

# **IUPAC: A Report of IUPAC Commission I.2 on Thermodynamics**

## **Standard quantities in chemical thermodynamics. Fugacities, activities, and equilibrium constants for pure and mixed phases**

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The standard chemical potential  $\mu^\circ$  is defined in terms of a specified state of aggregation, a specified temperature, a specified standard pressure, and if needed a specified standard molality. Other standard thermodynamic quantities are obtained from  $\mu^\circ$  by the usual mathematical operations. In this report, definitions are given for the standard chemical potential of a substance B in the following states of aggregation: gas phase whether pure or mixed, condensed phase whether pure or mixed, solvent substance in solution, and solute substance in solution. Fugacities and virial coefficients are introduced in the discussion of gases, and activity, activity coefficients, and osmotic coefficients in the discussion of condensed phases. Expressions for the standard equilibrium constant are given for a reaction in the gas phase, in a condensed phase, and in a solution. Some numerical examples are given and relative magnitudes of various terms in the equations are illustrated.

### **1. Introduction**

The chemical potential  $\mu_B$  of a substance B in a mixture of substances B, C, ... is related to the Gibbs energy  $G$  of the mixture by

$$\mu_B = (\partial G / \partial n_B)_{T, p, n_C \neq B}, \quad (1)$$

where  $T$  is the thermodynamic temperature,  $p$  is the pressure, and  $n_B, n_C, \dots$  are the amounts of substance of B, C, ... For a pure substance B, the chemical potential  $\mu_B^*$  is given by

$$\mu_B^* = G^* / n_B = G_m^*, \quad (2)$$

where  $G_m^*$  is the molar Gibbs energy, and where the superscript \* attached to a symbol denotes the property of the pure substance. The superscripts  $^\circ$  or  $^\ominus$  attached to a symbol may be used to denote a standard thermodynamic quantity. Consequently, the standard chemical potential of substance B may be written as  $\mu_B^\circ$

or  $\mu_B^\ominus$ , but in this document the superscript  $^\ominus$  has been used to denote standard thermodynamic quantities.

The word **mixture** is used to describe a gaseous or liquid or solid phase containing more than one substance, when the substances are all treated in the same way. It is sometimes convenient to treat a mixture of constant composition, such as air, as a pure substance.

The word **solution** is used to describe a liquid or solid phase containing more than one substance, when for convenience one (or more) substance, which is called the **solvent**, is treated differently from the other substances, which are called **solutes**. When, as is often but not necessarily the case, the sum of the mole fractions of the solutes is small compared with unity, the solution is called a **dilute solution**. In the definitions given in Section 5, the solvent substance is denoted by A and the solute substances by B, C,  $\dots$ .† The superscript  $^\infty$  attached to the symbol for a property of a solution denotes the property in the limit of infinite dilution.

The nomenclature and notation used in this document are consistent with the recommendations given in “Quantities, Units and Symbols in Physical Chemistry” (the IUPAC “Green Book”),<sup>(1)</sup> and in ISO 31 Part 8.<sup>(2)</sup>

## 2. Standard thermodynamic quantities

Absolute values of some thermodynamic quantities, such as chemical potential and entropy, are not needed in applications of thermodynamics. Only changes in these quantities, such as may result from variation of temperature or pressure or composition, can be measured or calculated. It is therefore important to define a reference value of each thermodynamic quantity to which such changes may be referred. Suitable reference values, known as **standard thermodynamic quantities**, can be derived if the standard chemical potential is defined.

The **standard chemical potential** is defined in terms of a specified state of aggregation, a specified temperature, and an arbitrarily chosen but specified standard pressure. For a solute substance in a solution {see section 5(a)} an arbitrarily chosen but specified standard molality is also required. Other standard thermodynamic quantities are obtained from the standard chemical potential by the usual mathematical operations.

For example, the standard molar entropy  $S_B^\ominus$  of a substance B is defined by  $S_B^\ominus(z, T) = -d\mu_B^\ominus(z, T)/dT$ , where  $z$  denotes the state of aggregation.

For a specified standard pressure, and in the case of a solute for a specified standard molality, all standard thermodynamic quantities are independent of pressure and composition, but they remain functions of temperature. Consequently, standard thermodynamic quantities are often required at many temperatures. For quantities such as  $\{H_m^\ominus(T) - H_m^\ominus(T')\}$  and  $\Delta_r H_m^\ominus(T')$  it is generally convenient to choose the reference temperature  $T'$  as one of a relatively small number of values: 0, 273.15 K, 293.15 K, or 298.15 K. The most widely used values are  $T' = 0$ , for

†The solvent is often denoted by the subscript 1 and the solutes by the subscripts 2, 3,  $\dots$ .

example with the quantity  $\{H_m^\circ(T) - H_m^\circ(0)\}$ , and  $T' = 298.15\text{ K}$ , for example with the quantity  $\Delta_f H_m^\circ(298.15\text{ K})$ . Because the quantity “298.15 K” occurs so frequently, workers may find it convenient to use a special symbol for this quantity. No recommendation is made for a special symbol for “298.15 K” but  $T_a$ ,  $T'$ , and  $T^\circ$  have all been used. Since  $T^\circ$  would mean a standard temperature in general, the use of  $T^\circ$  to mean exclusively 298.15 K is strongly discouraged. The value of any reference temperature for standard thermodynamic quantities should *always* be stated when the numerical value of any standard thermodynamic quantity is given and any special symbol used to denote the quantity “298.15 K” should *always* be defined.

The **standard pressure** is denoted by  $p^\circ$  and the recommended value for standard thermodynamic quantities is  $10^5\text{ Pa}$ .<sup>†</sup> Most standard quantities published before 1982 used a standard pressure of “1 standard atmosphere”, a non-SI unit equal to 101 325 Pa. If the value of  $p^\circ$  changes from  $p_1^\circ$  to  $p_2^\circ$  then the difference in the standard chemical potential (see sections 3 to 5) is given by one of the following three equations.

(a) For a gaseous (g) substance B, whether pure or mixed,

$$\mu_B^\circ(\text{g}, T, p_2^\circ) - \mu_B^\circ(\text{g}, T, p_1^\circ) = RT \cdot \ln(p_2^\circ/p_1^\circ). \quad (3)$$

(b) For a substance B in a condensed phase (cd), whether pure or mixed, or for a solvent substance B in a solution,

$$\mu_B^\circ(\text{cd or solvent}, T, p_2^\circ) - \mu_B^\circ(\text{cd or solvent}, T, p_1^\circ) = \int_{p_1^\circ}^{p_2^\circ} V_B^* \cdot dp, \quad (4)$$

where  $V_B^* = V_B^*(\text{cd}, T, p)$  is the molar volume of pure substance B.

