# NLS and Sharp Interpolation Estimates Michael Weinstein, Commun. Math. Phys. 87, 567-576 (1983)

Sebastian Gherghe

April 3 2019

#### Introduction: Prelude to the NLS

Finding the "best constant" of various inequalities usually has some geometric or analytic significance. As an example:

## Theorem (Isoperimetric Inequality)

Consider a domain  $D \in \mathbb{R}^n$  with boundary  $\partial D$ . Then,

$$V^{n-1} \leqslant n^{-n} \omega_n^{-1} A^n$$

- V is the volume (Lebesgue Measure) of D
- $\omega_n$  is the volume of the n-ball of radius 1
- A is the "surface area" of D, generalized by by  $M_{n-1}(\partial D)$ , the (n-1)-dimensional Minkowski Content of  $\partial D$ .

Equality is obtained only when D is the n-ball!

#### Introduction: Prelude to the NLS

How does this geometric fact concern inequalities? Consider the following Sobolev-type Embedding theorem:

#### **Theorem**

Consider  $u \in C_0^{\infty}(\mathbb{R}^N)$ . Then

$$||u||_{L^{\frac{n}{n-1}}}\leqslant C(n)||\nabla u||_{L^1}$$

and the best constant is the same as in the Isoperimetric Inequality, ie.  $C(n)_{sharp} = n^{-n}\omega_n^{-1}$ .

#### Introduction

Weinstein's 1983 paper presents a relationship between the sharp constant for a classical interpolation inequality by Nirenberg and Gagliardo and a criterion for global existence for the Nonlinear Schrödinger Equation:

$$i\frac{\partial\phi}{\partial t} + \Delta\phi + |\phi|^{2\sigma}\phi = 0 \tag{NLS}$$

where  $x \in \mathbb{R}^N$ ,  $t \in \mathbb{R}^+$  and initial conditions  $\phi(x,0) = \phi_0(x)$  in the  $L^2$ -critical case  $N\sigma = 2$ .

#### Introduction

#### This presentation will tackle three things:

- I. Sharpest Constant for the (classical) Gagliardo-Nirenberg Interpolation Inequality and the ground state solution of the NLS: Solution of a Variational Problem
- II. A sufficient condition for global existence of  $H^1$  solutions of the NLS: Global Existence for the IVP in the  $L^2$  critical case
- III. (time-permitting) comments on stability and finite blow-up

Consider a particular Gagliardo-Nirenberg Interpolation Inequality:

$$||f||_{2\sigma+2}^{2\sigma+2} \le C_{\sigma,N}^{2\sigma+2} ||\nabla f||_2^{\sigma N} ||f||_2^{2+\sigma(2-N)}$$

for  $0<\sigma<\frac{2}{N-2}$  and  $N\geqslant 2$ . This is a particular case of the general inequality:

$$||D^{j}f||_{p} \leqslant C||D^{m}f||_{r}^{\alpha}||f||_{q}^{1-\alpha}$$

for some function  $f: \mathbb{R}^N \to \mathbb{R}$  with  $p=2\sigma+2, \ m=1, \ q=r=2$ , which fixes  $\alpha=n(\frac{1}{2}-\frac{1}{2\sigma+2})$ 

In order to find the sharpest constant  $C_{\sigma,N}$ , we define the functional:

$$J^{\sigma,N}(f) = \frac{\|\nabla f\|_2^{\sigma N} \|f\|_2^{2+\sigma(2-N)}}{\|f\|_{2\sigma+2}^{2\sigma+2}}$$

Now this turns into a minimization problem!

We will show that the minimum is attained at some  $H^1$  function  $\psi^*$ . We obtain explicit forms for  $C_{\sigma,N}$  and the equation it solves by applying the Euler-Lagrange equations to  $J^{\sigma,N}$ .

Note that by the original estimate  $J^{\sigma,N}$  is defined on  $H^1(\mathbb{R}^N)$  for  $0<\sigma<\frac{2}{N-2}$ .

