Prerequisites

Leibnitz rule for constants $-\infty < a, b < +\infty$:

$$\frac{d}{dx}\left(\int_{a}^{b} f(x,t)dt\right) = \int_{a}^{b} \frac{\partial}{\partial x} f(x,t)dt$$

Additionally, let at least one of the following conditions hold:

- (i) f(x,t) is measurable and nonnegative
- (ii) $\int_a^b |f(x,t)| dt$ is finite

then we can switch the order of integration according to Tonelli/Fubini, respectively.

Note: In our case, at least one of (i), (ii) is nearly always satisfied. Think about why.

Lastly, we need (*) Theorem 1 concerning uniform convergence from these notes:

http://www.math.ucla.edu/~tao/resource/general/131bh.1.03s/week45.pdf

Let D be the domain of \mathbf{x} .

$$E[y(\mathbf{x}) - t] = \int_{t_1}^{t_2} \int_{D} (y(\mathbf{x}) - t)^2 p(\mathbf{x}, t) d\mathbf{x} dt$$

$$= \underbrace{\int_{D} \int_{t_1}^{t_2} (y(\mathbf{x}) - t)^2 p(\mathbf{x}, t) dt d\mathbf{x}}_{Fubini/Tonelli}$$

$$= \int_{D} p(\mathbf{x}) \int_{t_1}^{t_2} (y(\mathbf{x}) - t)^2 p(t|\mathbf{x}) dt d\mathbf{x}$$

Minimizing the loss cumulative loss for all **t** equals minimizing the loss for each t_i separately(**). Note: Inside the interior integral, **x** is constant. Let $y(\mathbf{x}) = z$:

$$\frac{\partial}{\partial z} \int_{t_1}^{t_2} (z - t)^2 p(t | \mathbf{x}) dt = \underbrace{\int_{t_1}^{t_2} \frac{\partial}{\partial z} (z - t)^2 p(t | \mathbf{x}) dt}_{Leibnitz \, rule}$$
$$= 2 \int_{t_1}^{t_2} (z - t) p(t | \mathbf{x}) dt$$

We can now solve for z:

$$2\int_{t_{1}}^{t_{2}} (z-t)p(t|\mathbf{x})dt = 0$$

$$\Leftrightarrow z\int_{t_{1}}^{t_{2}} p(t|\mathbf{x})dt = \int_{t_{1}}^{t_{2}} tp(t|\mathbf{x})dt$$

$$\Leftrightarrow y(\mathbf{x}) = E[t|\mathbf{x}]$$

According to the Leibnitz rule, this only holds for finite limits t_1, t_2 . To extend this proof to the infinite domain, we construct the sequence:

$$f'_n = \frac{\partial}{\partial z} \int_{-n}^{n} (z - t)^2 p(t|\mathbf{x}) dt$$

Because probabilities sum to one, if n tends to infinity, $p(t|\mathbf{x})$ vanishes for most t_i . The $(z-t)^2$ will not compensate that, as we required $E[y(\mathbf{x})-t]$ to be finite earlier.

Accordingly, $\lim_{n\to\infty} f_n{'}$ converges to g uniformly. Due to (*), the functions f_n converge uniformly to f, with f'=g.

Spoken plainly, this means if we have an infinite domain \mathbb{R} , we can approximate the solution arbitrarily close by increasing $[t_1, t_2]$.

(**) Only because the x_i are independent. In our situation this is the case, otherwise we would also have to integrate over all possibilities $p(x|x_i, ... x_0)$.

Auxiliary calculation

Using the Jacobian integral substitution, the area of an infinitesimal d-dimensional volume element is

$$\begin{split} d^n V &= \left| \det \frac{\partial (x_i)}{\partial (r,\phi_j)} \right| dr d\theta_1 \dots d\theta_{n-1} \\ d^n V &= \left| \det \frac{\partial (x_i)}{\partial (r,\phi_j)} \right| dr d\theta_1 \dots d\theta_{n-1} \\ &= r^{n-1} \underbrace{\sin^{n-2} \theta_1 \sin^{n-3} \theta_2 \cdot \dots \cdot \sin \theta_{n-2}}_{g(\theta_1,\dots,\theta_{n-1})} dr d\theta_1 \dots d\theta_{n-1} \\ &= r^{n-1} g(\mathbf{\theta}) dr d\theta_1 \dots d\theta_{n-1} \end{split}$$

Above can be seen here:

https://en.wikipedia.org/wiki/N-sphere#Spherical coordinates

Furthermore: We have $||\mathbf{x}||^2 = r^2$ for all $\mathbf{0}$, because the squared length of a coordinate point is r^2 .

The surface area $S_D = S_{n-1}$ of an n dimensional sphere with radius r is denoted by

$$S_{n-1} = \frac{dV_n(R)}{dR} = nC_n R^{n-1}$$

where $V_n(R)$ denotes the volume of a sphere of radius R.

