Boring but important disclaimers:

If you are not getting this from the GitHub repository or the associated Canvas page (e.g. CourseHero, Chegg etc.), you are probably getting the substandard version of these slides Don't pay money for those, because you can get the most updated version for free at

https://github.com/julianmak/academic-notes

The repository principally contains the compiled products rather than the source for size reasons.

- Associated Python code (as Jupyter notebooks mostly) will be held on the same repository. The source data however might be big, so I am going to be naughty and possibly just refer you to where you might get the data if that is the case (e.g. JRA-55 data). I know I should make properly reproducible binders etc., but I didn't...
- ▶ I do not claim the compiled products and/or code are completely mistake free (e.g. I know I don't write Pythonic code). Use the material however you like, but use it at your own risk.
- As said on the repository, I have tried to honestly use content that is self made, open source or explicitly open for fair use, and citations should be there. If however you are the copyright holder and you want the material taken down, please flag up the issue accordingly and I will happily try and swap out the relevant material.

OCES 2003 : Descriptive Physical Oceanography

(a.k.a. physical oceanography by drawing pictures)

Lecture 16: Dynamics 2 (waves and dynamic mechanisms)

Outlook of the next few lectures

Dynamics important, next few lectures on

- waves (this Lec. + 16, 18) and instabilities (Lec. 17)
 - → because waves are easier to talk about without maths...

Highlight gross features (i.e. those that can be drawn...)

- ▶ how to describe waves (Lec. 15)
- types of waves (Lec. 16)
 - \rightarrow consequence + leading to instabilities
- ▶ instabilities (Lec. 17)
 - → parcel-type (mechanistic) arguments for instability
- ▶ tides (particularly as internal gravity waves) (Lec. 18)

Outline

- gravity waves
 - \rightarrow gravity/buoyancy as restoring mechanism (\sqrt{gH})
- ▶ inertial waves
 - \rightarrow Coriolis as restoring mechanism (f)
 - \rightarrow e.g. Rossby waves, Kelvin waves
- ▶ inertial-gravity + internal waves $(\sqrt{gH} \text{ or } N, \text{ and } f)$
 - \rightarrow extra depth dimension to deal with
 - \rightarrow Brunt–Väisälä or buoyancy frequency N
- propagation mechanism (Rossby wave example)
 - → kinematic argument with **vorticity**

Key terms: buoyancy frequency, gravity waves, inertial waves, Rossby waves, Kelvin waves, vorticity inversion



Recap: what goes down must come up

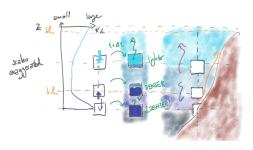


Figure: Schematic of the diffusive upwelling.

▶ diapycnal mixing contribute upwelling, strongest in boundary layers
 → broad diffusive boundary intensified upwelling

what causes the bounary intensification of κ_d ? dynamics!

- at the surface, lots of things... (convection, waves, Langmuir turbulence etc.)
- ▶ at the bottom, probably tidal conversion (Lec. 18) \rightarrow internal gravity waves (Lec. 16) \rightarrow shear instabilities (Lec. 17)



Recap: waves and dispersion relation

- waves are ubiquitous physical features
 - \rightarrow depends on physics
- wave described by the dispersion relation $\omega = \mathcal{F}(k)$
 - \rightarrow physics dictate the form of \mathcal{F}









Figure: Examples of systems supporting waves. All figures from Wikipedia except the cello one



Figure: Gravity waves with signal at the sea surface (as darker and lighter bands). Taken at HKUST.

- ► (linear) waves can interfere with each other
 - → constructive or destructive
 - → interference can lead to steepening and breaking ("becoming"

Recap: wave propagation

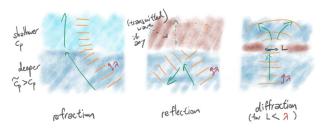


Figure: Schematic of refraction, reflection (and transmission), and diffraction. See Lec. 15.

b phase speed (in a direction) and group velocity as (note $\omega = \mathcal{F}(k)$)

$$c_{p,x} = \frac{\omega}{k}, \qquad c_{g,x} = \frac{\partial \omega}{\partial k}$$

- → individual wave vs. wavepacket behaviour
- \rightarrow contribute to wave phenomenon (e.g. refraction from

$$c_p = c_p(x)$$



Wave steepening and breaking



Figure: Schematic of mixing by (irreversible) wave breaking, with contours reconnecting leading to e.g. diapycnal mixing.

