

(1) Non-dimensional numbers

Let's manipulate LHS & RHS to arrive at a N-D representation

$$\frac{h_x \kappa}{K} = \boxed{Nu_x = 0.332 \ Pr^{\frac{1}{3}} Re_x^{\frac{1}{2}}} \quad (Pr > 0.6)$$

Local Nusselt Number: $\frac{\text{convection in fluid}}{\text{conduction in fluid}} = \frac{h_x K_f}{\kappa} = \frac{\text{cond. resistance in fluid}}{\text{conv. resistance}}$

Advantage of this representation is that, to test this relation, we only need to vary Pr & Re_x ($\sim 10 \times 10$ data points)

$$\frac{h_x}{\kappa} \xrightarrow{\downarrow} \frac{h_x \kappa}{\kappa^2} \xrightarrow{\downarrow} \frac{h_x}{\kappa}$$

In experiments,
this means either

- ① Vary α and vary U_∞ ✓
- or
- ② Vary α and vary λ \times (two N-D in RHS vary)
- or
- ③ Vary α and vary κ \times (LHS ND also varies)

Option 1 is the most convenient route to test the above numerical result numerically.

Most exp. use $Nu = \text{const.} \times Re_l^m Pr^n$ + calculate these unknowns.

Let us return to understanding results from analysis of T-B-L



Further numerical results reveal that

$$\delta_t = \frac{\delta}{Pr^{1/3}} = \frac{4.91 x}{Re_x^{1/2} Pr^{1/3}}$$

All of these are only valid for laminar
regime
 $f \quad Pr > 0.6$ (low to moderate Re_x)

* Average H-T coefficient

We generally use average: $\bar{h} = \frac{1}{L} \int_0^L h(x) dx$

$$\bar{Nu} = \frac{\bar{h} L}{k_f} = f(Re_L, Pr)$$

Laminar Flow over Plate: $Nu_x = 0.332 Pr^{1/2} Re_x^{1/2}$

$Re < 10^5$

$$\bar{Nu} = 0.664 Pr^{1/3} Re_L^{1/2}$$

$$Pr > 0.6$$

Q. You are performing H.T experiment on a cylinder
 If $L = 0.15 \text{ m}$ $\rightarrow 100 \text{ m/s}$ $T_\infty = 37^\circ\text{C}$ ($R_{cyl} \gg \delta_v$)
~~|||||~~ $T_s = 325^\circ\text{C}$
 $\dot{Q} = 1600 \text{ J}$ is lost every second.

As a lazy experimentalist, how will you calculate \dot{Q} for

$$L = 0.30 \text{ m} \quad \rightarrow 50 \text{ m/s} \quad T_\infty = 37^\circ\text{C}$$
~~|||||~~ $T_s = 425^\circ\text{C}$
 $\dot{Q} = ?$

Ans. $\overline{\dot{Q}} = ? = \overline{h} A_s (T_s - T_\infty)$

(Exp 1.) $\overline{h} A_s = \frac{1600}{325 - 37} = 5.56 \text{ W/K}$

Then using this $\overline{h} A_s$ for
 Exp 2. will give wrong result because \overline{h} is different.
 \Downarrow

Do we know h & U_∞ relation? $h_x = 0.332 \Pr^{\frac{1}{3}} K_f \sqrt{\frac{U_\infty \delta}{M L}}$

\downarrow average

$$\overline{h} = 0.664 \Pr^{\frac{1}{3}} K_f \sqrt{\frac{U_\infty \delta}{M L}} \rightarrow ?$$

Exp 1. [5.56 ? ? ?]

Exp 2. [$\overline{h}_2 = 0.664 \Pr^{\frac{1}{3}} K_f \sqrt{\frac{U_\infty \delta}{M L_2}}$]

Unknown
But same
as Exp 1

$$0.664 \left(\Pr^{\frac{1}{3}} K_f \sqrt{\frac{\delta}{M}} \right) = \overline{h}_1 \sqrt{\frac{L_1}{U_{\infty 1}}} = \overline{h}_2 \sqrt{\frac{L_2}{U_{\infty 2}}}$$

