Prerequisites

Leibnitz rule for constants $-\infty < a, b < +\infty$:

$$\frac{d}{dx}\left(\int_{a}^{b} f(x,t)dt\right) = \int_{a}^{b} \frac{\partial}{\partial x} f(x,t)dt$$

Additionally, let at least one of the following conditions hold:

- (i) f(x,t) is measurable and nonnegative
- (ii) $\int_a^b |f(x,t)| dt$ is finite

then we can switch the order of integration according to Tonelli/Fubini, respectively.

Note: In our case, at least one of (i), (ii) is nearly always satisfied. Think about why.

Lastly, we need (*) Theorem 1 concerning uniform convergence from these notes:

http://www.math.ucla.edu/~tao/resource/general/131bh.1.03s/week45.pdf

Let D be the domain of \mathbf{x} .

$$E[y(\mathbf{x}) - t] = \int_{t_1}^{t_2} \int_{D} (y(\mathbf{x}) - t)^2 p(\mathbf{x}, t) d\mathbf{x} dt$$

$$= \underbrace{\int_{D} \int_{t_1}^{t_2} (y(\mathbf{x}) - t)^2 p(\mathbf{x}, t) dt d\mathbf{x}}_{Fubini/Tonelli}$$

$$= \int_{D} p(\mathbf{x}) \int_{t_1}^{t_2} (y(\mathbf{x}) - t)^2 p(t|\mathbf{x}) dt d\mathbf{x}$$

Minimizing the loss cumulative loss for all **t** equals minimizing the loss for each t_i separately(**). Note: Inside the interior integral, **x** is constant. Let $y(\mathbf{x}) = z$:

$$\frac{\partial}{\partial z} \int_{t_1}^{t_2} (z - t)^2 p(t | \mathbf{x}) dt = \underbrace{\int_{t_1}^{t_2} \frac{\partial}{\partial z} (z - t)^2 p(t | \mathbf{x}) dt}_{Leibnitz \ rule}$$
$$= 2 \int_{t_1}^{t_2} (z - t) p(t | \mathbf{x}) dt$$

We can now solve for z:

$$2\int_{t_{1}}^{t_{2}} (z-t)p(t|\mathbf{x})dt = 0$$

$$\Leftrightarrow z\int_{t_{1}}^{t_{2}} p(t|\mathbf{x})dt = \int_{t_{1}}^{t_{2}} tp(t|\mathbf{x})dt$$

$$\Leftrightarrow y(\mathbf{x}) = E[t|\mathbf{x}]$$

According to the Leibnitz rule, this only holds for finite limits t_1 , t_2 . To extend this proof to the infinite domain, we construct the sequence:

$$f'_n = \frac{\partial}{\partial z} \int_{-n}^{n} (z - t)^2 p(t|\mathbf{x}) dt$$

Because probabilities sum to one, if n tends to infinity, $p(t|\mathbf{x})$ vanishes for most t_i . The $(z-t)^2$ will not compensate that, as we required $E[y(\mathbf{x})-t]$ to be finite earlier.

Accordingly, $\lim_{n\to\infty} f_n'$ converges to g uniformly. Due to (*), this means the functions f_n converge uniformly to f, with f'=g.

Spoken plainly, this means if we have an infinite domain \mathbb{R} , we can approximate the solution arbitrarily close by increasing $[t_1, t_2]$.

(**) Only because the x_i are independent. In our situation this is the case, otherwise we would also have to integrate over all possibilities $p(x|x_i, ... x_0)$.

Auxiliary calculation

The area of a single infinitesimal d-dimensional piece of $f(r, \mathbf{\theta})$ is $r^{d-1}d\theta_1 \cdot ... \cdot d\theta_{d-1} \cdot dr$. This is trivially an d-dimensional extension of the two-dimensional case shown below:





Additionally, to convert a function $f(r, \theta)$ from hyperspherical coordinates into cartesian coordinates $f(\mathbf{x})$, we use the following trigonometric conversion:

$$\mathbf{x} = \begin{pmatrix} x_1 \\ x_2 \\ \dots \\ x_{d-1} \\ x_d \end{pmatrix} = \begin{pmatrix} r \cos \theta_1 \\ r \sin \theta_1 \cos \theta_2 \\ \dots \\ r \sin \theta_1 \sin \theta_2 \sin \theta_3 \dots \sin \theta_{d-3} \sin \theta_{d-2} \cos \theta_{d-1} \\ r \sin \theta_1 \sin \theta_2 \sin \theta_3 \dots \sin \theta_{d-3} \sin \theta_{d-2} \sin \theta_{d-1} \end{pmatrix},$$

i.e.

 $f(\mathbf{x}) = f(r\cos\theta_1, r\sin\theta_1\cos\theta_2, \dots, r\sin\theta_1\sin\theta_2\sin\theta_3 \dots \sin\theta_{d-3}\sin\theta_{d-2}\sin\theta_{d-1}).$ Lastly, let

$$\begin{split} K(\boldsymbol{\theta}) &= \cos\theta_1^{\ 2} + (\sin\theta_1\cos\theta_2)^2 + \cdots \\ &\quad + (\sin\theta_1\sin\theta_2\sin\theta_3 \dots \sin\theta_{d-3}\sin\theta_{d-2}\sin\theta_{d-1})^2. \end{split}$$

Note: θ describe points on the unit d-sphere, so it is no surprise that $||K(\theta)||^2 = 1$ for all θ , because the radius of the unit sphere is 1.

