

# analisi: your Swiss Army Knife of molecular dynamics analysis

Available on linux's conda-forge:

```
conda config --add channels conda-forge    #if not already done
conda config --set channel_priority strict #if not already done
conda install analisi
```

Azure

Variant

Status

linux\_64\_mpimpichpython3.6.\_\_\_\_cpython  
linux\_64\_mpimpichpython3.7.\_\_\_\_cpython  
linux\_64\_mpimpichpython3.8.\_\_\_\_cpython  
linux\_64\_mpimpichpython3.9.\_\_\_\_cpython  
linux\_64\_mpiopenmpipython3.6.\_\_\_\_cpython  
linux\_64\_mpiopenmpipython3.7.\_\_\_\_cpython  
linux\_64\_mpiopenmpipython3.8.\_\_\_\_cpython  
linux\_64\_mpiopenmpipython3.9.\_\_\_\_cpython

Name	Downloads	Version	Platforms	Repository
recipe <b>analisi</b>	downloads <b>756</b>	conda-forge <b>v0.1.10</b>	platform <b>linux-64</b>	analisi

or by compiling from source following this section

Fastest possible example usage:

## command line example

- mean square displacement

```
./analisi -i tests/lammps.bin -Q > output_file
```

- spherical harmonics correlation functions

```
./analisi -i tests/lammps.bin -Y 4 -F 0.7 3.5 > output_file
```

- $g(r)$ , with time lags too

```
./analisi -i tests/lammps.bin -g 200 -F 0.7 3.5 > output_file
```

- green kubo autocorrelation integrals

```
./analisi -l tests/gk_integral.txt -H -a 'c_flux[1]' 'c_vcm[1][1]' > output_file
```

and many others...

## python example

```
#read trajectory
import numpy as np
pos = np.load( 'tests/data/positions.npy')
vel = np.load( 'tests/data/velocities.npy')
box = np.load( 'tests/data/cells.npy')
types = np.zeros(pos.shape[1], dtype = np.int32)
types[-16:-8]=1
types[-8:]=2
print('position array shape is {}'.format(pos.shape))
print('first cell matrix is {}'.format(box[0]))

#create trajectory object
import pyanalisi
analisi_traj = pyanalisi.Trajectory(pos, vel, types, box, True, False)

#do the calculation that you want
msd=pyanalisi.MeanSquareDisplacement(analisi_traj,10,4000,4,True,False,False)
msd.reset(3000)
msd.calculate(0)

result=np.array(msd,copy=False)

#other calculations
#...
#...

Heat transport coefficient calculation: correlation functions and gk integral for a multicomponent
fluid example

import numpy as np
with open(filepath_tests + '/data/gk_integral.dat', 'r') as flog:
    headers = flog.readline().split()
log = np.loadtxt(filepath_tests + '/data/gk_integral.dat', skiprows=1)

import pyanalisi
traj = pyanalisi.ReadLog(log, headers)
gk = pyanalisi.GreenKubo(analisi_log, '',
    1, ['c_flux[1]', 'c_vcm[1][1]'],
    False, 2000, 2, False, 0, 4, False, 1, 100)
gk.reset(analisi_log.getNtimesteps()-2000)
gk.calculate(0)
result = np.array(gk,copy=False)
```

## note

If you use this program and you like it, spread the word around and give it credits! Implementing stuff that is already implemented can be a waste of time. And why don't you try to implement something that you like inside it? You will get for free MPI parallelization and variance calculation,

that are already implemented in a very generic and abstracted way.

## Description

This is a framework for computing averages on molecular dynamics trajectories and block averages. An mpi parallelization over the blocks is implemented in a very abstract and general way, so it is easy to implement a new calculation that can be parallelized with no effort using MPI. The code has two interfaces: a command line interface that is parallelized with MPI and threads, and a python interface that is parallelized on threads only. With such many parallelization levels more accurate averages can be performed, as described in details in the next sections.

Features:

- python interface (reads numpy array)
- command line interface (reads binary lammps-like files or time series in column format)
- multithreaded
- command line interface has MPI too (for super-heavy calculations)
- command line calculates variance of every quantity and every function (in python you can do it by yourselves with `numpy.var` )

Calculations:

- $g$  of  $r$  ( and time too!)
- vibrational spectrum (this is nothing special)
- histogram of number of neighbours
- mean square displacement
- green kubo integral of currents
- multicomponent green kubo time domain formula (MPI here can be useful...)
- spherical harmonics number density time correlation analysis (MPI here can be useful...)
- atomic position histogram
- and more ...
- ...

