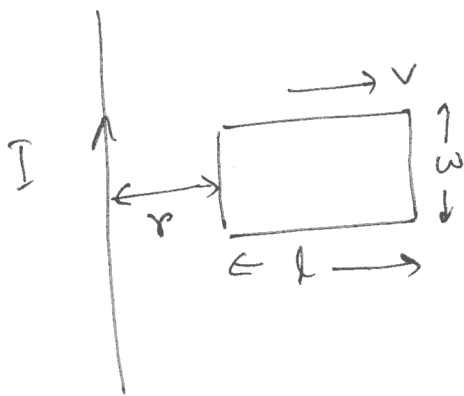


Assignment 7 (sols.)



Method 1:

$$|\mathcal{E}| = v\omega(B_1 - B_2)$$

$$B_1 = \frac{\mu_0 I}{2\pi r}, \quad B_2 = \frac{\mu_0 I}{2\pi(r+l)}$$

$$r = 15 \text{ cm} = 0.15 \text{ m}; \quad l = 10 \text{ cm} = 0.1 \text{ m}$$

$$v = 5 \text{ m/sec}; \quad w = 8 \text{ cm} = 0.08 \text{ m}, \quad I = 100 \text{ A}$$

$$|\mathcal{E}| = 5 \times 0.08 \cdot \frac{\mu_0}{2\pi} \left(\frac{100}{0.15} - \frac{100}{0.25} \right)$$

$$= 0.4 \times 2 \cdot 10^{-7} \times 10^2 \cdot \frac{0.25 - 0.15}{0.15 \times 0.15}$$

$$= 2.13 \times 10^{-5} \text{ volts.}$$

Method 2: $\phi = \int \vec{B} \cdot d\vec{a} = \int_r^{r+l} \frac{\mu_0 I}{2\pi x} \omega dx$ Note:
 $x \equiv x(t)$
 $r \equiv r(t)$

$$= \frac{\mu_0 I \omega}{2\pi} \ln x \Big|_r^{r+l} = \frac{\mu_0 I \omega}{2\pi} \ln \left(\frac{r+l}{r} \right)$$

$$\mathcal{E} = - \frac{d\phi}{dt} = - \frac{\mu_0 I \omega}{2\pi} \cdot \frac{r}{r+l} \left(\frac{1}{r} \frac{dr}{dt} - \frac{1}{r+l} \frac{dr}{dt} \right)$$

$$= - \frac{\mu_0 I \omega}{2\pi} \cdot \frac{r}{r+l} \left(\frac{v}{r} - \frac{v(r+l)}{r^2} \right)$$

$$= \frac{\mu_0 I \omega}{2\pi} \frac{dv}{v(r+l)} = v\omega \left(\frac{\mu_0 I}{2\pi r} - \frac{\mu_0 I}{2\pi(r+l)} \right)$$

$$= v\omega(B_1 - B_2) \quad (\text{Same as above})$$

Direction of \mathcal{E} will be in a direction to drive current such that flux decreases. Since flux is downward & decreases, induced current is clockwise.

The current in the loop at a time t is $I(t) = \frac{\mathcal{E}(t)}{R}$ where R is a non-negligible resistance in the loop.

This current creates a field ~~and~~ of its own and therefore a flux linked with it. Since \mathcal{E} is changing with time, so does this flux. This creates an additional induced emf \mathcal{E}' which we have ignored. We need to therefore find R so that, $\mathcal{E} \gg \mathcal{E}'$ & we can safely ignore \mathcal{E}' .