- growing waves by instability
 - \rightarrow convective and/or shear (see Lec. 17)
 - → mixing of material **across** isopycnals after reconnection, leading to diapycnal mixing
- feedback onto MOC (see Lec. 14)

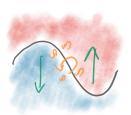


Figure: Velocity shear from waves can lead to mixing.

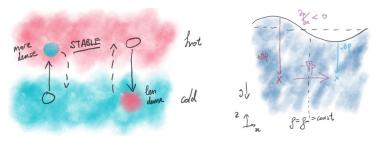


Figure: Gravity as restoring force. Pictures adapted from ones used in Lec. 7.

- deviation from resting isopycnal experiences restoring force from gravity (buoyancy)
 - → left case: internal isopycnal (as an **isotherm**)
 - \rightarrow right case: sea surface is the isopycnal
- ▶ weak damping, restoring force, overshoots ⇒ oscillatory motion (up and down in this case)

For simplicity, consider a **homogeneous** (i.e. $\rho = \text{const}$) layer of fluid (cf. Lec. 11 + 12, right case in previous Figure) until we get to internal waves

For simplicity, consider a **homogeneous** (i.e. $\rho = \text{const}$) layer of fluid (cf. Lec. 11 + 12, right case in previous Figure) until we get to internal waves

In this instance, dispersion relation for gravity waves is given by (without derivation)

$$\omega^2 = gk \tanh(kH)$$
 \Rightarrow $\omega = \pm \sqrt{gk \tanh(kH)}$

- ► H is water depth, and tanh = hyperbolic tangent, goes from -1 to 1
 - \rightarrow note symmetry in both directions (the \pm sign)
 - ($\omega \leq 0$ cases are just **shifts** in the **phase**)

$$\omega = \pm \sqrt{gk \tanh(kH)}$$

▶ for **deep** water waves ($kH \gg 1$) and **shallow** water waves ($kH \ll 1$),

$$\omega_{\mathrm{deep}} = \pm \sqrt{gk}, \qquad \omega_{\mathrm{shallow}} = \pm k\sqrt{gH}$$

- $\rightarrow kH \gg 1$ so $tanh(kH) \rightarrow 1$
- $ightarrow kH \ll 1$ with $anh(kH) pprox kH + O((kH)^3)$ (do a Taylor expansion)
- deep water waves are depth-independent and dispersive
- ▶ shallow water waves are slower in shallow waters $(c_p \sim \sqrt{H})$ and non-dispersive $(c_p = c_g)$



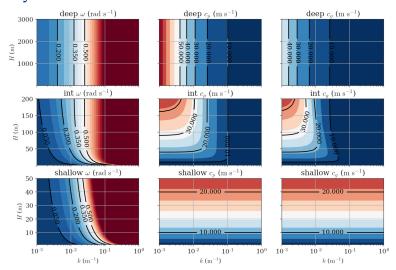
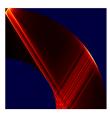


Figure: Water wave ω , c_p and c_g over (k,H) space, with k shown on a log axis. k chosen so wavelengths are roughly between 50 m to 5 km (recall $k=2\pi/\lambda$). Also note the transitions from shallow to intermediate to deep are really to do with $kH\sim H/\lambda$. See waves.ipynb.

Inertial waves

Inertial waves has the Coriolis "force" act as the restoring force

- generic for rotating systems
 - \rightarrow planetary interiors, stars, galactic disks
- mostly arise in context of internal waves
 - → at surface buoyancy effects can dominate
- ► limited in frequency by inertial frequency f_0 (cf. Coriolis parameter)
 - → revisit later when talking about inertia-gravity waves



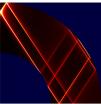


Figure: Inertial wave attractors in a homogeneous planetary interior at different tidal forcings. From Gordon Ogilvie (2009, Mon. Not. Royal. Astro. Soc).