(c) For a solute substance B in a solution,

$$\mu_B^\circ(\text{solute}, T, p_2^\circ) - \mu_B^\circ(\text{solute}, T, p_1^\circ) = \int_{p_1^\circ}^{p_2^\circ} V_B^\infty \cdot dp, \quad (5)$$

where  $V_B^\infty = V_B^\infty(\text{solute}, T, p)$  is the partial molar volume of the solute B at infinite dilution.

Consequently with  $p_1^\circ = 101\,325\text{ Pa}$  and  $p_2^\circ = 10^5\text{ Pa}$  the differences in standard thermodynamic quantities for condensed phases will nearly always be negligible compared with the uncertainties in current values, but for gases the differences, though small, might not be negligible. For example if  $p^\circ$  decreases from 101 325 Pa to  $10^5\text{ Pa}$  then, for a gaseous substance, the standard chemical potential decreases by about  $0.013 \cdot RT$  and the standard molar entropy increases by about  $0.013 \cdot R$ . The **standard molar enthalpy**  $H_B^\circ(\text{g}, T)$  of a gas B, and its derivatives with respect to temperature, do not depend in any way on the value chosen for the standard pressure  $p^\circ$ . Instances may occur when it is desirable to adopt other values for  $p^\circ$  but in any case the value adopted for  $p^\circ$  should *always* be stated when the numerical value of any standard thermodynamic quantity is given.

<sup>†</sup>**Editorial comment.** In *this Journal* no value for the standard pressure  $p^\circ$  is particularly recommended; 101 325 Pa is at present used at least as often as  $10^5\text{ Pa}$ . Other values may also be used, but the value used should always be declared.

Notwithstanding the above, the word “**normal**” in for example “normal boiling temperature” continues to mean the value at a pressure of 101 325 Pa.

The **standard molality** is denoted by  $m^\circ$  and for standard thermodynamic quantities the recommended value, which is used almost exclusively, is  $1 \text{ mol}\cdot\text{kg}^{-1}$ . Nevertheless, the value of  $m^\circ$  adopted should always be stated when the numerical value of any standard thermodynamic quantity of a solute is given. If the amount-of-substance concentration  $c_B$  is used in place of molality then the value adopted for the **standard amount-of-substance concentration**  $c^\circ$  should always be stated when the numerical value of any standard thermodynamic quantity of a solute is given; the recommended value for standard thermodynamic quantities is  $c^\circ = 1 \text{ mol}\cdot\text{dm}^{-3}$ .

In sections 3 to 5, definitions are given for the standard chemical potential of a substance B in the following states of aggregation: gas phase whether pure or mixed, condensed phase whether pure or mixed, solvent substance in a solution, and solute substance in a solution. The term “standard state” is often used to imply the values of the standard thermodynamic functions that are taken as a reference, but, for gaseous substances and for solute substances in solution, the standard chemical potential does not correspond to conditions that can be achieved experimentally.

However, the difference  $(\mu_B - \mu_B^\circ)$ , between the chemical potential of substance B and its standard chemical potential for the same temperature and state of aggregation, can always be calculated and measured using the definitions of sections 3 to 5. Standard methods may be used to determine more general differences in chemical potential.

Sections 3 to 5 also give definitions for additional quantities that represent differences in chemical potential that are especially important for a particular state of aggregation. Consequently, fugacities and virial coefficients are introduced in the discussion of gases so that the dependence of the chemical potential on pressure can be represented succinctly. For condensed phases and solutions (where the dependence of the chemical potential on composition is often more important than the dependence on pressure) activities, activity coefficients, and osmotic coefficients are introduced. All of these quantities may be determined experimentally or calculated from appropriate theories.

### 3. Gases and gaseous mixtures

The partial pressure  $p_B$  of a substance B in a gaseous mixture is a quantity with the same dimension as pressure defined by

$$p_B = y_B p, \quad (6)$$

where  $y_B$  is the mole fraction of B in the gaseous mixture and  $p$  is the pressure. Equation (6) is the definition of partial pressure recommended by IUPAC<sup>(1)</sup> and has the important property that  $\sum_B p_B = p$ . However, for a real gas,  $y_B p$  could not be obtained from the rate of change of momentum through collisions at a wall of known

area and, in this sense,  $y_B p$  is not a true pressure. To avoid ambiguity, partial pressure  $p_B$  is not used again in this document.

The standard chemical potential  $\mu_B^\circ$  of a gas B, whether a pure substance or a component of a gaseous mixture, is defined by

$$\mu_B^\circ(g, T) = \mu_B(g, T, p, y) - RT \cdot \ln(y_B p / p^\circ) - \int_0^p (V_B - RT/p) \cdot dp, \quad (7)$$

where  $y$  denotes the set of mole fractions  $y_B, y_C, \dots$  of B, C,  $\dots$ ,  $p^\circ$  is the standard pressure, and  $V_B = V_B(g, T, p, y)$  is the partial molar volume of B. The integral in equation (7) results from the pressure dependence of the chemical potential and is of central importance in the study of gases; for a perfect gas mixture the integrand in equation (7) is zero.

The fugacity  $f_B$  of the gas B in a gaseous mixture containing mole fractions  $y_B, y_C, \dots$  of the substances B, C,  $\dots$  is a quantity with the same dimension as pressure defined by

$$f_B = y_B p \cdot \exp \left[ \int_0^p \{ (V_B / RT) - (1/p) \} \cdot dp \right]. \quad (8)$$

The symbol  $\tilde{p}$  is sometimes used in place of  $f$  to denote fugacity especially when the discussion also involves the activity coefficients (see section 4) of substances in a liquid mixture.

It follows from the definitions of  $\mu_B^\circ$  and of  $f_B$  that the chemical potential  $\mu_B$  of a gas B in a gaseous mixture is given by

$$\mu_B(g, T, p, y) = \mu_B^\circ(g, T) + RT \cdot \ln(f_B / p^\circ). \quad (9)$$

The ratio  $f_B / y_B p$ , sometimes denoted by the symbol  $\phi_B$ , may be called the fugacity coefficient of the gas B. It follows from the definition of  $f_B$  that

$$\lim_{p \rightarrow 0} (f_B / y_B p) = \lim_{p \rightarrow 0} \phi_B = 1. \quad (10)$$

A perfect gas mixture of B, C,  $\dots$  is defined as one for which

$$pV = (n_B + n_C + \dots)RT.$$

Consequently,  $f_B = y_B p, f_C = y_C p, \dots$ , and  $\phi_B = \phi_C = \dots = 1$  for a perfect gas mixture.

