



# Metropolis

---

Davide Saccardo

17 Luglio 2018

Università di Trento

# Table of contents

1. Introduzione
2. Metodi
3. Analisi dati
4. Conclusione

# Introduzione

---

- E' il Machine Learning applicabile per calcolare l'energia di ground state di un Condensato di Bose-Einstein (BEC)?
- Possiamo imparare qualcosa in più sulla Restricted Boltzmann Machine?

1. Metodi Variational Monte Carlo per calcolare l'energia
2. Funzione d'onda di prova rappresentata con Restricted Boltzmann Machine
3. Parametri variazionali modificati con Stochastic Gradient Descent

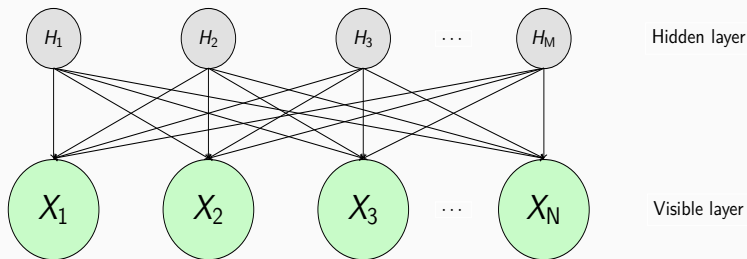
- Set of algorithms to give computers the ability to learn to recognize images (e.g. plants, handwritten numbers) or data distribution
- Huge success in the last years: Google Translate, ...
- Recent applications in physics with many-body problems such as Ising Model and Heisenberg model<sup>1</sup>, with encouraging results

---

<sup>1</sup>Giuseppe Carleo and Matthias Troyer. “Solving the quantum many-body problem with artificial neural networks”. In: *Science* 355.6325 (2017), pp. 602–606. ISSN: 0036-8075. DOI: 10.1126/science.aag2302. eprint: <http://science.sciencemag.org/content/355/6325/602.full.pdf>. URL: <http://science.sciencemag.org/content/355/6325/602>.

# Machine Learning

We represent our trial wave-function with a set of artificial Neural Networks. We call this way of representing the trial wave-function, Neural-Network Quantum State (NQS). The architecture selected is the Restricted Boltzmann Machine. The goal of the RBM is to learn the probability distribution of the simulated system. In quantum systems the probability distribution is the wave-function  $\Psi$ .



The joint probability distribution of the RBM is defined as

$$F_{RBM}(\mathbf{X}, \mathbf{H}) = \frac{1}{Z} \exp(-E(\mathbf{X}, \mathbf{H})),$$

where  $\mathbf{X}$  is the visible nodes and  $\mathbf{H}$  is the hidden nodes. The quantity  $Z$  represents the partition function or normalization constant of the system.

The quantity  $E(\mathbf{X}, \mathbf{H})$  is the function that specifies the relation between the visible and hidden nodes. It is called the energy of the node configuration. The choice of  $E(\mathbf{X}, \mathbf{H})$  is the heart of what sort of RBM we have.



There are several types of RBM. Since, in our case, the visible nodes need to take continuous values, we choose the Gaussian-Binary RBM:

$$E(\mathbf{X}, \mathbf{H}) = \sum_i^N \frac{(X_i - a_i)^2}{2\sigma^2} - \sum_j^M b_j H_j + \sum_{ij}^{N,M} \frac{X_i w_{ij} H_j}{\sigma^2},$$

To represent the wave-function, we use the so-called "marginal PDF" found by summing over all the hidden nodes:

$$\Psi(\mathbf{X}) = \sum_H F_{RBM}(\mathbf{X}, \mathbf{H}) = \frac{1}{Z} \sum_H \exp(-E(\mathbf{X}, \mathbf{H})).$$

Setting in the Gaussian-Binary RBM gives the final result

$$\Psi(\mathbf{X}) = \frac{1}{Z} \exp \left[ - \sum_i^N \frac{(X_i - a_i)^2}{2\sigma^2} \right] \prod_j^M \left( 1 + \exp \left[ b_j + \sum_i^N \frac{X_i w_{ij}}{\sigma^2} \right] \right).$$

# Condensato di Bose-Einstein

- Stato della materia in cui gas diluiti di bosoni vengono sottoposti a una transizione di fase quando vengono raffreddati fino a basse temperature ( $T \rightarrow 0$  K). La maggior parte dei bosoni condensa nel ground state. Questo causa nuovi fenomeni quantistici come superfluidità e coerenza di fase.
- Predetto da A. Einstein nel 1925 seguendo il lavoro del fisico indiano S. N. Bose (1924) sulla statistica dei bosoni.
- Verificati sperimentalmente nel 1995 dal team di Cornell e Wieman a Boulder con  $^{87}\text{Rb}$  e da Ketterle al MIT con  $^{23}\text{Na}$  raffreddati a temperature di 100 nK attraverso laser cooling e evaporative cooling in trappole magneto-ottiche. Premio Nobel 2001.

