

四、电磁波的传播

• I. 真空中的波动方程

Derived from Maxwell E.g.

$$\text{eq } \nabla \times (\nabla \times \vec{E}) = \nabla \times \left(-\frac{\partial \vec{B}}{\partial t} \right) = -\frac{\partial}{\partial t} (\nabla \times \vec{B}) = -\frac{\partial}{\partial t} (\mu_0 \epsilon_0 \frac{\partial \vec{E}}{\partial t})$$

$$\nabla \times (\nabla \times \vec{E}) = \nabla (\nabla \cdot \vec{E}) - \nabla^2 \vec{E} = -\nabla^2 \vec{E}$$

$$\Rightarrow \nabla^2 \vec{E} - \mu_0 \epsilon_0 \frac{\partial^2 \vec{E}}{\partial t^2} = 0$$

same for \vec{B} .

$$\nabla^2 \vec{B} - \mu_0 \epsilon_0 \frac{\partial^2 \vec{B}}{\partial t^2} = 0$$

where " $\nabla^2 f - \frac{1}{v^2} \frac{\partial^2 f}{\partial t^2} = 0$ " ... wave equation

\downarrow

$$c = \frac{1}{\sqrt{\mu_0 \epsilon_0}} \equiv \text{const.}$$

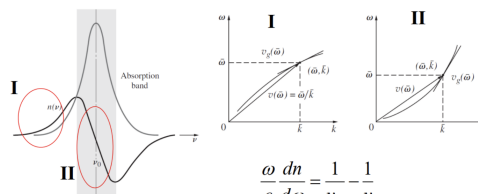
Solution (plane form) $A \cdot e^{(\vec{k} \cdot \vec{r} - \omega t)}$

• 2. 介质的色散

ϵ, μ are " f/ω " - dependent.

we learn in Optics:

Normal Dispersion & Anomalous Dispersion



(I) $v_g < v$ (II) $v_g > v$

$$\frac{\omega}{c} \frac{dn}{d\omega} = \frac{1}{v_g} - \frac{1}{v}$$

$$\Rightarrow \omega \frac{dn}{d\omega} = \frac{c}{v_g} - \frac{c}{v} = \underbrace{n_g}_{\text{Group index of refraction}} - n$$

Group index of refraction

• § 时谐电磁波 (单色波)

1) Def: $\vec{E}(\vec{x}, t) = \vec{E}(\vec{x}) e^{-i\omega t}$

$\vec{B}(\vec{x}, t) = \vec{B}(\vec{x}) e^{-i\omega t}$ $\vec{B} = \mu \vec{H}$

2) Into Maxwell E.g.

$$\left\{ \begin{array}{l} \nabla \cdot \vec{E} = 0 \\ \nabla \times \vec{E} = i\omega \mu \vec{H} \end{array} \right.$$

$$\left\{ \begin{array}{l} \nabla \cdot \vec{H} = 0 \\ \nabla \times \vec{H} = -i\omega \epsilon \vec{E} \end{array} \right.$$

Operator:

$\frac{\partial}{\partial t} \Rightarrow -i\omega$

$$\nabla \times (\nabla \times \vec{E}) = i\omega \mu \nabla \times \vec{H}$$

$$-\nabla^2 \vec{E} + \nabla(\nabla \cdot \vec{E}) = \omega^2 \mu \epsilon \vec{E}$$

Helmholtz E.g. $\left\{ \begin{array}{l} \nabla^2 \vec{E} + k^2 \vec{E} = 0 \quad (\nabla \cdot \vec{E} = 0) \\ \vec{B} = \frac{1}{\omega} \nabla \times \vec{E} \end{array} \right.$

Similarly: $\nabla^2 \vec{H} + k^2 \vec{H} = 0$

$$\vec{E} = \frac{1}{\omega \epsilon} \nabla \times \vec{H} \quad \vec{H} = -\frac{1}{\omega \mu} \nabla \times \vec{E}$$

3) Helmholtz E.g. " $\nabla^2 \vec{E} + k^2 \vec{E} = 0$ " where $k = \omega \sqrt{\mu \epsilon}$