The fugacity of a substance B in a liquid mixture may be defined by an equation similar to (8) but with liquid mole fraction  $x_B$  in place of gaseous mole fraction  $y_B$  and the partial molar volume  $V_B = V_B(l, T, p, x)$  of the substance B in the liquid mixture in place of  $V_B(g, T, p, y)$ .

When  $y_B = 1$ , the definitions given in this section for a gas B in a gaseous mixture reduce to those for a pure gas. Consequently for a pure gas B, the fugacity  $f_B^*$  is defined by

$$f_B^* = p \cdot \exp \left[ \int_0^p \{ (V_B^* / RT) - (1/p) \} \cdot dp \right], \quad (11)$$

where  $V_B^* = V_B^*(g, T, p)$  is the molar volume of the pure gas B, and the chemical potential  $\mu_B^*$  is given by

$$\mu_B^*(g, T, p) = \mu_B^\circ(g, T) + RT \cdot \ln(f_B^*/p^\circ). \quad (12)$$

It follows from the definitions of  $f_B^*$  that

$$\lim_{p \rightarrow 0} (f_B^*/p) = 1. \quad (13)$$

A pure perfect gas B is defined as one for which  $pV = n_B RT$ . Consequently  $f_B^* = p$  for a pure perfect gas B.

The fugacity of a pure liquid B may be defined by an equation similar to (11) but with the molar volume  $V_B^* = V_B^*(l, T, p)$  of the pure liquid B in place of  $V_B^*(g, T, p)$ .

It is often convenient to express the thermodynamic properties of a gas in terms of virial coefficients. At low to moderate pressures, the molar volume  $V_m$  of a gaseous mixture may be obtained from

$$V_m = (RT/p) + B(y) + \cdots, \quad (14)$$

where the second virial coefficient  $B(y)$  of the mixture is defined by

$$B(y) = \lim_{p \rightarrow 0} (V_m - RT/p). \quad (15)$$

For a binary mixture  $\{(1-y)A + yB\}$ , the second virial coefficient is given by

$$B(y) = (1-y)^2 B_{AA} + 2y(1-y) B_{AB} + y^2 B_{BB}, \quad (16)$$

where  $B_{AA}$  and  $B_{BB}$  are the second virial coefficients for pure A and B, and  $B_{AB}$  results from interactions of A with B. Alternatively,

$$B(y) = (1-y) B_{AA} + y B_{BB} + y(1-y) \delta_{AB}, \quad (17)$$

where

$$\delta_{AB} = 2B_{AB} - (B_{AA} + B_{BB}). \quad (18)$$

The partial molar volume  $V_B$  for component B in the mixture is given by

$$V_B = (RT/p) + B_{BB} + (1-y)^2 \delta_{AB} + \cdots, \quad (19)$$

and, since all virial coefficients are independent of pressure, the fugacity  $f_B$  and chemical potential  $\mu_B$  of B in the mixture are given by

$$\ln(f_B/y_B p) = [\{B_{BB} + (1-y)^2 \delta_{AB}\} p / RT] + \cdots, \quad (20)$$

$$\mu_B(g, T, p, y) = \mu_B^\circ(g, T) + RT \cdot \ln(y_B p / p^\circ) + p \{B_{BB} + (1-y)^2 \delta_{AB}\} + \cdots. \quad (21)$$

When  $y = 1$ , these equations reduce to those for a pure gas at low to moderate pressures. The use of some of these equations is illustrated by the following numerical example.

At 373.15 K values for the second virial coefficients of  $\text{CH}_3(\text{CH}_2)_2\text{CH}_3$  and

TABLE 1. Partial molar volumes  $V_B$ , fugacities  $f_B$ , and differences ( $\mu_B - \mu_B^\circ$ ) in chemical potential for  $\{0.8000\text{CH}_3(\text{CH}_2)_2\text{CH}_3 + 0.2000\text{CH}_3(\text{CH}_2)_4\text{CH}_3\}$  at a temperature of 373.15 K and a pressure of 100.000 kPa for  $p^\circ = 100.000$  kPa

$\frac{\delta_{AB}}{\text{cm}^3 \cdot \text{mol}^{-1}}$	B	$\frac{B_{BB}}{\text{cm}^3 \cdot \text{mol}^{-1}}$	$\frac{V_B - RT/p}{\text{cm}^3 \cdot \text{mol}^{-1}}$	$\frac{f_B}{\text{kPa}}$	$\frac{\mu_B - \mu_B^\circ}{\text{J} \cdot \text{mol}^{-1}}$
153	$\text{CH}_3(\text{CH}_2)_2\text{CH}_3$	-422	-416	78.935	-733.90
153	$\text{CH}_3(\text{CH}_2)_4\text{CH}_3$	-1031	-933	19.408	-5086.6
0	$\text{CH}_3(\text{CH}_2)_2\text{CH}_3$	-422	-422	78.919	-734.53
0	$\text{CH}_3(\text{CH}_2)_4\text{CH}_3$	-1031	-1031	19.346	-5096.5
0	$\text{CH}_3(\text{CH}_2)_2\text{CH}_3$	0	0	80.000	-692.32
0	$\text{CH}_3(\text{CH}_2)_4\text{CH}_3$	0	0	20.000	-4993.4

$\text{CH}_3(\text{CH}_2)_4\text{CH}_3$  are  $-422 \text{ cm}^3 \cdot \text{mol}^{-1}$  and  $-1031 \text{ cm}^3 \cdot \text{mol}^{-1}$  and, for mixtures of  $\text{CH}_3(\text{CH}_2)_2\text{CH}_3$  and  $\text{CH}_3(\text{CH}_2)_4\text{CH}_3$ ,  $\delta_{AB}$  is  $153 \text{ cm}^3 \cdot \text{mol}^{-1}$ .

