

Time-Ordering Invariance Under Monotone Clock Changes

Addendum to LTQG: $\sigma = \log(\tau/\tau_0)$

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Context

Let $\tau \in I \subset \mathbb{R}$ with I an interval and let $\sigma = F(\tau)$ be a *monotone* C^1 change of clock with $F' : I \rightarrow (0, \infty)$. Consider a (possibly time-dependent) Schrödinger generator $H(\tau)$ on a Hilbert space \mathcal{H} . Write $U_\tau(t_2, t_1)$ for the τ -time propagator solving

$$i \partial_\tau U_\tau(\tau, \tau_0) = H(\tau) U_\tau(\tau, \tau_0), \quad U_\tau(\tau_0, \tau_0) = \mathbf{1}, \quad (1)$$

and define the σ -time generator by the usual chain rule

$$\tilde{H}(\sigma) := \frac{d\tau}{d\sigma} H(\tau(\sigma)) = \frac{H(\tau(\sigma))}{F'(\tau(\sigma))}. \quad (2)$$

Denote by $U_\sigma(\sigma_2, \sigma_1)$ the corresponding σ -time propagator.

Main Lemma (Time-Ordering Invariance)

Lemma 1 (Chronological invariance under monotone clocks). *If F is strictly increasing, then for any $\tau_1 \leq \tau_2$ we have*

$$\mathcal{T}_\tau \left[\exp \left(-i \int_{\tau_1}^{\tau_2} H(\tau) d\tau \right) \right] = \mathcal{T}_\sigma \left[\exp \left(-i \int_{\sigma_1}^{\sigma_2} \tilde{H}(\sigma) d\sigma \right) \right], \quad (3)$$

with $\sigma_j = F(\tau_j)$. In particular, $U_\tau(\tau_2, \tau_1) = U_\sigma(\sigma_2, \sigma_1)$.

Proof. Strict monotonicity implies that the map $\tau \mapsto \sigma = F(\tau)$ preserves order: $\tau_a < \tau_b \iff \sigma_a < \sigma_b$. Hence the chronological partitions used to define Dyson expansions can be transported bijectively. Start from the Dyson series in τ :

$$U_\tau(\tau_2, \tau_1) = \mathbf{1} + \sum_{n \geq 1} (-i)^n \int_{\tau_1 \leq t_n \leq \dots \leq t_1 \leq \tau_2} H(t_1) \cdots H(t_n) dt_1 \cdots dt_n. \quad (4)$$

Change variables $s_k = F(t_k)$; then $dt_k = \frac{dt_k}{ds_k} ds_k = \frac{1}{F'(t_k)} ds_k$ and the ordered simplex is mapped to $\sigma_1 \leq s_n \leq \dots \leq s_1 \leq \sigma_2$ by monotonicity. Each factor transforms as $H(t_k) dt_k = \tilde{H}(s_k) ds_k$ by definition of \tilde{H} . Therefore every n -simplex integral equals the corresponding σ -time Dyson term, yielding the identity. \square

Existence & Uniqueness Hypotheses

Assumption 1 (Bounded case). $H(\tau)$ is strongly measurable and $\sup_{\tau \in I} \|H(\tau)\| < \infty$. Then the Dyson series converges in operator norm and defines a unique unitary propagator.

Assumption 2 (Kato Class (unbounded case)). $H(\tau)$ is self-adjoint on a dense domain D (independent of τ), $\tau \mapsto H(\tau)\psi$ is continuous $\forall \psi \in D$, and the family is stable in the sense of Kato (domain invariance and suitable bounds). Under these hypotheses the evolution family U_τ exists and is unique; the reparameterization $\tilde{H}(\sigma) = \frac{d\tau}{d\sigma} H(\tau(\sigma))$ preserves the same class, so U_σ exists and equals U_τ by Lemma 1.

Theorem 1 (Unitary equivalence under monotone reparameterization). *Under either set of hypotheses above, $U_\tau(\tau_2, \tau_1) = U_\sigma(\sigma_2, \sigma_1)$ for $\sigma_j = F(\tau_j)$. Consequently, spectra of Heisenberg-evolved observables and transition probabilities coincide for τ - and σ -descriptions.*

Remark 1 (Application to LTQG). For $F(\tau) = \log(\tau/\tau_0)$ with $\tau > 0$ we have $\frac{d\tau}{d\sigma} = \tau$ and hence $\tilde{H}(\sigma) = \tau(\sigma) H(\tau(\sigma))$. The lemma shows that the reparameterization leaves all physical predictions invariant while making multiplicative early-time structure additive in σ .