#### **Theorem**

For 
$$0 < \sigma < \frac{2}{N-2}$$
,  $\alpha = \frac{1}{N-2}$ 

$$\alpha \equiv \inf_{u \in H^1(\mathbb{R}^N)} J^{\sigma,N}(u)$$

is attained at a function  $\psi$  with the following properties

- 1)  $\psi$  is positive and a function of |x| alone
- 2)  $\psi \in H^1(\mathbb{R}^N) \cap C^\infty(\mathbb{R}^N)$
- 3)  $\psi$  is a solution of:

$$\frac{\sigma N}{2}\Delta\psi - (1 + \frac{\sigma}{2}(2 - N))\psi + \psi^{2\sigma + 1} = 0$$

of minimal L<sup>2</sup> norm (the ground state), and in addition

$$\alpha = \frac{\|\psi\|_2^{2\sigma}}{\sigma + 1}$$

#### Ideas behind the proof:

- (1) show that positive minimizers satisfy the same minimization problem, and a "symmetrization" argument to show the same for radial minimizers (Strauss 1977)
- (2) Compactness lemma gives compact embedding of radial functions, obtain a convergent minimizing sequence!
- (3) Apply the Euler-Lagrange equations to  $J^{\sigma,N}$

#### Lemma (Compactness Lemma)

For  $0 < \sigma < \frac{2}{N-2}$ , the embedding

$$H^1_{radial}(\mathbb{R}^N) o L^{2\sigma+2}(\mathbb{R}^N)$$

is compact.

#### Proof.

Proof follows from the interpolation estimate:

$$||u||_{2\sigma+2}^{2\sigma+2} \leqslant C||u||_{H^1}^{\sigma N}||u||_2^{2+\sigma(2-N)}$$

for  $0 < \sigma < \frac{2}{N-2}$  and  $u \in H^1(\Omega)$  where  $\Omega$  is a bounded domain.

We extend to  $u \in H(\mathbb{R}^N)$  if we can show that a bounded sequence in  $H(\mathbb{R}^N)_{radial}$  is uniformly small at infinity. This follows by Strauss' Radial Lemma.

#### Lemma (Radial Lemma 1)

Let  $N \ge 2$ . Every radial function  $u \in H^1$  is almost everywhere equal to a function U(x), continuous for  $x \ne 0$ , such that

$$|U(x)| \leqslant c|x|^{\frac{1-N}{2}} ||u||_{H^1}$$

for  $|x| \geqslant 1$ , and where c depends only on n.

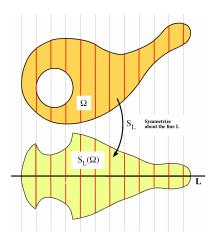
#### Sketch of 1).

Argue that if the minimization problem has a solution, then it also has a solution that is non-negative and radial.

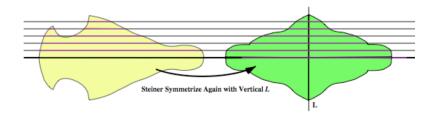
Let u be a solution. Argue that  $u^+(x) := \max(u(x), 0)$  solves the same minimization problem (observe that  $\nabla(u^+) = (\nabla u)^+$  and that  $u^+ \in H^1$ ).

To assume that u is radial we will argue using Steiner Symmetrization, imitating Strauss 1977 who in turn extended the classical version of this technique (Polya and Szegö 1951).

"Symmetrize" a given bounded domain with "nice" boundary about a hyperplane that passes through the origin.



Applying the symmetrization multiple times...