For the mixture  $\{0.8000\text{CH}_3(\text{CH}_2)_2\text{CH}_3 + 0.2000\text{CH}_3(\text{CH}_2)_4\text{CH}_3\}$  at pressure  $p = 100.000$  kPa, equation (17) gives

$$\begin{aligned} B(y)/(\text{cm}^3 \cdot \text{mol}^{-1}) &= (-0.8000 \cdot 422 - 0.2000 \cdot 1031 + 0.2000 \cdot 0.8000 \cdot 153) \\ &= -519. \end{aligned}$$

For the butane in the mixture ( $y_A = 1 - y = 0.8000$ ), where A denotes  $\text{CH}_3(\text{CH}_2)_2\text{CH}_3$ , equations (19) and (8) give

$$(V_A - RT/p)/(\text{cm}^3 \cdot \text{mol}^{-1}) = (-422 + 0.2000^2 \cdot 153) = -416,$$

$$\frac{f_A}{\text{kPa}} = 0.8000 \cdot 100.000 \cdot \exp \left\{ \frac{-416 \cdot 10^{-6} \text{ m}^3 \cdot \text{mol}^{-1} \cdot 10^5 \text{ Pa}}{8.31451 \text{ J} \cdot \text{K}^{-1} \cdot \text{mol}^{-1} \cdot 373.15 \text{ K}} \right\}$$

$$= 80.000 \cdot \exp(-0.01341) = 78.935,$$

and, for  $p^\circ = 100.000$  kPa, equation (9) gives

$$\begin{aligned} \mu_A(\text{g}, 373.15 \text{ K}, 10^5 \text{ Pa}, 0.8000) - \mu_A^\circ(\text{g}, 373.15 \text{ K}) \\ = (8.31451 \text{ J} \cdot \text{K}^{-1} \cdot \text{mol}^{-1} \cdot 373.15 \text{ K}) \cdot \ln(78.935 \text{ kPa}/100.000 \text{ kPa}) \\ = -733.90 \text{ J} \cdot \text{mol}^{-1}. \end{aligned}$$

Similarly, for the hexane in the mixture ( $y_B = y = 0.2000$ ), where B denotes  $\text{CH}_3(\text{CH}_2)_4\text{CH}_3$ ,

$$(V_B - RT/p)/(\text{cm}^3 \cdot \text{mol}^{-1}) = (-1031 + 0.8000^2 \cdot 153) = -933,$$

$$\frac{f_B}{\text{kPa}} = 0.2000 \cdot 100.000 \cdot \exp \left\{ \frac{-933 \cdot 10^{-6} \text{ m}^3 \cdot \text{mol}^{-1} \cdot 10^5 \text{ Pa}}{8.31451 \text{ J} \cdot \text{K}^{-1} \cdot \text{mol}^{-1} \cdot 373.15 \text{ K}} \right\}$$

$$= 20.000 \cdot \exp(-0.03007) = 19.408,$$

and, for  $p^\circ = 100.000$  kPa,

$$\begin{aligned} \mu_B(\text{g}, 373.15 \text{ K}, 10^5 \text{ Pa}, 0.2000) - \mu_B^\circ(\text{g}, 373.15 \text{ K}) \\ = (8.31451 \text{ J} \cdot \text{K}^{-1} \cdot \text{mol}^{-1} \cdot 373.15 \text{ K}) \cdot \ln(19.408 \text{ kPa}/100.000 \text{ kPa}) \\ = -5086.6 \text{ J} \cdot \text{mol}^{-1}. \end{aligned}$$

The results of these calculations are summarized in table 1 together with those obtained when the mixture is treated as an ideal mixture of imperfect gases ( $\delta_{AB} = 0$ ) and when the mixture is treated as a perfect gas mixture ( $f_A = y_A p$  and  $f_B = y_B p$  or at these conditions  $B_{AA} = B_{BB} = \delta_{AB} = 0$ ).

#### 4. Pure condensed phases and mixed condensed phases

For a substance B in a condensed phase which is a pure substance or a component in a mixture the **standard chemical potential**  $\mu_B^\circ$  is defined by

$$\mu_B^\circ(\text{cd}, T) = \mu_B^*(\text{cd}, T, p) + \int_p^{p^\circ} V_B^* \cdot dp, \quad (22)$$

where  $p^\circ$  is the standard pressure and  $V_B^* = V_B^*(\text{cd}, T, p)$  is the molar volume of pure B. The integral in equation (22) represents the pressure dependence of the chemical potential and often, especially when  $p \approx p^\circ$ , the integral is small and the last term in equations (22) and (24) is neglected.

For example, if  $|p^\circ - p| < 100 \text{ kPa}$ , then the contribution of the integral to equations (22) and (24) is usually less than about  $10 \text{ J} \cdot \text{mol}^{-1}$ , which is small compared with  $(\mu_B - \mu_B^\circ) \approx -1700 \text{ J} \cdot \text{mol}^{-1}$  for an ideal binary mixture with  $x_B = 0.5$  at ambient temperature.

The definition (22) of the standard chemical potential of a substance B in a condensed phase is identical with the definition in section 5(b) of the standard chemical potential of a solvent substance A in a solution.

The **activity coefficient**  $f_B$  of a substance B in a liquid or solid mixture containing mole fractions  $x_B, x_C, \dots$  of the substances B, C,  $\dots$  is a quantity of dimension 1 defined in terms of the chemical potential  $\mu_B$  of B in the mixture by

$$RT \cdot \ln(x_B f_B) = \mu_B(\text{cd}, T, p, x) - \mu_B^*(\text{cd}, T, p), \quad (23)$$

where  $x$  denotes the set of mole fractions  $x_B, x_C, \dots$ .

It follows from the definitions of  $\mu_B^\circ$  and  $f_B$  that the **chemical potential**  $\mu_B$  of the substance B in the mixture is given by

$$\mu_B(\text{cd}, T, p, x) = \mu_B^\circ(\text{cd}, T) + RT \cdot \ln(x_B f_B) - \int_p^{p^\circ} V_B^* \cdot dp. \quad (24)$$

The **(relative) activity**  $a_B$  of a substance B in a liquid or solid mixture is a quantity of dimension 1 defined in terms of the chemical potential  $\mu_B$  by

$$RT \cdot \ln a_B = \mu_B - \mu_B^*. \quad (25)$$

It follows from the definitions of  $f_B$  and  $a_B$  that

$$a_B = x_B f_B, \quad (26)$$

and that

$$\lim_{x_B \rightarrow 1} a_B = 1. \quad (27)$$



A mixture of substances B, C,  $\dots$  is called an **ideal mixture** when  $a_B = x_B$ ,  $a_C = x_C$ ,  $\dots$  or  $f_B = 1$ ,  $f_C = 1$ ,  $\dots$ .

