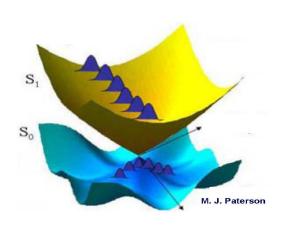
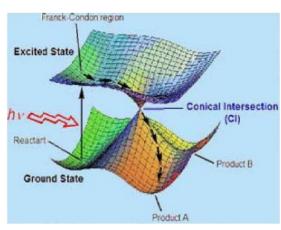
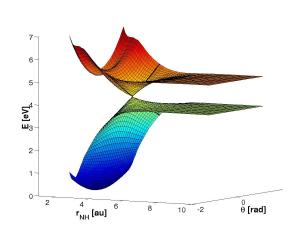
A brief introduction to JADE package

L. Du, Z. Lan*

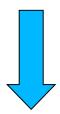
Nonadiabatic dynamics







Born-Oppenheimer approximation breaks down!!



A self-consistent treatment of nuclear and electronic degrees of freedom is needed.

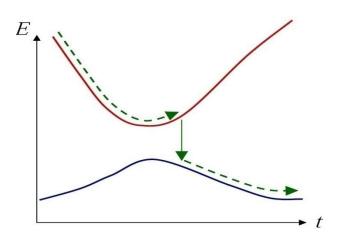
Surface hopping method

Methodology

A swarm of independent trajectories is considered.

Each trajectory evolves on a single potential energy surface.

Sudden hops between different potential energy surfaces are allowed.



Expansion of the electronic wavefunction

$$\Phi(\mathbf{r}, \mathbf{R}, t) = \sum_{i} c_i(t) \phi_i(\mathbf{r}, \mathbf{R})$$

The electronic wavefunction is expanded as a superposition of independent adiabatic states.

Quantum amplitudes determine the time-dependent population of state *i* at time *t*

Substitution into the timedependent Schrödinger equation yields a set of coupled equations for the quantum amplitudes.

$$i\hbar \frac{dc_j(t)}{dt} = c_j(t)\epsilon_j - i\hbar \sum_i c_i(t)\dot{\mathbf{R}} \cdot \mathbf{d}_{ji}$$

J. C. Tully, *J. Chem. Phys.* **93**, 1061 (1990).

nonadiabatic coupling

Tully's Surface Hopping Approach

$$i\hbar \frac{\partial \Phi(\mathbf{r}, \mathbf{R}, t)}{\partial t} = H_e \Phi(\mathbf{r}, \mathbf{R}, t).$$

$$\Phi(\mathbf{r}, \mathbf{R}, t) = \sum_{i} c_i(t)\phi_i(\mathbf{r}, \mathbf{R}),$$

$$i\hbar \frac{dc_j(t)}{dt} = \sum_i c_i(t) [H_{ji} - i\hbar \dot{\mathbf{R}} \cdot \mathbf{d}_{ji}],$$

transition probability

$$P_{ij} = -\frac{2\int_t^{t+\Delta t} dt \left[\hbar^{-1} \operatorname{Im}(c_i^* c_j \epsilon_i \delta_{ji}) - \operatorname{Re}(c_i^* c_j \dot{\mathbf{R}} \cdot \mathbf{d}_{ji})\right]}{|c_i(t)|^2}.$$

$$H_{ji} \equiv \int d\mathbf{r} \phi_j^*(\mathbf{r}, \mathbf{R}) \left[-\frac{\hbar}{2} \sum_l \frac{1}{m_l} \nabla_{\mathbf{r}_l}^2 + V_{rR}(\mathbf{r}, \mathbf{R}) \right] \phi_i(\mathbf{r}, \mathbf{R}),$$

$$\mathbf{d}_{ji} \equiv \int d\mathbf{r} \phi_j^*(\mathbf{r}, \mathbf{R}) [\nabla_{\mathbf{R}} \phi_i(\mathbf{r}, \mathbf{R})].$$

$$\frac{\partial}{\partial t} = \frac{\partial}{\partial \mathbf{R}} \frac{d\mathbf{R}}{dt}$$

$$F_{ij}(t) = \int d\mathbf{r} \phi_j^*(\mathbf{r}, \mathbf{R}) \frac{\partial \phi_i(\mathbf{r}, \mathbf{R})}{\partial t}$$

$$H_{ji} = \epsilon_i \delta_{ji}$$
.