Inertial-Gravity waves

In reality Coriolis and buoyancy effects both contribute

- ▶ large-scale and/or slow \Rightarrow Coriolis important (because Ro \ll 1), classify as inertial waves
 - \rightarrow e.g. Rossby waves
- ightharpoonup small-scale and/or fast \Rightarrow Coriolis unimportant
 - \rightarrow e.g. internal gravity waves
- somewhere in between? Poincaré or inertia-gravity waves

$$\omega = \pm \sqrt{f_0^2 + gH(k_x^2 + k_y^2)}$$

- \rightarrow for $gH(k_x^2 + k_y^2) \gg f_0$, recover gravity waves
- \rightarrow for $gH(k_x^2 + k_y^2) \ll f_0$, recover inertial waves



Rossby deformation radius

$$\omega = \pm \sqrt{f_0^2 + gH(k_x^2 + k_y^2)}$$

boundary between gravity and inertial regimes is roughly

$$L_d = \frac{\sqrt{gH}}{f_0}$$

- the Rossby deformation radius (for shallow water system)
 - \rightarrow roughly also the boundary where geostrophic approximation should hold (see Lec. 8 + 13)
 - \rightarrow estimates in a few slides

introduce a useful quantity

$$N = \sqrt{-\frac{g}{\rho_0} \frac{\partial \rho}{\partial z}}$$

- ► Brunt–Väisälä or buoyancy frequency (units: s⁻¹)
 - $\rightarrow N^2$ normally used
 - \rightarrow note $\partial \rho / \partial z < 0$ for stable stratification, i.e.

$$N^2 > 0$$

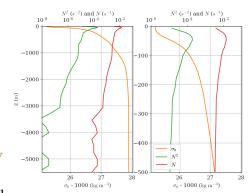


Figure: σ_0 (see Lec. 6) and the associated N^2 and N. $N^2 \ll 1$ means weakly stratified (weak density gradients), whilst $N^2 < 0$ shows unstable stratification (none in this case, but see Lec. 17). See plot-eos.ipynb.

simplistic view (!): $\sqrt{gH} \rightarrow N$

Generally then, internal inertia-gravity waves described by

$$\omega = \pm \sqrt{\frac{f_0^2 k_z^2 + N^2 k_x^2}{k_x^2 + k_z^2}} \approx \pm \sqrt{f_0^2 + \frac{N^2 k_x^2}{k_z^2}} \quad \text{(for } |k_z| \gg |k_x|\text{)}$$

- ► atmosphere and ocean has $N/f_0 = O(10^1 \text{ to } 10^2)$
 - \rightarrow so really we have gravity waves influenced by rotation
 - → refer to them here as internal waves

Generally then, internal inertia-gravity waves described by

$$\omega = \pm \sqrt{\frac{f_0^2 k_z^2 + N^2 k_x^2}{k_x^2 + k_z^2}} \approx \pm \sqrt{f_0^2 + \frac{N^2 k_x^2}{k_z^2}} \quad \text{(for } |k_z| \gg |k_x|\text{)}$$

- ► atmosphere and ocean has $N/f_0 = O(10^1 \text{ to } 10^2)$
 - \rightarrow so really we have gravity waves influenced by rotation
 - → refer to them here as internal waves
- ▶ note that $|f_0| \le |\omega| \le |N|$
 - \rightarrow since $0 \le k_{x,z}^2/(k_x^2 + k_z^2) \le 1$
 - \rightarrow frequency is **much lower** than gravity waves

Generally then, internal inertia-gravity waves described by

$$\omega = \pm \sqrt{\frac{f_0^2 k_z^2 + N^2 k_x^2}{k_x^2 + k_z^2}} \approx \pm \sqrt{f_0^2 + \frac{N^2 k_x^2}{k_z^2}} \quad \text{(for } |k_z| \gg |k_x|\text{)}$$

- ► atmosphere and ocean has $N/f_0 = O(10^1 \text{ to } 10^2)$
 - \rightarrow so really we have gravity waves influenced by rotation
 - → refer to them here as internal waves
- ▶ note that $|f_0| \le |\omega| \le |N|$
 - \rightarrow since $0 \le k_{x,z}^2/(k_x^2 + k_z^2) \le 1$
 - → frequency is **much lower** than gravity waves

internal tides to be seen as internal waves (Lec. 18)