\Downarrow

$$\overline{h}_2 = \overline{h}_1 \times \sqrt{\frac{L_1 \times U_{\infty 2}}{L_2 \times U_{\infty 1}}}$$

$$\text{Ans. } \bar{Q} = ? = \bar{h} A_s (T_s - T_\infty)$$

Expt. 1.

$$\bar{h} A_s = \frac{1600}{325-37} = 5.56 \text{ W/K}$$

Then using this $\bar{h} A_s$ for
Expt. 2. will give wrong result because \bar{h} is different.
↓

Do we know \bar{h} & U_∞ relation? $h_x = 0.332 \Pr^{\frac{1}{3}} k_f \sqrt{\frac{U_\infty \delta}{M \alpha}}$

↓ average

$$\bar{h} = 0.664 \Pr^{\frac{1}{3}} k_f \sqrt{\frac{U_\infty \delta}{M L}}$$

$$\text{Expt. 1. } \left[\begin{array}{c} \downarrow \\ 5.56 \end{array} \right]$$

$$\left[\begin{array}{c} \downarrow \\ ? \end{array} \right] \left[\begin{array}{c} \downarrow \\ ? \end{array} \right] \left[\begin{array}{c} \downarrow \\ ? \end{array} \right] \left[\begin{array}{c} \downarrow \\ ? \end{array} \right]$$

$$\text{Expt. 2. } \left[\begin{array}{c} \downarrow \\ ? \end{array} \right] \bar{h}_2 = 0.664 \Pr^{\frac{1}{3}} k_f \sqrt{\frac{U_\infty \delta}{M L_2}}$$

Fluid prop.
identical to Expt. 1

↓

$$0.664 \left(\Pr^{\frac{1}{3}} k_f \sqrt{\frac{\delta}{M}} \right) = \bar{h}_1 \sqrt{\frac{L_1}{U_{\infty 1}}} = \bar{h}_2 \sqrt{\frac{L_2}{U_{\infty 2}}}$$

$$\bar{h}_2 = \bar{h}_1 \times \sqrt{\frac{L_1}{L_2} \times \frac{U_{\infty 2}}{U_{\infty 1}}}$$

$$\bar{h}_2 A_1 = \left(\bar{h}_1 \bar{A}_1 \right) \times \sqrt{\frac{1}{4}} \Rightarrow \bar{h}_2 A_1 = \frac{\bar{h}_1 \bar{A}_1}{2}$$

↓

$$\bar{h}_2 \frac{A_1}{A_2} A_2 = \frac{\bar{h}_1 \bar{A}_1}{2}$$

0.5

↓

$$\bar{h}_2 A_2 = \bar{h}_1 A_1$$

$$\bar{Q}_2 = \bar{h}_2 A_2 (T_s - T_\infty) \quad \checkmark$$

$$\bar{h}_2 A_1 = \underbrace{\bar{h}_1 \bar{A}_1}_{\downarrow} \times \sqrt{\frac{1}{4}} \Rightarrow \bar{h}_2 A_1 = \frac{\bar{h}_1 \bar{A}_1}{2}$$

$$\bar{h}_2 \frac{A_1}{A_2} A_2 = \frac{\bar{h}_1 \bar{A}_1}{2}$$

\downarrow
0.5 \Downarrow

$$\bar{h}_2 A_2 = \bar{h}_1 A_1$$

\Downarrow

$$\dot{Q}_2 = \bar{h}_2 A_2 (T_s - T_\infty) \quad \checkmark$$

Alternative: Work with Non-dim. relation

$$\bar{h} = 0.664 Pr^{\frac{1}{3}} k_f \sqrt{\frac{SU_\infty}{ML}} \Rightarrow \bar{Nu} = 0.664 Pr^{\frac{1}{3}} Re_L^{\frac{1}{2}}$$

Now, what is changed in Exp.: Re_L ? No. $(\frac{U_\infty}{L})_1 = (\frac{U_\infty}{L})_2$
 Pr ? Obviously not

\Downarrow

$$(\bar{Nu})_1 = (\bar{Nu})_2$$

\Downarrow

$$\frac{\bar{h}_1 L_1}{k_f} = \frac{\bar{h}_2 L_2}{k_f} \Rightarrow \bar{h}_2 = \frac{\bar{h}_1 L_1}{L_2}$$