Additional quantities are sometimes used to describe the (gas + liquid) equilibrium of a mixture. The **Henry's law constant**  $k_{H,B}$  of a substance B in a liquid mixture is a quantity with the same dimension as pressure defined by

$$k_{H,B} = \lim_{x_B \rightarrow 0} (\tilde{p}_B/x_B) = (\partial \tilde{p}_B / \partial x_B)_{x_B=0}, \quad (28)$$

where  $\tilde{p}_B$  is the fugacity of the substance B in the *gaseous* mixture and  $x_B$  is the mole fraction of B in the *liquid* mixture; here the alternative<sup>(1)</sup> symbol  $\tilde{p}$  has been used to denote fugacity to avoid any confusion with the use of  $f$  to denote activity coefficient. It is important to note that Henry's law applies to a substance B at high dilution ( $x_B \rightarrow 0$ ) in a liquid mixture. By contrast, with the composition limit  $x_B \rightarrow 1$ , equation (28) becomes

$$\lim_{x_B \rightarrow 1} (\tilde{p}_B/x_B) = (\partial \tilde{p}_B / \partial x_B)_{x_B=1} = \tilde{p}_B^*, \quad (29)$$

where  $\tilde{p}_B^*$  is the fugacity of the pure gas B. Since equations (28) and (29) refer to a (gas + liquid) equilibrium, the fugacities  $\tilde{p}_B$  and  $\tilde{p}_B^*$  are evaluated at the vapour pressure of the mixture or of the pure fluid B. If the pure gas B is a perfect gas, then  $\tilde{p}_B = p_B^g$ , where  $p_B^g$  is the vapour pressure of pure B, and equation (29) becomes a statement of **Raoult's law**.

In general, the activities and activity coefficients defined in this section for a substance in a mixed condensed phase are functions of temperature, pressure, and composition.

## 5. Solutions

### (A) SOLUTE SUBSTANCE IN A SOLUTION

The **standard chemical potential**  $\mu_B^\circ$  of a solute substance B in a solution (especially a dilute liquid solution) of B, C,  $\dots$  in a solvent substance A is defined by

$$\mu_B^\circ(\text{solute}, T) = \{\mu_B(\text{solute}, T, p, m) - RT \cdot \ln(m_B/m^\circ)\}^\infty + \int_p^{p^\circ} V_B^\infty \cdot dp, \quad (30)$$

where  $m$  denotes the set of molalities  $m_B, m_C, \dots$  of the solutes B, C,  $\dots$  in the solvent A,  $m^\circ$  is the standard molality,  $^\infty$  denotes the limit at infinite dilution,  $p^\circ$  is the standard pressure, and  $V_B^\infty = V_B^\infty(\text{solute}, T, p)$  is the partial molar volume of the solute B at infinite dilution. In equation (30), both the terms inside  $\{\dots\}^\infty$  diverge as  $m_B \rightarrow 0$ , but their difference remains finite. The integral in equation (30) results from the pressure dependence of the chemical potential and often, especially when  $p \approx p^\circ$ , the integral is small and the last term in equations (30) and (32) is neglected.

For example, if  $|p^\circ - p| < 100 \text{ kPa}$ , then the contribution of the integral to equations (30) and (31) is usually less than about  $10 \text{ J} \cdot \text{mol}^{-1}$  which is small

compared with  $(\mu_B - \mu_B^\circ) \approx -5700 \text{ J} \cdot \text{mol}^{-1}$  for an ideal dilute solution with  $m/m^\circ = 0.1$  at ambient temperature.

The **activity coefficient**  $\gamma_B$  of a solute B in a solution (especially a dilute liquid solution) containing molalities  $m_B, m_C, \dots$  of solutes B, C,  $\dots$  in a solvent A is a quantity of dimension 1 defined in terms of the chemical potential  $\mu_B$  by

$$RT \cdot \ln(m_B \gamma_B / m^\circ) = \mu_B - \{\mu_B - RT \cdot \ln(m_B / m^\circ)\}^\infty. \quad (31)$$

It follows from the definitions of  $\mu_B^\circ$ (solute) and  $\gamma_B$  that the **chemical potential**  $\mu_B$  of the solute B in a solution containing molalities  $m_B, m_C, \dots$  of solutes B, C,  $\dots$  in a solvent A is given by

$$\mu_B(\text{solute}, T, p, m) = \mu_B^\circ(\text{solute}, T) + RT \cdot \ln(m_B \gamma_B / m^\circ) - \int_p^{p^\circ} V_B^\infty \cdot dp. \quad (32)$$

The name activity coefficient with the symbol  $\gamma_{c,B}$  may be used for the quantity defined in a way similar to  $\gamma_B$  but with amount-of-substance concentration  $c_B$  in place of molality  $m_B$  and a standard amount-of-substance concentration  $c^\circ$  in place of the standard molality  $m^\circ$ .

The **activity**  $a_B$  of a solute B in a solution (for example in a dilute liquid solution) containing molalities  $m_B, m_C, \dots$  of solutes B, C,  $\dots$  in a solvent A, is a quantity of dimension 1 defined in terms of the chemical potential  $\mu_B$  by

$$RT \cdot \ln a_B = \mu_B - \{\mu_B - RT \cdot \ln(m_B / m^\circ)\}^\infty. \quad (33)$$

The name relative activity is sometimes used instead of the name activity for this quantity.

It follows from the definitions of  $\gamma_B$  and  $a_B$  that

$$a_B = \gamma_B m_B / m^\circ, \quad (34)$$

and that

$$(a_B m^\circ / m_B)^\infty = 1. \quad (35)$$

The name activity with the symbol  $a_{c,B}$  may be used for the quantity defined in a way similar to  $a_B$  but with amount-of-substance concentration  $c_B$  in place of molality  $m_B$  and a standard amount-of-substance concentration  $c^\circ$  in place of the standard molality  $m^\circ$ .

A solution of solutes B, C,  $\dots$  in a solvent A is called an **ideal-dilute solution** when  $a_B = m_B / m^\circ$ ,  $a_C = m_C / m^\circ$ ,  $\dots$  or  $\gamma_B = 1$ ,  $\gamma_C = 1$ ,  $\dots$ .

## (B) SOLVENT SUBSTANCE IN A SOLUTION

The **standard chemical potential**  $\mu_A^\circ$  for the solvent A in a liquid or solid solution is defined by

$$\mu_A^\circ(\text{solvent}, T) = \mu_A^*(\text{cd}, T, p) + \int_p^{p^\circ} V_A^* \cdot dp, \quad (36)$$

where  $p^\circ$  is the standard pressure and  $V_A^* = V_A^*(\text{cd}, T, p)$  is the molar volume of pure A. The integral in equation (36) represents the pressure dependence of the chemical potential and often, especially when  $p \approx p^\circ$ , the integral is small and the last term in equation (36) is neglected. This definition of the standard chemical potential of the solvent A in a solution is identical with the definition in section 4 of the standard chemical potential of any substance B in a condensed phase.