## Building from source

Dependencies:

- C++17 compatible compiler (GCC 7+, clang 6+, intel 19.0.1+, source )
- cmake
- linux (mmap - maybe this dependency can be removed with some additional work)
- FFTW3 (included in the package)
- Eigen3 (included in the package)
- Boost (included in the package)
- Mpi (optional)
- libxdrfile (for gromacs file conversion – included in the package)
- python (optional)

### MPI build (why not?)

```
mkdir build
cd build
```

```
cmake ../ -DCMAKE_CXX_COMPILER=mpicxx -DCMAKE_C_COMPILER=mpicc -DUSE_MPI=ON
make
```

## non-MPI build (shame on you!)

```
mkdir build
cd build
cmake ../
make
```

After compiling the code you will find the executable `analisi` and the shared library `pyanalisi.*.so` (with a part of the name that depends on your particular python installation) in the building folder. To install the python library, simply copy it in your site-library folder, for example by executing the command

```
cp pyanalisi.*.so "`python3 -m site --user-site`/pyanalisi.so"
```

Be careful to choose the correct python executable, the same that you used to compile the library. This task can be performed as well by the command `make install/local` inside a python virtual environment, for example a miniconda installation.

To test that the library was compiled correctly and that there are no regressions, you can run the (small) test suite with the command

```
make test
```

## additional cmake options

- `-DBUILD_TESTS=OFF` skips the building of the C++ tests
- `-DSYSTEM_FFTW3=ON` use system's FFTW3 library
- `-DSYSTEM_EIGEN3=ON` use installed system's eigen3 library
- `-DSYSTEM_BOOST=ON` use detected system's boost library
- `-DSYSTEM_XDRFILE=ON` use detected system's xdrfile library
- `-DPYTHON_EXECUTABLE=/path/to/python` use this python installation to build the python interface
- `-DPYTHON_INTERFACE=OFF` don't build any python interface (the building process will not depend on python)

## Documentation

This document is better rendered in the pdf version. [Link to the pdf version.](#)

### Command line interface

The command line utility is able to read only binary trajectory files in the LAMMPS format, specified with the command line option `-i input_file`, or a time series in a column formatted text file with a header, specified with the command line option `-l input_file`. The trajectory file can be generated in many ways: - by LAMMPS :-) - by using the command line utility with the command line options `-i input_file -binary-convert output_file`, where `input_file` is the name of a plain text trajectory in the format:

```

[natoms]
[xlo] [xhi]
[ylo] [yhi]
[zlo] [zhi]
[id_1] [type_1] [x_1] [y_1] [z_1] [vx_1] [vy_1] [vz_1]
.      .      .      .      .      .      .      .
.      .      .      .      .      .      .      .
.      .      .      .      .      .      .      .
[id_natoms] [type_natoms] [x_natoms] [y_natoms] [z_natoms] [vx_natoms] [vy_natoms] [vz_natoms]
.
.
.

```

That is: for every step you have to provide the number of atoms, low and high coordinates of the orthorombic cell, and then for every atom its id, type id, positions and velocities. - by using the command line utility with the command line options `-i input_file -binary-convert-gromacs output_file typefile` and a gromacs trajectory (you have to provide the xdr library in the building process) - by using the python interface: see section Buffer protocol interface

### Command line interface common arguments

Where meaningful, the following arguments are shared by all calculation performed by the command line utility `-i input_file` specify input trajectory binary file. See lammps documentation for the documentation of how to produce a binary dump from lammps. The order of the dumped atomic quantities **must be id type xu yu zu vx vy vz** `-l input_file` specify input time series in text column format with headers `-N thread_number` specify the number of threads to use where parallelization is implemented with threads `-B number_of_block` when the code calculates variances of the quantity, it splits the trajectory in many blocks. With this option you can specify the number of blocks. If MPI parallelization is used, since different blocks are calculated in parallel, you may want to specify a number of blocks that is a multiple of the number of MPI processes. In this way, at the final iteration over the blocks, all processes will be busy. `-s timestep_number_skip` usually neighbour configurations in the trajectory are strongly correlated. With this option you specify how far must be two consecutive configuration used to calculate averages. `-S timestep_number_length` specify how long must be the function that the code calculates, measured in number of timesteps. (for example the length of the MSD plot)

## Python interface

### Creating a trajectory object

You can create a trajectory python object to be used in this library in two ways: - start from python arrays generated by other code - use a LAMMPS binary file that you have on the filesystem. This is the same file that is used by the command line interface.

