_1 + 1 &= 0 \Rightarrow x_1 = -\frac{1}{2} \\ \mathcal{O}(\epsilon^2) : -3x_1^2 + 3x_2 - x_2 &= 0 \Rightarrow x_2 = \frac{3}{8} \end{aligned}$$

$$\begin{aligned} x_0^3 - x_0 - \epsilon &= (\epsilon x_1 + \epsilon^2 x_2 + \mathcal{O}(\epsilon^3))^3 - (\epsilon x_1 + \epsilon^2 x_2 + \mathcal{O}(\epsilon^3)) + \epsilon \\ &= \mathcal{O}(\epsilon^3) - (\epsilon x_1 + \epsilon^2 x_2 + \mathcal{O}(\epsilon^3)) + \epsilon \\ \mathcal{O}(1) : &\text{none} \\ \mathcal{O}(\epsilon) : -x_1 + 1 &= 0 \Rightarrow x_1 = 1 \\ \mathcal{O}(\epsilon^2) : x_2 &= 0 \end{aligned}$$

Putting all this together we get the approximation root formulas

$$\begin{aligned} x_{+1} &= 1 - \frac{\epsilon}{2} - \frac{3\epsilon^2}{8} \\ x_{-1} &= -1 - \frac{\epsilon}{2} + \frac{3\epsilon^2}{8} \\ x_0 &= \epsilon \end{aligned}$$

$$(c) \epsilon x^3 - x + 1 = 0$$

We have a singularly perturbed system so we will apply scaling by letting  $x = \delta y$  to get  $\epsilon \delta^3 y^3 - \delta y + 1 = 0$ . Now we must balance

$$\text{Balance I and II : } \epsilon \delta^3 \sim \delta \Rightarrow \delta \sim \frac{1}{\sqrt{\epsilon}}$$

$$\text{Balance I and III : } \epsilon \delta^3 \sim 1 \Rightarrow \delta \sim \frac{1}{\sqrt[3]{\epsilon}}$$

$$\text{Balance II and III : } \delta \sim 1$$

where balancing I and II is what we want since we arrive at the equation  $y^3 - y + \sqrt{\epsilon} = 0$ . Using this formula we get the Taylor Series expansions

$$\begin{aligned} x_{+1} &= 1 + \sqrt{\epsilon}x_1 + \epsilon x_2 \\ x_{-1} &= -1 + \sqrt{\epsilon}x_1 + \epsilon x_2 \\ x_0 &= \sqrt{\epsilon}x_1 + \epsilon x_2 \end{aligned}$$

Plugging this back into the equation we get

$$\begin{aligned} y_{+1}^3 - y_{+1} + \sqrt{\epsilon} &= (1 + \sqrt{\epsilon}x_1 + \epsilon x_2)^3 - (1 + \sqrt{\epsilon}x_1 + \epsilon x_2) + \sqrt{\epsilon} \\ &= (1 + 3\sqrt{\epsilon}x_1 + \epsilon(3x_2 + 3x_1^2) + \mathcal{O}(\epsilon^{3/2})) - (1 + \sqrt{\epsilon}x_1 + \epsilon x_2 + \mathcal{O}(\epsilon^{3/2})) + \sqrt{\epsilon} \\ \mathcal{O}(1) : 1 - 1 &= 0 \\ \mathcal{O}(\sqrt{\epsilon}) : x_1 + 1 &= 0 \Rightarrow x_1 = -1 \\ \mathcal{O}(\epsilon) : 2x_2 + 3x_1^2 &= 0 \Rightarrow x_2 = -\frac{3}{2} \end{aligned}$$

$$\begin{aligned} y_{-1}^3 - y_{-1} + \sqrt{\epsilon} &= (-1 + \sqrt{\epsilon}x_1 + \epsilon x_2)^3 - (-1 + \sqrt{\epsilon}x_1 + \epsilon x_2) + \sqrt{\epsilon} \\ &= (1 + 3\sqrt{\epsilon}x_1 + \epsilon(3x_2 - 3x_1^2) + \mathcal{O}(\epsilon^{3/2})) - (-1 + \sqrt{\epsilon}x_1 + \epsilon x_2 + \mathcal{O}(\epsilon^{3/2})) + \sqrt{\epsilon} \\ \mathcal{O}(1) : 1 - 1 &= 0 \\ \mathcal{O}(\sqrt{\epsilon}) : x_1 + 1 &= 0 \Rightarrow x_1 = -1 \\ \mathcal{O}(\epsilon) : 2x_2 - 3x_1^2 &= 0 \Rightarrow x_2 = \frac{3}{2} \end{aligned}$$