## Theorem (Ljusternik-Gross Sphericalization Theorem)

Let  $\Omega$  be a non-empty compact set and  $\mathbb G$  the family of all multiple symmetrizations (finite composition of symmetrizations) of  $\Omega$ . Then there is a subsequence  $\Omega_n \subset \mathbb G$  and a closed ball  $\overline B$  with the same volume as  $\Omega$  and

$$\Omega_n o \overline{B}$$
 as  $n o \infty$ 

Define

$$D = \{(x, t) \in \mathbb{R}^{N+1} : 0 \leqslant t \leqslant u(x)\}$$

Let  $D^*$  be the symmetrization of D around the hyperplane  $x_1 = 0$ . Then,  $D^*$  is of the form

$$D = \{(x,t) \in \mathbb{R}^{N+1} : 0 \leqslant t \leqslant u^*(x)\}$$

and we argue that  $u^*$  is also a solution of the same problem. So by successive choice of hyperplanes, we get a (non-negative) radial solution.

## Sketch of 2).

Since  $J^{\sigma,N}(u) \geqslant 0$  there exists a minimizing sequence  $u_{\nu} \in H^1(\mathbb{R}^N) \cap L^{2\sigma+2}(\mathbb{R}^N)$ . We assume this  $u_{\nu}$  to be positive and radial.

Define the scaling

$$u^{\lambda,\mu}(x) \equiv \mu u(\lambda x)$$

and fix values of  $\lambda$  and  $\mu$  as follows.

$$\lambda_{\nu} = \frac{\|u_{\nu}\|_2}{\|\nabla u_2\|_2}$$

and

$$\mu_{\nu} == \frac{\|u_{\nu}\|_{2}^{\frac{N}{2}-1}}{\|\nabla u_{2}\|_{2}^{\frac{N}{2}}}$$

Then the sequence  $\psi_{\nu}(x)=u^{\lambda_{\nu},\mu_{\nu}}(x)$  with the following properties:

- (a)  $\psi_{\nu} \geqslant 0$ ,  $\psi_{\nu} = \psi_{\nu}(|x|)$
- (b)  $\psi_{\nu} \in H^1(\mathbb{R}^N)$
- (c)  $\|\psi_{\nu}\|_{2} = 1$  and  $\|\nabla\psi_{\nu}\|_{2} = 1$
- (d)  $J^{\sigma,N}(\psi_{\nu})\downarrow \alpha$  as  $\nu\to\infty$

Since the sequence  $\psi_{\nu}$  is bounded in  $H^1(\mathbb{R}^N)$  some subsequence has a weak  $H^1$  limit  $\psi^*$  (Recall the Banach-Alaoglu Theorem). We can use the Compactness Lemma to take  $\psi_{\nu}$  strongly convergent to  $\psi^*$  in  $L^{2\sigma+2}(\mathbb{R}^N)$  for  $0<\sigma<\frac{2}{N-2}$ .

By weak convergence we have  $\|\psi^*\|_2 \leqslant 1$  and  $\|\nabla \psi^*\|_2 \leqslant 1$ . So using the definition of the operator  $J^{\sigma,N}$ :

$$\alpha \leqslant J^{\sigma,N}(\psi^*)$$

$$\leqslant \frac{1}{\int |\psi^*|^{2\sigma+2} dx}$$

$$= \lim_{\nu \uparrow \infty} J(\psi_n u) = \alpha$$

It follows that

$$\|\psi^*\|_2^{\sigma N} \|\psi^*\|_2^{2+\sigma(2-N)} = 1$$

so then  $\|\psi^*\|=\|\nabla\psi^*\|_2=1$  and we have strong convergence of  $\psi_{\nu}\to\psi^*$  in  $H^1$ .

## Sketch of 3).

Applying the Euler-Lagrange equation to  $J^{\sigma,N}$ :

$$\left. \frac{d}{d\epsilon} \right|_{\epsilon=0} J^{\sigma,N}(\psi^* + \epsilon \eta) = 0$$

for all  $\eta \in C_0^\infty(\mathbb{R}^N)$ . Using  $\|\psi^*\| = \|\nabla \psi^*\|_2 = 1$  we obtain from the E-L equation:

$$\frac{\sigma N}{2} \Delta \psi^* - (1 + \frac{\sigma}{2} (2 - N)) \psi^* + \alpha (\sigma + 1) (\psi^*)^{2\sigma + 1} = 0$$

