analisi: your Swiss Army Knife of molecular dynamics analysis

NOTE: after the merge request to conda forge will be accepted conda-forge will be available.

Available on linux's conda-forge:

or by compiling from source following this section

Fastest possible example usage:

command line example

- ullet mean square displacement
- ./analisi -i tests/lammps.bin -Q > output_file
 - spherical harmonics correlation functions
- ./analisi -i tests/lammps.bin -Y 4 -F 0.7 3.5 > output_file
 - g(r), with time lags too
- ./analisi -i tests/lammps.bin -g 200 -F 0.7 3.5 > output_file
 - green kubo autocorrelation integrals
- ./analisi -l tests/gk_integral.txt -H -a 'c_flux[1]' 'c_vcm[1][1]' > output_file and many others...

python example

```
#read trajectory
import numpy as np
pos = np.load( 'tests/data/positions.npy')
vel = np.load( 'tests/data/velocities.npy')
box = np.load( 'tests/data/cells.npy')
types = np.zeros(pos.shape[1], dtype = np.int32)
types[-16:-8]=1
types[-8:]=2
print('position array shape is {}'.format(pos.shape))
print('first cell matrix is {}'.format(box[0]))
#create trajectory object
import pyanalisi
analisi_traj = pyanalisi.Trajectory(pos, vel, types, box,True, False)
#do the calculation that you want
msd=pyanalisi.MeanSquareDisplacement(analisi_traj,10,4000,4,True,False,False)
msd.reset(3000)
```

```
msd.calculate(0)
result=np.array(msd,copy=False)
#other calculations
#...
#...
Heat transport coefficient calculation: correlation functions and gk integral for a multicomponent
fluid example
import numpy as np
with open(filepath_tests + '/data/gk_integral.dat', 'r') as flog:
    headers = flog.readline().split()
log = np.loadtxt(filepath_tests + '/data/gk_integral.dat', skiprows=1)
import pyanalisi
traj = pyanalisi.ReadLog(log, headers)
gk = pyanalisi.GreenKubo(analisi_log,'',
       1,['c flux[1]','c vcm[1][1]'],
       False, 2000, 2, False, 0, 4, False, 1, 100)
gk.reset(analisi log.getNtimesteps()-2000)
gk.calculate(0)
result = np.array(gk,copy=False)
```

note

If you use this program and you like it, spread the word around and give it credits! Implementing stuff that is already implemented can be a waste of time. And why don't you try to implement something that you like inside it? You will get for free MPI parallelization and variance calculation, that are already implemented in a very generic and abstracted way.

Description

This is a framework for computing averages on molecular dynamics trajectories and block averages. An mpi parallelization over the blocks is implemented in a very abstract and general way, so it is easy to implement a new calculation that can be parallelized with no effort using MPI. The code has two interfaces: a command line interface that is parallelized with MPI and threads, and a python interface that is parallelized on threads only. With such many parallelization levels more accurate averages can be performed, as described in details in the next sections.

Features:

- python interface (reads numpy array)
- command line interface (reads binary lammps-like files or time series in column format)
- multithreaded
- command line interface has MPI too (for super-heavy calculations)
- command line calculates variance of every quantity and every function (in python you can do it by yourselves with numpy.var)

Calculations:

- g of r (and time too!)
- vibrational spectrum (this is nothing special)
- histogram of number of neighbours
- mean square displacement
- green kubo integral of currents
- multicomponent green kubo time domain formula (MPI here can be useful...)
- spherical harmonics number density time correlation analysis (MPI here can be useful...)
- atomic position histogram
- and more ...
- . . .

Building from source

Dependencies:

- C++17 compatible compiler (GCC 7+, clang 6+, intel 19.0.1+, source)
- cmake
- linux (mmap maybe this dependency can be removed with some additional work)
- FFTW3 (included in the package)
- Eigen3 (included in the package)
- Boost (included in the package)
- Mpi (optional)
- libxdrfile (for gromacs file conversion optional)
- python (optional)

MPI build (why not?)

```
mkdir build
cd build
cmake ../ -DCMAKE_CXX_COMPILER=mpicxx -DCMAKE_C_COMPILER=mpicc -DUSE_MPI=ON
make
```

non-MPI build (shame on you!)

```
mkdir build
cd build
cmake ../
make
```

After compiling the code you will find the executable analisi and the shared library pyanalisi.*.so (with a part of the name that depends on your particular python installation) in the building folder. To install the python library, simply copy it in your site-library folder, for example by executing the command

```
cp pyanalisi.*.so "`python3 -m site --user-site`/pyanalisi.so"
```

Be careful to choose the correct python executable, the same that you used to compile the library. This task can be performed as well by the command make install/local inside a python virtual environment, for example a miniconda installation.