$$\begin{aligned} y_0^3 - y_0 + \sqrt{\epsilon} &= (\sqrt{\epsilon}x_1 + \epsilon x_2 + \mathcal{O}(\epsilon^{3/2}))^3 - (\sqrt{\epsilon}x_1 + \epsilon x_2 + \mathcal{O}(\epsilon^{3/2})) + \sqrt{\epsilon} \\ &= \sqrt{\epsilon}(-x_1 + 1) - \epsilon x_2 + \epsilon^{3/2}x_1 + \epsilon^2(x_1^2x_2 + x_1x_2) \\ \mathcal{O}(1) : \text{none} \\ \mathcal{O}(\sqrt{\epsilon}) : -x_1 + 1 &= 0 \Rightarrow x_1 = 1 \\ \mathcal{O}(\epsilon) : x_2 &= 0 \Rightarrow x_2 = 0 \end{aligned}$$

Putting all this together we get the following

$$\begin{aligned}x_{+1} &= \delta y_{+1} = \epsilon^{-1/2} - \frac{1}{2} - \frac{3}{8}\epsilon^{1/2} \\x_{-1} &= \delta y_{-1} = \epsilon^{-1/2} + \frac{1}{2} + \frac{3}{8}\epsilon^{1/2} \\x_0 &= \delta y_0 = 1 + \epsilon\end{aligned}$$

$$(d) (1 - \epsilon)x^2 - 2x + 1 = 0$$

Trying to do our regular Taylor series expansion, we let

$$x \approx x_0 + \epsilon x_1 + \epsilon^2 x_2$$

we end up with

$$(1 - \epsilon)(x_0 + \epsilon x_1 + \epsilon^2 x_2)^2 - 2(x_0 + \epsilon x_1 + \epsilon^2 x_2) + 1$$

Collecting terms we get

$$\begin{aligned}\mathcal{O}(1) : x_0^2 - 2x_0 + 1 &= 0 \Rightarrow x_0 = 1 \\ \mathcal{O}(\epsilon) : 2x_0x_1 + x_0^2 - 2x_1 &= 0 \Rightarrow 1 = 0\end{aligned}$$

where we arrive at a contradiction. This is due to the repeated roots, but if we look at the roots with respect to the perturbation we get

$$x = \frac{1 \pm \sqrt{\epsilon}}{1 - \epsilon}$$

where we see that it scales with respect to  $\sqrt{\epsilon}$ . Expanding with respect to this we get

$$(1 - \epsilon)(x_0 + \sqrt{\epsilon}x_1 + \epsilon x_2)^2 - 2(x_0 + \sqrt{\epsilon}x_1 + \epsilon x_2) + 1 = 0$$

Collect terms again we get

$$\begin{aligned}\mathcal{O}1 : x_0^2 - 2x_0 + 1 &= 0 \Rightarrow x_0 = 1 \\ \mathcal{O}(\sqrt{\epsilon}) : 2x_0x_1 - 2x_1 &= 0 \Rightarrow \text{Inconclusive} \\ \mathcal{O}(\epsilon) : -x_0^2 + x_1^2 + 2x_0x_2 - 2x_2 &= 0 \Rightarrow x_1 = \pm 1\end{aligned}$$

From this we have that the perturbed roots are

$$x_{\pm 1} = 1 \pm \sqrt{\epsilon} + \mathcal{O}(\epsilon)$$

$$(e) \quad \epsilon(x^2 + x) + 1 = 0$$

This system is singularly perturbed so we must do scaling ( $x = \delta y$ ) and balancing

$$\epsilon \delta^2 y^2 + \epsilon \delta y + 1 = 0$$

Balancing term I with term II we get  $\delta \sim 1$  which leads to the same problem as we originally had. Balancing term I and III we get  $\delta \sim \frac{1}{\sqrt{\epsilon}}$  which works out. We end up getting

$$y^2 + \sqrt{\epsilon}y + 1 = 0$$

It is important to note if we try balancing term II and III we get  $\delta \sim \frac{1}{\epsilon}$  which leads to another singular perturbed system. Applying the usual expansion (w.r.t to  $\sqrt{\epsilon}$  instead of  $\epsilon$ ) we get

$$(y_0 + \sqrt{\epsilon}y_1 + \epsilon y_2)^2 + \sqrt{\epsilon}(y_0 + \sqrt{\epsilon}y_1 + \epsilon y_2) + 1$$