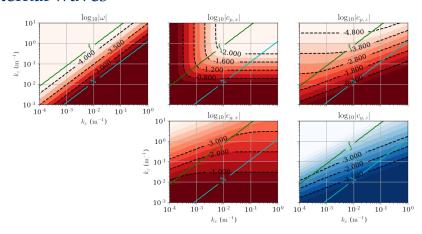


Figure: Inertial-gravity waves (with the $k_z\gg k_x$ approximation) ω , $c_{p,x}$, $c_{p,y}$, $c_{g,x}$ and $c_{p,y}$ as a log-log plot in (k_x,k_z) space, with $f=5\times 10^{-5}$ and $N=3\times 10^{-3}$ (oceanic relevant values). The contours denote the exponent x of $|10^x|$ and the colour shading denotes the sign (more blue = more negative actual values rather than exponents, more red = more positive actual values rather than exponents, since k_x and k_z is chosen to be positive, everything except $c_{g,z}$ is positive. Contours of f and N plotted with an offset plotted to show the boundary beyond which everything is either gravity waves or inertial oscillations. See waves . ipynb.

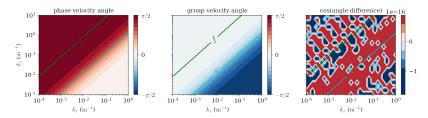


Figure: Inertial-gravity waves (with the $k_z \gg k_x$ approximation) phase velocity c_p angles and group velocity c_g angles (in radians, relative to the horizontal, and note $\pi/2 = 90^{\circ}$). The final panel shows $c_p \cdot c_g = |c_p||c_g|\cos\theta$ (which is zero up to rounding errors). Contours of f and N plotted with an offset plotted as in the previous diagram. See waves . i.pynb.

note that, for inertial-gravity waves (left as a bonus exercise),

$$c_p \cdot c_g = 0$$

ightarrow i.e. phase and group velocities are **perpendicular** to each other (see Lec. 4)



Deformation radius

 boundary given again by the Rossby deformation radius (for the continuously stratified case)

$$L_d = \frac{NH}{f}$$

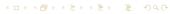
- \rightarrow $L_{d,atmos} = O(1000 \text{ km})$, scale of cyclones and anti-cyclones, i.e. weather systems form (synoptic structures)
- \rightarrow $L_{d,ocean} = O(50 \text{ km})$, scale of ocean eddies
- ▶ latitude (through *f*) and *H* dependent
 - \rightarrow smaller L_d for **high** latitudes and **shallow** regions
 - → consequence for **geostrophic approximation?** (e.g. shelves and coasts, see Lec. 21 + 22)

Deformation radius

 boundary given again by the Rossby deformation radius (for the continuously stratified case)

$$L_d = \frac{NH}{f}$$

- \rightarrow $L_{d,atmos} = O(1000 \text{ km})$, scale of cyclones and anti-cyclones, i.e. weather systems form (synoptic structures)
- \rightarrow $L_{d,ocean} = O(50 \text{ km})$, scale of ocean eddies
- ▶ latitude (through *f*) and *H* dependent
 - \rightarrow smaller L_d for **high** latitudes and **shallow** regions
 - \rightarrow consequence for **geostrophic approximation?** (e.g. shelves and coasts, see Lec. 21 + 22)
- internal L_d defined analogously (normally smaller than above)



Kelvin waves (more on this in Lec. 18, 21 + 22)

A type of boundary wave

- ► need *f* and a **boundary**
 - \rightarrow could be land (coastal Kelvin waves) (see Lec. 18, 21 + 22)
 - \rightarrow could be a wave guide (e.g. equator where f changes sign, equatorial Kelvin waves) (see OCES 4001, El-Niño, QBO etc.)
- needs f but propagates at the gravity wave speed, with

$$\omega = k\sqrt{gH}$$

- → non-dispersive
- → fairly fast (gravity wave speed)

Kelvin waves (more on this in Lec. 18, 21 + 22)

A type of boundary wave

- need f and a boundary
 - \rightarrow could be land (coastal Kelvin waves) (see Lec. 18, 21 + 22)
 - \rightarrow could be a wave guide (e.g. equator where f changes sign, equatorial Kelvin waves) (see OCES 4001, El-Niño, QBO etc.)
- needs f but propagates at the gravity wave speed, with

$$\omega = k\sqrt{gH}$$

- → non-dispersive
- \rightarrow fairly fast (gravity wave speed)
- NOTE the lack of ±!

