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Introduction

There are different models for the electromagnetic momentum, stress, and force density inside a material (see Table 1). Some experiments have taken steps towards resolving their validity.

- The 1973 experiment by Ashkin and Dziedzic [1] showed that pulses directed onto an air-water interface caused it to bulge towards air (the lower index region) regardless of the illumination direction.
- A similar experiement by Casner et al. [2] showed that a beam directed onto a fluid-fluid interface caused it to bulge downwards, also towards the lower index region.

These observation suggest the Minkowski form but this has not been firmly established [3, 4].

We study radiation pressure on a fluid interface using a coupled electromagnetic and fluid dynamic simulator that solves Maxwell's equations, electromagnetic momentum continuity, and the Navier Stokes equation. Force densities of Abraham, Minkowski, and Einstein-Laub all predict fluid deformations consistent with experiment. Force densities of Amperian and Chu however, predict polarization-dependent deformation inconsistant with experiments.

$egin{aligned} \mathbf{Methodology} \end{aligned}$

Maxwell's equations are solved using FDTD. The force density is calculated from momentum continuity given the stess tensor \bar{T} and momentum density \vec{G} from the postulate sets listed in Table 1. Fluid deformation is modelled using the Navier Stokes equation.

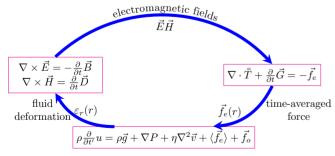


Figure 1. Simulation framework used to study radiation pressure in fluid media. Because $t \ll t'$, the time-averaged (as oppossed to instantaneous) force density is used to calculate fluid deformation.

Form	\vec{G}	$ec{S}$	$ar{ar{T}}$
Minkowski	$\vec{D} imes \vec{B}$	$c^2 \vec{D} \times \vec{B}$	$(\vec{D}\cdot\vec{E}+\vec{B}\cdot\vec{H})\bar{ar{I}}/2-\vec{D}\vec{E}-\vec{B}\vec{H}$
Abraham	$\vec{E} \times \vec{H}/c^2$	$ec{E} imes ec{H}$	$(\vec{D} \cdot \vec{E} + \vec{B} \cdot \vec{H})\bar{\bar{I}}/2 - \vec{D}\vec{E} - \vec{B}\vec{H}$
Einstein-Laub	$\vec{E} \times \vec{H}/c^2$	$\vec{E} imes \vec{H}$	$(\varepsilon_o \vec{E}^2 + \mu_o \vec{H}^2) \bar{\bar{I}}/2 - \vec{D}\vec{E} - \vec{B}\vec{H}$
Amperian	$\varepsilon_o \vec{E} \times \vec{B}$	$\vec{E} \times \vec{B}/\mu_o$	$(\varepsilon_o \vec{E}^2 + \mu_o^{-1} \vec{B}^2) \bar{\bar{I}}/2 - \varepsilon_o \vec{E} \vec{E} - \mu_o^{-1} \vec{B} \vec{B}$
Chu	$\vec{E} \times \vec{H}/c^2$	$ec{E} imes ec{H}$	$(\varepsilon_o \vec{E}^2 + \mu_o \vec{H}^2) \bar{\bar{I}}/2 - \varepsilon_o \vec{E} \vec{E} - \mu_o \vec{H} \vec{H}$

Table 1. Five formulations of electrodynamics used here to model radiation pressure on fluid interfaces.

Results

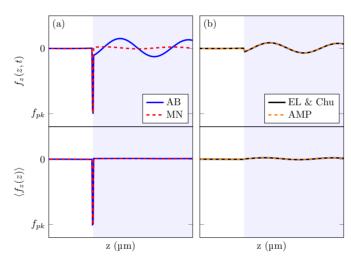


Figure 3. Degeneracy of the time-averaged force densities predicted by different formulation of electrodynamics. Here we compare the (top row) instantaneous and (bottom row) time-averaged force densities for the case of monochromatic wave of 500 nm wavelength incident from air into water.

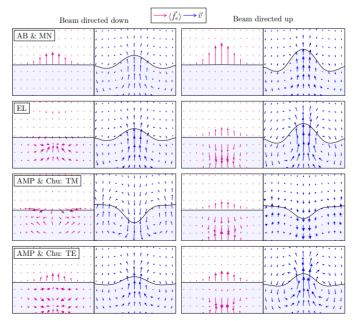


Figure 2. Simulation of the Ashkin-Dziedzic experiment using five electrodynamic models for illumination from above and below (magenta arrows) depict time-averaged force density and the (blue-arrows) depict the velocity field of the fluid.

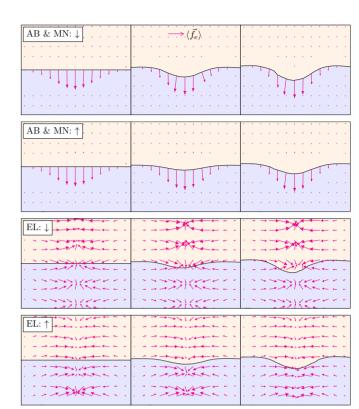


Figure 4. Simulation of the Casner et al. experiment using the three models validated in Figure 2. Magenta arrows depict time-averaged force density.

Conclusions

- Radiation pressure on fluid interfaces cannot distinguish betweeen Abraham and Minkowski models.
- The Abraham, Minkowski, and Einstein-Laub models all correctly describe observations, but present different mechanisms for bulge formation. The first two rely interface forces, the latter relies on body forces.
- The Amperian and Chu models are inconsistent with experiments.

References

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- [4] Astrath N G C, Malacarne L C, Baesso M L, Lukasievicz G V B and Bialkowski S E. Unravelling the effects of radiation forces in water. *Nat. Commun.*, 5:023826, 2014.