

■ p.39 Riemann tensor: 4つの基本的性質

$$(\nabla_a \nabla_b - \nabla_b \nabla_a) w_c \equiv R_{abc}{}^d w_d \quad (1)$$

左辺で w_c は任意の dual vector であり、derivative の交換子は (1,3)型 tensor を定義するので右辺のように表すことができる。

明らかに

$$(\natural 1) \quad R_{bac}{}^d = -R_{abc}{}^d \quad (2)$$

(3.2.17),(3.2.14):

$$\begin{aligned} \nabla_{[a} \nabla_b w_{c]} &= \partial_{[a} (\nabla_b w_{c]}) - \Gamma_{[a}{}^d{}_{[b} (\nabla_{d|} w_{c]}) - \Gamma_{[a}{}^d{}_{[b} (\nabla_{c]} w_d) \\ &= \partial_{[a} (\nabla_b w_{c]}) \quad \because \Gamma_{ab}^c = \Gamma_{ba}^c \\ &= \partial_{[a} \partial_b w_{c]} - \partial_{[a} (\Gamma_{bc]}^d w_d) \\ &= 0 \end{aligned}$$

$$\begin{aligned} \therefore 0 &= 2 \nabla_{[a} \nabla_b w_{c]} \\ &= \nabla_{[a} \nabla_b w_{c]} - \nabla_{[b} \nabla_a w_{c]} \\ &= R_{[abc]}{}^d w_d \\ (\natural 2) \quad \therefore R_{[abc]}{}^d &= 0 \end{aligned} \quad (3)$$

(3.2.15):

$$\begin{aligned} 0 &= (\nabla_a \nabla_b - \nabla_b \nabla_a) g_{cd} = R_{abc}{}^e g_{ed} + R_{abd}{}^e g_{ce} \\ &= R_{abcd} + R_{abdc} \\ (\natural 3) \quad \therefore R_{abdc} &= -R_{abcd} \end{aligned} \quad (4)$$

(3.2.20): (3.2.14),(3.2.15) より

$$0 = 3 R_{[abc]d} = R_{abcd} + R_{cabd} + R_{cbad} \quad (5)$$

$$(5) \text{ で } a \leftrightarrow d \text{ の置換 } 0 = R_{dbca} + R_{cdba} + R_{bcd a} = R_{b\cancel{d}ac} - R_{cdab} - R_{b\cancel{c}ad}$$

$$(5) \text{ で } b \leftrightarrow d \text{ の置換 } 0 = R_{adcb} + R_{cadb} + R_{dcab} = -R_{a\cancel{d}bc} - R_{c\cancel{a}bd} - R_{cdab}$$

$$(5) \text{ で } c \leftrightarrow d \text{ の置換 } \times -1 \quad 0 = -R_{abdc} - R_{dabc} - R_{bdac} = R_{abcd} + R_{c\cancel{d}ba} - R_{b\cancel{d}ac}$$

$$\begin{aligned} \text{上の四式の和を取ると} \quad 0 &= 2(R_{abcd} - R_{cdab}) \\ \therefore R_{abdc} &= \therefore R_{cdab} \end{aligned} \quad (6)$$

(3.2.24) Bianchi identity:

$$(\natural 4) \quad \nabla_{[a} R_{bc]d}{}^e = 0 \quad (7)$$

Riemann tensorの自由度勘定:

(\natural 1) ~ (\natural 4) の性質は (時) 空間の次元に関係が無い。 n 次元であるとする、(\natural 1),(\natural 3)より自由度は ${}_n C_2 \times {}_n C_2$. 更に(\natural 2)は ${}_n C_3 \times n$ 個の恒等式なので、Riemann tensor の自由度は少なくとも

$$\begin{aligned} {}_n C_2 \times {}_n C_2 - {}_n C_3 &= n^2(n-1)^2/4 - n^2(n-1)(n-2)/6 \Rightarrow \frac{n}{\text{自由度}} \left| \begin{array}{ccc} 2 & 3 & 4 \\ 1 & 6 & 20 \end{array} \right. \\ &= n^2(n^2-1)/12 \end{aligned}$$

にまで減る。 $n=4$ であれば $16 \times 15/12 = \mathbf{20}$ となる。 加えてBianchi恒等式(\natural 4)がある。

■ Ricci tensor, Ricci scalar

(3.2.26): (5)より

$$\begin{aligned} 0 &= (R_{abcd} + R_{cabd} + R_{bcad}) g^{bd} \\ &= R_{ac} + 0 - R_{ca} \\ \therefore R_{ac} &= R_{ca} \end{aligned}$$

(3.2.28): Weyl tensor

$$\begin{aligned} B_{abcd} g^{bd} &\equiv g_{a[c} g_{d]b} g^{bd} = \frac{1}{2} (g_{ac} g_{db} - g_{ad} g_{cb}) g^{bd} \\ &= \frac{1}{2} (n g_{ac} - g_{ad} g_c^d) \\ &= \frac{(n-1)}{2} g_{ac} \\ A_{abcd} g^{bd} &\equiv (g_{a[c} R_{d]b} - g_{b[c} R_{d]a}) g^{bd} = \frac{1}{2} (g_{ac} R_{db} - g_{ad} R_{cb} + g_{bd} R_{ca} - g_{bc} R_{da}) g^{bd} \\ &= \frac{1}{2} (g_{ac} R - R_{ca} + n R_{ca} - R_{ac}) \\ &= \frac{n-2}{2} R_{ac} + \frac{1}{2} g_{ac} R \\ \underbrace{(R_{abcd} - \alpha A_{abcd} - \beta B_{abcd} R)}_{\equiv C_{abcd}} g^{bd} &= \left\{ 1 - \alpha \frac{(n-2)}{2} \right\} R_{ac} - \frac{\alpha + \beta(n-1)}{2} g_{ac} R \\ \alpha = \frac{2}{n-2}, \beta = -\frac{\alpha}{n-1} &= -\frac{2}{(n-1)(n-2)} \Rightarrow = 0 \end{aligned}$$

(3.2.32) : Einstein tensor

Bianchi 恒等式(7)は

$$\begin{aligned} \nabla_a R_{bcd}^e + \nabla_b R_{cad}^e + \nabla_c R_{abd}^e &= 0 \\ \text{上式} \times g^{bd} &\Rightarrow \nabla_a R_c^e - \nabla_b R_a^b{}^e - \nabla_c R_a^e = 0 \\ \text{上式} \times g_{be}, \quad b \leftrightarrow c &\Rightarrow \nabla_a R_{bc} - \nabla_b R_{ac} + \nabla_d R_{abc}^d = 0 \end{aligned}$$

$$\text{上式} \times g^{bc} \Rightarrow \nabla_a R - 2 \nabla_b R_a^b = 0$$

$$\therefore G_{ab} \equiv R_{ab} - \frac{1}{2} g_{ab} R \Rightarrow \nabla^a G_{ab} = 0$$