

The computation aspects of the equivalent-layer technique: review and perspective

Diego Takahashi 1,* , André L. A. Reis 2 , Vanderlei C. Oliveira Jr. 1 and Valéria C. F. Barbosa 1

Correspondence*: Valéria C.F. Barbosa valcris@on.br

¹Observatório Nacional, Department of Geophysics, Rio de Janeiro, Brasil

² Universidade do Estado do Rio de Janeiro, Department of Applied Geology, Rio de Janeiro, Brasil

1 FUNDAMENTALS

- 2 Let d be a $D \times 1$ vector, whose i-th element d_i is the observed potential field at the position (x_i, y_i, z_i) ,
- 3 $i \in \{1:D\}$, of a topocentric Cartesian system with x, y and z axes pointing to north, east and down,
- 4 respectively. Consider that d_i can be satisfactorily approximated by a harmonic function

$$f_i = \sum_{j=1}^{P} g_{ij} p_j , \quad i \in \{1 : D\} ,$$
 (1)

- 5 where, p_i represents the scalar physical property of a virtual source (i.e., monopole, dipole, prism) located
- 6 at $(x_j, y_j, z_j), j \in \{1 : P\}$ and

$$g_{ij} \equiv g(x_i - x_j, y_i - y_j, z_i - z_j), \quad z_i < \min\{z_j\}, \quad \forall i \in \{1 : D\},$$
 (2)

- 7 is a harmonic function, where $\min\{z_j\}$ denotes the minimum z_j , or the vertical coordinate of the shallowest
- 8 virtual source. These virtual sources are called equivalent sources and they form an equivalent layer. In
- 9 matrix notation, the potential field produced by all equivalent sources at all points (x_i, y_i, z_i) , $i \in \{1 : D\}$,
- 10 is given by:

$$\mathbf{f} = \mathbf{G}\mathbf{p} \,, \tag{3}$$

- where \mathbf{p} is a $P \times 1$ vector with j-th element p_j representing the scalar physical property of the j-th
- 12 equivalent source and G is a $D \times P$ matrix with element g_{ij} given by equation 2.
- 13 The equivalent-layer technique consists in solving a linear inverse problem to determine a parameter
- 14 vector p leading to a predicted data vector f (equation 3) sufficiently close to the observed data vector d,
- whose i-th element d_i is the observed potential field at (x_i, y_i, z_i) . The notion of *closeness* is intrinsically
- related to the concept of vector norm (e.g., Golub and Van Loan, 2013, p. 68) or measure of length (e.g.,
- 17 Menke, 2018, p. 41). Because of that, almost all methods for determining p actually estimate a parameter
- vector $\tilde{\mathbf{p}}$ minimizing a length measure of the difference between \mathbf{f} and \mathbf{d} (see subsection 2.1). Given an
- 19 estimate \tilde{p} , it is then possible to compute a potential field transformation

$$\mathbf{t} = \mathbf{A}\tilde{\mathbf{p}} \,, \tag{4}$$

where t is a $T \times 1$ vector with k-th element t_k representing the transformed potential field at the position

21
$$(x_k, y_y, z_k), k \in \{1:T\}$$
, and

$$a_{kj} \equiv a(x_k - x_j, y_k - y_j, z_k - z_j), \quad z_k < \min\{z_j\}, \quad \forall k \in \{1:T\},$$
 (5)

22 is a harmonic function representing the kj-th element of the $T \times P$ matrix A.

23 1.1 Spatial distribution and total number of equivalent sources

- There is no well-established criteria to define the optimum number P or the spatial distribution of the
- 25 equivalent sources. We know that setting an equivalent layer with more (less) sources than potential-field
- 26 data usually leads to an underdetermined (overdetermined) inverse problem (e.g., Menke, 2018, p. 52–53).
- 27 Concerning the spatial distribution of the equivalent sources, the only condition is that they must rely on a
- 28 surface that is located below and does not cross that containing the potential field data. Soler and Uieda
- 29 (2021) present a practical discussion about this topic.