**using Python arrays: the buffer protocol interface** You must have 4 arrays. In general the interface supports any object that supports python's buffer protocol. This include, among others, numpy arrays. Suppose you have a trajectory of `t` timesteps and `n` atoms. You need the following arrays: - position array, of shape `(t,n,3)` - velocity array, of shape `(t,n,3)` - cell array, of shape `(t,3,3)` only diagonal matrices (orthorombic cells) are supported at the moment, or if a lammps formatted cell is provided `(t,6)` The lammps format simply list the low and high

coordinates of the orthorombic cell, as in the header of each timestep that you find in the lammps trajectory format (shown in command line interface section ) - types array, of shape (n) (integer array)

In general no particular units of measure are required, and the output will reflect the units that you used in the input. The calculations that the program does are reported exactly in the following sections. From those formulas you can deduce the units of the output given the units that you used in the input.

Then you must decide if you want that the coordinates are rewrapped inside the cell or not. At the moment only orthorombic cells are supported in all calculations but those that need only unwrapped coordinates, like MSD.

The lammps format for the cell is [x\_lo, x\_hi, y\_lo, y\_hi, z\_lo, z\_hi]: you have to provide only the coordinates boundaries.

The syntax for creating the object is

```
import pyanalisi
analisi_traj = pyanalisi.Trajectory(positions_array,
                                    velocity_array,
                                    types_array,
                                    box_array,
                                    use_matrix_or_lammps_cell_format,
                                    wrap_atomic_coordinates)
```

where `use_matrix_or_lammps_cell_format` is `True` if usual matrix format for the cell is given and `False` if a lammps formatted cell is provided and `wrap_atomic_coordinates` is `True` if you want to wrap all the atomic coordinates inside the cell.

You can write a LAMMPS binary trajectory (that can be used by the command line interface with `mpi`, for example) by calling

```
analisi_traj.write_lammps_binary('output_path', start_timestep, end_timestep)
```

where `start_timestep` is the first timestep index to dump (indexes start from 0) and `end_timestep` is the first timestep that will not be written. If `end_timestep == -1`, it will dump everything till the end of the trajectory. This is a very convenient way of moving heavy computations on a cluster where MPI can be used, or more in general to convert a generic trajectory format in the format used by the command line tool. For example

```
#read trajectory. It can come from everywhere
import numpy as np
pos = np.load( 'tests/data/positions.npy') #shape (N_timesteps, N_atoms, 3)
vel = np.load( 'tests/data/velocities.npy') #shape (N_timesteps, N_atoms, 3)
box = np.load( 'tests/data/cells.npy') #shape (N_timesteps, 3, 3)
types = np.load( 'tests/data/types.npy') #shape (N_atoms), dtype=np.int32

#create trajectory object and dump to file
import pyanalisi
analisi_traj = pyanalisi.Trajectory(pos, vel, types, box,
                                   True, # matrix format for the box array
                                   False # don't wrap the coordinates
                                   )
```

```

analisi_traj.write_lammps_binary("output_filename.bin"
                                , 0, # starting timestep
                                -1  # last timestep:
                                )    # -1 dumps full trajectory

```

**LAMMPS binary trajectory interface** This interface is a little more complicated, since it was designed for computing block averages of very big files. It can read the same files that the command line program reads. The object is created with

```
analisi_traj = pyanalisi.Traj('path_of_binary_file')
```

Then you have to call some more functions, BEFORE calling the compute routines :

```

analisi_traj.setWrapPbc(True) #optional: if you want to wrap positions inside the cell
analisi_traj.setAccessWindowSize(size_in_number_of_steps_of_the_reading_block)
analisi_traj.setAccessStart(first_timestep_to_read)

```

The first line is to set the pbc wrapping (don't use this if you have to compute the MSD!). Since the access to the trajectory is done in blocks, the second line sets the size of the reading block. The bigger the block the bigger the memory allocated to store the positions and the velocities. The last call sets the index of the first timestep that is going to be read, and then read it. In this moment the selected chunk of the trajectory is loaded into your memory. At this point you are able to call the needed compute routine, making sure that it is not going to read past the last timestep of the block. Later you can call again `setAccessStart` to load a different part of the trajectory and compute the quantity again. Only the data of the current block is stored in the memory, and if eventually there is some overlap with the previous block, the data is copied from memory and the overlapped part is not read again from the filesystem.