The field created by the current, $B' \sim \frac{\mu_0 I(t)}{2\pi \cdot L}$

where L is some length ~ 10 cm.

$$\begin{aligned} \therefore \phi' &\sim B' \cdot \text{area of loop} = \frac{\mu_0 I(t) \times 0.08 \times 0.1}{2\pi \times 0.1} \\ &= 2 \times 10^{-7} \times 0.08 \times I(t) \end{aligned}$$

$$\therefore \phi' \sim 0.16 \times 10^{-5} \times I(t)$$

A typical time scale in the problem would be,

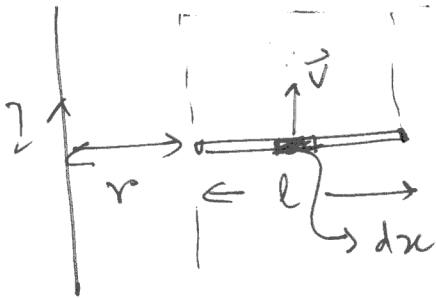
$$\tau' \sim \frac{l}{v} = \frac{0.1}{5} \text{ sec}$$

$$\begin{aligned} \therefore \mathcal{E}' &\sim \frac{\phi'}{\tau'} = \frac{0.16 \times 10^{-5} \times 5 \times I(t)}{0.1} = 8 \times 10^{-5} I(t) \\ &= 8 \times 10^{-5} \frac{\mathcal{E}(t)}{R} \end{aligned}$$

$$\therefore R \sim \left(8 \times 10^{-5} \frac{\epsilon}{\epsilon'} \right) \Omega$$

\therefore even for $R \gg 10^{-5} \Omega$ we will have $\epsilon \gg \epsilon'$.

2.



Let the ~~rod~~ conducting rod sweep an area in time t .

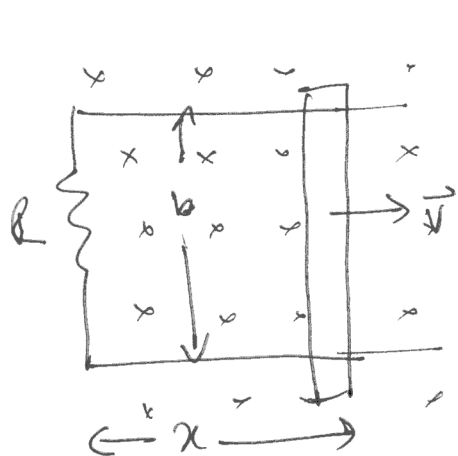
Area swept by an element dx of rod, at distance x from current carrying wire, is given as $\underbrace{vt}_{\text{vertical length}} dx = da$.

Magnetic field at a distance x from current carrying wire, $B = \frac{\mu_0 I}{2\pi x}$.

$$\begin{aligned} \therefore \phi &= \int B da = \int_r^{\text{rel}} \frac{\mu_0 I}{2\pi x} vt dx \\ &= \frac{\mu_0 I}{2\pi} vt \cdot \ln\left(\frac{\text{rel}}{r_0}\right) \end{aligned}$$

$$\therefore |\epsilon| = \frac{d\phi}{dt} = \frac{\mu_0 I}{2\pi} v \ln\left(\frac{\text{rel}}{r_0}\right)$$

3.



$\times \Rightarrow$ denotes magnetic field direction

$\therefore \vec{B} = -B\hat{z}$ pointing into the page.

$$\vec{v} = v\hat{x}$$

The rod moves ~~is~~ to the right in the \hat{x} direction due to some external force.

Let at a time, t , the crossbar of mass m is at a distance x as shown.

Then it ~~is~~ forms closed loop and during its motion, area of loop increases with time.

With the magnetic field acting downwards, ~~mag~~ flux increases with time.

Flux, $\phi = BA = Bbx$

$$\therefore \mathcal{E} = -\frac{d\phi}{dt} = -\frac{d}{dt}(Bbx) = -Bb\frac{dx}{dt} = -Bbv$$

This induced emf will cause current in the loop in the counterclockwise direction so that it opposes the increase in flux by giving rise to a magnetic field to oppose the external magnetic field \vec{B} .

Now, induced current, $I = \frac{|\mathcal{E}|}{R} = \frac{Bbv}{R}$

∴ Magnetic force experienced by bar,

$$\vec{F}_B = I \vec{l} \times \vec{B} = I b \hat{y} \times (-B \hat{z}).$$

$$= -I b B (\hat{y} \times \hat{z}) = -I b B \hat{x}$$

$$= - \frac{B^2 b^2 v}{R} \hat{x}$$

→ opposite to \vec{v} .

∴ For the bar to keep moving at a constant speed, an external force needs to be present so that,

$$\vec{F}_{ext} = - \vec{F}_B = \frac{B^2 b^2 v}{R} \hat{x}.$$

Suppose at $t=0$, speed of rod $= v_0$ & the external agent stops pushing. Then,

$$F_B = - \frac{B^2 b^2 v}{R} = m a = m \frac{dv}{dt}.$$

$$\Rightarrow \frac{dv}{v} = - \frac{B^2 b^2}{mR} dt = - \frac{1}{\tau} dt, \quad \tau = \frac{mR}{B^2 b^2}$$

Integrating,

$$v(t) = v_0 e^{-t/\tau}.$$

∴ Speed decreases exponentially in the absence of external force doing work.

~~Now distance~~

$$x = \int_0^\infty v(t) dt = \int_0^\infty v_0 e^{-t/\tau} dt = v_0 \tau.$$
$$= \frac{mR}{B^2 b^2} v_0.$$

Power delivered by \vec{F}_{ext} = power dissipated in resistor

$$\therefore P = \vec{F}_{\text{ext}} \cdot \vec{v} = \frac{B^2 b^2 v}{R} \cdot v = \left(\frac{B b v}{R} \right)^2 = \frac{\mathcal{E}^2}{R} = I^2 R.$$

4. (a) $\phi = BA = (B_0 + bt) \pi a^2 = \pi (B_0 + bt) a^2$
At $t=0$, $\phi = \pi B_0 a^2$.

(b) $\mathcal{E} = - \frac{d\phi}{dt} = - \frac{d}{dt} (BA) = \cancel{\pi a^2 \frac{dB}{dt}}$
 $= -A \frac{dB}{dt} = -\pi a^2 \cdot \frac{d}{dt} (B_0 + bt)$
 $= -\pi a^2 \cdot b$

(c) $I = \frac{|\mathcal{E}|}{R} = \frac{\pi a^2 b}{R}$

(d) $P = I^2 R = \left(\frac{\pi b a^2}{R} \right)^2 R = \frac{(\pi b a^2)^2}{R}$

5. Faraday's law : $\oint \vec{E} \cdot d\vec{s} = - \frac{d}{dt} \int \vec{B} \cdot d\vec{a}$
(integral form)

Inside solenoid, $B = \mu_0 n I$

$$\therefore \phi = \int \vec{B} \cdot d\vec{a} = B \cdot \pi r^2 = \mu_0 n I \pi r^2$$
$$= \mu_0 n (ct) \pi r^2 \quad (\because I = ct)$$

\therefore

∴ Faraday's law gives magnitude of tangential component of \vec{E} as,

$$E_{\theta} \cdot 2\pi r = \frac{d}{dt} (\mu_0 n C t + \pi r^2)$$

$$E_{\theta} \cdot 2\pi r = \mu_0 n C \cancel{r^2}$$

$$\Rightarrow E_{\theta} = \frac{\mu_0 n C r}{2}$$

Outside solenoid, $\phi = (\mu_0 n C t) \pi R^2$

where R : radius of solenoid.

This happens because, field is zero outside.

$$\therefore E_{\theta} \cdot 2\pi r = \frac{d}{dt} (\mu_0 n C t + \pi R^2)$$

$$E_{\theta} \cdot 2\pi r = \mu_0 n C \cancel{r^2}$$

$$\Rightarrow E_{\theta} = \frac{\mu_0 n C R^2}{2r}$$

To check that these results are correct, use,

$$\vec{\nabla} \times \vec{E} = - \gamma \frac{\vec{B}}{\gamma t}$$

$$\& \quad \vec{\nabla} \times \vec{E} = \hat{z} \frac{1}{r} \frac{\partial}{\partial r} (r E_{\theta})$$

in cylindrical coordinates for this problem. This is the only term in curl which survives.