From a theoretical point of view, the equivalent layer reproducing a given potential field data set cannot cross the true gravity or magnetic sources. This condition is a consequence of recognizing that the equivalent layer is essentially an indirect solution of a boundary value problem of potential theory (e.g., Roy, 1962; Zidarov, 1965; Dampney, 1969; Li et al., 2014; Reis et al., 2020). In practical applications, however, there is no guarantee that this condition is satisfied. Actually, its is widely known from practical experience (e.g., Gonzalez et al., 2022) that the equivalent-layer technique works even for the case in which the layer cross the true sources.

Regarding the depth of the equivalent layer, Dampney (1969) proposed a criterion based on horizontal data sampling, suggesting that the equivalent-layer depth should be between two and six times the horizontal grid spacing, considering evenly spaced data. However, when dealing with a survey pattern that has unevenly spaced data, Reis et al. (2020) adopted an alternative empirical criterion. According to their proposal, the depth of the equivalent layer should range from two to three times the spacing between adjacent flight lines. The criteria of Dampney (1969) and Reis et al. (2020) are valid for planar equivalent layers. Cordell (1992) have proposed and an alternative criterion for scattered data that leads to an undulating equivalent layer. This criterion have been slightly modified by Guspí et al. (2004), Guspí and Novara (2009) and Soler and Uieda (2021), for example, and consists in setting one equivalent source below each datum at a depth proportional to the horizontal distance to the nearest neighboring data points. Soler and Uieda (2021) have compared different strategies for defining the equivalent sources depth for the specific problem of interpolating gravity data, but they have not found significant differences between them.

49 **1.2 Matrix** G

37

38

39 40

41

42 43

44

45

46

47

48

Generally, the harmonic function g_{ij} (equation 2) is defined in terms of the inverse distance between the observation point (x_i, y_i, z_i) and the j-th equivalent source at (x_j, y_j, z_j) ,

$$\frac{1}{r_{ij}} \equiv \frac{1}{\sqrt{(x_i - x_j)^2 + (y_i - y_j)^2 + (z_i - z_j)^2}},$$
(6)

52 or by its partial derivatives of first and second orders, respectively given by

$$\partial_{\alpha} \frac{1}{r_{ij}} \equiv \frac{-(\alpha_i - \alpha_j)}{r_{ij}^3} , \quad \alpha \in \{x, y, z\} , \tag{7}$$

53 and

$$\partial_{\alpha\beta} \frac{1}{r_{ij}} \equiv \begin{cases} \frac{3(\alpha_i - \alpha_j)^2}{r_{ij}^5} , & \alpha = \beta ,\\ \frac{3(\alpha_i - \alpha_j)(\beta_i - \beta_j)}{r_{ij}^5} - \frac{1}{r_{ij}^3} , & \alpha \neq \beta , \end{cases} \quad \alpha, \beta \in \{x, y, z\} . \tag{8}$$

In this case, the equivalent layer is formed by punctual sources representing monopoles or dipoles (e.g., Dampney, 1969; Emilia, 1973; Leão and Silva, 1989; Cordell, 1992; Oliveira Jr. et al., 2013; Siqueira et al., 55 2017; Reis et al., 2020; Takahashi et al., 2020; Soler and Uieda, 2021; Takahashi et al., 2022). Another 56 common approach consists in not defining g_{ij} by using equations 6–8, but other harmonic functions 57 58 obtained by integrating them over the volume of regular prisms (e.g., Li and Oldenburg, 2010; Barnes and Lumley, 2011; Li et al., 2014; Jirigalatu and Ebbing, 2019). There are also some less common approaches 59 defining the harmonic function g_{ij} (equation 2) as the potential field due to plane faces with constant physical property (Hansen and Miyazaki, 1984), doublets (Silva, 1986) or by computing the double 61 62 integration of the inverse distance function with respect to z (Guspí and Novara, 2009).

A common assumption for most of the equivalent-layer methods is that the harmonic function g_{ij} 63 64 (equation 2) is independent on the actual physical relationship between the observed potential field and their true sources (e.g., Cordell, 1992; Guspí and Novara, 2009; Li et al., 2014). Hence, g_{ij} can be 65 defined according to the problem. The only condition imposed to this function is that it decays to zero 66 as the observation point (x_i, y_i, z_i) goes away from the position (x_j, y_i, z_j) of the j-th equivalent source. 67 However, several methods use a function g_{ij} that preserves the physical relationship between the observed 68 69 potential field and their true sources. For the case in which the observed potential field is gravity data, g_{ij} 70 is commonly defined as a component of the gravitational field produced at (x_i, y_i, z_i) by a point mass or 71 prism located at (x_j, y_j, z_j) , with unit density. On the other hand, g_{ij} is commonly defined as a component of the magnetic induction field produced at (x_i, y_i, z_i) by a dipole or prism located at (x_i, y_i, z_i) , with unit magnetization intensity, when the observed potential field is magnetic data. 73