To test that the library was compiled correctly and that there are no regressions, you can run the (small) test suite with the command

make test

additional cmake options

- -DBUILD_TESTS=OFF skips the building of the C++ tests
- -DSYSTEM_FFTW3=ON try to detect system's FFTW3 library
- -DPYTHON_EXECUTABLE=/path/to/python use this python installation to build the python interface
- -DPYTHON_INTERFACE=0FF don't build any python interface (the building process will not depend on python)

Documentation

This document is better rendered in the pdf version. Link to the pdf version.

Command line interface

The command line utility is able to read only binary trajectory files in the LAMMPS format, specified with the command line option -i input_file, or a time series in a column formatted text file with a header, specified with the command line option -l input_file. The trajectory file can be generated in many ways: - by LAMMPS:-) - by using the command line utility with the command line options -i input_file -binary-convert output_file, where input_file is the name of a plain text trajectory in the format:

That is: for every step you have to provide the number of atoms, low and high coordinates of the orthorombic cell, and then for every atom its id, type id, positions and velocities. - by using the command line utility with the command line options -i input_file -binary-convert-gromacs output_file typefile and a gromacs trajectory (you have to provide the xdr library in the building process) - by using the python interface: see section Buffer protocol interface

Command line interface common arguments

Where meaningful, the following arguments are shared by all calculation performed by the command line utility - -i input_file specify input trajectory binary file. See lammps documentation for the documentation of how to produce a binary dump from lammps. The order of the dumped

atomic quantities **must be id type xu yu zu vx vy vz** --l input_file specify input time series in text column format with headers - -N thread_number specify the number of threads to use where parallelization is implemented with threads - -B number_of_block when the code calculates variances of the quantity, it splits the trajectory in many blocks. With this option you can specify the number of blocks. If MPI parallelization is used, since different blocks are calculated in parallel, you may want to spacify a number of blocks that is a multiple of the number of MPI processes. In this way, at the final iteration over the blocks, all processes will be busy. - -s timestep_number_skip usually neighbour configurations in the trajectory are strongly correlated. With this option you specify how far must be two consecutive configuration used to calculate averages. - -S timestep_number_length specify how long must be the function that the code calculates, measured in number of timesteps. (for example the length of the MSD plot)

Python interface

Creating a trajectory object

You can create a trajectory python object to be used in this library in two ways: - start from python arrays generated by other code - use a LAMMPS binary file that you have on the filesystem. This is the same file that is used by the command line interface.

using Python arrays: the buffer protocol interface You must have 4 arrays. In general the interface supports any object that supports python's buffer protocol. This include, among others, numpy arrays. Suppose you have a trajectory of t timesteps and n atoms. You need the following arrays: - position array, of shape (t,n,3) - velocity array, of shape (t,n,3) - cell array, of shape (t,3,3) only diagonal matrices (orthorombic cells) are supported at the moment, or if a lammps formated cell is provided (t,6) The lammps format simply list the low and high coordinates of the orthorombic cell, as in the header of each timestep that you find in the lammps trajetory format (shown in command line interface section) - types array, of shape (n) (integer array)

In general no particular units of measure are required, and the output will reflect the units that you used in the input. The calculations that the program does are reported exactly in the following sections. From those formulas you can deduce the units of the output given the units that you used in the input.

Then you must decide if you want that the coordinates are rewrapped inside the cell or not. At the moment only orthorombic cells are supported in all calculations but those that need only unwrapped coordinates, like MSD.

The lammps format for the cell is [x_lo, x_hi, y_lo, y_hi, z_lo, z_hi]: you have to provide only the coordinates boundaries.

The syntax for creating the object is

where use_matrix_or_lammps_cell_format is True if usual matrix format for the cell is given and False if a lammps formatted cell is provided and wrap_atomic_coordinates is True if you want to wrap all the atomic coordinates inside the cell.