The **osmotic coefficient**  $\phi$  of the solvent A in a solution (especially a dilute liquid solution) containing molalities  $m_B, m_C, \dots$  of solutes B, C,  $\dots$  is a quantity of dimension 1 defined in terms of chemical potential  $\mu_A$  of the solvent A in the solution by

$$\phi = (\mu_A^* - \mu_A)/(RTM_A \cdot \Sigma_B m_B), \quad (37)$$

where  $\mu_A^*$  is the chemical potential of the pure solvent A at the same temperature and pressure, the sum over the solutes B, C,  $\dots$  is denoted by  $\Sigma_B$ , and  $M_A$  is the molar mass of the solvent A.

For an **ideal dilute solution** as defined in sections 5(a) it can be shown that  $\phi = 1$ .

It follows from the definitions of  $\mu_A^\circ$  and  $\phi$  that the **chemical potential**  $\mu_A$  of the solvent A in the solution is given by

$$\mu_A(\text{solvent}, T, p, m) = \mu_A^\circ(\text{solvent}, T) - RT\phi M_A \cdot \Sigma_B m_B - \int_p^{p^\circ} V_A^* \cdot dp, \quad (38)$$

where  $m$  denotes the set of molalities  $m_B, m_C, \dots$ . Often, especially when  $p \approx p^\circ$ , the integral in equation (38) is small and the last term is neglected. However, in practice the integral is not always negligible compared with  $RT\phi M_A \cdot \Sigma_B m_B$ , although  $|\mu_A - \mu_A^\circ|$  for the solvent substance in solution is usually very much less than  $|\mu_B - \mu_B^\circ|$  for the solutes (especially for an electrolyte solute) in the same solution.

For example, if  $p^\circ = 100 \text{ kPa}$  and  $p \ll p^\circ$ , as is the case for isopiestic measurements with aqueous solutions near ambient temperature, then the contribution of the integral in equations (36) and (38) is about  $2 \text{ J} \cdot \text{mol}^{-1}$  which is comparable with  $(\mu_A - \mu_A^\circ) \approx -4.5 \text{ J} \cdot \text{mol}^{-1}$  but small compared with  $(\mu_B - \mu_B^\circ) \approx -5700 \text{ J} \cdot \text{mol}^{-1}$  for an ideal dilute solution with  $\Sigma_B m_B = 0.1 \text{ mol} \cdot \text{kg}^{-1}$ .

The **activity**  $a_A$  of the solvent A in a solution (especially a dilute liquid solution) containing molalities  $m_B, m_C, \dots$ , of solutes B, C,  $\dots$  is a quantity of dimension 1 defined in terms of the chemical potential  $\mu_A$  of the solvent A in the solution by

$$RT \cdot \ln a_A = \mu_A - \mu_A^* = -RT\phi M_A \cdot \Sigma_B m_B. \quad (39)$$

This definition of the activity of the solvent A in a solution is identical with the definition in section 4 of the activity of any substance in a mixture.

In general, the activities, activity coefficients, and osmotic coefficients defined in this section for solutes and solvents are functions of temperature, pressure, and composition.

## 6. Composition variables and activity coefficients

Sometimes the same mixed condensed phase will have been studied by one group of workers as a mixture and by another group as a solution, and the results must be adjusted before the different measurements can be compared. The mole fraction  $x_A$  of the solvent A in a solution containing molalities  $m_B = n_B/n_A M_A$ ,  $m_C, \dots$  of the solutes B, C,  $\dots$  is given by

$$x_A = 1/(1 + M_A \cdot \sum_B m_B), \quad (40)$$

where  $M_A$  is the molar mass of the solvent A. The mole fraction  $x_B$ , molality  $m_B$ , and amount-of-substance concentration  $c_B$  of the (solute) substance B are related to each other by

$$x_B = m_B M_A / (1 + M_A \cdot \sum_B m_B), \quad (41)$$

$$m_B = x_B / x_A M_A, \quad (42)$$

and

$$c_B / m_B = \rho / (1 + \sum_B m_B M_B) = \rho - \sum_B c_B M_B, \quad (43)$$

where  $M_B$  is the molar mass of the solute B and  $\rho$  is the density of the solution.

For a given mixed condensed phase, the activity coefficient  $f_B$  of the substance B with mole fraction  $x_B$  in the mixture, the activity coefficient  $\gamma_B$  of the solute B with molality  $m_B$  in the solution, and the activity coefficient  $\gamma_{c,B}$  of the solute B with amount-of-substance concentration  $c_B$  in the solution will, in general, be different. These activity coefficients are related to each other by

$$f_B / f_B^\infty = \gamma_B (1 + M_A \cdot \sum_B m_B) = \gamma_{c,B} (\rho + M_A \cdot \sum_B c_B - \sum_B c_B M_B) / \rho_A^*, \quad (44)$$

$$\gamma_B = (f_B / f_B^\infty) x_A, \quad (45)$$

and

$$\gamma_B = \gamma_{c,B} (\rho - \sum_B c_B M_B) / \rho_A^* = \gamma_{c,B} (c_B / m_B \rho_A^*), \quad (46)$$

where  $\rho_A^*$  is the density of the pure (solvent) substance A and  $f_B^\infty$  denotes the activity coefficient of the substance B in the mixture at infinite dilution when  $x_A \rightarrow 1$ . Note that whereas  $\gamma_B^\infty = 1$ , in general  $f_B^\infty \neq 1$ ; this is a direct consequence of the differences between the definition of the standard chemical potential of a solute in a solution and the definition of the standard chemical potential of a substance in a mixture. However, the chemical potential  $\mu_B$  of the substance B in a mixed condensed phase is the same whether the phase is treated as a mixture with the composition described in terms of mole fraction, or as a solution with the composition described in terms of molality, or as a solution with the composition described in terms of amount-of-substance concentration.

## 7. Equilibrium constant

The standard equilibrium constant  $K^\circ$  for a chemical reaction is defined by

$$RT \cdot \ln\{K^\circ(T)\} = -\sum_B \nu_B \mu_B^\circ(z, T) = -\Delta_r G_m^\circ(T), \quad (47)$$

where  $z$  denotes the state of aggregation and  $\nu_B$  denotes the stoichiometric number of the substance  $B$  in the chemical reaction:

$$0 = \sum_B \nu_B B. \quad (48)$$

$\Delta_r G_m^\circ$  is the change in the standard molar Gibbs energy due to the reaction. It follows from this definition that  $K^\circ$  and  $\Delta_r G_m^\circ$  are, like all standard thermodynamic quantities, independent of pressure and composition, and that  $K^\circ$  is a quantity of dimension 1.