Collecting terms we get

$$\begin{aligned} \mathcal{O}(1) : y_0^2 + 1 &= 0 \Rightarrow y_0 = \pm i \\ \mathcal{O}(\sqrt{\epsilon}) : 2y_0y_1 + y_0 &= 0 \Rightarrow y_1 = \frac{1}{2} \\ \mathcal{O}(\epsilon) : y_1^2 + 2y_0y_2 + y_1 &= 0 \Rightarrow y_2 = \frac{i}{8} \end{aligned}$$

Therefore fore the perturbed roots we get

$$y_{\pm} = \pm i - \frac{1}{2}\sqrt{\epsilon} \pm \frac{i}{8} + \mathcal{O}(\epsilon^{3/2})$$



$$(f) \ x^2 - 1 = \epsilon x$$

Applying our regular Taylor series expansion we get

$$x \approx x_0 + \epsilon x_1 + \epsilon^2 x_2$$

and plugging back into the equation

$$(x_0 + \epsilon x_1 + \epsilon^2 x_2)^2 - \epsilon(x_0 + \epsilon x_1 + \epsilon^2 x_2) - 1$$

Collecting terms we get

$$\mathcal{O}(1) : x_0^2 - 1 = 0 \Rightarrow x_0 = \pm 1$$

$$\mathcal{O}(\epsilon) : 2x_0x_1 - x_0 = 0 \Rightarrow x_1 = \frac{1}{2}$$

$$\mathcal{O}(\epsilon^2) : x_1^2 + 2x_0x_2 - x_1 = 0 \Rightarrow x_2 = \pm \frac{1}{8}$$

Therefore the perturbed roots are

$$x_{\pm 1} = \pm 1 + \frac{1}{2}\epsilon \pm \frac{1}{8}\epsilon^2$$

$$(g) \ x^2 - 1 = \epsilon e^x$$

Taking the log of both sides we get

$$\log(x^2 - 1) = \log(\epsilon) + x \quad \Rightarrow \quad x = \log(x^2 - 1) + \log\left(\frac{1}{\epsilon}\right)$$

for which we can apply fixed point iteration. Taking when only the  $\epsilon$  term is present to be our initial guess we get

$$\begin{aligned} x_1 &= \log\left(\frac{1}{\epsilon}\right) \\ x_2 &= \log\left(\log\left(\frac{1}{\epsilon}\right)^2 - 1\right) + \log\left(\frac{1}{\epsilon}\right) \end{aligned}$$

We can see from the above that if we keep expanding we get  $\mathcal{O}(\log(\frac{1}{\epsilon})^3)$ . The above iteration converges as  $\epsilon \rightarrow 0$ .

$$(h) \ x^2 - 4 = \epsilon \ln x$$

Just like question 3(a), we will use fixed point iteration

$$\begin{aligned} x_n &= 2 + \frac{\epsilon \ln(x)}{x + 2} \Rightarrow \\ x_1 &= 2 \\ x_2 &= 2 + \frac{\epsilon \ln(2)}{4} \\ x_3 &= 2 + \frac{\epsilon \ln\left(2 + \frac{\epsilon \ln(2)}{4}\right)}{4 + \frac{\epsilon \ln(2)}{4}} \\ &= 2 + \frac{\epsilon}{4} \left(1 - \frac{\epsilon \ln(2)}{16}\right) (\ln(2) + \ln(1 + \frac{\epsilon \ln(2)}{8})) \end{aligned}$$

The other perturbed root is near zero. We can see that the above iteration converges as  $\epsilon \rightarrow 0$ .