$$i\hbar \frac{dc_j(t)}{dt} = c_j(t)\epsilon_j - i\hbar \sum_i c_i(t)\dot{\mathbf{R}} \cdot \mathbf{d}_{ji}.$$

$$P_{ij} = \frac{2\int_t^{t+\Delta t} dt \operatorname{Re}(c_i^* c_j \dot{\mathbf{R}} \cdot \mathbf{d}_{ji})}{|c_i(t)|^2}.$$

$$\rho_{kl} \equiv c_k c_l^*$$

$$\frac{d\rho_{ll}}{dt} = i\hbar^{-1} \left[\left(-H_{kl}^c + i\hbar \mathbf{F}_{kl}^c \cdot \mathbf{v}^c \right) \rho_{lk} - \left(-H_{lk}^c + i\hbar \mathbf{F}_{lk}^c \cdot \mathbf{v}^c \right) \rho_{kl} \right]$$
$$= -2\hbar^{-1} H_{kl} \operatorname{Im}(\rho_{kl}) - 2 \mathbf{F}_{kl}^c \cdot \mathbf{v}^c \operatorname{Re}(\rho_{kl}).$$

$$P_{l\to k} = \max\left[0, \frac{-2\Delta t}{\rho_{ll}} \operatorname{Re}(\rho_{kl}) \mathbf{F}_{kl}^c \cdot \mathbf{v}^c\right] \text{(adiabatic)},$$

$$P_{l \to k} = \max \left[0, \frac{2\Delta t}{\hbar \rho_{ll}} \operatorname{Im}(\rho_{kl}) W_{lk}^{c} \right] (diabatic)$$

Nonadiabatic dynamics at TDDFT level

Why?

The only pratical method for large systems.

- Efficiency and accuracy (?)
- Available package
- Map to the CIS form easily
- Potential Possiblilty on GPU computation

Challenging?

- Nonadiabatic couplings
 Numerical v.s. analytical ones
- Charge-transfer problems
- Single-reference problems: S0/S1 Crossing
- Other problems: double excitations...

TD-DFT

- how to get non-adiabatic coupling vector
- numerical

$$F_{ij}\left(t + \frac{\Delta t}{2}\right) = \frac{\int d\mathbf{r}\phi_j^*(t)\phi_i(t + \Delta t) - \int d\mathbf{r}\phi_j^*(t + \Delta t)\phi_i(t)}{2\Delta t}.$$

$$\left|\Psi_{i}\right\rangle = \sum_{k} c_{k}^{i} \left|\Phi_{i}^{CSF}\right\rangle \qquad \begin{pmatrix} \mathbf{A} & \mathbf{B} \\ \mathbf{B} & \mathbf{A} \end{pmatrix} \begin{pmatrix} \mathbf{X} \\ \mathbf{Y} \end{pmatrix} = \boldsymbol{\omega} \begin{pmatrix} 1 & \mathbf{0} \\ \mathbf{0} & -1 \end{pmatrix} \begin{pmatrix} \mathbf{X} \\ \mathbf{Y} \end{pmatrix}$$

- J. Pittner et al. Chem. Phys. 356 (2009) 147–152
- I. Tavernelli et al. J. Mol. Struct. Theochem 914 (2009) 22-29

Numerical nonadiabatic couplings

AO Overlap:
$$<\mu_p^{AO}(t)|\mu_q^{AO}(t+\delta t)>$$

MO Overlap:
$$<\kappa_m^{MO}(t)|\kappa_n^{MO}(t+\delta t)>$$

CSF Overlap (Overlap of Slater determinant)

$$<\psi_k^{CSF}(t)|\psi_l^{CSF}(t+\delta t)>$$

Overlap of Electronic Wavefunctions

$$<\phi_j(t)|\phi_i(t+\delta t)>$$

Nonadiabatic couplings

$$F_{ji} = \frac{1}{\delta t} < \phi_j(t) | \phi_i(t + \delta t) > 0$$

Algorithms

- 1. Initialization of velocity, gradients, and quantum amplitudes
- 2. Time propagation of coordinates and velocities on the selected PES.
- 3. Computation of energies, gradients, and nonadiabatic coupling vectors of all relevant states at new position and velocity.
- 4. Time propagation of quantum amplitudes and computation of hopping probabilities.
- 5. Random number generation and comparison with hopping probabilities
- 6. If hopping is rejected: inversion of velocity.
- 7. If hopping is performed: velocity adjustment and update of the active PES for molecular dynamics.
- 8. Back to point 2.