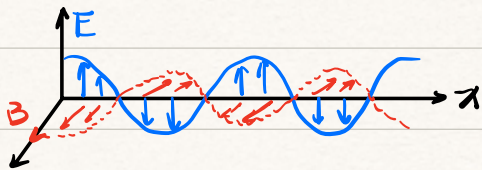
on condition: $\nabla \cdot \vec{E} = 0$

solution: $\vec{E} = \vec{E}_0 e^{i(kx - \omega t)}$
[plane wave]

• 4 时谐平面电磁波

• 振幅比 (ratio of Amplitude)

$$\frac{|\vec{E}|}{|\vec{B}|} = \begin{cases} \frac{1}{\sqrt{\epsilon\mu}} = v & (\text{普通}) \\ \frac{1}{\sqrt{\mu_0\epsilon_0}} = c & (\text{vacuum}) \end{cases}$$



$$\vec{E} = E_0 \vec{e}_x e^{i(kx - \omega t)}$$

$k = \frac{2\pi}{\lambda}$ 角波数

• 平面电磁波 $\left\{ \begin{array}{l} \text{横波} \\ \vec{E} \perp \vec{B} \dots \perp \vec{k} \quad \text{eg } i\vec{k} \cdot \vec{E} = \nabla \cdot \vec{E} = 0 \\ \vec{E} \times \vec{B} \rightarrow \vec{k} \text{ 方向} \quad \vec{E}, \vec{B} \text{ 同相} \quad E = cB \end{array} \right.$

electromagnetic wave

$$\nabla \times \vec{E} = -i\omega \vec{B} = i\vec{k} \times \vec{E} \Rightarrow \vec{B} = \frac{1}{\omega} \vec{k} \times \vec{E}$$

$$E = v \cdot B \Leftrightarrow \vec{B} = \frac{1}{v} \hat{k} \times \vec{E}$$

• Energy

(1) 时谐电磁波

Energy density $W = \frac{1}{2} (\epsilon E^2 + \frac{1}{\mu} B^2)$

E flow density $\vec{S} = \vec{E} \times \vec{H}$

$$\bar{W} = \frac{1}{T} \int_0^T W dt = \frac{1}{2} \text{Re} (\vec{E}^* \cdot \epsilon \vec{E}) = \frac{1}{2} \epsilon E^2$$

$$\bar{S} = \frac{1}{T} \int_0^T S dt = \frac{1}{2} \text{Re} (\vec{E}^* \times \vec{H}) = \frac{1}{2} \sqrt{\frac{\epsilon}{\mu}} E^2 \vec{e}_k$$

(2) 全反射

$$\vec{S}_{\text{reflex}} \stackrel{\text{in value}}{=} \vec{S}_{\text{input}}$$

• 5. 电磁波在界面上的折射与反射

we learn in Optics:

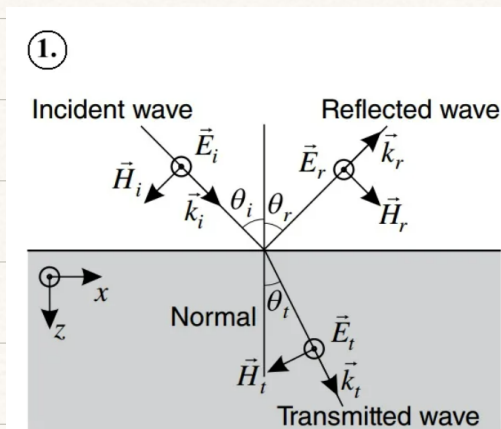
Reflection: $\theta' = \theta$

Refraction: (Snell Law)

$$\frac{\sin \theta}{\sin \theta''} = \frac{n_2}{n_1}$$

... derived from
Huygen's Principle

Fresnel's Law (2 types)



(1) 入射 \vec{E} 垂直于纸面

$$\frac{E'}{E} = \frac{\sqrt{\epsilon_1} \cos \theta - \sqrt{\epsilon_2} \cos \theta''}{\sqrt{\epsilon_1} \cos \theta + \sqrt{\epsilon_2} \cos \theta''} = -\frac{\sin(\theta - \theta'')}{\sin(\theta + \theta'')}$$

$$\frac{E''}{E} = \frac{2\sqrt{\epsilon_1} \cos \theta}{\sqrt{\epsilon_1} \cos \theta + \sqrt{\epsilon_2} \cos \theta''} = \frac{2 \cos \theta \sin \theta''}{\sin(\theta + \theta'')}$$