This can be seen here:

http://scipp.ucsc.edu/~haber/ph116A/volume_11.pdf

There we can also find the identity (Eq.7):

$$\begin{split} nC_n &= \int \dots \int d\Omega_{n-1} \\ (*) &= \int \dots \int g(\mathbf{\theta}) d\theta_1 d\theta_2 \dots d\theta_{n-1} \end{split}$$

$$\begin{split} p(r,\theta_{1},\theta_{2},\dots,\theta_{n-1}) &= \int_{0}^{r} \int_{0}^{2\pi} \dots \int_{0}^{\pi} \underbrace{(2\pi)^{-\frac{D}{2}} e^{-\frac{1}{2}r^{2}}}_{pdf\ normal\ dist} \underbrace{r^{d-1} \sin^{d-2}\theta_{1} \sin^{d-3}\theta_{2} \cdot \dots \cdot \sin\theta_{d-2}\, d\theta_{1} \dots d\theta_{d-1} dr}_{infinitismal\ area} \\ &= \int_{0}^{r} (2\pi)^{-\frac{D}{2}} e^{-\frac{1}{2}r^{2}} r^{d-1} \int_{0}^{2\pi} \dots \int_{0}^{\pi} \underbrace{g(\mathbf{\theta}) d\theta_{1} \dots d\theta_{n-1}}_{=nc_{n} \cdot 1, acc. to.(*)} dr \\ &= \int_{0}^{r} (2\pi)^{-\frac{D}{2}} e^{-\frac{1}{2}r^{2}} r^{d-1} S_{D} \, dr \end{split}$$

Ergo
$$p(r)dr = S_D r^{d-1} (2\pi)^{-\frac{D}{2}} e^{-\frac{1}{2}r^2}$$
.

Now we are looking for the maximum density $\max_{r} p(r)$:

$$\frac{d}{dr} \left[\log r^{d-1} + \log e^{-\frac{1}{2}r^2} \right] = \frac{(d-1)r^{d-2}}{r^{d-1}} - r = 0.$$

$$\Leftrightarrow$$
 $(d-1)=r^2$.

Because radii are non-negative, we have a maximum at $\sqrt{d-1}$.

Now if we set $||\mathbf{x}|| = \sqrt{d-1}$, we get

$$\frac{p(\mathbf{x})}{p(\mathbf{0})} = \frac{(2\pi)^{-\frac{D}{2}}e^{-\frac{1}{2}\|\mathbf{x}\|^2}}{(2\pi)^{-\frac{D}{2}}e^{-\frac{1}{2}\|\mathbf{0}\|^2}} = e^{-\frac{1}{2}(d-1)}.$$

4A-3.

Let $L(\mathbf{w})$:

$$L(\mathbf{w}) = \frac{1}{2} \sum_{n=1}^{N} (t_n - \mathbf{w}^T \boldsymbol{\phi}(x_n))^2 + \frac{\lambda}{2} \mathbf{w}^T \mathbf{w}$$

Derivative of $\ln L(\mathbf{w})$ with respect to w_i :

$$\frac{\partial}{\partial w_i} \ln L(\mathbf{w}) = \sum_{n=1}^N \phi_i(x_n) \left(t_n - \mathbf{w}^T \boldsymbol{\phi}(x_n) \right) + \lambda w_i$$

Conversion to matrix/vector operations:

$$\sum_{n=1}^{N} \phi_i(x_n) \left(t_n - \mathbf{w}^T \boldsymbol{\phi}(x_n) \right) + \lambda w_i = \operatorname{col}_i(\boldsymbol{\Phi})^T (\mathbf{t} - \boldsymbol{\Phi} \mathbf{w}) + \lambda w_i$$

Generalized for all w:

$$\frac{\partial}{\partial \mathbf{w}} \ln L(\mathbf{w}) = \mathbf{\Phi}^T (\mathbf{t} - \mathbf{\Phi} \mathbf{w}) + \lambda \mathbf{w}$$

Setting zero and solving for w:

$$\Phi^{T}(\mathbf{t} - \Phi \mathbf{w}) + \lambda \mathbf{w} = 0$$

$$\Leftrightarrow \Phi^{T}\mathbf{t} + (-\Phi^{T}\Phi + \lambda \mathbf{I})\mathbf{w} = 0$$

$$\Leftrightarrow \mathbf{w} = (-\lambda \mathbf{I} + \Phi^{T}\Phi)^{+}\Phi^{T}\mathbf{t}$$

As usual, ${\bf A}^+$ denotes the Penrose pseudo inverse. Because λ can be an arbitrary normative factor, it is also possible to write:

$$\mathbf{w} = (\lambda \mathbf{I} + \mathbf{\Phi}^T \mathbf{\Phi})^+ \mathbf{\Phi}^T \mathbf{t}$$

