## Problem 2

Nagaoka bags

January 28-29, 2017

In this problem we face Hall effect-like resistivity density tensor. In the steady state there is no charge density  $\nabla E = 0$ , so we have Laplace equation  $\nabla \phi = E$ 

$$\nabla^2 \phi = 0 \tag{1}$$

Potential difference is between x = 0 and x = a sides of square. Also, let's shift the y axis so corresponding square side changes from [0, a] to [-a/2, a/2].

## 1 Boundary conditions

Boundary conditions are

$$\phi(x, a/2) = \phi_0$$

$$\phi(x, -a/2) = -\phi_0$$
(2)

and we should satisfy the requirment that no current leave through x=0 and x=a edges

$$E_x(0,y) = \lambda E_y(0,y)$$

$$E_x(a,y) = \lambda E_y(a,y),$$

$$\lambda = \frac{\rho_{xy}}{\rho_{xx}}$$
(3)

Since  $E = -\nabla \phi$ 

$$\frac{\partial \phi}{\partial x}(0, y) = \lambda \frac{\partial \phi}{\partial y}(0, y) 
\frac{\partial \phi}{\partial x}(a, y) = \lambda \frac{\partial \phi}{\partial y}(a, y)$$
(4)

## 2 Solution

One of possible solutions to (1) is linear function  $\phi(x,y) = (ax+b)(cy+d)$ . However, only

$$\phi(x,y) = \frac{2\phi_0}{a} \left[ \lambda \left( x + b \right) + y \right] \tag{5}$$

satisfies boundary conditions (4).

Then we assume separable solutions  $\phi(x,y) = X(x)Y(y)$ . (1) splits into

$$\frac{d^2X}{dx^2} = -k^2X, \frac{d^2Y}{dy^2} = -k^2Y \tag{6}$$

From boundaries (4)  $X_0'Y = \lambda Y'X_0$   $(X_0isforx = 0, X' = \frac{dX}{dx})$ . Then differentiate  $X_0'Y' = \lambda Y''X_0$  and substitute into (6)

$$X_0'Y' = \lambda k^2 Y X_0 \tag{7}$$

Eliminating Y' we have real k value

$$k^2 = \left(\frac{X_0'}{\lambda X_0}\right)^2 \tag{8}$$

That results in

$$\phi_k(x,y) = (A_k \cos kx + B_k \sin kx) \left( C_k e^{kY} + D_k e^{-kY} \right) \tag{9}$$

Applying the (4) boundary conditions

$$\frac{B_k}{A_k} = \lambda \frac{C_k e^{kY} - D_k e^{-kY}}{C_k e^{kY} + D_k e^{-kY}} \tag{10}$$

Either  $C_k = 0, B_k = -\lambda A_k$  or  $D_k = 0, B_k = \lambda A_k$ .

Most general is linear combination

$$\phi_k(x,y) = R_k \left( \cos kx + \lambda \sin kx \right) e^{ky} + S_k \left( \cos kx - \lambda \sin kx \right) e^{-ky},$$

$$R_k = A_k C_k, S_k = A_k D_k$$
(11)

From boundaries (4) we have  $\sin ka = 0 \to k_n = \frac{n\pi}{L}$  and

$$\phi(x,y) = \frac{2\phi_0}{a} \left[ \lambda (x+b) + y \right] + \sum_{n=1,2,3,...} \cos k_n x \left( R_{kn} e^{k_n y} + S_{kn} e^{-k_n y} \right) + \lambda \sin k_n x \left( R_{kn} e^{k_n y} - S_{kn} e^{-k_n y} \right)$$
(12)

From symmetry of equations (1), (4) we find  $S_n = (-1)^{n+1}R_n$  and

$$\phi(x,y) = \frac{2\phi_0}{a} \left[ \lambda \left( x - \frac{a}{2} \right) + y \right] + \sum_{m=1,3..} T_m \left[ \cos \left( \frac{m\pi}{a} x \right) \cosh \left( \frac{m\pi}{a} y \right) + \lambda \sin \left( \frac{m\pi}{a} x \right) \sinh \left( \frac{m\pi}{a} y \right) \right] + \sum_{n=2,4} U_n \left[ \cos \left( \frac{n\pi}{a} x \right) \cosh \left( \frac{n\pi}{a} y \right) + \lambda \sin \left( \frac{n\pi}{a} x \right) \sinh \left( \frac{n\pi}{a} y \right) \right]$$
(13)

where  $T_m = 2R_m$  and  $U_n = 2R_n$  (m odd, n even). Applying boundary conditions at  $y = \pm \frac{a}{2}$  one can achieve

$$T_{m} = \frac{8\phi_{0}\lambda}{\pi^{2}\cosh(m\pi/2)} - \frac{4\lambda}{\pi\cosh(m\pi/2)} \sum_{n=2,4} U_{n}\cosh\left(\frac{n\pi}{2}\right) \frac{n}{n^{2} - m^{2}}$$
(14)

$$\frac{-4\lambda}{\pi \sinh(n\pi/2)} \sum_{m=1,3} T_m \sinh\left(\frac{m\pi}{2}\right) \frac{m}{m^2 - n^2}$$
 (15)

To get solutions for  $I_y$  and R we should appropriate partial derivatives of  $\phi$  and get  $J_{x,y}$  from E

$$J_x = \frac{1}{\rho_{xx}} \sum_{m=1,3,..} T_m \frac{m\pi}{a} \sin\left(\frac{m\pi x}{a}\right) \cosh\left(\frac{m\pi y}{a}\right) - \frac{1}{\rho_{xx}} \sum_{n=2,4,..} U_n \frac{n\pi}{a} \sin\left(\frac{n\pi x}{a}\right) \sinh\left(\frac{n\pi y}{a}\right)$$
 (16)

$$J_{y} = -\frac{2\phi_{0}}{a\rho_{xx}} - \frac{1}{\rho_{xx}} \sum_{m=1,3} T_{m} \frac{m\pi}{a} \cos\left(\frac{m\pi x}{a}\right) \sinh\left(\frac{m\pi y}{a}\right) - \frac{1}{\rho_{xx}} \sum_{m=2,4} U_{n} \frac{n\pi}{a} \cos\left(\frac{n\pi x}{a}\right) \cosh\left(\frac{n\pi y}{a}\right)$$
(17)

In order to get total  $I_y$  we should count integral  $I_y = \int_0^a J_y dx$ . All harmonics except the constant first part of  $J_y$  will give zero impact.  $I_y = \frac{2\phi_0}{\rho_{xx}}$ . Remember that potential difference is  $\Phi_0 = \phi_0 - (-\phi_0) = 2\phi_0$ . Finally,

$$I_y = \frac{\Phi_0}{\rho_{xx}} \tag{18}$$

$$R = \rho_{xx} \tag{19}$$

It is often considered in textbooks that magnetic force (in Hall effect magnetic force results in  $\rho_{xy} \neq 0$ ) is balanced by induced Hall potential and the current flows only parallel to the y axis. It causes the same answer for  $I_y$  and R. However, in general J has both non-zero components everywhere except the boundary.