## MSBM Project Documentation

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## 1 PINN Architecture

We are solving the 2D MSBM equations at steady state along the x-direction, so the solution varies only along the y-direction. Furthermore, the y-direction velocity is zero at steady state. Thus, our only input is y, and our solution is u and  $\phi$ , where u is the x-direction velocity and  $\phi$  is the particle volume fraction.

We use two separate PINNs for u and  $\phi$ , respectively, which are denoted in the style of [1] as

$$u(\theta_u, y) = \Gamma_u(\mathcal{M}_u(\theta_u, I_u(y))),$$

$$\phi(\theta_{\phi}, y) = \Gamma_{\phi}(\mathcal{M}_{\phi}(\theta_{\phi}, I_{\phi}(y))).$$

The input transformations are  $I_u(y)$  and  $I_{\phi}(y)$  for each PINN, where

$$I_u(y) = \frac{2y}{H} - 1$$

normalizes the input such that  $I_u(y) \in [-1, 1]$ ; this range is beneficial for the choice of activation function, which are discussed later.

As for  $I_{\phi}(y)$ , each script will specify using either a Fourier expansion or a Gaussian expansion:

$$I_{\phi}(y) = a \cdot \left[ \sin(2\pi b_i^{\top} | \frac{2y}{H} - 1 |), \cos(2\pi b_i^{\top} | \frac{2y}{H} - 1 |) \right]$$

where a is a scalar hyperparameter and  $b_i$  is a learnable parameter, or

$$I_{\phi}(y) = \alpha_i \cdot \exp\left(-\beta_i \left( (\kappa_i + 0.1) \left| \frac{2y}{H} - 1 \right| \right)^{(\gamma_i + 1)} \right),$$

where the input is still similarly normalized, but then passed through a modified Gaussian expansion for learnable parameters  $\alpha_i$ ,  $\beta_i$ ,  $\gamma_i$ ,  $\kappa_i > 0$ ,  $i \in [1, N_{\phi}]$ , where  $N_{\phi}$  is the number of neurons for  $\phi$ 's PINN.

The modified Gaussian expansion was chosen because its solution for y is similar in shape to the expected solution of  $\phi$ , and it is modified by including  $(\gamma_i + 1)$  (rather than a simple square) and  $(\kappa_i + 0.1)$ ; this allows for a high-frequency solution for  $\phi$ , but only where high-frequencies are expected. The  $(\cdot + 1)$  and  $(\cdot + 0.1)$  are included to prevent singularities; the chosen values 1 and 0.1 seem to be arbitrary, and the model will perform well so long as they are non-zero and positive. The modified Gaussian expansion is more problem-specific and stable than a Fourier expansion.

As for the core neural networks themselves,  $\mathcal{M}_u(\theta_u, I_u(y))$  and  $\mathcal{M}_\phi(\theta_\phi, I_\phi(y))$  are

$$\mathcal{M}(\theta_u, I_u(y)) = \sigma \big( W_u^{(L_u)} \, \sigma \big( W_u^{(L_u-1)} \cdots \sigma \big( W_u^{(1)} \cdot I_u(y) + b_u^{(1)} \big) \cdots + b_u^{(L_u-1)} \big) + b_u^{(L_u)} \big),$$

$$\mathcal{M}(\theta_{\phi}, I_{\phi}(y)) = \sigma \big( W_{\phi}^{(L_{\phi})} \, \sigma \big( W_{\phi}^{(L_{\phi}-1)} \cdots \sigma \big( W_{\phi}^{(1)} \cdot I_{\phi}(y) + b_{\phi}^{(1)} \big) \cdots + b_{\phi}^{(L_{\phi}-1)} \big) + b_{\phi}^{(L_{\phi})} \big).$$

where the number of layers for each PINN is denoted as  $L_u$  and  $L_{\phi}$ , and the weights and biases are collectively denoted as  $\theta_u = \{W_u^{(l)}, b_u^{(l)}\}_{l=1}^{L_u}$  and  $\theta_{\phi} = \{W_{\phi}^{(l)}, b_{\phi}^{(l)}\}_{l=1}^{L_{\phi}}$ . The weights and biases are in this case matrices with sizes that depend on the number of neurons  $N_u$  and  $N_{\phi}$ , respectively. Neurons contribute to PINN width, while layers contribute to PINN depth. As for the activation function,  $\sigma$ , it is the same hyperbolic tangent function for both PINNs.