Rescaling we let  $\psi^*=[\alpha(\sigma+1)]^{-\frac12\sigma}\psi$  and we obtain a solution the equation in Theorem B with

$$\alpha = \frac{\|\psi\|_2^{2\sigma}}{(\sigma+1)}$$



#### **Theorem**

For 
$$0 < \sigma < \frac{2}{N-2}$$
,

$$\alpha \equiv \inf_{u \in H^1(\mathbb{R}^N)} J^{\sigma,N}(u)$$

is attained at a function  $\psi$  with the following properties

- 1)  $\psi$  is positive and a function of |x| alone
- 2)  $\psi \in H^1(\mathbb{R}^N) \cap C^\infty(\mathbb{R}^N)$
- 3)  $\psi$  is a solution of:

$$\frac{\sigma N}{2}\Delta\psi - (1 + \frac{\sigma}{2}(2 - N))\psi + \psi^{2\sigma + 1} = 0$$

of minimal L<sup>2</sup> norm (the ground state), and in addition

$$\alpha = \frac{\|\psi\|_2^{2\sigma}}{\sigma + 1}$$

Also obtain two corollaries:

## Corollary (Corollary B.1)

The best (smallest) constant for which the classical interpolation estimate holds is given by

$$C_{\sigma,N} = \left(\frac{\sigma+1}{\|\psi\|_2^{2\sigma}}\right)^{\frac{1}{2\sigma+2}}$$

where  $\psi$  is the ground state of the NLS.

#### Corollary (Corollary B.2)

Let  $0 < \sigma < \frac{2}{N-2}$ . Then, the equation

$$\Delta u - u + u^{2\sigma + 1} = 0$$

has a positive, radial solution of class  $H^1(\mathbb{R}^N)$ .

Building on the results of Ginibre and Velo for global existence:

## Theorem (Theorem 3.1)

Let  $\phi_0 \in H^1(\mathbb{R}^N)$ . Then:

(i) If  $0<\sigma<\frac{2}{N}$ , then there exists a unique solution  $\phi\in C([0,\infty];H^1(\mathbb{R}^N))$  of the IVP (NLS) in the sense of the equivalent integral equation:

$$\phi = U(t - t_0)\phi - i \int_{t_0}^t U(t - s)(|\phi|^{2\sigma}\phi)(s)d2s$$

where U(t) is the Schrödinger Kernel.

## Theorem (Theorem 3.1)

- (ii) If  $\sigma = \frac{2}{N}$ , then for  $\|\phi_0\|_2$  sufficiently small, the conclusion of (i) holds.
- (iii) As long as  $\phi(x,t)$  remains in  $H^1(\mathbb{R}^N)$ , the quantities

$$\mathcal{N}(\phi) \equiv \int |\phi(x,t)|^2 dx$$

and

$$\mathcal{H}(\phi) \equiv \int (|
abla \phi(x,t)|^2 - rac{1}{\sigma+1}|\phi(x,t)|^{2\sigma+2})$$

are constants in time.

Remark If  $\sigma \geqslant \frac{2}{N}$ , solutions may develop singularities in finite time.



## Theorem (Theorem A)

Let  $\phi_0 \in H^1(\mathbb{R}^N)$ . For  $\sigma = \frac{2}{N}$ , a sufficient condition for global existence in the IVP is:

$$\|\phi_0\|_2 \le \|\psi\|_2$$

where  $\psi$  is a positive solution of the equation

$$\Delta u - u + u^{\frac{4}{N}+1} = 0$$

of minimal  $L^2$  norm (the ground state), and  $\psi e^{\frac{it}{2}}$  is an exact solution of the IVP.

The idea behind the proof of Theorem A:

In the theorem of Ginibre and Velo  $(0 < \sigma < \frac{2}{N-2})$  they show that the length T of the interval of existence  $[t_0, t_0 + T]$  can be taken to depend only on  $\|\phi(t_0)\|_{H^1}$ .