# Gross-Pitaevskii equation

Negli esperimenti sono stati utilizzati gas diluiti (e.g. atomi alcalini) e non uniformi perchè confinati in trappole magneto-ottiche. L'equazione di riferimento è l'equazione di Gross-Pitaevskii (GP):

$$i\hbar \frac{\partial \Psi(\mathbf{r}, t)}{\partial t} = \left( -\frac{\hbar^2}{2m} \nabla^2 + V(\mathbf{r}) + g|\Psi(\mathbf{r}, t)|^2 \right) \Psi(\mathbf{r}, t)$$

dove  $i = \sqrt{-1}$ ,  $\hbar$  è la costante di Planck ridotta,  $t$  è il tempi,  $\Psi$  è la funzione d'onda del condensato,  $\mathbf{r}$  rappresenta la posizione dei bosoni e  $V$  è il potenziale esterno che intrappola i bosoni. La quantità  $g$  è la costante di accoppiamento

$$g = \frac{4\pi\hbar^2 a}{m}$$

dove  $a$  is la lunghezza di scattering dell'onda s. Confronteremo le sue soluzioni in [2] con i nostri risultati.

La condizione di diluizione è data dal parametro  $na^3$ , dove  $n$  è la densità del sistema. Quando il parametro è piccolo, GP funziona molto bene. Tuttavia ci sono esperimenti in cui il parametro supera tale valore. A qual punto è importante studiare il sistema con un approccio many-body. In letterature, troviamo molti papers in cui vengono utilizzati metodi Monte Carlo nel range diluito fino alla densità dell' $^4\text{He}$ . Per esempio:

1. in [3], Dubois e Glyde usano Variational Monte Carlo (VMC);
2. in [4], Giorgini et al. usano Diffusion Monte Carlo (DMC);
3. in [5], Grüter et al. usano the path-integral Monte Carlo.

In this thesis, we use ML and we compare the results with GP equation since we consider a dilute system. Nevertheless, out of the dilute range, we note that ML results should be compared to the Monte Carlo ones which we have briefly mentioned above.

# Metodi

---

## II sistema

We consider a gas of  $N_p$  bosons in spherical and elliptical harmonic oscillator potentials. The interaction between bosons is modeled by the hard-core model as in [3]. The Hamiltonian of the system is

$$H = \sum_{i=1}^{N_p} \left[ -\frac{1}{2} \frac{\hbar^2}{m} \nabla_i^2 + V_{\text{ext}}(\mathbf{r}_i) \right] + \sum_{i \neq k} V_{\text{int}}(r_{ik}), \quad (1)$$

where  $V_{\text{ext}}(\mathbf{r})$  is the harmonic oscillator potential given by

$$V_{\text{ext}}(\mathbf{r}) = \begin{cases} \frac{1}{2} m \omega_{ho}^2 r^2 & \text{Spherical} \\ \frac{1}{2} m [\omega_{ho}^2 (x^2 + y^2) + \omega_z^2 z^2] & \text{Elliptical} \end{cases}$$

and  $V_{\text{int}}(r_{ik})$  is the hard-shell interaction potential

$$V_{\text{int}}(r_{ik}) = \begin{cases} 0, & r_{ik} > a \\ \infty, & r_{ik} < a. \end{cases}$$

The quantity  $r_{ik} = |\mathbf{r}_i - \mathbf{r}_k|$  represents the distance between particle  $i$  and  $k$ , while  $a$  is the size of the interaction between particles.

## No interaction - spherical trap

When we set the interaction to be zero ( $V_{int} = 0$ ), we are left with a harmonic oscillator potential, where we consider a spherical shape for simplicity. In this case the solutions are known analytically. In general the energy is given by  $E_n = \hbar\omega_{ho}(n + \frac{1}{2})$ . The ground state is

$$E(N_p, D) = \frac{1}{2}DN_p \hbar\omega_{ho}.$$

where  $D$  is the dimension of the system and  $N_p$  is the number of particles. This case is useful to benchmark our code at the beginning.

## No interaction - elliptic trap

The interaction is still null ( $V_{int} = 0$ ), the trap now is considered to be elliptic.