74 The main challenge in the equivalent-layer technique is the computational complexity associated with 75 handling large datasets. This complexity arises because the sensitivity matrix G (equation 3) is dense 76 regardless of the harmonic function g_{ij} (equation 2) employed. In the case of scattered potential-field 77 data, the structure of G is not well-defined, regardless of the spatial distribution of the equivalent sources. However, in a specific scenario where (i) each potential-field datum is directly associated with a single 78 79 equivalent source located directly below it, and (ii) both the data and sources are based on planar and regularly spaced grids, Takahashi et al. (2020, 2022) demonstrate that G exhibits a block-Toeplitz Toeplitz-80 81 block (BTTB) structure. In such cases, the product of G and an arbitrary vector can be efficiently computed using a 2D fast Fourier transform as a discrete convolution.

2 LINEAR INVERSE PROBLEM OF EQUIVALENT-LAYER TECHNIQUE

83 2.1 General formulation

A general formulation for almost all equivalent-layer methods can be achieved by first considering that the $P \times 1$ parameter vector \mathbf{p} (equation 3) can be reparameterized into a $Q \times 1$ vector \mathbf{q} according to:

$$\mathbf{p} = \mathbf{H} \mathbf{q} \,, \tag{9}$$

where **H** is a $P \times Q$ matrix. The predicted data vector **f** (equation 3) can then be rewritten as follows:

$$\mathbf{f} = \mathbf{G} \,\mathbf{H} \,\mathbf{q} \,. \tag{10}$$

- Note that the original parameter vector \mathbf{p} is defined in a P-dimensional space whereas the reparameterized parameter vector \mathbf{q} (equation 9) lies in a Q-dimensional space. For convenience, we use the terms P-space
- 89 and Q-space to designate them.
- In this case, the problem of estimating a parameter vector $\tilde{\mathbf{p}}$ minimizing a length measure of the difference
- 91 between f (equation 3) and d is replaced by that of estimating an auxiliary vector \tilde{q} minimizing the goal
- 92 function

$$\Gamma(\mathbf{q}) = \Phi(\mathbf{q}) + \mu \,\Theta(\mathbf{q}) \,, \tag{11}$$

93 which is a combination of particular measures of length given by

$$\Phi(\mathbf{q}) = (\mathbf{d} - \mathbf{f})^{\top} \mathbf{W}_d (\mathbf{d} - \mathbf{f}) , \qquad (12)$$

94 and

$$\Theta(\mathbf{q}) = (\mathbf{q} - \bar{\mathbf{q}})^{\top} \mathbf{W}_{a} (\mathbf{q} - \bar{\mathbf{q}}) , \qquad (13)$$

- 95 where the regularization parameter μ is a positive scalar controlling the trade-off between the data-misfit
- 96 function $\Phi(\mathbf{q})$ and the regularization function $\Theta(\mathbf{q})$; \mathbf{W}_d is a $D \times D$ symmetric matrix defining the relative
- 97 importance of each observed datum d_i ; \mathbf{W}_q is a $Q \times Q$ symmetric matrix imposing prior information on \mathbf{q} ;
- 98 and $\bar{\bf q}$ is a $Q \times 1$ vector of reference values for $\bf q$ that satisfies

$$\bar{\mathbf{p}} = \mathbf{H}\,\bar{\mathbf{q}}\,,\tag{14}$$

- 99 where $\bar{\mathbf{p}}$ is a $P \times 1$ vector containing reference values for the original parameter vector \mathbf{p} .
- After obtaining an estimate $\tilde{\mathbf{q}}$ for the reparameterized parameter vector \mathbf{q} (equation 9), the estimate $\tilde{\mathbf{p}}$ for
- 101 the original parameter vector (equation 3) is computed by

$$\tilde{\mathbf{p}} = \mathbf{H}\,\tilde{\mathbf{q}}\,. \tag{15}$$

The reparameterized vector $\tilde{\mathbf{q}}$ is obtained by first computing the gradient of $\Gamma(\mathbf{q})$,

$$\nabla\Gamma(\mathbf{q}) = -2\mathbf{H}^{\mathsf{T}}\mathbf{G}^{\mathsf{T}}\mathbf{W}_{d}(\mathbf{d} - \mathbf{f}) + 2\mu\mathbf{W}_{q}(\mathbf{q} - \bar{\mathbf{q}}).$$
(16)