### Creating a time series object

Before analyzing a time series, for example to calculate the integral of the autocorrelation function, it is necessary to create a time series object. This object will hold the data that a different function can analyze. It is created with the following:

```
time_series = pyanalisi.ReadLog(data_array, headers_array)
```

where `data_array` and `header_array` are python objects that support the buffer protocol interface, for example numpy arrays, or a plain python list. The `data_array` is a two dimensional array where the first index is the timestep index and the second index is the column index. The `header_array` object must describe the columns that are stored in the bidimensional floating point array `data_array` using Python Strings.

### MSD

Given a trajectory  $\mathbf{x}_t$  where  $i \in \{1, \dots, N_{atoms}\}$  is the atomic index and  $t$  is the timestep index, defining the center of mass position of the atomic species  $j$  at the timestep  $t$  as

$${}^jcm_t = \frac{1}{N_j} \sum_{i|type(i)=j}^i \mathbf{x}_t$$

where  $N_j$  is the number of atoms of the specie  $j$ , the code computes the following

$$MSD_t^{typej} = \frac{1}{N_{typej}} \sum_{i|type(i)=typej} \frac{1}{N_{ave}} \sum_{l=1}^{N_{ave}} |{}^i\mathbf{x}_{t+l} - {}^i\mathbf{x}_l|^2$$

$$MSDcm_t^{typej} = \frac{1}{N_{ave}} \sum_{l=1}^{N_{ave}} |{}^{typej}cm_{t+l} - {}^{typej}cm_l|^2$$

If the option `--mean-square-displacement-self` is provided in the command line or in the python interface the documented argument is set to `True`, the atomic mean square displacement for each atomic species is calculated in the reference system of the center of mass of that particular atomic specie. That is, in this case the following is computed:

$$MSD_t^{typej} = \frac{1}{N_{typej}} \sum_{i|type(i)=typej} \frac{1}{N_{ave}} \sum_{l=1}^{N_{ave}} |({}^i\mathbf{x}_{t+l} - {}^{typej}cm_{t+l}) - ({}^i\mathbf{x}_l - {}^{typej}cm_l)|^2$$

$$MSDcm_t^{typej} = \frac{1}{N_{ave}} \sum_{l=1}^{N_{ave}} |{}^{typej}cm_{t+l} - {}^{typej}cm_l|^2$$

In the output you have many columns, one for each of the  $N_{types}$  atomic species, first the block of the atomic MSD and then eventually the block of the center of mass MSD if asked to compute. The center of mass MSD is computed only if the command line option `-Q` is provided or the documented argument is set to `True` in the python interface constructor. The output is the following:

$$\{MSD_t^1, \dots, MSD_t^{N_{types}}, MSDcm_t^1, \dots, MSDcm_t^{N_{types}}\}$$

In the command line output after each column printed as described in the line above you will find the variance calculated with a block average over the specified number of blocks.

The options that you can use for this calculation, in summary, are:

`./analisi -i lammp_binary [-N number of threads] [-S stop timestep] [-s skip every] [-B number of blocks] -q | -Q [-mean-square-displacement-self]`

where any option of `-q`, `-Q`, `--mean-square-displacement-self` activates the calculation of the MSD.

To use this calculation in the python interface, you have to first create a trajectory object (that can be both from a lammps binary or from python arrays objects - described in the section creating a trajectory object), then you have to create the following instance if the numpy array trajectory is used:

```
msd_calculation = pyanalisi.MeanSquareDisplacement(
    trajectory_instance, # Trajectory instance
    skip, # in calculating averages skip this amount of timestep,
        # as in -s option
    tmax, # size of the MSD(t) plot in number of timesteps
    nthreads, # number of parallel threads to use in the computation
    center_of_mass, # if True calculates center of mass displacement too
    use_cm_reference, # use the reference system of the center of mass
        # of all atoms of the same type
    debug_flag # if True produces some dump files...
)
```



if the lammps binary trajectory instance is used, you must use the `pyanalisi.MeanSquareDisplacement_lammps`, with exactly the same arguments. The calculation will be exactly the same. After initializing the object, the calculation can be initialized with

```
msd_calculation.reset(block_size)
```

that sets the number of steps that will be used to calculate the average over the trajectory. If `tmax` is greater than `block_size`, then `tmax` will be setted to `block_size`. The calculation is started with:

```
msd_calculation.calculate(first_timestep)
```

where `first_timestep` is the index of the first timestep that will be used to calculate the averages. After the calculation over this block is finished, the result can be collected in a numpy array with:

```
result = np.array(msd_calculation, copy=False)
```

In general the object `msd_calculation` supports the buffer protocol interface, so to get the result you can use anything that can interface with the buffer.