Let us introduce lengths in unit of  $a_{ho} = \sqrt{\hbar/(m\omega_{ho})}$ ,  $r \rightarrow r/a_{ho}$  and energy in units of  $\hbar\omega_{ho}$ . The Hamiltonian can be rearranged as

$$H = \sum_{k=1}^{N_p} \frac{\hbar\omega_{ho}}{2} \left( -a_{ho}^2 \nabla_k^2 + a_{ho}^2 \hbar \left[ x_k^2 + y_k^2 + \frac{\omega_z^2}{\omega_{ho}^2} z_k^2 \right] \right).$$

We set  $\lambda = \omega_z/\omega_{ho}$ , we get

$$H = \sum_{k=1}^{N_p} \frac{1}{2} \left( -\nabla_k^2 + V_{ext}(\mathbf{r}_k) \right) \quad (2)$$

where  $V_{ext} = x_k^2 + y_k^2 + \lambda^2 z_k^2$ . As in [2], we set  $\lambda = \sqrt{8}$ . In [3], the energy of non-interacting bosons in this trap is shown to be

$$\frac{E}{N} \rightarrow E_{ho} = \hbar\omega_{ho} \left( 1 + \frac{\lambda}{2} \right) = 2.414 \hbar\omega_{ho}.$$



## Interaction - elliptic trap

At this point, we turn on the interaction  $V_{int} \neq 0$  and we consider an elliptic trap. The Hamiltonian is

$$H = \sum_{k=1}^{N_p} \frac{1}{2} \left( -\nabla_k^2 + V_{ext}(\mathbf{r}_k) \right) + \sum_{k < i}^{N_p} V_{int}(\mathbf{r}_k, \mathbf{r}_i). \quad (3)$$

where  $V_{ext} = x_k^2 + y_k^2 + \lambda^2 z_k^2$  and  $\lambda = \sqrt{8}$  as above.

This frame uses the `allsmallcaps` titleformat.

### Potential problems

As this titleformat also uses smallcaps you face the same problems as with the `smallcaps` titleformat. Additionally this format can cause some other problems. Please refer to the documentation if you consider using it.

As a rule of thumb: Just use it for plaintext-only titles.

This frame uses the `allcaps` titleformat.

## Potential Problems

This titleformat is not as problematic as the `allsmallcaps` format, but basically suffers from the same deficiencies. So please have a look at the documentation if you want to use it.

## **Analisi dati**

---

The theme provides sensible defaults to  
`\emph{emphasize}` text, `\alert{accent}` parts  
or show `\textbf{bold}` results.

becomes

The theme provides sensible defaults to *emphasize* text, **accent** parts or  
show **bold** results.

# FONT FEATURE TEST

- Regular
- *Italic*
- SMALLCAPS
- **Bold**
- ***Bold Italic***
- **SmallCaps**
- Monospace
- *Monospace Italic*
- Monospace Bold
- *Monospace Bold Italic*

## Items

- Milk
- Eggs
- Potatos

## Enumerations

1. First,
2. Second and
3. Last.

## Descriptions

**PowerPoint** Meeh.  
**Beamer** Yeeeha.

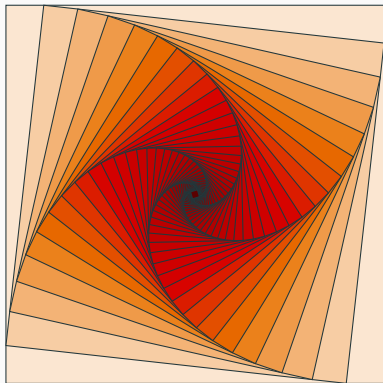
- This is important



- This is important
- Now this

- This is important
- Now this
- And now this

- This is really important
- Now this
- And now this



**Figure 1:** Rotated square from [texample.net](http://texample.net).

**Table 1:** Largest cities in the world (source: Wikipedia)

City	Population
Mexico City	20,116,842
Shanghai	19,210,000
Peking	15,796,450
Istanbul	14,160,467

Three different block environments are pre-defined and may be styled with an optional background color.

## **Default**

Block content.

## **Alert**

Block content.

## **Example**

Block content.

## **Default**

Block content.