Introduction to the codes

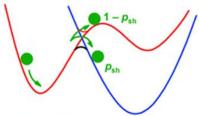
Excited State Method

CASSCF, ADC(2), TDDFT,CIS,TDDFTB

$$\hat{H}_{e} \chi(\mathbf{r}, \mathbf{R}) = E_{e}(\mathbf{R}) \chi(\mathbf{r}, \mathbf{R})$$

$$HC = SCE$$

$$\begin{pmatrix} \mathbf{A} & \mathbf{B} \\ \mathbf{B} & \mathbf{A} \end{pmatrix} \begin{pmatrix} \mathbf{X} \\ \mathbf{Y} \end{pmatrix} = \boldsymbol{\omega} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} \mathbf{X} \\ \mathbf{Y} \end{pmatrix}$$



Surface Hopping Dynamics

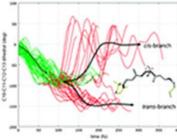
$$i\hbar \frac{dc_{j}(t)}{dt} = \sum_{i} c_{i}(t) [H_{ji} - i\hbar \mathbf{R} \cdot \mathbf{d}_{ji}]$$
$$\mathbf{d}_{ji} \equiv \int d\mathbf{r} \phi_{j}^{*}(\mathbf{r}, \mathbf{R}) [\nabla_{\mathbf{R}} \phi_{i}(\mathbf{r}, \mathbf{R})]$$

$$P_{ij} = \frac{2 \int_t^{t+\Delta t} dt \operatorname{Re}(c_i^* c_j \dot{\mathbf{R}} \cdot \mathbf{d}_{ji})}{|c_i(t)|^2}$$





Code development



Analysis

code management

- version management: git
- code management: Makefile

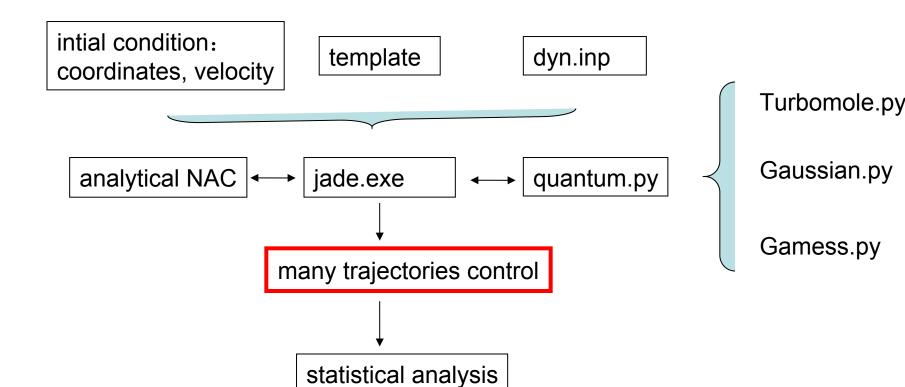
Code development and feature

- On-the-fly dynamics surface hopping dynamics
- the velocity-Verlet algorithm;
- fewest switches algorithms
- Langevin thermostat
- Interfaces
- TURBOMOLE (CIS, TD-DFT, RI-CC2, ADC(2))
- GAUSSIAN (CIS, TD-DFT)
- GAMESS (CIS TD-DFT)
- Q-CHEM (in progress)
- Molpro (in progress)
- Initial conditions
- Wigner distribution
- enviroment effects
- Langevin dynamics
- effective fragment potential (interface with Gamess-US & Q-CHEM) in process
- Output and analysis
- Statistical analysis of results
- transition density analysis tools

how-to: install

- See INSTALL file in doc/ directory
- To compile the Fortran90 programs of the JADE package, go to the src/ directory.
- go to src directory and run make; make install;
- if the compilation is successful
- then, set the environment variable \$JADE_HOME
- export JADE_HOME=/home/jade-package
- source \$JADE_HOME/bin/JADERC

code



how-to: initial conditions

wigner sampling

```
$SH_HOME/wigner
```

- 1. frequency calcuation: gaussian/turbomole/mndo prepare.py
- 2. sampling process sampling.py
- 3. filter a list of sample geometries mdfilter.py

