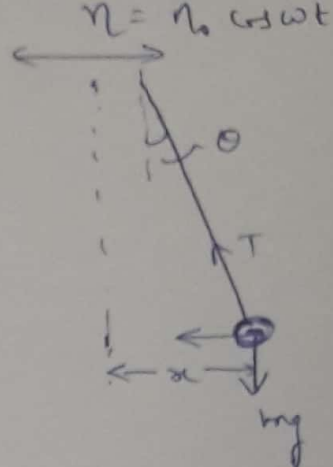


**Q-1**

(a) Pls follow derivation in class to arrive at

$$\ddot{x} + \omega_0^2 x + \gamma \dot{x} = \omega_0^2 \eta_0 \cos \omega t$$



(b) Steady state:

$$x = A \cos(\omega t - \delta)$$

$$A = \frac{\omega_0^2 \eta_0}{\sqrt{(\omega_0^2 - \omega^2)^2 + (\omega \gamma)^2}}$$

$$\tan \delta = \frac{\omega \gamma}{\omega_0^2 - \omega^2}$$

(c) At exact resonance ( $\omega \simeq \omega_0$ ) for  $Q \gg 1$

$$A_{\max} = Q \eta_0$$

$$Q = N \cdot x = 50 \times 3.14 = 157.0$$

$$\Rightarrow A_{\max} \simeq 157 \text{ mm.}$$

**Q-2**

Power dissipated against frictional force

(a)  $P = \frac{dW}{dt} = -F \cdot v$  (work done against b v)  
 $= + b v^2$

Instantaneous power  $P(t) = + b v^2$

(b)  $x = A \cos(\omega t - \delta)$

$$\Rightarrow \dot{x} = -A\omega \sin(\omega t - \delta)$$

Mean power  $\langle P \rangle = b \langle v^2 \rangle = b (A\omega)^2 \langle \sin^2(\ ) \rangle$   
 $= \frac{1}{2} b (A\omega)^2$

Q.3: Given  $\frac{dE(t)}{dt} = P(t) = \frac{ke^2}{c^3} a^2$   
 (a)  $c \leftarrow \text{constant}$

$$P(t) = c a^2(t)$$

for damped H.O.:  $E(t) = E_0 e^{-\gamma t}$

$$\Rightarrow \frac{dE}{dt} = -\gamma E$$

Cycle avg. loss  $\frac{d}{dt} \langle E(t) \rangle = -\gamma \langle E(t) \rangle = -c \langle a^2(t) \rangle$

→ for 1d SHM:

$$X(t) = X_0 \cos \omega t; \quad \omega = 2\pi\nu$$

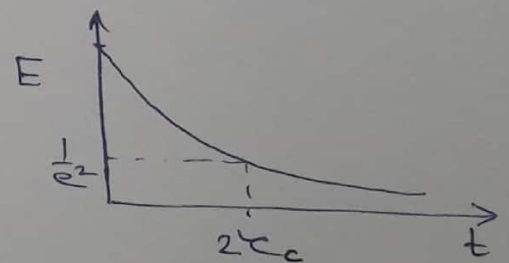
$$\ddot{x} = a(t) = -\omega^2 X_0 \cos \omega t$$

$$\langle a^2(t) \rangle = \omega^4 X_0^2 \langle \cos^2 \omega t \rangle = \frac{1}{2} \omega^4 X_0^2$$

→ Coeff. of damping

$$\gamma = \frac{c \langle a^2(t) \rangle}{\langle E \rangle} = \frac{c \left( \frac{1}{2} \omega^4 X_0^2 \right)}{\frac{1}{2} m \omega^2 X_0^2} = \frac{c \omega^2}{m}$$

$$\boxed{\gamma = \frac{c \omega^2}{m}}$$



⇒ Energy decays to  $\frac{1}{2} E_0$  in  $\tau_c = \frac{1}{\gamma}$  sec

(b) Q-factor:  $Q = \frac{\omega_0}{\gamma} = \frac{m}{c \omega_0}$

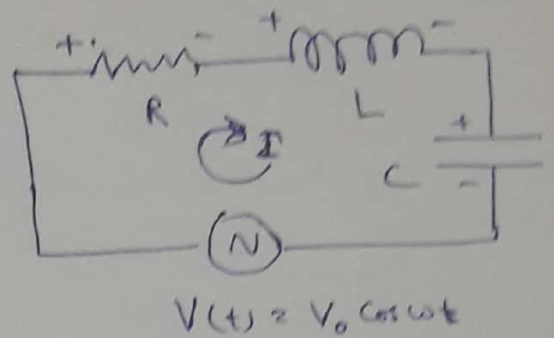
Plugging in values  $Q \sim 10^5$

⇒ It is very good oscillator.

Note: Recall, this radiative loss of ~~charge~~ energy was an argument going against stability of H-atom.

Q-4

Using voltage drop across  
R, L, C



$$IR + L \frac{dI}{dt} + \frac{Q}{C} = V_0 \cos \omega t$$

$$\Rightarrow \text{Use } I = \frac{dQ}{dt} \quad \& \quad \omega_0^2 = \frac{1}{LC}$$

$$\gamma = \frac{R}{L}$$

$$\ddot{Q} + \omega_0^2 Q + \gamma \dot{Q} = \frac{V_0}{L} \cos \omega t$$

(b) In steady state

$$Q = Q_0 \cos(\omega t - \delta)$$

$$Q_0 = \frac{\omega_0^2 (V_0/L)}{\sqrt{(\omega_0^2 - \omega^2)^2 + (\omega\gamma)^2}} = \frac{V_0}{\omega \sqrt{R^2 + (\omega L - 1/\omega C)^2}}$$

$$\tan \delta = \frac{\omega\gamma}{\omega_0^2 - \omega^2} = \frac{1}{R} (\omega L - \frac{1}{\omega C})$$

(c) Power delivered to the ckt by the voltage source

$$P(t) = V(t) \cdot I(t)$$

$$V(t) = V_0 \cos \omega t$$

$$I(t) = I_0 \sin(\omega t - \delta)$$

Max power is consumed near resonance  $\omega \approx \omega_0$

$$P(t) = V_0 I_0 \cos \omega t \sin(\omega t - \delta)$$

$$\langle P(t) \rangle = \frac{1}{T} \int_0^T P(t) dt = \frac{V_0^2}{2Z} \overbrace{\cos \phi}^{\text{power factor}} ;$$

$$Z = \sqrt{R^2 + (\omega L - \frac{1}{\omega C})^2}$$

$$\langle P \rangle = I_{\text{rms}}^2 R = \frac{1}{2} \frac{V_o^2}{R}$$

$\Rightarrow$  All the power will be dissipated in the resistor  $R$