The centered sphere is described by $B_0(r) \coloneqq \{x_1^2 + \dots + x_d^2 \le r^2 : x_i \in \mathbb{R}\}$.

Armed with this knowledge, $P(B_0(r))$ becomes:

$$\begin{split} \int_0^r \int_0^{2\pi} \dots \int_0^\pi \underbrace{(2\pi)^{-\frac{D}{2}} e^{-\frac{1}{2}||\mathbf{x}||^2}}_{pdf \; normal \; dist.} \underbrace{r^{d-1} d\theta_1 \dots d\theta_{d-1} dr}_{infinitismal \; area} &= \int_0^r \int_0^{2\pi} \dots \int_0^\pi (2\pi)^{-\frac{D}{2}} e^{-\frac{1}{2}r^2} \underbrace{K(\mathbf{\theta})}_{=1} r^{d-1} d\theta_1 \dots d\theta_{d-1} dr, \\ &= \int_0^r r^{d-1} (2\pi)^{-\frac{D}{2}} e^{-\frac{1}{2}r^2} \underbrace{\int_0^{2\pi} \dots \int_0^\pi d\theta_1 \dots d\theta_{d-1} dr}_{Surface \; Area \; unit \; n-sphere \; S_D} \\ &= \int_0^r S_D r^{d-1} (2\pi)^{-\frac{D}{2}} e^{-\frac{1}{2}r^2} dr. \end{split}$$

Ergo
$$p(r)dr = S_D r^{d-1} (2\pi)^{-\frac{D}{2}} e^{-\frac{1}{2}r^2}.$$

Now we are looking for the maximum density $\max_{x} p(r)$:

$$\frac{d}{dr} \left[\log r^{d-1} + \log e^{-\frac{1}{2}r^2} \right] = \frac{(d-1)r^{d-2}}{r^{d-1}} - r = 0.$$

$$\Leftrightarrow$$
 $(d-1)=r^2$.

Because radii are non-negative, we have a maximum at $\sqrt{d-1}$.

Now if we set $||\mathbf{x}|| = \sqrt{d-1}$, we get

$$\frac{p(\mathbf{x})}{p(0)} = \frac{(2\pi)^{-\frac{D}{2}} e^{-\frac{1}{2} ||\mathbf{x}||^2}}{(2\pi)^{-\frac{D}{2}} e^{-\frac{1}{2} ||\mathbf{0}||^2}} = \frac{e^{-\frac{1}{2}(d-1)}}{e^{-\frac{1}{2}}} = e^{-\frac{d}{2}}.$$

4A-3.

Let $L(\mathbf{w})$:

$$L(\mathbf{w}) = \frac{1}{2} \sum_{n=1}^{N} (t_n - \mathbf{w}^T \boldsymbol{\phi}(x_n))^2 + \frac{\lambda}{2} \mathbf{w}^T \mathbf{w}$$

Derivative of $\ln L(\mathbf{w})$ with respect to w_i :

$$\frac{\partial}{\partial w_i} \ln L(\mathbf{w}) = \sum_{n=1}^N \phi_i(x_n) \left(t_n - \mathbf{w}^T \boldsymbol{\phi}(x_n) \right) + \lambda w_i$$

Conversion to matrix/vector operations:

$$\sum_{n=1}^{N} \phi_i(x_n) \left(t_n - \mathbf{w}^T \boldsymbol{\phi}(x_n) \right) + \lambda w_i = \operatorname{col}_i(\boldsymbol{\Phi})^T (\mathbf{t} - \boldsymbol{\Phi} \mathbf{w}) + \lambda w_i$$

Generalized for all w:

$$\frac{\partial}{\partial \mathbf{w}} \ln L(\mathbf{w}) = \mathbf{\Phi}^T (\mathbf{t} - \mathbf{\Phi} \mathbf{w}) + \lambda \mathbf{w}$$

Setting zero and solving for w:

$$\Phi^{T}(\mathbf{t} - \Phi \mathbf{w}) + \lambda \mathbf{w} = 0$$

$$\Leftrightarrow \Phi^{T}\mathbf{t} + (-\Phi^{T}\Phi + \lambda \mathbf{I})\mathbf{w} = 0$$

$$\Leftrightarrow \mathbf{w} = (-\lambda \mathbf{I} + \Phi^{T}\Phi)^{+}\Phi^{T}\mathbf{t}$$

As usual, A^+ denotes the Penrose pseudo inverse.

Because λ can be an arbitrary normative factor, it is also possible to write:

$$\mathbf{w} = (\lambda \mathbf{I} + \mathbf{\Phi}^T \mathbf{\Phi})^+ \mathbf{\Phi}^T \mathbf{t}$$

Pictures from Python:

