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DEPARTMENT OF PHYSICS

DEGREE IN PHYSICS – LAUREA IN FISICA

THESIS

Phase-matching method by temperature control of partially degenerate Four-Wave Mixing fenomena

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Abstract: studying degenerate Four-Wave Mixing fenomena in a rectangular silicon-based waveguide. Analyze its behaviour at different temperatures to maximize phase-match. Developing a project to build an actual prototype.

Contents

1 Introduction

1.1 Motivation

In my scholastic career I approached the study of optics in an elective course of my third year of undergraduate academic degree. I became very interested in the subject, in particular because of the potential developement of photonics for communication and computing. I discussed different topics with the professors and some researchers. One of the researcher, who would later become my advisor, suggested me a few interesting topics in guided-wave optics which could be studied as a undergraduate thesis.

I chose to study the temperature behaviour of silicon waveguides. The reason of the choice is that there are always discrepancies between the design of silicon objects and their actual production by the factory. The aim of the study is to exploit behaviour of silicon to correct these possible production defects of the materials.

The study is focused on a silicon-based waveguide with a silicon (Si) core of rectangular section and a silica (SiO_2) cladding. The purpose of the waveguide is to generate light through a partial degenerate Four-Wave Mixing phenomenon.

1.2 Partially degenerate Four-Wave Mixing

Four-Wave Mixing (FWM) is a nonlinear optical effect due to the susceptibility of silicon. Silicon is an isotropic material and therefore second order susceptibility $\chi^{(2)}$ (almost) vanishes. Third order susceptibility $\chi^{(3)}$, on the other hand, is non-zero and is responsible among the other effects for FWM, which is indeed a third order phenomenon.

In general FWM implies the mix of four waves of different frequency. Partial degenerate FWM, instead, is a specific case in which two waves have the same frequency and are considered as one (pump). Therefore it is about the pump wave and a second weaker one (signal) which generate a third wave (idler) of frequency different from that of the former two.

2 Electromagnetic waves in dielectric media

To understand electromagnetic phenomena the starting point is the well known Maxwell's equations:

$$\nabla \times \mathbf{E} = -\frac{\partial \mathbf{B}}{\partial t} \tag{1a}$$

$$\nabla \times \mathbf{H} = \mathbf{J} + \frac{\partial \mathbf{D}}{\partial t} \tag{1b}$$

$$\nabla \cdot \mathbf{D} = \rho_{\mathrm{f}} \tag{1c}$$

$$\nabla \cdot \mathbf{B} = 0 \tag{1d}$$

where $\mathbf{E}$ and $\mathbf{H}$ are electric and magnetic field vectors and $\mathbf{D}$ and $\mathbf{B}$ are corresponding electric and magnetic flux densities. $\mathbf{J}$ is the current density vector and ρ_{f} is the free charge density. In dielectric media, such as silicon, we can also consider $\mathbf{J} = 0$ and $\rho_{\mathrm{f}} = 0$.

$\mathbf{D}$ and $\mathbf{B}$ are related to the electromagnetic fields through the relations:

$$\mathbf{D} = \varepsilon_0 \mathbf{E} + \mathbf{P} \tag{2a}$$

$$\mathbf{H} = \mu_0 \mathbf{H} + \mathbf{M} \tag{2b}$$

where ε_0 is the vacuum permittivity, μ_0 is the vacuum permeability, and $\mathbf{P}$ and $\mathbf{M}$ are the induced electric and magnetic polarizations repectively. For a nonmagnetic medium, as in our case, $\mathbf{M} = 0$.

Solution in linear dielectric medium

A linear dielectric medium is characterized by a linear relation between the polarization density and the electric field:

$$\mathbf{P} = \varepsilon_0 \chi \mathbf{E} \tag{3}$$

where χ is the susceptibility of the medium and $\varepsilon_r \equiv 1 + \chi$ is the relative dielectric constant.

Taking the curl of Eq. (1a) and substituting in it Eq. (1c) it leads to the Helmholtz equation:

$$\nabla^2 \mathbf{E} - \frac{1}{c^2} \frac{\partial^2 \mathbf{E}}{\partial t^2} = \mu_0 \frac{\partial^2 \mathbf{P}}{\partial t^2} \tag{4}$$

where $c = 1/\sqrt{\mu_0 \varepsilon_0}$ is the speed of light in vacuum.

Defining the refractive index $n \equiv \sqrt{\varepsilon_r}$ of the material the previous equation (4) and using the definition of $\mathbf{P}$ as in Eq. (3)

$$\nabla^2 \mathbf{E} - \frac{n^2}{c^2} \frac{\partial^2 \mathbf{E}}{\partial t^2} = 0 \tag{5}$$

Moreover, n represent the ratio between the speed of light in vacuum and in the material $n = c/v$ and it depends on the frequency of the EM wave.

The general solution of Eq. (4) is a transverse electromagnetic (TEM) wave, which electric and magnetic fields are perpendicular to the propagation direction. Without losing generality, we could assume $\hat{z}$ as the propagation direction and electric field of the form:

$$\mathbf{E} = \frac{1}{2} \hat{x}\, E \exp[\mathrm{i}(kz - \omega t)] + c.c. \tag{6}$$

with E amplitude of the electric field, $k = 2\pi/\lambda$ the wavevector and $\omega = 2\pi f$ the angular frequency. λ and f are relatively the wavelength and the frequency of the electromagnetic wave. The magnetic field has a similar equation.

2.1 Silicon waveguides

In linear optics the technology for trasmitting a light wave by confining it in a finite space is the guided-wave optics. The instruments employed for achieving such purpose are called waveguides.

In a ray-optics picture, silicon waveguides are dielectric waveguides and make use of the interface between two media with different refractive index. Precisely they exploit the phenomenon of total internal reflection: if a propagating wave reaches a boundary between two mediums, one with higher refractive index (n_H) and another with lower refractive index (n_L), with an angle of incidence greater than the *critical angle*, then the wave is completely reflected in the first medium. Angles are defined relative to the normal vector of the interface and the critical angle can be obtained by the Snell law:

$$\theta_C = \arcsin\left(\frac{n_L}{n_H}\right)$$

Silicon waveguides are composed by core and cladding **(?)**. The former is the higher refractive index medium in which the wave is confined, the latter is the lower refractive index medium which surrounds the core, thus creating the interface. Theoretically the cladding could also be vacuum.

Wave equation and modes

Considering an ideal dielectric waveguide (no absorption) with the purpose of transmitting EM waves in the $\hat{z}$ direction, the total field distribution is a sum of continous TEM plane waves. Defining the amplitude of the fields as $E(\mathbf{r}, t) = F(x, y)A(z, t)$ and $k = \beta$ and substituting them in Eq. (6) leads to the following equation:

$$\mathbf{E} = \frac{1}{2}\hat{x}\, F(x, y)A(z, t) \exp[\mathrm{i}(\beta z - \omega t)] + c.c.$$

The previous equation represents what is called a "transverse electric" (TE) wave. A second family of solution could be acquired by defining the sum of TEM wave with the magnetic field always perpendicular to the propagation direction. This is called a "transverse magnetic" (TM) wave.

Is then possible to study the propagation of the wave along the waveguide using the Helmholtz equation (4) or (5), obtaining the next equations:

$$\frac{\partial^2 F(x, y)}{\partial x^2} + \frac{\partial^2 F(x, y)}{\partial y^2} + \left[k_0^2 n^2(x, y) - \beta^2\right] F(x, y) = 0 \tag{7a}$$

$$\frac{\partial^2 \tilde{A}(z, \omega)}{\partial z^2} + 2\mathrm{i}\beta \frac{\partial \tilde{A}(z, \omega)}{\partial z} + 3n^2 k_0^2 \tilde{A}(z, \omega) = 0 \quad ! \neq ! \quad 2\mathrm{i}\beta_0 \frac{\partial \tilde{A}(z, \omega)}{\partial z} + (\tilde{\beta}^2 - \beta_0^2)\tilde{A}(z, \omega) = 0 \tag{7b}$$

where k_0 is the wavevector in vacuum and $n = n(x, y)$ is a function of the position (mainly x and y for a straight waveguide).

The first equation, once solved, leads to the field distribution $F(x, y)$ in the plane perpendicular to the waveguide axis ($\hat{z}$). Its solution is strictly linked to the geometry of the waveguide and to the dependence of the refractive index n to the position in the plane.

The second equation (7b), on the other hand, define the change of field distribution along the $\hat{z}$ axis, or the waveguide axis. Imposing to this equation to be zero means that we force the field distribution F not to change. Indeed not all field distributions are transmitted through the waveguide without energy loss **(?)** and thus without changing F. Those field distributions which maintain the same transverse distribution and polarization at all locations along the waveguide axis are called modes. These modes can be categorized in two main groups, depending on the polarization of the wave: transverse electric (TE) modes and transverse magnetic (TM) modes, which have respectively the electric and magnetic field transverse to the waveguide axis.

Solving for β, we obtain that only a discrete set of values β_m verify the equation. Replacing those values in Eq. (7a), leads to the field distribution F(x,y) to be defined. Therefore both TE and TM modes can be again classified by their order, that is proportional to the number of maxima of the modulus of the electric and magnetic fields in the core section. The first order has only one maximum at the center of the core, the second order has two maxima placed simmetrically from the center and so on.

Moreover it is noteworthy that the field distributions of even orders are even functions and those of odd orders are odd functions.

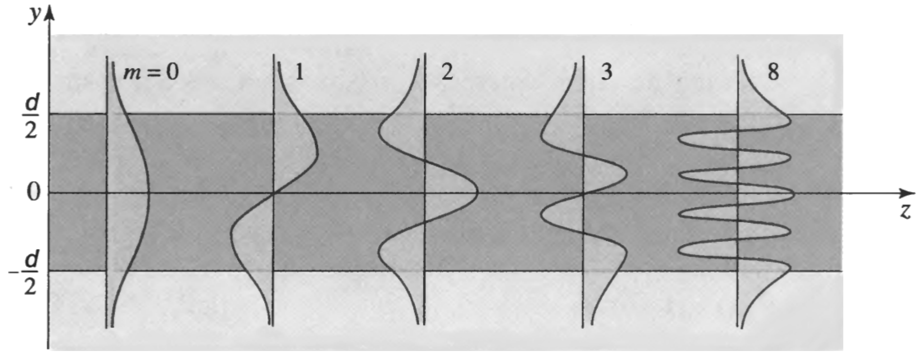

Figure 1: graphical rapresentation of field distribution of 1D modes

Propagation constant and effective index

A very important parameter of each mode is its *propagation constant* β, defined as the component of the wavevector k in the waveguide axis:

$$\beta \equiv k|_z \tag{8}$$

Higher-order modes travel with smaller propagation constants. For a 1D slab waveguide β has a simple relation with the bounce angles θ_m, which are quantized between 0 and $\bar{\theta}_C = \pi - \theta_C$:

$$\beta_m = n_H k_0 \cos\theta_m$$

From the propagation constant, we can also define the refractive index relative to the propagation mode, called *effective index*:

$$\beta = k_0 n_{eff} \quad \Rightarrow \quad n_{eff} \equiv \frac{\beta}{k_0} \tag{9}$$

Its value is between the higher index of the core and lower one of the cladding. Again, for the 1D slab waveguide case, the effective index can be rewritten as in the following formula:

$$n_{eff} \equiv n_H \cos\theta_m$$

2.2 Rectangular dielectric waveguides

In general, in a rectangular dielectric waveguide the classification of modes need two indexes because we have two degrees of freedom. In Figure (2.2) is *** a graphical representation of the distribution of the fields in modulus.

For our purpose, due to the thin geometry of the core, modes with the higher effective index are fundamental modes in the height axis. The first row is of our interest. Only one significant index is then left to describe the order of modes.

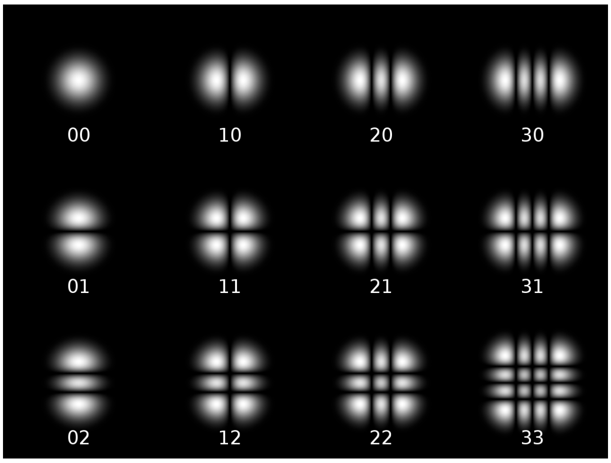

Figure 2: graphical rapresentation of field distribution of 2D modes

3 Non-linear optics

A nonlinear dielectric medium is characterized by a nonlinear relation between **P** and **E**, such as

$$\mathbf{P} = \varepsilon_0 \left(\chi^{(1)} \cdot \mathbf{E} + \chi^{(2)} : \mathbf{E}^2 + \chi^{(3)} \vdots \mathbf{E}^3 + \cdots \right) = \mathbf{P}_{\mathrm{L}} + \mathbf{P}_{\mathrm{NL}} \tag{10}$$

which originates from the behaviour of bound electrons.

In everyday condition linear effects are much larger than nonlinear ones. The relation between **P** and **E** becomes nonlinear when E has a value comparable to interatomic electric fields, which are typically $\sim 10^5 - 10^8\ \mathrm{V\,m^{-1}}$.

For an isotropic medium the second order term is zero (in the dipole approximation) and this is our scenario. Thus the dominant nonlinearity is the third order and the material is called a Kerr medium.

3.1 Third order effect: Four-Wave Mixing

Many processes are result of third order nonlinearities: Third-Harmonic Generation (THG), Kerr Effect, Cross-phase modulation (XPM), Self-phase modulation (SPM) and Four-Wave Mixing (FWM).

The parametric process of Four-Wave Mixing originates from the third order nonlinear response of a material to an electromagnetic field and involves interaction among four optical waves.

The nonlinear component of the induced electric polarization vector, considering significant only the third-order term, is

$$\mathbf{P}_{NL} = \varepsilon_0 \chi^{(3)} \vdots \mathbf{E}^3 \tag{11}$$

Considering a silicon waveguide with a rectangular core section, to study the propagation of four wave, the total electric field and be expressed as

$$\mathbf{E} = \frac{1}{2} \hat{x} \sum_{j=1}^{4} E_j \exp[i(\beta_j z - \omega_j t)] + c.c. \tag{12}$$

Substituting Eq. (11) in Eq. (12) we obtain $8^3 = 512$ terms. We can then write the non linear terms $\mathbf{P}_{\rm NL}$ of the induced electric polarization as

$$\mathbf{P}_{\rm NL} = \frac{1}{2}\hat{x}\sum_{j=1}^{4} P_j \exp[i(\beta_j z - \omega_j t)] + c.c. \tag{13}$$

where for example P_4 would be

$$\begin{aligned} P_4 \quad &= \tfrac{3\varepsilon_0}{4}\chi^3_{xxxx}\left[|\mathrm{E}_4|^2\mathrm{E}_4 2(|\mathrm{E}_1|^2 + |\mathrm{E}_2|^2 + |\mathrm{E}_3|^2)\mathrm{E}_4 \right. \\ &\left. + 2\mathrm{E}_1\mathrm{E}_2\mathrm{E}_3\mathrm{e}^{i\phi_+} + 2\mathrm{E}_1\mathrm{E}_2\mathrm{E}_3^*\mathrm{e}^{i\phi_-}\right] \end{aligned}$$

with ϕ_+ and ϕ_- defined as

$$\begin{aligned} \phi_+ \quad &= (\beta_1 + \beta_2 + \beta_3 - \beta_4 +)z - (\omega_1 + \omega_2 + \omega_3 - \omega_4)t \\ \phi_- \quad &= (\beta_1 + \beta_2 - \beta_3 - \beta_4 +)z - (\omega_1 + \omega_2 - \omega_3 - \omega_4)t \end{aligned}$$

Replacing Eq. (13) in Helmholtz equation (4) leads to eight coupled equations: four of them are equivalent of Eq. (7a) and describe the field distribution of the mode of each wave. The remaining four are the equivalent of Eq. (7b) and describe how the field change while travelling through the waveguide. Those, considering $E_j(\mathbf{r}) = F_j(x, y)A_j(z)$ where F_j is the spatial distribution of the mode, can be written as:

$$\frac{dA_1}{dz} = \frac{\mathrm{i}n_L\omega_1}{c}\left[\left(f_{11}|A_1|^2 + 2\sum_{k\neq 1} f_{1k}|A_k|^2\right)A_1 + 2f_{1234}A_2^*A_3A_4\mathrm{e}^{+\mathrm{i}\Delta k_z}\right] \tag{14a}$$

$$\frac{dA_2}{dz} = \frac{\mathrm{i}n_L\omega_2}{c}\left[\left(f_{22}|A_2|^2 + 2\sum_{k\neq 2} f_{2k}|A_k|^2\right)A_2 + 2f_{2134}A_1^*A_3A_4\mathrm{e}^{+\mathrm{i}\Delta k_z}\right] \tag{14b}$$

$$\frac{dA_3}{dz} = \frac{\mathrm{i}n_L\omega_3}{c}\left[\left(f_{33}|A_3|^2 + 2\sum_{k\neq 3} f_{3k}|A_k|^2\right)A_3 + 2f_{3412}A_1A_2A_4^*\mathrm{e}^{-\mathrm{i}\Delta k_z}\right] \tag{14c}$$

$$\frac{dA_4}{dz} = \frac{\mathrm{i}n_L\omega_4}{c}\left[\left(f_{44}|A_4|^2 + 2\sum_{k\neq 4} f_{4k}|A_k|^2\right)A_4 + 2f_{4321}A_1A_2A_3^*\mathrm{e}^{-\mathrm{i}\Delta k_z}\right] \tag{14d}$$

where

$$f_{ij} = \frac{\langle |F_i|^2|F_j|^2\rangle}{\sqrt{\langle |F_i|^2\rangle\langle |F_j|^2\rangle}} \tag{15}$$

and

$$f_{ijkl} = \frac{\langle F_i^*F_j^*F_kF_l\rangle}{\sqrt{\langle |F_i|^2\rangle\langle |F_j|^2\rangle\langle |F_k|^2\rangle\langle |F_l|^2\rangle}} \tag{16}$$

are the overlap integrals of the different modes, where the brackets mean an integration on a plane $\langle g(x,y)\rangle = \int\int_{-\inf}^{\inf} g(x,y)\mathrm{dxdy}$. The physical interpretation of such quantities is how much different modes are superimposed/**??**/overlapped with each others.

Approximate solution

The previous equations are quite general (they include SPM, XPM in addition to FWM) and require a numerical approach to be solved exactly. For this reason, we propose an approximate solution to understand the physical implication. The first approximation is called “undepleted pump” and consist of assuming the pump waves to be much more intense than the other waves. As a further simplification, we assume that all overlap integrals are nearly the same

$$f_{ijkl} \approx f_{ij} \approx \frac{1}{A_{eff}} \qquad i,j,k,l = 1,2,3,4 \tag{17}$$

From this definition follows naturally the nonlinear parameter definition

$$\gamma_j = \frac{n_2 \omega_j}{c A_{eff}} \approx \gamma \tag{18}$$

where γ is the average value (because frequencies have relatively small differencies).

Using the last two equations and the undepleted pump approximation, (14a) and (14b) are easily solved to obtain the pump fields

$$A_1(z) = A_1(0)\exp[\mathrm{i}\gamma(P_1 + 2P_2)z] \tag{19a}$$
$$A_2(z) = A_2(0)\exp[\mathrm{i}\gamma(P_2 + 2P_1)z] \tag{19b}$$

where $P_j = |A_j(z=0)|^2$.

Using the same approximations on (14c) and (14d) and replacing Eqs (19) where needed we obtain two linear coupled equations for signal and idler:

$$\frac{dA_3}{dz} = 2\mathrm{i}\gamma[(P_1 + P_2)A_3 + A_1(0)A_2(0)\mathrm{e}^{-\mathrm{i}\phi}A_4^*] \tag{20a}$$
$$\frac{dA_4}{dz} = 2\mathrm{i}\gamma[(P_1 + P_2)A_4 + A_1(0)A_2(0)\mathrm{e}^{-\mathrm{i}\phi}A_3^*] \tag{20b}$$

where $\phi = [\Delta k - 3\gamma(P_1 + P_2)]z$.

These equations can be solved introducing $B_j = A_j \exp[-2\mathrm{i}\gamma(P_1 + P_2)z]$ for $j = 3,4$ giving their general solution as

$$B_j(z) = (a_j \mathrm{e}^{gz} + b_j \mathrm{e}^{-gz})\exp(-i\kappa z/2)$$

with $\kappa = \Delta k + \gamma(P_1 + P_2)$ *effective phase mismatch* and $g = \sqrt{4\gamma^2 P_1 P_2 - \kappa^2/4}$ the *parametric gain.*

The solutions of Eqs (20a) and (20b) are therefore

$$A_3(z) = \left(a_3 \mathrm{e}^{gz} + b_3 \mathrm{e}^{-gz}\right) \exp\left[-\frac{\mathrm{i}}{2}(\Delta k - 3\gamma(P_1 + P_2))z\right] \tag{21a}$$

$$A_4(z) = \left(a_4 \mathrm{e}^{gz} + b_4 \mathrm{e}^{-gz}\right) \exp\left[-\frac{\mathrm{i}}{2}(\Delta k - 3\gamma(P_1 + P_2))z\right] \tag{21b}$$

****

**** è necessario verificare conservazione di energia e momento

In order to be efficient, the parametric process needs to verify both the conservation of energy and momentum.

$$\omega_1 + \omega_2 = \omega_3 + \omega_4$$

$$\Delta k = \beta_3 + \beta_4 - \beta_1 - \beta_2$$

****

$$\Omega_S = \omega_1 + \omega_3 = \omega_4 - \omega_1$$

where we assume $\omega_3 < \omega_4$.

$$\Omega_S = |\omega_P - \omega_I| = |\omega_S - \omega_P|$$

We can define a length scale, known as the coherence length, using $L_{coh} = 2\pi/|\Delta\kappa|$, where $\Delta\kappa$ is the maximum value of the effective phase mismatch that can be tolerated. Moreover significant FWM occurs if $L_{sam} < L_{coh}$ where L_{sam} is the sample length.

4 Computational models

To understand the behaviour of the waveguide, it is necessary to first know how waves of different frequency are transmitted through the waveguide. We chose to simulate the modes with temperature and frequency as a parameters.

Our aim was to develop a model to work with an infrared laser pump with $\lambda_p = 1.55\,\mu\text{m}$. The reason for this choice is that silicon is transparent at these wavelength and also because of the availability of lasers that emit at the specific wavelength.

We narrowed our simulations to a frequency gap centered in $\omega_p = 2\pi\nu_p = 2\pi c/\lambda_p$. The gap chosen was precisely $[1.45\,\mu\text{m}, 1.65\,\mu\text{m}]$ or the equivalent $[182\,\text{THz}, 207\,\text{THz}]$. The temperature gap, on the other hand, was chosen to be $[0\,\text{K} - 400\,\text{K}]$ over ambient temperature $T_{amb} = 293.15\,\text{K} = 20\,°\text{C}$, that is $[293.15\,\text{K} - 693.15\,\text{K}]$.

4.1 Simulation

To obtain definite solutions, we had to assign the materials and a geometry to the problem first.

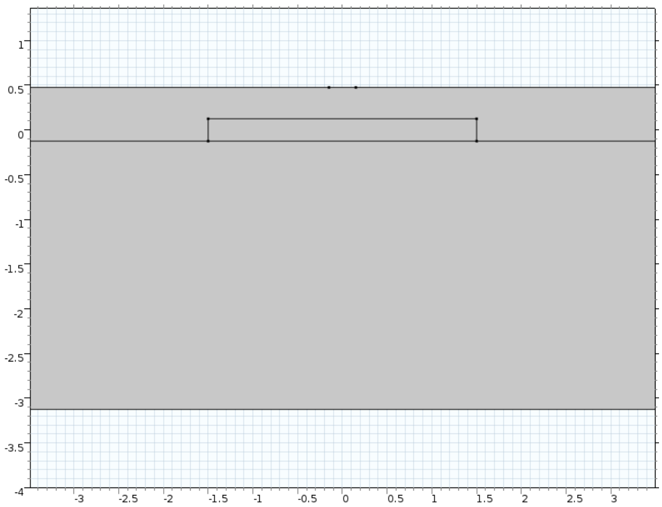

Figure 3: Section of the rectangular waveguide

We modelized a dielectric waveguide with a silicon (Si) core $3\,\mu\text{m}$ wide and $0.25\,\mu\text{m}$ high. The cladding was set to be silica (SiO2) and coated the top of the core with a thickness of 0.35 µm. For both materials was defined a refractive index dependent on the wavelength $n = n(\lambda)$, as given by the empirical relationship called *Sellmeier equation.*

$$n^2(\lambda) = 1 + \sum_i \chi_{0i} \frac{\lambda^2}{\lambda^2 - \lambda_i^2} \tag{22}$$

where the weights are defined experimentally and vary depending on the material.

We positioned a thermal reservoir, a "*heater*", 0.3 µm wide at the top of the cladding, exactly over the center of the waveguide core and set to mantain a specific temperature in the gap, given by the configuration chosen. To define the density ρ in [kgm^3], the thermal conductivity κ_{TH} in [W m^{-1} K^{-1}] and the heat capacity C_P in [J kg^{-1} K^{-1}] of the different materials.

Materials	Properties			
	ρ [kgm^3]	κ_{TH} [W m^{-1} K^{-1}]	C_P [J kg^{-1} K^{-1}]	TOC [K^{-1}]
Silicon (Si)	2329	130	700	1.87×10^{-4}
Silica (SiO_2)	2203	1.38	703	0

Table 1: Thermal properties of silicon and silica glass

Moreover, in order to be influenced by the different temperatures, we had to add to the definition of refractive index of the silicon a term dependent on the temperature. For our purpose this dependence can easily be considered linear:

$$n = n(\lambda) + \mathrm{TOC} \cdot T \tag{23}$$

where $\mathrm{TOC}|_{\mathrm{Si}} = 1.87 \times 10^{-4}\,\mathrm{K}^{-1}$[3] is called Termo-Optic Coefficient. The coefficient of silica glass is negligible and coefficient was therefore considered zero $\mathrm{TOC}|_{\mathrm{SiO_2}} \approx 0$.

4.1.1 Numerical solution

Because the degree of freedom of the problem were too many, we resorted to automated algorithms which gave us a numerical, and therefore approximate, solution.

Heat transfer: stationary solution

The first step was solving the equation for the heat transfer to obtain a stationary solution for each temperature in the gap. That means to know the local temperature for each position in the core and in the cladding, for each initial setup.

EM mode analysis

The second step was to find the different modes of the waveguide for each set of parameters, such as temperature and wavelength.

We emploied an automated algorithm to solve, in the frequency domain, the Helmholtz equation on a fine grid in the waveguide section. Each point of evaluation took into account the effective index of the material, slightly changed by the local temperature obtained by the stationary solution.

This algorithm was set to find the eight modes which effective index n_{eff} was closest to $n_{eff} = 4$ in absolute value.

4.1.2 Data output

We obtained the field distribution and a series of parameters as our data output for each parametric configuration. In particular the parameters were the effective index n_{eff}, the effective area A_{eff} and the nonlinear parameter γ, for each mode and for each wavelength and temperature.

4.2 Data classification

The next step was to classify these data, depending on the type (TE or TM) and order (1,2,3...) of the modes. We achieved this by analyzing a sample taken at $y = 0$ of the field distributions. The ratio between the maxima values of the x-component and y-component of the electric field allowed to distinguish the mode type: TE or TM. Those fields which had the maximum component in the y direction were categorized as TE modes while those which had the maximum component in the x direction as TM modes. The number of maxima of the modulus of the field showed us the order of the mode.

We obtained then a well organized data matrix filled with the value of effective index n_{eff}, effective area A_{eff} and nonlinear parameter γ for each combination of parameters.

4.3 Data processing

4.3.1 Polynomial fit

Starting from raw data consisting of the effective index for each parametric configuration, we fitted them with a polinomial function of frequency $\omega = 2\pi\nu$ or λ and temperature T, obtaining:

$$n_{eff} = f_{mt,\,mo}\,(\omega, T) = f'_{mt,\,mo}\,(\lambda, T)$$

where *mt* and *mo* means respectively *mode type* (TE or TM) and *mode order* (1, 2, 3, ...)

Specifically, we used a polynomial expression of fifth order in wavelength and second order in temperature "`poly52`".

4.3.2 Order combination

Due to the intrinsic symmetry of the problem (and thus of the integrals) not all combinations of orders can produce FWM. Only those in which the sum of the orders is an even number are interesting. To narrow the possible sets, we chose only full TE or full TM combinations.

We eventually obtained a list of combinations and for each one we evaluated the phase mismatch.

4.3.3 Conservation of energy and momentum

As it is, the phase mismatch is a function of the wavelenghts of pump, signal and idler. Imposing the conservation of energy,

$$2\omega_P = \omega_S + \omega_I \qquad \Longrightarrow \qquad \omega_I = 2\omega_P - \omega_S \tag{24}$$

we obtained the only frequency permitted for ω_I.

We replaced its value in the conservation of momentum formula, giving us the phase mismatch as a function only of signal wavelength and temperature of the heater, for each set of orders.

Then, starting from the fits of the effective index at different wavelengths and temperatures, we evaluated the phase mismatch of all sets. The evaluation was made on a grid of temperature and signal wavelength arbitrarily dense.

4.3.4 Selection of sets

Once evaluated all the points on the grid for all the sets available, we selected only the interesting combinations. The critical parameter on which the choice was based on is the coherence length: only those sets which had a maximum value of coherence length greater than a supposed sample length L_{sam} =1 cm was chosen.

$$L_{coh} = \frac{2\pi}{|\Delta k|} \geq L_{sam}$$

Again, we had a list of the only combination interesting for our purpose.

It is important to note that, in general, each combination could have more that one maximum of coherence length L_{coh} in wavelength for each temperature value.

4.3.5 Data aggregation

After the selection of sets, we evaluated the dependence of each maximum on temperature.

4.4 Second heater configuration

In order to achieve a greater effect from the different temperature, we tried to move the heater from over the center of the core to the side. This operation was intended to increase the difference of behaviour between low and high orders enabling thus the creation of new peaks in the coherence length.

4.5 Different geometries

5 Results

6 Conclusions

6.1 Improvements

6.2 Summary

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