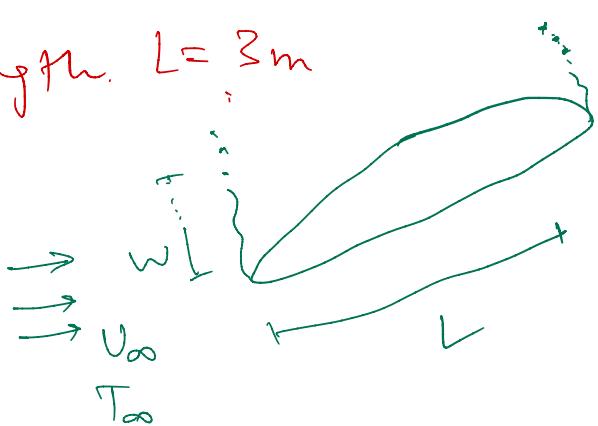
$$\dot{Q}_2 = \bar{h}_2 A_2 (T_{s_2} - T_\infty) \times \frac{A_1}{A_2}$$

$$\dot{Q}_2 = [\bar{h}_1 A_1] \frac{L_1}{L_2} \times \frac{A_2}{A_1} (T_{s_2} - T_\infty) = [\bar{h}_1 A_1] (T_{s_2} - T_\infty)$$

Q. A large object of char. length $L = 3\text{ m}$
f width $= 2\text{ m}$.

$$V_\infty = 7\text{ m/s}$$

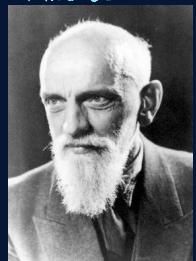
$$T_\infty = 20^\circ\text{C}$$



Total drag force on plate is 1 N. What is \bar{h} ?

Ans

What's the relation?



* Mom. f Heat Transfer Analogies

Consider the following Non-dim. G-E^s

Mom. Transfer: $u' \frac{\partial u'}{\partial x'} + v' \frac{\partial u'}{\partial y'} = \frac{1}{Re_L} \frac{L^2}{\delta_v^2} \frac{\partial^2 u'}{\partial y'^2}$

↓ using same char. scales as before (only $y_{ch} = \delta_t$ instead of δ_v)

Energy Transfer: $u' \frac{\partial \theta'}{\partial x'} + v' \frac{\partial \theta'}{\partial y'} = \frac{1}{Re_L Pr} \frac{L^2}{\delta_t^2} \frac{\partial^2 \theta'}{\partial y'^2}$

↓

If $Pr = 1$

$\delta_t = \delta_v$ & The governing Eq (and B.C) are identical for mom. & energy transfer

↓

$$\left. \frac{\partial u'}{\partial y'} \right|_{y'=0} = \left. \frac{\partial \theta'}{\partial y'} \right|_{y'=0}$$

↓

We found earlier:

$$T_w = M \left. \frac{du}{dy} \right|_{y=0} = C_f \left(\frac{8U_\infty^2}{2} \right)$$

↓
skin friction coeff.
↓
 $0.664 Re_x^{-1/2}$

Dim. less.

$$\left. \frac{du'}{dy'} \right|_{y=0} = C_f \frac{8U_\infty \delta}{M}$$

↓

$$\left. \frac{\partial \theta'}{\partial y'} \right|_{y=0} = 0.332 Pr^{1/3} \sqrt{\frac{U_\infty S}{M x}}$$

$$Nu_x = \frac{h x}{k_f} = \left[- \frac{C_f \left. \frac{\partial T}{\partial y} \right|_{y=0}}{T_s - T_\infty} \right] \frac{x}{\delta} = + \left(\left. \frac{\partial \theta'}{\partial y'} \right|_{y=0} \right) \frac{x}{\delta}$$

Non-dim.

$$\left. \frac{\partial \theta'}{\partial y'} \right|_{y=0} = Nu_x \left(\frac{\delta}{x} \right)$$

$$Nu_x = 0.332 Pr^{1/3} Re_x^{1/2}$$

$$\frac{C_f}{2} \left(\frac{\beta U_\infty}{M} \right) \delta = N u_x \left(\frac{\delta}{x} \right)$$



$$C_f \frac{Re_x}{2} = Nu_x \rightarrow \text{Analogy - I :}$$

Reynolds Analogy

(for $Pr \approx 1$)

Advantages :

⊕ No need to conduct exp. to vary thermal properties.