- 103 Then, by considering that $\nabla\Gamma(\tilde{\mathbf{q}})=\mathbf{0}$ (equation 16), where $\mathbf{0}$ is a vector of zeros, as well as adding and
- 104 subtracting the term $(\mathbf{H}^{\top}\mathbf{G}^{\top}\mathbf{W}_{d}\mathbf{G}\,\mathbf{H})\,\bar{\mathbf{q}}$, we obtain

$$\tilde{\boldsymbol{\delta}}_{q} = \mathbf{B}\,\boldsymbol{\delta}_{d}\,,\tag{17}$$

105 where

$$\tilde{\mathbf{q}} = \tilde{\boldsymbol{\delta}}_q + \bar{\mathbf{q}} \,, \tag{18}$$

106

$$\boldsymbol{\delta}_d = \mathbf{d} - \mathbf{G} \,\mathbf{H} \,\bar{\mathbf{q}} \,, \tag{19}$$

107

$$\mathbf{B} = \left(\mathbf{H}^{\mathsf{T}} \mathbf{G}^{\mathsf{T}} \mathbf{W}_{d} \mathbf{G} \mathbf{H} + \mu \mathbf{W}_{q}\right)^{-1} \mathbf{H}^{\mathsf{T}} \mathbf{G}^{\mathsf{T}} \mathbf{W}_{d},$$
 (20)

108 or, equivalently (Menke, 2018, p. 62),

$$\mathbf{B} = \mathbf{W}_q^{-1} \mathbf{H}^{\top} \mathbf{G}^{\top} \left(\mathbf{G} \mathbf{H} \mathbf{W}_q^{-1} \mathbf{H}^{\top} \mathbf{G}^{\top} + \mu \mathbf{W}_d^{-1} \right)^{-1} . \tag{21}$$

- 109 Evidently, we have considered that all inverses exist in equations 20 and 21.
- The $Q \times D$ matrix B defined by equation 20 is commonly used for the case in which D > Q, i.e., when
- there are more data than parameters (overdetermined problems). In this case, we consider that the estimate
- 112 \tilde{q} is obtained by solving the following linear system for δ_q (equation 18):

$$\left(\mathbf{H}^{\top}\mathbf{G}^{\top}\mathbf{W}_{d}\mathbf{G}\mathbf{H} + \mu\mathbf{W}_{q}\right)\tilde{\boldsymbol{\delta}}_{q} = \mathbf{H}^{\top}\mathbf{G}^{\top}\mathbf{W}_{d}\boldsymbol{\delta}_{d}.$$
(22)

- 113 On the other hand, for the cases in which D < Q (underdetermined problems), matrix B is usually defined
- according to equation 21. In this case, the general approach involves estimating \tilde{q} in two steps. The first
- 115 consists in solving a linear system for a dummy vector, which is subsequently used to compute \tilde{q} by a

matrix-vector product as follows:

$$\left(\mathbf{G}\mathbf{H}\mathbf{W}_{q}^{-1}\mathbf{H}^{\top}\mathbf{G}^{\top} + \mu\mathbf{W}_{d}^{-1}\right)\mathbf{u} = \boldsymbol{\delta}_{d}$$

$$\tilde{\boldsymbol{\delta}}_{q} = \mathbf{W}_{q}^{-1}\mathbf{H}^{\top}\mathbf{G}^{\top}\mathbf{u}$$
(23)

where u is a dummy vector. After obtaining $\tilde{\delta}_q$ (equations 22 and 23), the estimate \tilde{q} is computed with