$$(i) \ x^2 - 2\epsilon x - \epsilon = 0$$

Trying to do our regular Taylor series expansion, we let

$$x \approx x_0 + \epsilon x_1 + \epsilon^2 x_2$$

we end up with

$$(x_0 + \epsilon x_1 + \epsilon^2 x_2)^2 - 2\epsilon(x_0 + \epsilon x_1 + \epsilon^2 x_2) - \epsilon$$

Collecting terms we get

$$\mathcal{O}(1) : x_0^2 = 0 \Rightarrow x_0 = 0$$

$$\mathcal{O}(\epsilon) : 2x_0x_1 + -2x_0 - 1 \Rightarrow -1 = 0$$

where we arrive at a contradiction. This is due to the repeated roots, but if we look at the roots with respect to the perturbation we get

$$x = 1 \pm 2\sqrt{\epsilon^2 + \epsilon} \approx 1 \pm 2\sqrt{\epsilon}$$

where we see that it scales with respect to  $\sqrt{\epsilon}$ . Expanding with respect to this we get

$$(x_0 + \sqrt{\epsilon}x_1 + \epsilon x_2)^2 - 2\epsilon(x_0 + \sqrt{\epsilon}x_1 + \epsilon x_2) - \epsilon = 0$$

Collect terms again we get

$$\mathcal{O}1 : x_0^2 = 0 \Rightarrow x_0 = 0$$

$$\mathcal{O}(\sqrt{\epsilon}) : 2\sqrt{\epsilon}x_0x_1 = 0 \Rightarrow \text{Inconclusive}$$

$$\mathcal{O}(\epsilon) : 2\epsilon x_0x_2 + x_1^2 + 2x_0 - 1 = 0 \Rightarrow x_1 = \pm 1$$

$$\mathcal{O}(\epsilon^{3/2}) : 2x_1x_2 - 2x_1 = 0 \Rightarrow x_2 = 1$$

From this we have that the perturbed roots are

$$x = \pm\sqrt{\epsilon} + \epsilon$$

$$(j) \epsilon x^5 + x^2 - 2x + 1 = 0$$

This is singularly perturbed so we apply scaling by letting  $x = \delta y$  and we get

$$\epsilon \delta^5 y^5 + \delta^2 y^2 - 2\delta y + 1 = 0$$

Balancing the term I against term II we get  $\delta \sim e^{-1/3}$ . It is important to note that when balancing term I against term II and term III we get

$$2^{5/4} y^5 + \frac{\sqrt{2} y^2}{\epsilon^{1/4}} + 2^{5/4} + \epsilon^{1/4} = 0, \quad y^5 + \epsilon^{-2/5} y^2 - 2\epsilon^{-1/5} y + 1 = 0$$

which blow up at small  $\epsilon$ . Also trying to balancing the other terms against each other leads to  $\delta$  being proportional to some scalar, which just leads to the original problem essentially. Now, plugging back in and expanding we get

$$(y_0 + \epsilon^{1/3} y_1 + \epsilon^{2/3} y_2)^5 + (y_0 + \epsilon^{1/3} y_1 + \epsilon^{2/3} y_2)^2 - 2\epsilon^{1/3} (y_0 + \epsilon^{1/3} y_1 + \epsilon^{2/3} y_2) + 1$$

Collecting terms we get

$$\begin{aligned} \mathcal{O}(1) : y_0^5 + y_2 = 0 &\Rightarrow y_0 = 0, -1 \text{ ( we will use } y_0 = 0) \\ \mathcal{O}(\epsilon^{1/3}) : y_0^4 y_1 + 2y_0 y_1 - 2y_0 = 0 &\Rightarrow y_1 = -\frac{2}{3} \\ \mathcal{O}(\epsilon^{2/3}) : 5y_0^4 y_2 + 10y_1^2 y_0^3 + 2y_0 y_1 + y_1^2 + y_1 + 1 &\Rightarrow y_2 = \frac{11}{9} \end{aligned}$$

Using this and plugging back into  $x = \delta y$  we get that the perturbed root is

$$x = \delta y = \frac{1}{\epsilon^{1/3}} \left( -1 - \frac{2}{3} \epsilon^{1/3} + \frac{11}{9} \epsilon^{2/3} \right) = -\frac{1}{\epsilon^{1/3}} - \frac{2}{3} + \frac{11}{9} \epsilon^{1/3}$$