The transition probability is calculated as P = (f/DE^2) / max(f/DE^2)

how-to: wigner sampling

```
[wigner]
n_atom = 6
n_{mode} = 12
label_random = 1
nr = 5000
nbin = 50 ! used to check random
label_read_vib = 1
label es output = 3
filename_es_output = freq.log
label displacement = 1
label dis wigner = 1
n_geom = 1000
label frozen = 0
number frozen = 2
list_frozen = 1, 2
```

how-to: template

- td-dft
- support: gaussian/turbomole/gamess
- \$SH_HOME/test

how-to: dyn.inp

- FORTRAN-NAMELIST-STYLE
- up to now, include three sections:
- &control
- &quantum
- &langevin
- •

how-to: dyn.inp

```
&control
dyn_method = 201,
ntime = 400,
dtime = 0.5,
ntime_ele = 100,
n sav stat = 1,
n sav_traj = 1,
qm_method = 11,
n \text{ state} = 3,
md state list = "1, 2, 3",
i state = 3,
seed random = -1,
cor_dec = 0.1,
label nac phase = 1,
label_reject_hops = 1,
hop e = 10,
label read velocity = 0,
label restart = 0,
```

```
&langevin
gamma0 = 0.1
temperature = 300.0
/
```

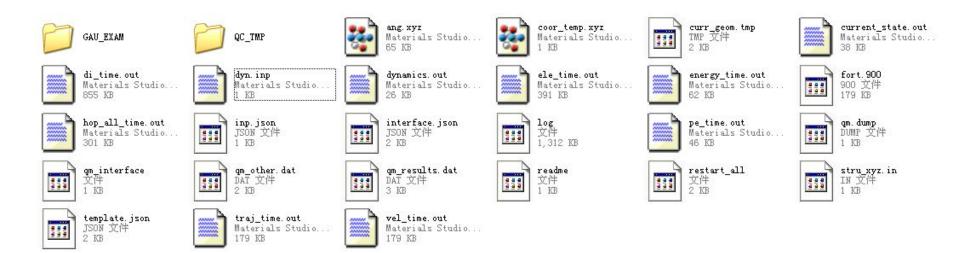
```
&quantum
qm_method = 11,
qm_package = 102,
ci_td_use_file_type = "log",
ci_assign_problem = "X+Y",
is_do_cis_casida = "yes",
/
```

restartable

- restartable
- set label_restart to 1

directory in dynamics

jade.exe > abc.log &



how-to: many-trajectory control

- write your own scripts....
- manyjob.py

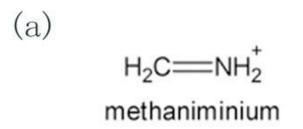
how-to: statistical analysis

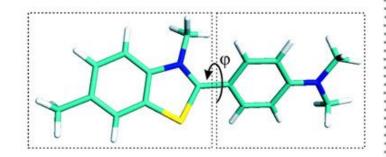
- see \$SH_HOME/contrib
- jobstatus.py
- jobfilter.py

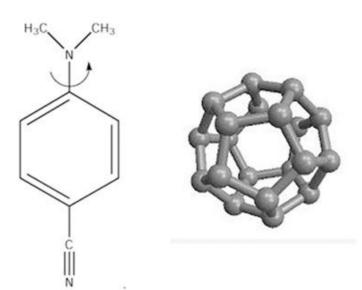
- structural analysis: geom_time.py
- energy analysis: pe_time.py
- population analysis: state_population.py
- spectrum analysis ...