equation 18. 118

2.2 Formulation without reparameterization

- Note that, for the particular case in which $\mathbf{H} = \mathbf{I}_P$ (equation 9), where \mathbf{I}_P is the identity of order P, 120
- $P=Q, \mathbf{p}=\mathbf{q}, \bar{\mathbf{p}}=\bar{\mathbf{q}}$ (equation 14) and $\tilde{\mathbf{p}}=\tilde{\mathbf{q}}$ (equation 15). In this case, the linear system (equations 22) 121
- and 23) is directly solved for

$$\tilde{\boldsymbol{\delta}}_p = \tilde{\mathbf{p}} - \bar{\mathbf{p}} \,, \tag{24}$$

instead of $\tilde{\boldsymbol{\delta}}_q$ (equation 18). 123

Linear system solvers 2.3 124

- 125 According to their properties, the linear systems associated with over and underdetermined problems
- (equations 22 and 23) can be solved by using direct methods such as LU, Cholesky or QR factorization, for 126
- example (Golub and Van Loan, 2013, sections 3.2, 4.2 and 5.2). These methods involve factorizing the 127
- linear system matrix in a product of "simple" matrices (i.e., triangular, diagonal or orthogonal). Here, we 128
- consider the *Cholesky factorization*, (Golub and Van Loan, 2013, p. 163). 129
- Let us consider a real linear system Mx = y, where M is a symmetric and positive definite matrix 130
- (Golub and Van Loan, 2013, p. 159). In this case, the Cholesky factorization consists in computing 131

$$\mathbf{M} = \mathcal{G}\mathcal{G}^{\top}, \tag{25}$$

- where \mathcal{G} is a lower triangular matrix called *Cholesky factor* and having positive diagonal entries. Given \mathcal{G} ,
- the original linear system is replaced by two triangular systems, as follows:

$$\mathbf{\mathcal{G}} \mathbf{s} = \mathbf{y}$$

$$\mathbf{\mathcal{G}}^{\top} \mathbf{x} = \mathbf{s}$$
(26)

- where s is a dummy vector. For the overdetermined problem (equation 22), 134
- $(\mathbf{H}^{\top}\mathbf{G}^{\top}\mathbf{W}_{d}\mathbf{G}\mathbf{H} + \mu\mathbf{W}_{q}), \mathbf{x} = \tilde{\boldsymbol{\delta}}_{q} \text{ and } \mathbf{y} = (\mathbf{H}^{\top}\mathbf{G}^{\top}\mathbf{W}_{d}\boldsymbol{\delta}_{d}).$ For the underdetermined problem (equation 23), $\mathbf{M} = (\mathbf{G}\mathbf{H}\mathbf{W}_{q}^{-1}\mathbf{H}^{\top}\mathbf{G}^{\top} + \mu\mathbf{W}_{d}^{-1}), \mathbf{x} = \mathbf{u} \text{ and } \mathbf{y} = \boldsymbol{\delta}_{d}.$
- 136
- The use of direct methods for solving large linear systems may be problematic due to computer (i) storage 137
- of large matrices and (ii) time to perform matrix operations. This problem may be specially complicated in 138
- equivalent-layer technique for the cases in which the sensitivity matrix G does not have a well-defined 139
- structure (sec. 1.2) 140
- These problems can be overcame by solving the linear system using an iterative method. These methods 141
- produce a sequence of vectors that typically converge to the solution at a reasonable rate. The main 142
- computational cost associated with these methods is usually some matrix-vector products per iteration. The 143
- conjugate gradient (CG) is a very popular iterative method for solving linear systems in equivalent-layer 144

- methods. This method was originally developed to solve systems having a square and positive definite
- 146 matrix. There are two adapted versions of the CG method. The first is called *conjugate gradient normal*
- 147 equation residual (CGNR) Golub and Van Loan (2013, sec. 11.3) or conjugate gradient least squares
- 148 (CGLS) (Aster et al., 2019, p. 165) and is used to solve overdetermined problems (equation 22). The second
- is called *conjugate gradient normal equation error* (CGNE) method Golub and Van Loan (2013, sec. 11.3)
- and is used to solve the underdetermined problems (equation 23). Algorithm 1 outlines the CGLS method
- applied to the overdetermined problem (equation 22).