$$\sigma = \tanh(\cdot) = \frac{e^{(\cdot)} - e^{-(\cdot)}}{e^{(\cdot)} + e^{-(\cdot)}}.$$

Lastly, the output transformations are  $\Gamma_u(\cdot)$  and  $\Gamma_{\phi}(\cdot)$ , where

$$\Gamma_u(\cdot) = (\cdot) \cdot (1 + I_{\phi}(y)) \cdot (1 - I_{\phi}(y)),$$

$$\Gamma_{\phi}(\cdot) = \frac{\phi_m}{1 + e^{-(\cdot)}} \cdot (1 + I_{\phi}(y)) \cdot (1 - I_{\phi}(y)),$$

which serve to enforce the zero Dirichlet boundary conditions as hard constraints for both u and  $\phi$ , and also bound the solution for  $\phi$  such that  $\phi \in [0, \phi_m]$ .

## 2 Loss Handling

The inverse problems involve solving for the velocity u and particle volume fraction  $\phi$ , and in some cases the lift force exponent  $\beta$ , using known synthetic or experimental data. First, u is computed using data points  $u_{\text{data}}(y_i)$  at specific spatial locations  $y_i$ . Subsequently,  $\phi$ , or both  $\phi$  and  $\beta$ , are determined using a combination of data and physics-based loss terms.

The loss function for each component is generalized as:

$$\mathcal{L}_{j} = m_{j} \sum_{i=1}^{M} \operatorname{mask}_{j}(\lambda_{j,i}) \cdot f_{j}(y_{i}),$$

where:

- j = 1, ..., L indexes the loss component, with L = 1 for scripts solving solely for u using  $u_{\text{data}}(y_i)$ , and L = 5 for scripts solving for  $\phi$  or both  $\phi$  and  $\beta$ .
- M is the number of collocation or data points  $y_i$ , where i = 1, ..., M.
- $m_j$  is a global scalar weight for the j-th loss term, balancing its contribution to the total loss.
- $\lambda_{j,i}$  is a local, self-adaptive weight for the *i*-th point and *j*-th loss term, enabling the model to focus on regions with higher residuals.
- $\operatorname{mask}_{j}(\lambda_{j,i}) = \operatorname{softplus}(\lambda_{j,i}) = \ln(1 + e^{\lambda_{j,i}})$  ensures non-negative weights, providing numerical stability. In some implementations,  $\operatorname{mask}_{j}(\lambda_{j,i}) = \lambda_{j,i}^{2}$  is used for stronger emphasis on high-error points.
- $f_j(y_i)$  is the residual or error term for the j-th loss component, evaluated at point  $y_i$ .

The specific loss terms  $f_j(y_i)$  depend on the problem and are defined as follows:

• Data loss for u (used when L=1):

$$f_{\text{data},u}(y_i) = (u(\theta_u, y_i) - u_{\text{data}}(y_i))^2,$$

where  $u(\theta_u, y_i)$  is the PINN prediction for velocity, and  $u_{\text{data}}(y_i)$  is the known data.

- Physics and constraint losses (used when solving for  $\phi$  or  $\phi$  and  $\beta$ , with L=5):
  - 1. Particle conservation:

$$f_{\text{particle}}(y_i) = (\nabla \cdot \mathbf{J}(y_i))^2$$
,

enforcing the divergence-free condition for the particle migration flux J.

2. Momentum balance (xy-component):

$$f_{\text{mom},xy}(y_i) = \left( \left( \nabla \cdot \mathbf{\Sigma} \right)_{xy} (y_i) \right)^2,$$

where  $\Sigma$  is the total stress tensor.

3. Momentum balance (yy-component):

$$f_{\text{mom},yy}(y_i) = \left( \left( \nabla \cdot \mathbf{\Sigma} \right)_{yy} (y_i) \right)^2.$$

4. Mass conservation:

$$f_{\text{mass}}(y_i) = \left(\frac{1}{M} \sum_{i=1}^{M} \phi(\theta_{\phi}, y_i) - \phi_{\text{avg}}\right)^2,$$

ensuring the average particle volume fraction matches the known  $\phi_{\rm avg}.$ 

5. Symmetry constraint:

$$f_{\text{sym},\phi}(y_i) = (\phi(\theta_{\phi}, y_i) - \phi(\theta_{\phi}, -y_i))^2,$$

enforcing symmetry of  $\phi$  across the channel centerline.