Only the knowledge of C_f is enough

Disadvantages :

⊖ Only works for $Pr \approx 1$

⊖ Implicit assumption of const. U_∞

⊖ Laminar R.L

Chilton-Colburn Analogy - II

Arbitrary Pr .

$$C_{f,x} = 0.664 Re_x^{-1/2} \quad \text{and} \quad Nu_x = 0.332 Pr^{1/3} Re_x^{1/2}$$

↓ ↓
Divide

$$C_{f,x} \frac{Re_x}{2} = Nu_x Pr^{-1/3}$$

Or use averaged versions:

$$\bar{C}_f = 1.32 Re_L^{-1/2}, \quad \bar{Nu} = 0.664 Pr^{1/2} Re_L^{1/2} \Rightarrow \frac{\bar{C}_f Re_x}{2} = \bar{Nu} Pr^{-1/3}$$

Advantages

- (+) \Pr is arbitrary ($\Pr > 0$, & though)
- (+) Surprisingly, works for Turbulent flows

Disadvantages

- (-) U_∞ still needs to const. so that $B-L \approx$ (that it's based on) are valid

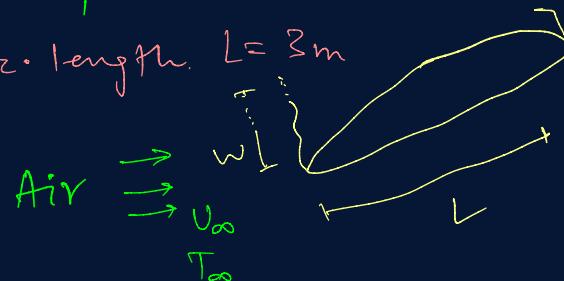
Let's Return to our example:

Q. A large object of char. length $L = 3\text{ m}$

fwidht = 2 m .

$$U_\infty = 7\text{ m/s}$$

$$T_\infty = 20^\circ\text{C}$$



Total drag force on plate is 1 N . What is h ?

Ans



Allen Colburn

We can now solve this by Colburn's blessing:

[Give me info. on mom. transfer,
I'll give you info. on heat transfer.]

$$\rightarrow \overline{C}_f \frac{Re_L}{2} = \overline{Nu} \Pr^{-1/3}$$

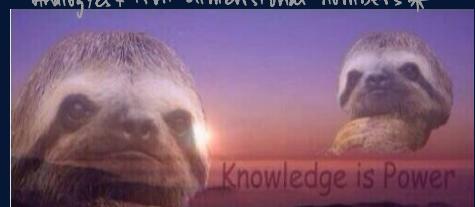
$$\text{Recall: } \tau_w = C_f \frac{\rho U_\infty^2}{2} \rightarrow \text{Units? } \frac{\text{kg}}{\text{m} \cdot \text{s}^2} = \frac{\text{N}}{\text{m}^2}$$



$$\text{Here, } F_{\text{total}} = \tau_w A_s = \overline{C}_f \frac{\rho U_\infty^2}{2} A_s$$

$$\left(\frac{F_{\text{total}}}{\rho U_\infty^2 A_s} \right) \times Re_L \downarrow = \underbrace{\overline{Nu}}_{\text{Find } h} \Pr^{-1/3}$$

When you finally understand the reason for Analogies of Non-dimensional numbers





Mom. Transfer Correlation for Turbulent flows

Laminar Flow:

$$F_{\text{Drag}} = \frac{c_f}{2} \rho U_\infty^2 + \quad \begin{array}{l} \text{Drag coefficient} \\ = \text{Friction coeff} \end{array}$$