## GreenKubo

Given  $M$  vector time series of length  $N$   ${}^m\mathbf{J}_t$ ,  $m \in \{1 \dots M\}$ ,  $t \in \{1 \dots N\}$ , implements an expression equivalent to the following formula:

$$\begin{aligned} {}^{ij}C_t &= \frac{1}{N_{ave}} \sum_{m=1}^{N_{ave}} \frac{1}{3} \sum_{c=1}^3 {}^{ij}J_m^c \cdot {}^{ij}J_{m+l}^c \\ {}^{ij}L_t &= \sum_{l=0}^t {}^{ij}C_l \\ {}^{ij}\bar{L}_t &= \frac{1}{t} \sum_{l=0}^t {}^{ij}C_l \cdot l \\ GK_t &= \frac{1}{{}^{00}[(L_t)^{-1}]} \\ \bar{GK}_t &= \frac{1}{{}^{00}[(L_t - \bar{L}_t)^{-1}]} \end{aligned}$$

but with the trapezoidal rule in place of the sums marked with \*. Note that  $L_t$  is a matrix. To get the correct units of measure, you have still to multiply all the quantities but the  $C_t$ s by the integration timestep.  $N_{ave}$  is the number of timesteps on which the code runs the average. Every quantity is written in the output in the following order:

$$\{{}^{00}C_t, {}^{00}L_t, {}^{00}\bar{L}_t, {}^{01}C_t, {}^{01}L_t, {}^{01}\bar{L}_t, \dots, {}^{MM}C_t, {}^{MM}L_t, {}^{MM}\bar{L}_t, GK_t, \bar{GK}_t\}$$

If the command line tool is used, the variance of the block average is automatically computed, and after each column you find its variance. Moreover you find in the output a useful description of the columns with their indexes.

## $\mathbf{g}(\mathbf{r}, t)$

Given a central atom of type  $j$  at timestep  $l$ , calculates the histogram of the minimum image's radial distance of atoms of type  $i$  at timestep  $l + t$ . The histogram can be of the same atom

(you now, it spreads around while time is passing) or of all atoms different from the original one. Everything is averaged over  $l$  and atoms of the same type. In particular, defining the index of a pair of atoms  $i$  and  $j$  and a timelag  $t$  at a timestep  $l$  as

$${}^{ij}h_t = \lfloor (|{}^j\mathbf{x}_{l+t} - {}^i\mathbf{x}_l|_{\text{min.image}} - r_{\text{min}})/dr \rfloor,$$

the histogram  $g(r, t)$  between specie  $I$  and  $J$  is defined as

$${}^{IJ}g_t^r = \sum_{i|type(i)=I} \sum_{j|type(j)=J} \delta({}^{ij}h_t, r)$$

where  $\delta(\cdot, \cdot)$  is the Kronecker delta. In practice the program, for each atoms, it adds 1 to the corresponding bin of the histogram.

For each  $t, I, J$ , with  $J \geq I$ , the position of the histograms in the memory is  $N_{types}(N_{types} + 1)/2 - (J + 1)(J + 2)/2 + I$  (0 is the first). Given this order of the histograms, the layout in the memory is the following:

$$\{ {}^0g_t^0, \dots, {}^0g_t^{r_{max}}, \dots, {}^{N_{types}(N_{types}+1)/2}g_t^0, \dots, {}^{N_{types}(N_{types}+1)/2}g_t^{r_{max}}, \dots \}$$

The layout of the command line output is a gnuplot-friendly one, where the output is organized in blocks, one for each  $t$ , separated by two blank lines, and every line corresponds to a histogram bin for every combination of  $I, J$ :

$$\begin{aligned} & \{ \quad {}^0g_0^0, \dots, \quad {}^{N_{types}(N_{types}+1)/2}g_0^0 \quad \} \\ & \{ \quad \vdots, \quad \vdots \quad \} \\ & \{ \quad {}^0g_0^{r_{max}}, \dots, \quad {}^{N_{types}(N_{types}+1)/2}g_0^{r_{max}} \quad \} \\ & \vdots \\ & \{ \quad {}^0g_{t_{max}}^0, \dots, \quad {}^{N_{types}(N_{types}+1)/2}g_{t_{max}}^0 \quad \} \\ & \{ \quad \vdots, \quad \vdots \quad \} \\ & \{ \quad {}^0g_{t_{max}}^{r_{max}}, \dots, \quad {}^{N_{types}(N_{types}+1)/2}g_{t_{max}}^{r_{max}} \quad \} \end{aligned}$$

where we collapsed the indexes  $I, J$  into a single number, the position order of the histogram. As usual in the command line output every column is an average over all the blocks and it is followed by the variance of the mean.