You can write a LAMMPS bynary trajectory (that can be used by the command line interface with mpi, for example) by calling

```
analisi_traj.write_lammps_binary('output_path', start_timestep, end_timestep)
```

where start_timestep is the first timestep index to dump (indexes start from 0) and end_timestep is the first timestep that will not be written. If end_timestep == -1, it will dump everything till the end of the trajectory. This is a very convenient way of moving heavy computations on a cluster where MPI can be used, or more in general to convert a generic trajectory format in the format used by the command line tool. For example

LAMMPS binary trajectory interface This interface is a little more complicated, since it was designed for computing block averages of very big files. It can read the same files that the command line program reads. The object is created with

```
analisi_traj = pyanalisi.Traj('path_of_binary_file')
```

Then you have to call some more functions, BEFORE calling the compute routines:

```
analisi_traj.setWrapPbc(True) #optional: if you want to wrap positions inside the cell
analisi_traj.setAccessWindowSize(size_in_number_of_steps_of_the_reading_block)
analisi_traj.setAccessStart(first_timestep_to_read)
```

The first line is to set the pbc wrapping (don't use this if you have to compute the MSD!). Since the access to the trajectory is done in blocks, the second line sets the size of the reading block. The bigger the block the bigger the memory allocated to store the positions and the velocities. The last call sets the index of the first timestep that is going to be read, and then read it. In this moment the selected chunk of the trajectory is loaded into your memory. At this point you are able to call the needed compute routine, making sure that it is not going to read past the last timestep of the block. Later you can call again setAccessStart to load a different part of the trajectory and compute the quantity again. Only the data of the current block is stored in the

memory, and if evenutally there is some overlap with the previous block, the data is copied from memory and the overlapped part is not read again from the filesystem.

Creating a time series object

Before analyzing a time series, for example to calculate the integral of the autocorrelation function, it is necessary to create a time series object. This object will hold the data that a different function can analyze. It is created with the following:

time_series = pyanalisi.ReadLog(data_array, headers_array)

where data_array and header_array are python objects that support the buffer protocol interface, for example numpy arrays, or a plain python list. The data_array is a two dimensional array where the first index is the timestep index and the second index is the column index. The header_array object must describe the columns that are stored in the bidimensional floating point array data_array using Python Strings.

MSD

Given a trajectory ${}^{\mathbf{i}}\mathbf{x_t}$ where $i \in \{1, \dots, N_{atoms}\}$ is the atomic index and t is the timestep index, defining the center of mass position of the atomic species j at the timestep t as

$${}^{j}cm_{t} = \frac{1}{N_{j}} \sum_{i|type(i)=j}^{i} \mathbf{x}_{t}$$

where N_j is the number of atoms of the specie j, the code computes the following

$$\begin{split} MSD_{t}^{typej} &= \frac{1}{N_{typej}} \sum_{i|type(i)=typej} \frac{1}{N_{ave}} \sum_{l=1}^{N_{ave}} \left| {}^{i}\mathbf{x}_{t+l} - {}^{i}\mathbf{x}_{l} \right|^{2} \\ MSDcm_{t}^{typej} &= \frac{1}{N_{ave}} \sum_{l=1}^{N_{ave}} \left| {}^{typej}cm_{t+l} - {}^{typej}cm_{l} \right|^{2} \end{split}$$

If the option --mean-square-displacement-self is provided in the command line or in the python interface the documented argument is set to True, the atomic mean square displacement for each atomic species is calculated in the reference system of the center of mass of that particular atomic specie. That is, in this case the following is computed:

$$\begin{split} MSD_{t}^{typej} &= \frac{1}{N_{typej}} \sum_{i|type(i)=typej} \frac{1}{N_{ave}} \sum_{l=1}^{N_{ave}} \left| (^{i}\mathbf{x}_{t+l} - ^{typej} cm_{t+l}) - (^{i}\mathbf{x}_{l} - ^{typej} cm_{l}) \right|^{2} \\ MSDcm_{t}^{typej} &= \frac{1}{N_{ave}} \sum_{l=1}^{N_{ave}} \left| ^{typej} cm_{t+l} - ^{typej} cm_{l} \right|^{2} \end{split}$$

In the output you have many columns, one for each of the N_{types} atomic species, first the block of the atomic MSD and then eventually the block of the center of mass MSD if asked to compute. The center of mass MSD is computed only if the command line option -Q is provided or the documented argument is set to True in the python interface constructor. The output is the following:

$$\{MSD_t^1, \dots, MSD_t^{N_{types}}, MSDcm_t^1, \dots, MSDcm_t^{N_{types}}\}$$

In the command line output after each column printed as described in the line above you will find the variance calculated with a block average over the specified number of blocks.

