## Phys 5405 HW 9 4.2 4.7(a,b) 4.8(a)

**4.2** A point dipole with dipole moment  $\mathbf{p}$  is located at the point  $\mathbf{x}_0$ . From the properties of the derivative of a Dirac delta function, show that for calculation of the potential  $\Phi$  or the energy of a dipole in an external field, the dipole can be described by an effective charge density

$$\rho_{\text{eff}}(\mathbf{x}) = -\mathbf{p} \cdot \nabla \delta(\mathbf{x} - \mathbf{x}_0)$$

**4.2** The potential from the dipole is given by

$$\Phi(\vec{x}) = \frac{1}{4\pi\epsilon_0} \int d^3x' \vec{p}(\vec{x}') \cdot \nabla' \left( \frac{1}{|\vec{x} - \vec{x}'|} \right) . \tag{1}$$

The distribution of the dipole is a delta function,

$$\vec{p}(\vec{x}') = \vec{p}\,\delta(\vec{x}' - \vec{x}_0) \ . \tag{2}$$

Plug it into the formula of potential and integrate by parts,

$$\Phi(\vec{x}) = -\frac{1}{4\pi\epsilon_0} \int d^3x' \vec{p} \cdot \nabla' \delta(\vec{x}' - \vec{x}_0) \frac{1}{|\vec{x} - \vec{x}'|} \equiv \frac{1}{4\pi\epsilon_0} \int d^3x' \frac{\rho_{\text{eff}}(\vec{x}')}{|\vec{x} - \vec{x}'|} . \tag{3}$$

Therefore, the dipole can be described by an effective charge density,

$$\rho_{\text{eff}}(\vec{x}) = -\vec{p} \cdot \nabla \delta(\vec{x} - \vec{x}_0) .$$
(4)

4.7 A localized distribution of charge has a charge density

$$\rho(\mathbf{r}) = \frac{1}{64\pi} r^2 e^{-r} \sin^2 \theta$$

- (a) Make a multipole expansion of the potential due to this charge density and determine all the nonvanishing multipole moments. Write down the potential at large distances as a finite expansion in Legendre polynomials.
- (b) Determine the potential explicitly at any point in space, and show that near the origin, correct to  $r^2$  inclusive,

$$\Phi(\mathbf{r}) \simeq \frac{1}{4\pi\epsilon_0} \left[ \frac{1}{4} - \frac{r^2}{120} P_2(\cos\theta) \right]$$

## 4.7 (a) The multipole expansion is

$$\Phi(\vec{x}) = \frac{1}{4\pi\epsilon_0} \sum_{l=0}^{\infty} \sum_{m=-l}^{l} \frac{4\pi}{2l+1} q_{lm} \frac{Y_{lm}(\theta,\phi)}{r^{l+1}} , \qquad (5)$$

with multipole moments,

$$q_{lm} = \int Y_{lm}^*(\theta, \phi) r^l \rho(\vec{x}) d^3 x$$

$$= \int Y_{lm}^*(\theta, \phi) r^l \left( \frac{1}{64\pi} r^2 e^{-r} \sin^2 \theta \right) r^2 \sin \theta dr d\theta d\phi$$
(6)

The spherical harmonics are related to the associated Legendre polynomials as

$$Y_{lm}(\theta,\phi) = \sqrt{\frac{2l+1}{4\pi} \frac{(l-m)!}{(l+m)!}} P_l^m(\cos\theta) e^{im\phi} .$$
 (7)

Consider the integral  $\int_0^{2\pi} e^{im\phi} d\phi$ , it vanishes for  $m \neq 0$ . Therefore,

$$q_{lm} = \delta_{m,0} \int \sqrt{\frac{2l+1}{4\pi}} P_l(\cos\theta) r^l \left(\frac{1}{64\pi} r^2 e^{-r} \sin^2\theta\right) r^2 \sin\theta dr d\theta d\phi$$
$$= \frac{2\pi}{64\pi} \sqrt{\frac{2l+1}{4\pi}} \delta_{m,0} \left(\int dr \, r^{l+4} e^{-r}\right) \left(\int d\theta \, P_l(\cos\theta) \sin^3\theta\right)$$
(8)

By the definition of Gamma functions,

$$\int_0^\infty dr \, r^{l+4} e^{-r} = \Gamma(l+5) = (l+4)!$$
 (9)

For another integral, we change the variable as  $x = \cos \theta$ , and we also make use of the identity,  $x^2 = \frac{2}{3}P_2(x) + \frac{1}{3}$ , then the integral is now given by

$$\int_{-1}^{1} dx \, P_l(x)(1-x^2) = \frac{2}{3} \int_{-1}^{1} dx \, P_l(x)(1-P_2(x)) = \frac{2}{3} \left( 2\delta_{l,0} - \frac{2}{5}\delta_{l,2} \right) . \tag{10}$$

Therefore, the only non-zero multipole moments are

$$q_{00} = \frac{1}{2} \frac{1}{\sqrt{\pi}} , \quad q_{20} = -3\sqrt{\frac{5}{\pi}} .$$
 (11)

Therefore, the potential is given by

$$\Phi = \frac{1}{4\pi\epsilon_0} \left[ \frac{1}{r} - \frac{6P_2(\cos\theta)}{r^3} \right] = \frac{1}{4\pi\epsilon_0} \left[ \frac{P_0(\cos\theta)}{r} - \frac{6P_2(\cos\theta)}{r^3} \right] . \tag{12}$$

- **4.8** A very long, right circular, cylindrical shell of dielectric constant  $\epsilon/\epsilon_0$  and inner and outer radii a and b, respectively, is placed in a previously uniform electric field  $E_0$  with its axis perpendicular to the field. The medium inside and outside the cylinder has a dielectric constant of unity.
- (a) Determine the potential and electric field in the three regions, neglecting end effects.