The Green Book<sup>(1)</sup> recommends the symbol  $K^\circ$ ,  $K^\ominus$ , or  $K$  for the equilibrium constant defined by equation (47); the name “thermodynamic equilibrium constant” with the symbol  $K$  has also been used by IUPAC.<sup>(3)</sup> ISO 31-8 recommends<sup>(2)</sup> the name “standard equilibrium constant” with the symbol  $K^\circ$  or  $K^\ominus$ . In this document the name **standard equilibrium constant** with symbol  $K^\circ$  ( $K^\ominus$  would be an equivalent symbol) has been used for the equilibrium constant defined by equation (47) to distinguish it from other equilibrium constants that are often used in practice although they are not constant at fixed temperature and need not be of dimension 1. In the following paragraphs, definitions are given for some of these equilibrium constants; they have been given the symbol  $K$  and their dependence on temperature, pressure, and composition has been explicitly indicated; where appropriate, the conditions under which  $K$  becomes constant at constant temperature and becomes a quantity of dimension 1 (that is the conditions under which  $K$  approaches  $K^\circ$ ) are noted. It is assumed that all quantities refer to the state of chemical equilibrium.

The **standard equilibrium constant for a reaction in the gas phase** is given by

$$K^\circ(g, T) = \Pi_B (f_B/p^\circ)^{\nu_B} \quad (49)$$

$$= \lim_{p \rightarrow 0} \Pi_B (y_B p/p^\circ)^{\nu_B}. \quad (50)$$

Equation (50) provides the most direct method of determining  $K^\circ(g, T)$  experimentally and is useful when, as is often the case, the fugacity  $f_B$  is not known for each gas  $B$  in the reaction mixture. In many instances the extrapolation to zero pressure is not carried out and, at least at low to moderate pressures, is not necessary. In practice the quantity

$$K_p(g, T, p, y) = \Pi_B (y_B p)^{\nu_B}, \quad (51)$$

which has the dimension of (pressure) $^{\sum_B \nu_B}$ , is often used. At constant temperature,  $K_p$  may be approximately constant, especially at low pressures where

$$\lim_{p \rightarrow 0} \{K_p(g, T, p, y)\} = K_f(g, T) = \Pi_B (f_B)^{\nu_B}, \quad (52)$$

which is independent of pressure but need not equal  $K^\circ$ .

The **standard equilibrium constant for a reaction in a condensed phase** is given by

$$K^\circ(\text{cd}, T) = \{\Pi_B(x_B f_B)^{v_B}\} \cdot \exp\left\{-\int_p^{p^\circ} \Sigma_B(v_B V_B^*/RT) \cdot dp\right\}, \quad (53)$$

where  $x_B$  and  $f_B$  are the mole fraction and activity coefficient of the substance B in the reaction mixture at equilibrium;  $V_B^*$  is the molar volume of pure substance B.

Often, especially when  $p \approx p^\circ$ , the pressure integral is small and the quantity

$$K_x(\text{cd}, T, p) = \Pi_B(x_B f_B)^{v_B}, \quad (54)$$

is a good approximation to  $K^\circ(\text{cd}, T)$ . If it is assumed that the reaction mixture is an ideal mixture then  $K_x$  reduces to

$$K_x(\text{cd}, T, p, x) = \Pi_B(x_B)^{v_B}, \quad (55)$$

which is often used in practice. However, since few real mixtures are even approximately ideal,  $K_x$  usually varies with composition and is often a poor approximation to  $K^\circ(\text{cd}, T)$ .

The **standard equilibrium constant for a reaction in a solution** of solutes B, C,  $\dots$  in a solvent A is given by

$$K^\circ(\text{solution}, T) = (a_A)^{v_A} \cdot \Pi_B(a_B)^{v_B} \cdot \exp\left\{-\int_p^{p^\circ} (v_A V_A^* + \Sigma_B v_B V_B^\infty)/RT\right\} \cdot dp, \quad (56)$$

where  $a$  denotes activity,  $V_A^*$  is the molar volume of the pure solvent, and  $V_B^\infty$  is the partial molar volume of the solute B at infinite dilution. Often, especially when  $p \approx p^\circ$  the pressure integral is small and the quantity

$$K_a(\text{solution}, T, p) = (a_A)^{v_A} \cdot \Pi_B(a_B)^{v_B} \quad (57)$$

$$= \exp(-v_A \phi M_A \Sigma_B m_B) \cdot \Pi_B(m_B \gamma_B / m^\circ)^{v_B}, \quad (58)$$

varies with pressure but may still be a good approximation to  $K^\circ$ . If the solvent factor  $\exp(-v_A \phi M_A \Sigma_B m_B)$ , which is often close to unity, is neglected then the resulting quantity

$$K_{m\gamma}(\text{solution}, T, p, m) = \Pi_B(m_B \gamma_B / m^\circ)^{v_B}, \quad (59)$$

varies with composition but may still be a good approximation to  $K^\circ$ . The quantity

$$K_m(\text{solution}, T, p, m) = \Pi_B(m_B)^{v_B}, \quad (60)$$

which has dimension  $(\text{molality})^{\Sigma_B v_B}$ , is often used in practice. Although  $K_m$  may be approximately constant, especially in a dilute solution at  $p \approx p^\circ$ , it need not equal  $K^\circ$ .

By contrast

$$K^\circ(\text{solution}, T) = \lim_{\Sigma_B m_B \rightarrow 0} \Pi_B(m_B / m^\circ)^{v_B}, \quad (p = p^\circ). \quad (61)$$

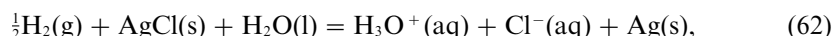
This equation provides the most direct method of determining  $K^\circ(\text{solution}, T)$  experimentally (provided  $p = p^\circ$ ) and is useful when, as is often the case, the osmotic

coefficient  $\phi$  for the solvent and the activity coefficients  $\gamma_B, \gamma_C, \dots$  for the solutes are not known.

If amount-of-substance concentration  $c_B$  is used in place of molality  $m_B$  with a standard amount-of-substance concentration  $c^\circ$  in place of the standard molality  $m^\circ$ , then equations (57) and (59) to (61) are modified by writing  $a_{c,B}$  in place of  $a_B$  and  $\gamma_{c,B}$  in place of  $\gamma_B$ ; as indicated in sections 4 and 5(b), the activity  $a_A$  of the solvent is not altered by changing the variable used to describe the composition of a mixture or solution. Since amount-of-substance concentration, unlike molality, varies with both temperature and pressure, it is not possible to define a standard equilibrium constant  $K^\circ(\text{solution}, T)$  for a reaction in solution when the composition of the solution is described in terms of concentration.