Kelvin waves (more on this in Lec. 18, 21 + 22)

 boundary introduces asymmetry in this case: general solution like

$$\eta \sim \mathrm{e}^{\pm f_0 y / \sqrt{gH}} \cos(kx - \omega t)$$

- \rightarrow take $y \le 0$ to be **boundary**, if $f_0 > 0$ (NH), need minus sign, and vice-versa
- → wave propagates cyclonically (same sign as f)
- ▶ taking $f_0 > 0$ (NH),

$$\eta \sim e^{-y/L_d} \cos(kx - \omega t),$$

so decay over the $L_d = \sqrt{gH}/f_0$

Rossby waves (more on this later)

A (particularly important) type of **inertial** wave

- requires a gradient in background vorticity
 - $\rightarrow \partial f/\partial y = \beta$ (planetary case)
 - ightarrow background flow $-\partial U/\partial y \sim \nabla \times u$ (see later and Lec. 17)
- dispersion relation given by (on β -plane)

$$\omega = -\frac{\beta k_x}{k_x^2 + k_y^2}$$

 \rightarrow note that Rossby waves propagate to the **west** (more generally, **retrograde** or against the mean flow) Since

$$c_{p,x} = \frac{\omega}{k_x} = -\frac{\beta}{k_x^2 + k_y^2} < 0,$$

and **long** waves ($k_x \ll 1$) are **fast**(er)



Rossby waves (more on this later)

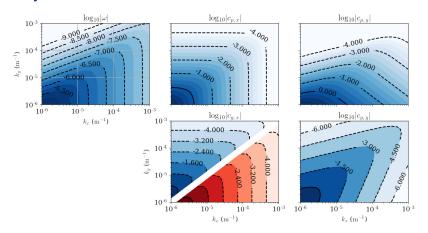


Figure: Rossby waves ω , $c_{p,x}$, $c_{p,y}$, $c_{g,x}$ and $c_{p,y}$ as a log-log plot in (k_x, k_y) space, with magnitude also as logs. The contours denote the exponent x of $|10^x|$ and the colour shading denotes the sign (more blue = more negative actual values, more red = more positive actual values); since k_x and k_y is chosen to be positive, everything except $c_{g,x}$ is negative. Choice of k_x and k_y correspond to wavelengths roughly between 6 km to 6000 km (Rossby waves are usually seen as planetary-scale waves). See waves.ipynb.

Propagation mechanism: Rossby waves

Rossby waves propagate west-ward (or, more generally, retrograde)

$$c_{p,x} = \frac{\omega}{k_x} = -\frac{\beta}{k_x^2 + k_y^2} < 0$$
why?

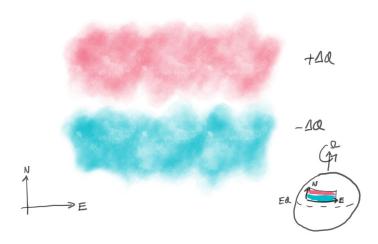
Propagation mechanism: Rossby waves

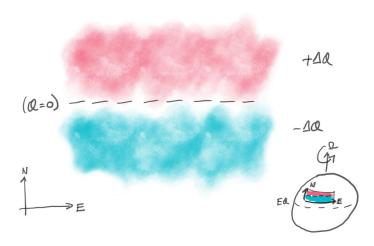
Rossby waves propagate west-ward (or, more generally, retrograde)

$$c_{p,x} = \frac{\omega}{k_x} = -\frac{\beta}{k_x^2 + k_y^2} < 0$$
 why?