## **Alert**

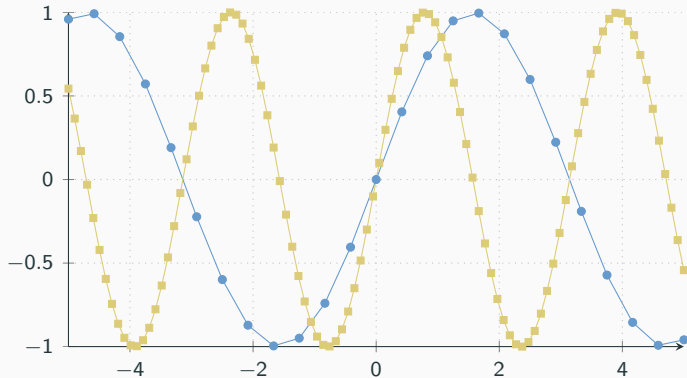
Block content.

## **Example**

Block content.

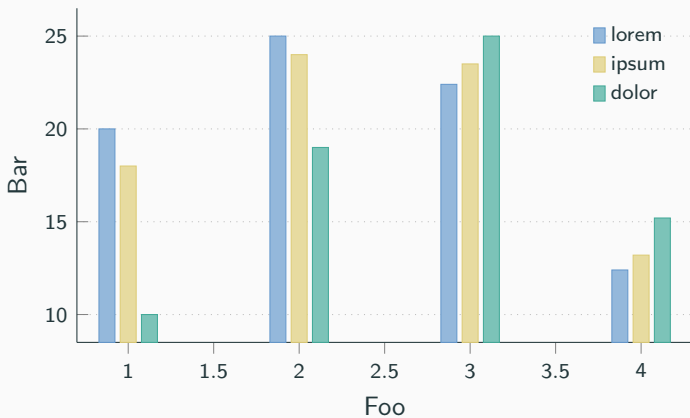
$$e = \lim_{n \rightarrow \infty} \left(1 + \frac{1}{n}\right)^n$$

# LINE PLOTS





# BAR CHARTS



*Veni, Vidi, Vici*

**metropolis** defines a custom beamer template to add a text to the footer. It can be set via

```
\setbeamertemplate{frame footer}{My custom footer}
```

# REFERENCES

Some references to showcase [allowframebreaks] [**knuth92**,  
**ConcreteMath**, **Simpson**, **Er01**, **greenwade93**]

# Conclusione

---

# SUMMARY

Get the source of this theme and the demo presentation from

`github.com/matze/mtheme`

The theme *itself* is licensed under a Creative Commons Attribution-ShareAlike 4.0 International License.



**Questions?**

# BACKUP SLIDES

Sometimes, it is useful to add slides at the end of your presentation to refer to during audience questions.

The best way to do this is to include the `appendixnumberbeamer` package in your preamble and call `\appendix` before your backup slides.

**metropolis** will automatically turn off slide numbering and progress bars for slides in the appendix.



## References

---



Giuseppe Carleo and Matthias Troyer. “Solving the quantum many-body problem with artificial neural networks”. In: *Science* 355.6325 (2017), pp. 602–606. ISSN: 0036-8075. DOI: 10.1126/science.aag2302. eprint: <http://science.sciencemag.org/content/355/6325/602.full.pdf>. URL: <http://science.sciencemag.org/content/355/6325/602>.



F. Dalfovo and S. Stringari. “Bosons in anisotropic traps: Ground state and vortices”. In: *Phys. Rev. A* 53 (4 Apr. 1996), pp. 2477–2485. DOI: 10.1103/PhysRevA.53.2477. URL: <https://link.aps.org/doi/10.1103/PhysRevA.53.2477>.

## REFERENCES II



J. L. DuBois and H. R. Glyde. “Bose-Einstein condensation in trapped bosons: A variational Monte Carlo analysis”. In: *Phys. Rev. A* 63 (2 Jan. 2001), p. 023602. DOI: 10.1103/PhysRevA.63.023602. URL: <https://link.aps.org/doi/10.1103/PhysRevA.63.023602>.



S. Giorgini, J. Boronat, and J. Casulleras. “Ground state of a homogeneous Bose gas: A diffusion Monte Carlo calculation”. In: *Phys. Rev. A* 60 (6 Dec. 1999), pp. 5129–5132. DOI: 10.1103/PhysRevA.60.5129. URL: <https://link.aps.org/doi/10.1103/PhysRevA.60.5129>.



Peter Grüter, David Ceperley, and Frank Laloë. “Critical Temperature of Bose-Einstein Condensation of Hard-Sphere Gases”. In: *Phys. Rev. Lett.* 79 (19 Nov. 1997), pp. 3549–3552. DOI: 10.1103/PhysRevLett.79.3549. URL: <https://link.aps.org/doi/10.1103/PhysRevLett.79.3549>.