In practice, many chemical reactions involve more than one phase, and the expressions for the standard equilibrium constant are necessarily more complicated than those given above for reactions in a single phase. Since it is not possible to consider all combinations, the general procedure is illustrated here by a single example.

The chemical reaction:



involves four phases and may be studied using a galvanic cell. The standard equilibrium constant contains a term from equation (49) for the gas ( $\text{H}_2$ ), a term from equation (53) for each of the pure solid phases ( $\text{Ag}$  and  $\text{AgCl}$ ), and terms from equation (56) for the solvent ( $\text{H}_2\text{O}$ ) and solute ( $\text{HCl}$ ). The complete expression is

$$K^\circ(T) = \{(a_{\text{HCl}}/a_{\text{H}_2\text{O}})(p^\circ/f_{\text{H}_2})^{1/2}\} \cdot \exp \left[ - \int_p^{p^\circ} \{(V_{\text{HCl}}^\infty - V_{\text{H}_2\text{O}}^* + V_{\text{Ag}}^* - V_{\text{AgCl}}^*)/RT\} \cdot dp \right],$$

and  $K^\circ(298.15 \text{ K}) \approx 5694$  when  $p^\circ = 100 \text{ kPa}$ . The contributions of the various factors to  $K^\circ$  are considered in some detail below; for this reaction, a change of 0.01 mV in the electromotive force (e.m.f.) of the galvanic cell corresponds to 0.04 per cent in  $K^\circ$ .

If the volumes of the condensed phases are independent of pressure, then the contribution of the integral to  $K^\circ$  is a factor

$$\exp\{-(p^\circ - p)(V_{\text{HCl}}^\infty - V_{\text{H}_2\text{O}}^* + V_{\text{Ag}}^* - V_{\text{AgCl}}^*)/RT\} \approx 1 + 6 \cdot 10^{-6}(p^\circ - p)/\text{kPa},$$

which is often negligible.

At the temperature 298.15 K, the second virial coefficient of  $\text{H}_2$  is about  $15 \text{ cm}^3 \cdot \text{mol}^{-1}$  and, using equation (20), the fugacity coefficient  $\phi_{\text{H}_2}$  for  $\text{H}_2$  is

$$\phi_{\text{H}_2} = f_{\text{H}_2}/(y_{\text{H}_2}p) = 1 + 6 \cdot 10^{-6}(p/\text{kPa}).$$

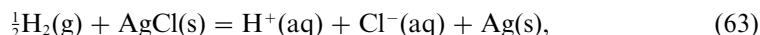
Consequently, the gas imperfections contribute a factor 1.0003 to  $K^\circ$  at a total pressure  $p = 100 \text{ kPa}$ . While this factor may be neglected for many measurements, the vapour pressure of the solution is close to that of pure water (about 3.2 kPa at 298.15 K) so  $y_{\text{H}_2} \approx 0.968$  at  $p = 100 \text{ kPa}$  and, when  $p^\circ = 100 \text{ kPa}$ ,

$$(f_{\text{H}_2}/p^\circ)^{1/2} \approx (y_{\text{H}_2}p/p^\circ)^{1/2} \approx 1 - 0.016,$$

which corresponds to 0.4 mV in the e.m.f. of the galvanic cell.

The activity of  $a_{\text{H}_2\text{O}}$  of  $\text{H}_2\text{O}$  differs from unity by about  $M_{\text{A}} \cdot \Sigma_{\text{B}} m_{\text{B}} = 2M_{\text{A}}m_{\text{HCl}}$  and contributes more than 0.04 per cent to  $K^\circ$  (or more than 0.01 mV to the e.m.f.) for molalities  $m_{\text{HCl}}$  of HCl greater than about  $0.01 \text{ mol} \cdot \text{kg}^{-1}$ . Consequently,  $a_{\text{H}_2\text{O}}$  should be considered for precise work except at high dilution.

Alternatively, the chemical reaction may be written



and the solvent no longer enters the expression for  $K^\circ$ . It should be noted that for a given solution,  $m(\text{H}^+) \neq m(\text{H}_3\text{O}^+)$  and that  $K^\circ$  is different for reactions (62) and (63).

The above discussion indicates that for many purposes the gas imperfections and the pressure dependence of the chemical potentials for the condensed phases may be neglected. Consequently, for the reaction (63), which also avoids  $a_{\text{H}_2\text{O}}$ , the standard equilibrium constant may be written

$$\begin{aligned} K^\circ(T) &= a_{\text{HCl}} \{p^\circ / (y_{\text{H}_2} p)\}^{1/2} = \{m_{\text{H}^+} m_{\text{Cl}^-} / (m^\circ)^2\} (\gamma_{\text{H}^+} \gamma_{\text{Cl}^-}) \{p^\circ / y_{\text{H}_2} p\}^{1/2} \\ &= (m_{\pm} \gamma_{\pm} / m^\circ)^2 \{p^\circ / y_{\text{H}_2} p\}^{1/2}, \end{aligned}$$

where  $m_{\pm}$  and  $\gamma_{\pm}$  are the mean ionic molality and mean ionic activity coefficient for HCl in the aqueous solution. As with most electrolyte solutions, the activity coefficients of HCl can be neglected only for very approximate calculations. For example,  $\gamma_{\pm} = 0.796$  for an aqueous HCl solution of molality  $0.1 \text{ mol} \cdot \text{kg}^{-1}$  and an equilibrium constant calculated with  $\gamma_{\pm} = 1$  differs from  $K^\circ$  by almost 40 per cent; even when  $m_{\pm} = 10^{-4} \text{ mol} \cdot \text{kg}^{-1}$ ,  $\gamma_{\pm}$  accounts for about 2.3 per cent of  $K^\circ$ .

## REFERENCES

1. IUPAC-Physical Chemistry Division: *Quantities, Units and Symbols in Physical Chemistry*. Blackwell Scientific Publications: Oxford. First edition **1988**, and second edition **1993**.
2. International Standard: *Quantities and units — Part 8: Physical chemistry and molecular physics*. ISO 31-8. **1992**.
3. IUPAC-Physical Chemistry Division: *Manual of Symbols and Terminology for Physicochemical Quantities and Units. Appendix IV — Notation of states and processes, significance of the word standard in chemical thermodynamics and remarks on commonly tabulated forms of thermodynamic functions*. *Pure Appl. Chem.* **1982**, 54, 1239–1250.