The options that you can use for this calculation, in summary, are:

./analisi -i lammp_binary [-N number of threads] [-S stop timestep] [-s skip every] [-B number of blocks] -q | -Q [-mean-square-displacement-self]

where any option of -q, -Q, --mean-square-displacement-self activates the calculation of the MSD.

To use this calculation in the python interface, you have to first create a trajectory object (that can be both from a lammps binary or from python arrays objects - described in the section creating a trajectory object), then you have to create the following instance if the numnpy array trajectory is used:

if the lammps binary trajectory instance is used, you must use the pyanalisi.MeanSquareDisplacement_lammps, with exactly the same arguments. The calculation will be exactly the same.

GreenKubo

Given M vector time series of length $N^{m}\mathbf{J}_{t}$, $m \in \{1...M\}$, $t \in \{1...N\}$, implements an expression equivalent to the following formula:

$$i^{j}C_{t} = \frac{1}{N_{ave}} \sum_{m=1}^{N_{ave}} \frac{1}{3} \sum_{c=1}^{3} i J_{m}^{c} \cdot {}^{j}J_{m+l}^{c}$$

$$i^{j}L_{t} = \sum_{l=0}^{t} {}^{ij}C_{l}$$

$$i^{j}\bar{L}_{t} = \frac{1}{t} \sum_{l=0}^{t} {}^{ij}C_{l} \cdot l$$

$$GK_{t} = \frac{1}{{}^{00}[(L_{t})^{-1}]}$$

$$G\bar{K}_{t} = \frac{1}{{}^{00}[(L_{t} - \bar{L}_{t})^{-1}]}$$

but with the trapezoidal rule in place of the sums marked with *. Note that L_t is a matrix. To get the correct units of measure, you have still to multiply all the quantities but the C_t s by the

integration timestep. N_{ave} is the number of timesteps on which the code runs the average. Every quantity is written in the output in the following order:

$$\{{}^{00}C_{t}, {}^{00}L_{t}, {}^{00}\bar{L}_{t}, {}^{01}C_{t}, {}^{01}L_{t}, {}^{01}\bar{L}_{t}, \dots, {}^{MM}C_{t}, {}^{MM}L_{t}, {}^{MM}\bar{L}_{t}, GK_{t}, G\bar{K}_{t}\}$$

If the command line tool is used, the variance of the block average is automatically computed, and after each column you find its variance. Moreover you find in the output a useful description of the columns with their indexes.

$\mathbf{g}(\mathbf{r},\mathbf{t})$

Given a central atom of type j at timestep l, calculates the histogram of the minimum image's radial distance of atoms of type i at timestep l + t. The histogram can be of the same atom (you now, it spreads arond while time is passing) or of all atoms different from the original one. Everything is averaged over l and atoms of the same type. In particular, defining the index of a pair of atoms i and j and a timelag t at a timestep l as

$$^{ij}h_t = |(|^j \mathbf{x}_{l+t} - ^i \mathbf{x}_l|_{\text{min.image}} - rmin)/dr|,$$

the histogram g(r,t) between specie I and J is defined as

$$^{I\,J}g_{t}^{r} = \sum_{i|type(i)=I} \sum_{j|type(j)=J} \delta(^{i\,j}h_{t},r)$$

where $\delta(\cdot, \cdot)$ is the Kronecker delta. In practice the program, for each atoms, it adds 1 to the corresponding bin of the histogram.

For each t, I, J, with $J \ge I$, the position of the histograms in the memory is $N_{types}(N_{types} + 1)/2 - (J+1)(J+2)/2 + I$ (0 is the first). Given this order of the histograms, the layout in the memory is the following:

$$\{{}^0g^0_t, \dots, {}^0g^{r_{max}}_t, \dots, {}^{N_{types}(N_{types}+1)/2}g^0_t, \dots, {}^{N_{types}(N_{types}+1)/2}g^{r_{max}}_t, \dots\}$$