Then if  $\phi(x,t)$  is a maximally defined solution on  $[t_0,t_{max}]$  we have two possibilities:

(i) 
$$t_{max} = +\infty$$

(ii) 
$$\lim_{t\uparrow t_{max}} \lVert \phi(t) \rVert_{H^1} = +\infty$$

In Ginibre and Velo's proof they use the invariants  $\mathcal{N}(\phi)$  and  $\mathcal{H}(\phi)$  to obtain an *a priori* bound of the type:

$$\|\phi(t)\|_{H^1(\mathbb{R}^N)} \leqslant C(\mathcal{N}, \mathcal{H})$$

We seek to imitate their proof by showing a particular version of the bound.

Sketch of Proof of Theorem A: Using the constants of motion and the interpolation estimate:

$$\|\nabla\phi(t)\|_{2}^{2} \leqslant \mathcal{H} + \frac{C_{\sigma,N}^{2\sigma+2}}{\sigma+1} \|\phi_{0}\|_{2}^{2+\sigma(2-N)} \|\nabla\phi(t)\|_{2}^{\sigma N}$$

If the case  $\sigma = \frac{2}{N}$ , we re-arrange:

$$(1 - \frac{C_N^{\frac{4}{N}}}{\frac{2}{N} + 1} \|\phi_0\|_2^{\frac{4}{N}}) \|\nabla\phi(t)\|_2^2 \leqslant \mathcal{H}$$

Using Corollary 1.1 (explicit expression for  $C_{\sigma,N}^{2\sigma+2}$ ), we obtain the estimate:

$$(1 - (\frac{\|\phi_0\|_2}{\|\psi\|_2})^{\frac{4}{N}})\|\nabla\phi(t)\|_2^2 \leqslant \mathcal{H}$$

Taking  $\|\phi_0\|_2 \leq \|\psi\|_2$ , we obtain a time-independent bound on  $\|\nabla\phi(t)\|_2^2$ .

We use the fact that the scaling  $f(x) \to \lambda^{\frac{1}{\sigma}} f(\lambda x)$  leaves the  $L^2$  norm of f unchanged when  $\sigma = \frac{2}{N}$ . Since  $\psi$  solves the E-L equation for  $J^{\sigma,N}$  in the critical case, then re-scaling  $\psi$  by  $\frac{1}{\sigma}$  yields the equation

$$\Delta\psi - \psi + \psi^{\frac{4}{N}+1} = 0$$



## Theorem (Theorem A)

Let  $\phi_0 \in H^1(\mathbb{R}^N)$ . For  $\sigma = \frac{2}{N}$ , a sufficient condition for global existence in the IVP is:

$$\|\phi_0\|_2 \le \|\psi\|_2$$

where  $\psi$  is a positive solution of the equation

$$\Delta u - u + u^{\frac{4}{N}+1} = 0$$

of minimal  $L^2$  norm (the ground state), and  $\psi e^{\frac{it}{2}}$  is an exact solution of the IVP.

First, recall some conservation laws:

#### **Theorem**

Let  $|x|\phi_0(x) \in L^2$ , and let  $\phi(x,t)$  be an  $H^1$  solution of the NLS for  $0 \le t \le T$ . Then, for  $0 \le t \le T$ :

(i) 
$$\frac{d}{dt} \int \{|x\phi - it\nabla\phi|^2 - \frac{t^2}{\sigma + 1}|\phi|^{2\sigma + 1}\} dx = t\frac{\sigma N - 2}{\sigma + 1} \int |\phi|^{2\sigma + 2} dx$$

(ii) 
$$\frac{d^2}{dt^2} \int |\phi|^2 |x|^2 dx = 2\mathcal{H}(\phi_0) + \frac{N}{\sigma+1} (\frac{2}{N} - \sigma) \int |\phi|^{2\sigma+2} dx$$

Remark Identity (i) is referred to as "pseudoconformal conservation law". (i) proved by Ginibre and Velo 1979, (ii) proved by Vlasov et al 1971.