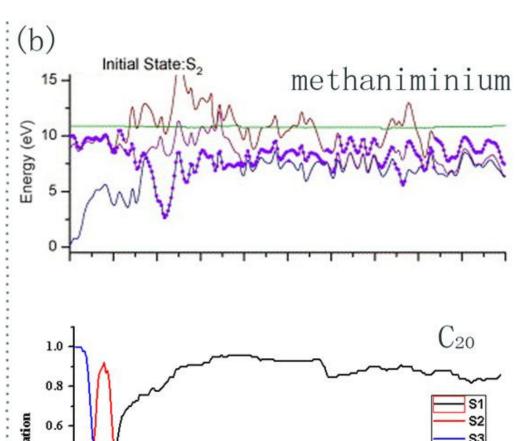
examples

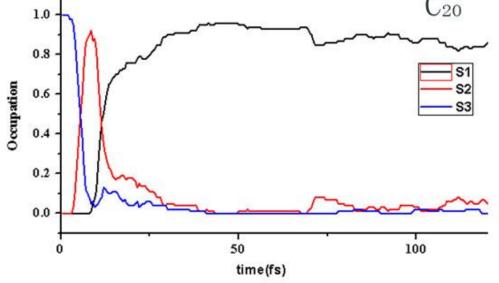
- simple examples
- \$SH_HOME/test



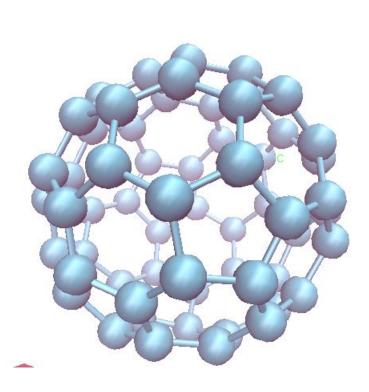




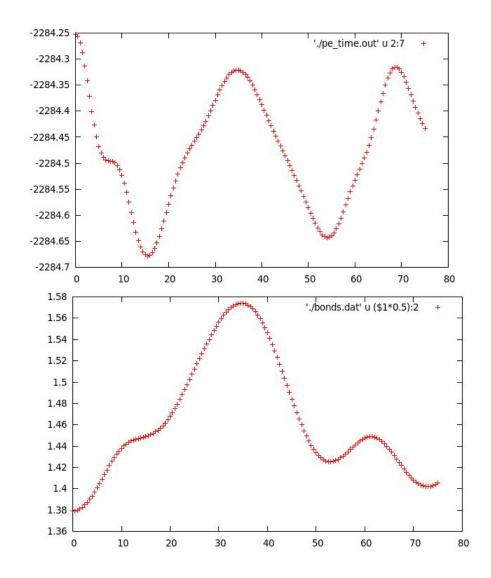


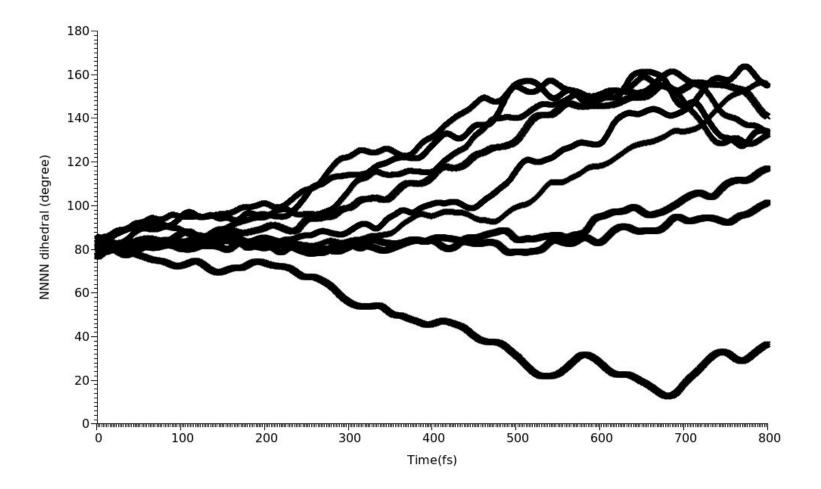


C60



turbomole; tddft/pbe; n_state=6; i_state=6





on-going

- analysis tools
- Q-CHEM interface..
- MCSCF with analytical NAC

- DFTB
- QM/MM
- open shell

Thanks