The layout of the command line output is a gnuplot-friendly one, where the output is organized in blocks, one for each t, separated by two blank lines, and every line corresponds to a histogram bin for every combination of I, J:

$$\left\{ \begin{array}{cccc} {}^{0}g_{0}^{0}, \ldots, {}^{N_{types}(N_{types}+1)/2}g_{0}^{0} & \right\} \\ \left\{ & \vdots, & \vdots, & \\ & \vdots, & \vdots, & \\ & {}^{0}g_{0}^{r_{max}}, \ldots, {}^{N_{types}(N_{types}+1)/2}g_{0}^{r_{max}} & \right\} \\ \vdots & & & \\ \left\{ & {}^{0}g_{t_{max}}^{0}, \ldots, {}^{N_{types}(N_{types}+1)/2}g_{t_{max}}^{0} & \right\} \\ \left\{ & \vdots, & \vdots, & \\ & {}^{0}g_{t_{max}}^{r_{max}}, \ldots, {}^{N_{types}(N_{types}+1)/2}g_{t_{max}}^{r_{max}} & \right\} \\ \end{array}$$

where we collapsed the indexes I, J into a single number, the position order of the histogram. As usual in the command line output every column is an average over all the blocks and it is followed by the variance of the mean.

Spherical harmonics correlations

Calculation procedure:

The implemented formula for the real spherical harmonics is the following:

$$Y_{\ell m} = \begin{cases} (-1)^m \sqrt{2} \sqrt{\frac{2\ell + 1}{4\pi} \frac{(\ell - |m|)!}{(\ell + |m|)!}} \ P_{\ell}^{|m|}(\cos \theta) \sin(|m|\varphi) & \text{if } m < 0 \\ \sqrt{\frac{2\ell + 1}{4\pi}} \ P_{\ell}^{m}(\cos \theta) & \text{if } m = 0 \\ (-1)^m \sqrt{2} \sqrt{\frac{2\ell + 1}{4\pi} \frac{(\ell - m)!}{(\ell + m)!}} \ P_{\ell}^{m}(\cos \theta) \cos(m\varphi) & \text{if } m > 0 \,. \end{cases}$$

Where $\cos \theta$, $\sin \varphi$, $\cos \varphi$ are calculated using cartesian components:

$$\cos \theta = \frac{z}{\sqrt{x^2 + y^2 + z^2}}$$
$$\cos \varphi = \frac{x}{\sqrt{x^2 + y^2}}$$
$$\sin \varphi = \frac{y}{\sqrt{x^2 + y^2}}$$

and $\sin |m|\varphi$, $\cos |m|\varphi$ are evaluated using Chebyshev polynomials with a recursive definition:

$$\begin{cases} \cos m\varphi &= 2\cos(m-1)\varphi\cos\varphi - \cos(m-2)\varphi\\ \sin m\varphi &= 2\cos\varphi\sin(m-1)\varphi - \sin(m-2)\varphi \end{cases}$$

and P_{ℓ}^{m} are the associated Legendre polynomials, calculated with the following set of recursive definition:

$$P_{\ell+1}^{\ell+1}(x) = -(2\ell+1)\sqrt{1-x^2}P_{\ell}^{\ell}(x)P_{\ell+1}^{\ell}(x) = x(2\ell+1)P_{\ell}^{\ell}(x)$$

that allows us to calculate every (ℓ, ℓ) and every $(\ell, \ell + 1)$ element of the (ℓ, m) values. Then we have the recursion to go up in ℓ , for any value of it:

$$P_{\ell}^{m} = \frac{x(2\ell-1)P_{\ell-1}^{m} - (\ell+m-1)P_{\ell-2}^{m}}{\ell-m}$$

The program, given a number ℓ_{max} and a triplet (x, y, z), is able to calculate every value of $Y_{\ell m}(x, y, z) \forall \ell \in [0, \ell_{max}]$ and for all allowed values of m with a single recursion. Let

$$y_{\ell m}^{Ij}(t) = \int_{V_I} \mathrm{d}\theta \mathrm{d}\varphi \mathrm{d}r \rho^j(r,\theta,\varphi,t) Y_{\ell m}(\theta,\varphi)$$

for some timestep t in some volume V_I taken as a the difference of two concentric spheres centered on the atom I of radius r_{inner} and r_{outer} , and where ρ_j is the atomic density of the species j. Since the densities are taken as sums of dirac delta functions, it is sufficient to evaluate the spherical harmonics functions at the position of the atoms. Then we calculate the following:

$$c_\ell^{Jj}(t) = \langle \frac{1}{N_J} \sum_{I \text{ is of type } J} \sum_{m=-\ell}^\ell y_{\ell m}^{Ij}(0) y_{\ell m}^{Ij}(t) \rangle$$

where $\langle \cdot \rangle$ is an average operator, and we do an additional average over all the N_J central atoms of the type J. The $\langle \cdot \rangle$ average is implemented as an average over the starting timestep.

Credits

Written by Riccardo Bertossa during his lifetime at SISSA