Key bits to the pictorial/parcel (cf. Lec 5 for temperature) argument:

- ► the initial wave conserves and carries vorticity (spini-ness, recall Lec. 4, 11, 12) into the external environments
 - → these are now vorticity anomalies
- vorticity anomalies induces a velocity/flow (because spini-ness)
- ▶ induced flow seen to self-advect the wave and move it to the West (retrograde in the general case)





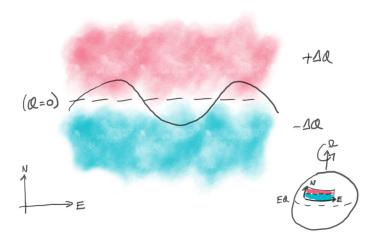
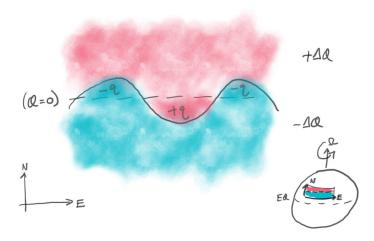
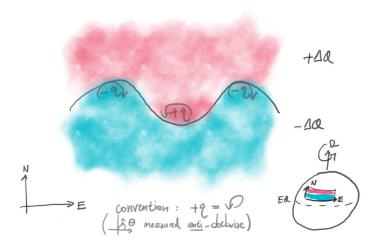
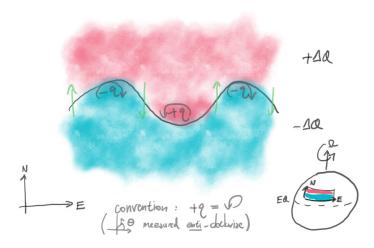


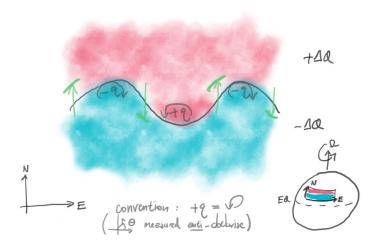
Figure: Rossby wave propagation schematic.

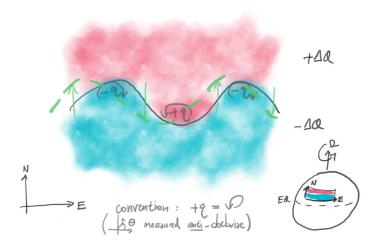




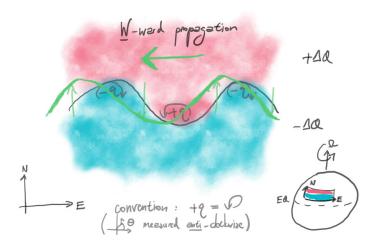














Summary

- gravity waves (gravity/buoyancy)
- inertial waves (Coriolis)
- inertial-gravity waves (general)
 - \rightarrow internal waves have

$$|f| \le |\omega| \le |N|$$

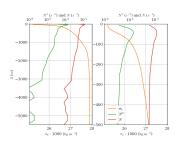


Figure: σ_0 (see Lec. 6) and the associated N^2 and N. See plot_eos.ipvnb.

► Brunt–Väisälä or buoyancy frequency *N*

$$N = \sqrt{-\frac{g}{\rho_0} \frac{\partial \rho}{\partial z}}$$

 \rightarrow measure of stratification strength (see also Lec. 17)



Summary

- parcel argument for west-ward
 Rossby wave propagation
 - → conservation of vorticity
 - → vorticity anomalies induces flow
 - $\rightarrow self\text{-}advecting$

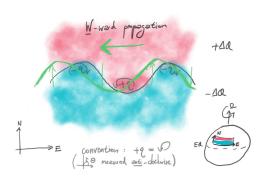


Figure: Rossby wave propagation schematic.

▶ generalisations exist (e.g. internal gravity waves in Harnik et al., 2008, J. Atmos. Sci)

Summary

- parcel argument for west-ward Rossby wave propagation
 - → conservation of vorticity
 - → vorticity anomalies induces flow
 - $\rightarrow self\text{-}advecting$

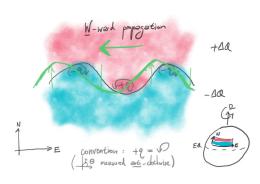


Figure: Rossby wave propagation schematic.

- peneralisations exist (e.g. internal gravity waves in Harnik et al., 2008, J. Atmos. Sci)
- ▶ two such waves interacting? (see Lec. 17)
 - → potential for instabilities