Planes ↓ & while taking off
↑ & while landing

For Blunt shapes :  → low pressure Region \Rightarrow Additional Pressure drag

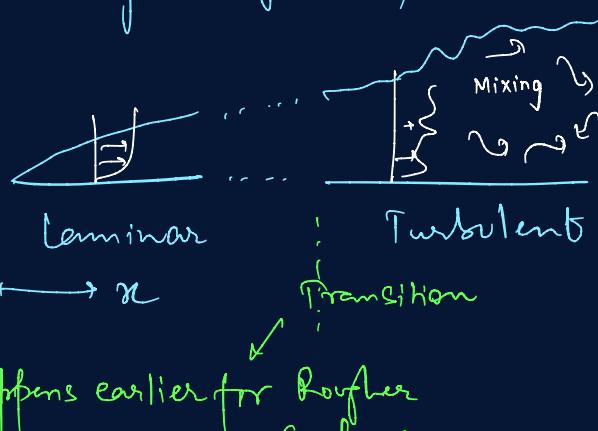
From Blasius numerical soln

$$\delta_v \approx \frac{5x}{Re_x^{1/2}},$$

$$C_{f,x} = 0.664 Re_x^{-1/2}$$

$$\bar{C}_f = 1.32 Re_L^{-1/2} \rightarrow \left(\frac{1}{L} \int_0^L C_{f,x} dx \right)$$

For long-enough surfaces :



for x such that $Re_x > 10^5$
Turbulent Reynolds

happens earlier for rougher surfaces

What do you expect? $\delta_v \propto n^?$ \Rightarrow Increased

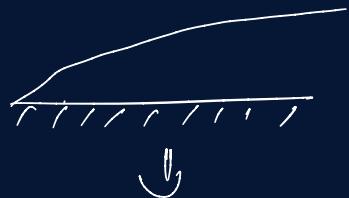
Turbulence:

$$\delta_t \approx ? = \frac{0.38x}{Re_x^{1/5}}, \quad C_{f,x} = 0.059 Re_x^{-1/5}$$

$$\bar{C}_f = 0.074 Re_L^{1/5}$$

Turbulent \bar{C}_f ↑ rapidly with roughness (δ_{laminar} , not much)

~~if~~



↓
Use Laminar Correlations

$$Re = 10^3$$

if

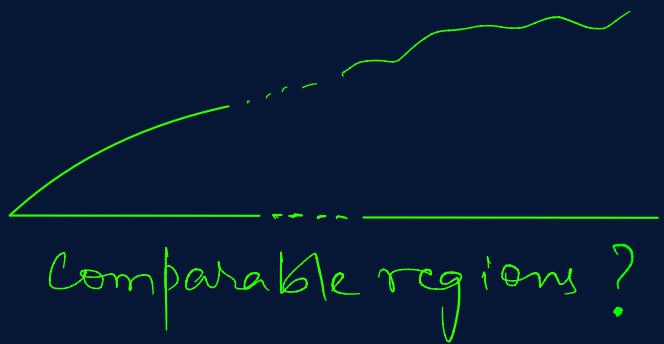
Laminar Region

Turbulent Region



↓
Use Turbulent Correlations

Q. What if



↓
What to use?

* H.T correlations for Turbulent flows

Laminar: From Blasius numerical solⁿ

$$\delta_t = \frac{\delta_v}{Pr^{1/3}}$$

$$Nu_x = \frac{h_x x}{K_f} = 0.332 Re_x^{1/2} Pr^{1/3} \quad (Pr > 0.6)$$

$$\bar{Nu} = 0.664 Re_x^{0.5} Pr^{1/2}$$

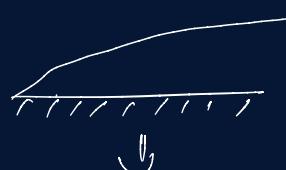
Turbulent : $Re > 10^5$

$$Nu_x = 0.0296 Re_x^{0.8} Pr^{1/3} \quad (Pr > 0.6)$$

$$\bar{Nu} = 0.037 Re_L^{0.8} Pr^{1/3}$$

Similarly, like before, use correlations wisely !

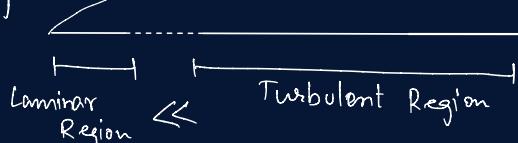
if



$$Re = 10^3$$

Use Laminar Correlations

if



Use Turbulent Correlations

* Sp1. Case: Uniform heat flux (Newmann BC for θ instead of Dirichlet)

Laminar: $Nu_x = 0.453 Re_x^{0.5} Pr^{1/3}$

$$Pr > 0.6$$

Turbulent: $Nu_x = 0.031 Re_x^{0.8} Pr^{1/3}$

Slightly increased Nu_x from isothermal case?