## Spherical harmonics correlations

### Calculation procedure:

The implemented formula for the real spherical harmonics is the following:

$$Y_{\ell m} = \begin{cases} (-1)^m \sqrt{2} \sqrt{\frac{2\ell+1}{4\pi} \frac{(\ell-|m|)!}{(\ell+|m|)!}} P_{\ell}^{|m|}(\cos \theta) \sin(|m|\varphi) & \text{if } m < 0 \\ \sqrt{\frac{2\ell+1}{4\pi}} P_{\ell}^m(\cos \theta) & \text{if } m = 0 \\ (-1)^m \sqrt{2} \sqrt{\frac{2\ell+1}{4\pi} \frac{(\ell-m)!}{(\ell+m)!}} P_{\ell}^m(\cos \theta) \cos(m\varphi) & \text{if } m > 0. \end{cases}$$

Where  $\cos \theta$ ,  $\sin \varphi$ ,  $\cos \varphi$  are calculated using cartesian components:

$$\begin{aligned}\cos \theta &= \frac{z}{\sqrt{x^2 + y^2 + z^2}} \\ \cos \varphi &= \frac{x}{\sqrt{x^2 + y^2}} \\ \sin \varphi &= \frac{y}{\sqrt{x^2 + y^2}}\end{aligned}$$

and  $\sin |m|\varphi$ ,  $\cos |m|\varphi$  are evaluated using Chebyshev polynomials with a recursive definition:

$$\begin{cases} \cos m\varphi &= 2 \cos(m-1)\varphi \cos \varphi - \cos(m-2)\varphi \\ \sin m\varphi &= 2 \cos \varphi \sin(m-1)\varphi - \sin(m-2)\varphi \end{cases}$$

and  $P_\ell^m$  are the associated Legendre polynomials, calculated with the following set of recursive definition:

$$P_{\ell+1}^{\ell+1}(x) = -(2\ell+1)\sqrt{1-x^2}P_\ell^\ell(x)P_{\ell+1}^\ell(x) = x(2\ell+1)P_\ell^\ell(x)$$

that allows us to calculate every  $(\ell, \ell)$  and every  $(\ell, \ell+1)$  element of the  $(\ell, m)$  values. Then we have the recursion to go up in  $\ell$ , for any value of it:

$$P_\ell^m = \frac{x(2\ell-1)P_{\ell-1}^m - (\ell+m-1)P_{\ell-2}^m}{\ell-m}$$

The program, given a number  $\ell_{max}$  and a triplet  $(x, y, z)$ , is able to calculate every value of  $Y_{\ell m}(x, y, z) \forall \ell \in [0, \ell_{max}]$  and for all allowed values of  $m$  with a single recursion. Let

$$y_{\ell m}^{Ij}(t) = \int_{V_I} d\theta d\varphi dr \rho^j(r, \theta, \varphi, t) Y_{\ell m}(\theta, \varphi)$$

for some timestep  $t$  in some volume  $V_I$  taken as a the difference of two concentric spheres centered on the atom  $I$  of radius  $r_{inner}$  and  $r_{outer}$ , and where  $\rho_j$  is the atomic density of the species  $j$ . Since the densities are taken as sums of dirac delta functions, it is sufficient to evaluate the spherical harmonics functions at the position of the atoms. Then we calculate the following:

$$c_\ell^{Jj}(t) = \langle \frac{1}{N_J} \sum_{I \text{ is of type } J} \sum_{m=-\ell}^{\ell} y_{\ell m}^{Ij}(0) y_{\ell m}^{Ij}(t) \rangle$$

where  $\langle \cdot \rangle$  is an average operator, and we do an additional average over all the  $N_J$  central atoms of the type  $J$ . The  $\langle \cdot \rangle$  average is implemented as an average over the starting timestep.

## Credits

Written by Riccardo Bertossa during his lifetime at SISSA