3 FLOATING-POINT OPERATIONS

- 152 Two important factors affecting the efficiency of a given matrix algorithm are the storage and amount of
- 153 required arithmetic. Here, we quantify this last factor associated with different computational strategies to
- solve the linear system of the equivalent-layer technique (section 6). To do it, we opted by counting *flops*,
- which are floating point additions, subtractions, multiplications or divisions (Golub and Van Loan, 2013,
- p. 12–14). This is a non-hardware dependent approach that allows us to do direct comparison between
- 157 different equivalent-layer methods. Most of the flops count used here can be found in Golub and Van Loan
- 158 (2013, p. 12, 106, 107 and 164).
- Let us consider the case in which the overdetermined problem (equation 22) is solved by Cholesky
- 160 factorization (equation 25 and 26) directly for the parameter vector $\tilde{\mathbf{p}}$ by considering the particular case in
- which $\mathbf{H} = \mathbf{I}_P$ (equation 9 and subsection 2.2), $\mu = 0$ (equation 11), $\mathbf{W}_d = \mathbf{I}_D$ (equation 12) and $\bar{\mathbf{p}} = \mathbf{0}$
- 162 (equation 14), where I_P and I_D are the identities of order P and D, respectively. Based on the information
- provided in table 1, the total number of flops can be determined by aggregating the flops required for
- 164 various computations. These computations include the matrix-matrix and matrix-vector products $\mathbf{G}^{\top}\mathbf{G}$
- and $G^{\top}d$, the Cholesky factor \mathcal{G} , and the solution of triangular systems. Thus, we can express the total
- 166 number of flops as follows:

$$f_{\text{Cholesky}} = 1/3D^3 + 2D^2 + 2(P^2 + P)D$$
. (27)

- 167 The same particular overdetermined problem can be solved by using the CGLS method (Algorithm 1).
- 168 In this case, we use table 1 again to combine the total number of flops associated with the matrix-vector
- and inner products defined in line 3, before starting the iteration, and the 3 saxpys, 2 inner products and 2
- matrix-vector products per iteration (lines 7-12). By considering a maximum number of iterations ITMAX,
- 171 we obtain

$$f_{\text{CGLS}} = 2P(D+1) + \text{ITMAX} [2P(2D+3) + 4D]$$
 (28)

- 172 The same approach used to deduce equations 27 and 28 is applied to compute the total number of flops for
- 173 the selected equivalent-layer methods discussed in section 6.

4 NUMERICAL STABILITY

- 174 All equivalent-layer methods aim at obtaining an estimate \tilde{p} for the parameter vector p (equation 3), which
- 175 contains the physical property of the equivalent sources. Some methods do it by first obtaining an estimate
- 176 \tilde{q} for the reparameterized parameter vector q (equation 9) and then using it to obtain \tilde{p} (equation 15).
- 177 The stability of a solution \tilde{p} against noise in the observed data is rarely addressed. Here, we follow the
- 178 numerical stability analysis presented in Sigueira et al. (2017).

For a given equivalent-layer method (section 6), we obtain an estimate $\tilde{\mathbf{p}}$ assuming noise-free potential-field data \mathbf{d} . Then, we create L different noise-corrupted data \mathbf{d}^{ℓ} , $\ell \in \{1:L\}$, by adding L different sequences of pseudorandom Gaussian noise to \mathbf{d} , all of them having zero mean. From each \mathbf{d}^{ℓ} , we obtain an estimate $\tilde{\mathbf{p}}^{\ell}$. Regardless of the particular equivalent-layer method used, the following inequality (Aster et al., 2019, p. 66) holds true:

$$\Delta p^{\ell} \le \kappa \, \Delta d^{\ell} \,, \quad \ell \in \{1 : L\} \,,$$
 (29)

84 where κ is the constant of proportionality between the model perturbation

$$\Delta p^{\ell} = \frac{\|\tilde{\mathbf{p}}^{\ell} - \tilde{\mathbf{p}}\|}{\|\tilde{\mathbf{p}}\|}, \quad \ell \in \{1 : L\},$$
(30)

185 and the data perturbation

$$\Delta d^{\ell} = \frac{\|\mathbf{d}^{\ell} - \mathbf{d}\|}{\|\mathbf{d}\|}, \quad \ell \in \{1 : L\},$$
(31)

with $\|\cdot\|$ representing the Euclidean norm. The constant κ acts as the condition number associated with the pseudo-inverse in a given linear inversion. The larger (smaller) the value of κ , the more unstable (stable) is the estimated solution. Equation 29 shows a linear relationship between the model perturbation Δp^{ℓ} and the data perturbation Δd^{ℓ} (equations 30 and 31). We estimate the κ (equation 29) associated with a given equivalent-layer method as the slope of the straight line fitted to the L points $(\Delta p^{\ell}, \Delta d^{\ell})$.