(2) 入射 \vec{E} 平行于纸面

$$\frac{E'}{E} = \frac{\sqrt{\epsilon_2} \cos \theta - \sqrt{\epsilon_1} \cos \theta''}{\sqrt{\epsilon_2} \cos \theta + \sqrt{\epsilon_1} \cos \theta''} = \frac{\tan(\theta - \theta'')}{\tan(\theta + \theta'')}$$

$$\frac{E''}{E} = \frac{2\sqrt{\epsilon_1} \cos \theta}{\sqrt{\epsilon_2} \cos \theta + \sqrt{\epsilon_1} \cos \theta''} = \frac{2 \cos \theta \sin \theta''}{\sin(\theta + \theta'') \cos(\theta - \theta'')}$$

偏振角 (Polarization angle): 当 $\theta_i + \theta_t = \frac{\pi}{2}$ 时, $r_{\parallel} = 0$, 此时的角度为 θ_p

临界角 (Critical angle) 光线从光密介质射向光疏介质 (内反射) 当入射角为某一数值时, 折射角等于 90° , 此入射角称临界角 $\theta_c = \arcsin \frac{n_t}{n_i}$.

Application.

(1) 半波损失 (电磁波垂直于界面方向入射, 反射光可能出现半波损失)

- \vec{E} 垂直于纸面入射: $R_{\perp} = \left(\frac{E'}{E}\right)_{\perp} = \frac{\sqrt{\epsilon_1} - \sqrt{\epsilon_2}}{\sqrt{\epsilon_1} + \sqrt{\epsilon_2}}$
- \vec{E} 平行于纸面入射: $R_{\parallel} = \left(\frac{E'}{E}\right)_{\parallel} = \frac{\sqrt{\epsilon_2} - \sqrt{\epsilon_1}}{\sqrt{\epsilon_1} + \sqrt{\epsilon_2}}$

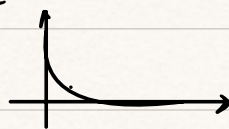
如果 $\sqrt{\epsilon_2} > \sqrt{\epsilon_1}$, $R_{\perp} < 0$ 表示 \vec{E}' 与入射方向相反, 这时产生了半波损失。

• 6. 有导体时电磁波的传播

preface $\rightarrow \rho(t) = 0 \quad \nabla \cdot \vec{J} + \frac{\partial \rho(t)}{\partial t} = 0 \quad \& \quad \vec{J} = \sigma \vec{E}$

• "Conductor" Gauss Law $\nabla \cdot \vec{E} = \frac{\rho}{\epsilon}$

then solution to PDE $\rightarrow \rho(t) = \rho_0 e^{-\frac{\sigma}{\epsilon} t}$

$\frac{\sigma}{\epsilon \omega} \gg 1$. 良导体 $\rho(t) \rightarrow 0$ like 

• Complex ϵ $\epsilon' = \epsilon + i \frac{\sigma}{\omega}$
不引起损耗. cost.

Helmholtz Eq. ϵ' into Maxwell Eq.

$\nabla^2 \vec{E} + k^2 \vec{E} = 0$, when $k = \omega \sqrt{\mu \epsilon'}$

$\vec{E}(\vec{x}) = \vec{E}_0 e^{i \vec{k} \cdot \vec{x}} \quad \vec{k} = \vec{\beta} + i \vec{\alpha}$

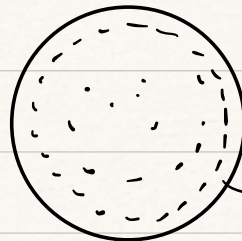
$\Rightarrow \vec{E} = \vec{E}_0 e^{-\vec{\alpha} \cdot \vec{x}} e^{i(\vec{\beta} \cdot \vec{x} - \omega t)}$

α 衰减 const β 传播 (phase const)

conditions $\begin{cases} \beta^2 - \alpha^2 = \omega^2 \mu \epsilon \\ 2 \vec{\beta} \cdot \vec{\alpha} = \omega \mu \sigma \end{cases}$

• Skin Effect

in (交变 \vec{E}/\vec{B})



$\vec{J} \rightarrow \text{skin}$

靠近. J 大.

• Penetrate depth

$E \downarrow (\frac{1}{e})$'s d.

$\delta = \frac{1}{\alpha} = \sqrt{\frac{2}{\omega \mu \sigma}}$

- 谐振腔

- 波导

- 截止频率、波长

To be updated~