Glassey 1977 proved a result on finite time blow up of solutions to the NLS in the case  $\sigma \geqslant \frac{2}{N}$ . Tsutsumi strengthened this:

## Theorem (Tsutsumi)

Let either

(*i*) 
$$\mathcal{H}(\phi_0) < 0$$

(ii) 
$$\mathcal{H}(\phi_0) = 0$$
 and  $\text{Im} \int x \cdot \overline{\phi_0} \nabla \phi_0 dx < 0$ 

(iii) 
$$\mathcal{H}(\phi_0) > 0$$
 and  $\text{Im} \int x \cdot \overline{\phi_0} \nabla \phi_0 dx \leqslant -2 \sqrt{\mathcal{H}(\phi_0)} \|x \phi_0\|$ 

Then, there exists a time  $0 < T < \infty$  such that

$$\lim_{t\uparrow T} \lVert \nabla \phi(t)\rVert_2 = +\infty$$

In the critical case  $\sigma = \frac{N}{2}$  the identity due to Vlasov reduces to:

$$\frac{d^2}{dt^2}\int |\phi|^2|x|^2dx=2\mathcal{H}(\phi_0)$$

If we consider particular solutions  $\Phi(x, t) = e^{it/2}R(x)$  where R(x) is an  $H^1$  function satisfying the equation:

$$\Delta u - u + u^{\frac{4}{N}+1} = 0$$

then we have that  $\mathcal{H}(R)=0$ . As a consequence Weinstein obtains the following instability result when  $\sigma=2/N$ , which expresses the "sharpness" of the condition in his global existence theorem.

## Theorem (Instability Theorem)

Let  $\sigma = 2/N$ . The nontrivial  $H^1$  solutions of

$$\Delta u - u + u^{\frac{4}{N}+1} = 0$$

are unstable for the nonlinear Schrödinger equation in the following sense:

Let  $R \in H^1(R \neq 0)$  solve the above equation. Then for any  $\delta > 0$ , there is a function  $\eta$ , with  $\|\eta - R\|_2 < \delta$  such that for  $\phi(x,t)$  the solution of the IVP with  $\phi_0 = \eta$  and:

$$\lim_{t\to T^-}\lVert \nabla \phi(t)\rVert_2=\infty$$

for some  $0 < T < \infty$ 

The following picture emerges in the critical case  $\sigma = 2/N$ :

(1) If  $\phi_0 \in H^1(\mathbb{R}^N)$  and  $\|\phi_0\|_2 < \|\psi\|_2$ , where  $\psi$  is the ground state of

$$\Delta u - u + u^{\frac{4}{N}+1} = 0$$

(ie. positive radial and  $H^1$  solution of minimal  $L^2$  norm), then the IVP has a unique global solution  $\phi(x,t)$  of class  $C([0,\infty):H^1(\mathbb{R}^N))$ .

- (2) If  $\mathcal{H}(\phi_0) < 0$  then the solution  $\phi(x,t)$  of the NLS blows up in finite time in  $H^1(\mathbb{R}^N)$ .
- (3) By Tsutsumi's Theorem  $\mathcal{H}(\phi_0) \geqslant 0$  is not sufficient for global existence.
- (4) If  $\|\phi_0\|_2 < \|\psi\|_2$ , then  $\mathcal{H}(\phi_0) \geqslant 0$

(5) If R is a nontrivial  $H^1$  solution of the above equation then  $Re^{it/2}$  is an exact solution of the NLS and  $\mathcal{H}(Re^{it/2})=0$ . These solutions are unstable in the sense of the Instability Theorem.

#### Conclusion

Weinstein obtains a sharp sufficient condition for global existence for the NLS

$$\frac{\partial \phi}{\partial t} + \Delta \phi + |\phi|^{2\sigma} \phi = 0$$

in the  $L^2$  critical case  $\sigma = \frac{2}{N}$ , in terms of an exact stationary solution of the NLS.

This condition is derived by investigating the the sharpest constant for a classic interpolation inequality of the type Gagliardo-Nirenberg.