5 NOTATION FOR SUBVECTORS AND SUBMATRICES

Here, we use a notation inspired on that presented by Van Loan (1992, p. 4) to represent subvectors and submatrices. Subvectors of **d**, for example, are specified by $\mathbf{d}[\mathbf{i}]$, where **i** is a list of integer numbers that "pick out" the elements of **d** forming the subvector $\mathbf{d}[\mathbf{i}]$. For example, $\mathbf{i} = (1, 6, 4, 6)$ gives the subvector $\mathbf{d}[\mathbf{i}] = [d_1 \ d_6 \ d_4 \ d_6]^{\top}$. Note that the list **i** of indices may be sorted or not and it may also have repeated indices. For the particular case in which the list has a single element $\mathbf{i} = (i)$, then it can be used to extract the *i*-th element $d_i \equiv \mathbf{d}[i]$ of **d**. Sequential lists can be represented by using the colon notation. We consider two types of sequential lists. The first has starting index is smaller than the final index and increment of 1. The second has starting index is greater than the final index and increment of -1. For example,

$$\mathbf{i} = (3:8) \Leftrightarrow \mathbf{d}[\mathbf{i}] = [d_3 \ d_4 \ \dots \ d_8]^{\top}$$

$$\mathbf{i} = (8:3) \Leftrightarrow \mathbf{d}[\mathbf{i}] = [d_8 \ d_7 \ \dots \ d_3]^{\top}$$

$$\mathbf{i} = (:8) \Leftrightarrow \mathbf{d}[\mathbf{i}] = [d_1 \ d_2 \ \dots \ d_8]^{\top}$$

$$\mathbf{i} = (3:) \Leftrightarrow \mathbf{d}[\mathbf{i}] = [d_3 \ d_4 \ \dots \ d_D]^{\top}$$

191 where D is the number of elements forming \mathbf{d} .

The notation above can also be used to define submatrices of a $D \times P$ matrix G. For example, $\mathbf{i} = (2, 7, 4, 6)$ and $\mathbf{j} = (1, 3, 8)$ lead to the submatrix

$$\mathbf{G[i,j]} = \begin{bmatrix} g_{21} & g_{23} & g_{28} \\ g_{71} & g_{73} & g_{78} \\ g_{41} & g_{43} & g_{48} \\ g_{61} & g_{63} & g_{68} \end{bmatrix}.$$

Note that, in this case, the lists i and j "pick out", respectively, the rows and columns of G that form the submatrix G[i, j]. The *i*-th row of G is given by the $1 \times P$ vector G[i, :]. Similarly, the $D \times 1$ vector G[:, j]represents the j-th column. Finally, we may use the colon notation to define the following submatrix:

$$\mathbf{i} = (2:5), \mathbf{j} = (3:7) \Leftrightarrow \mathbf{G}[\mathbf{i}, \mathbf{j}] = \begin{bmatrix} g_{23} & g_{24} & g_{25} & g_{26} & g_{27} \\ g_{33} & g_{34} & g_{35} & g_{36} & g_{37} \\ g_{43} & g_{44} & g_{45} & g_{46} & g_{47} \\ g_{53} & g_{54} & g_{55} & g_{56} & g_{57} \end{bmatrix},$$

which contains the contiguous elements of G from rows 2 to 5 and from columns 3 to 7.

COMPUTATIONAL STRATEGIES

- 193 The linear inverse problem of the equivalent-layer technique (section 2) for the case in which there are
- large volumes of potential-field data requires dealing with: 194
- (i) the large computer memory to store large and full matrices; 195
- (ii) the long computation time to multiply a matrix by a vector; and 196
- (iii) the long computation time to solve a large linear system of equations. 197
- Here, we review some strategies aiming at reducing the computational cost of the equivalent-layer technique. 198
- 199 We quantify the computational cost by using flops (section 3) and compare the results with those obtained
- 200 for Cholesky factorization and CGLS (equations 27 and 28). We focus on the overall strategies used by the
- 201 selected methods.

202 6.1 Moving window

- 203 The initial approach to enhance the computational efficiency of the equivalent-layer technique is
- commonly denoted moving window and involves first splitting the observed data d_i , $i \in \{1 : D\}$, into 204
- M overlapping subsets (or data windows) formed by D^m data each, $m \in \{1: M\}$. The data inside the 205
- m-th window are usually adjacent to each other and have indices defined by an integer list i^m having 206
- D^m elements. The number of data D^m forming the data windows are not necessarily equal to each other. 207
- Each data window has a $D^m \times 1$ observed data vector $\mathbf{d}^m \equiv \mathbf{d}[\mathbf{i}^m]$. The second step consists in defining 208
- a set of P equivalent sources with scalar physical property p_i , $j \in \{1 : P\}$, and also split them into M 209
- overlapping subsets (or source windows) formed by P^m data each, $m \in \{1 : M\}$. The sources inside the 210
- m-th window have indices defined by an integer list \mathbf{j}^m having P^m elements. Each source window has a 211
- $P^m \times 1$ parameter vector \mathbf{p}^m and is located right below the corresponding m-th data window. Then, each
- $\mathbf{d}^m \equiv \mathbf{d}[\mathbf{i}^m]$ is approximated by 213

$$\mathbf{f}^m = \mathbf{G}^m \mathbf{p}^m \,, \tag{32}$$

- where $G^m \equiv G[i^m, j^m]$ is a submatrix of G (equation 3) formed by the elements computed with equation 214
- 2 using only the data and equivalent sources located inside the window m-th. The main idea of the moving-215
- window approach is using the $\tilde{\mathbf{p}}^m$ estimated for each window to obtain (i) an estimate $\tilde{\mathbf{p}}$ of the parameter 216
- vector for the entire equivalent layer or (ii) a given potential-field transformation t (equation 4). The main 217
- advantages of this approach is that (i) the estimated parameter vector $\tilde{\mathbf{p}}$ or transformed potential field are 218
- 219 not obtained by solving the full, but smaller linear systems and (ii) the full matrix G (equation 3) is never
- 220 stored.

230

221 Leão and Silva (1989) presented a pioneer work using the moving-window approach. Their method requires a regularly-spaced grid of observed data on a horizontal plane z_0 . The data windows are defined by 222 square local grids of $\sqrt{D'} \times \sqrt{D'}$ adjacent points, all of them having the same number of points D'. The 223 equivalent sources in the m-th data window are located below the observation plane, at a constant vertical 224 distance Δz_0 . They are arranged on a regular grid of $\sqrt{P'} \times \sqrt{P'}$ adjacent points following the same 225 grid pattern of the observed data. The local grid of sources for all data windows have the same number 226 of elements P'. Besides, they are vertically aligned, but expands the limits of their corresponding data 227 windows, so that D' < P'. Because of this spatial configuration of observed data and equivalent sources, 228 we have that $G^m = G'$ (equation 32) for all data windows (i.e., $\forall m \in \{1 : M\}$), where G' is a $D' \times P'$ 229 constant matrix.

By omitting the normalization strategy used by Leão and Silva (1989), their method consists in directly 231 computing the transformed potential field t_c^m at the central point $(x_c^m, y_c^m, z_0 + \Delta z_0)$ of each data window 232 233 as follows:

$$t_c^m = (\mathbf{a}')^\top \mathbf{B}' \mathbf{d}^m, \quad m \in \{1 : M\},$$
 (33)

where a' is a $P' \times 1$ vector with elements computed by equation 5 by using all equivalent sources in the m-th window and only the coordinate of the central point in the m-th data window and

$$\mathbf{B}' = \left(\mathbf{G}'\right)^{\top} \left[\mathbf{G}' \left(\mathbf{G}'\right)^{\top} + \mu \mathbf{I}_{D'}\right]^{-1}$$
(34)

is a particular case of matrix B associated with underdetermined problems (equation 21) for the particular 236 case in which $\mathbf{H} = \mathbf{W}_q = \mathbf{I}_{P'}$ (equations 9 and 13), $\mathbf{W}_d = \mathbf{I}_{D'}$ (equation 12), $\bar{\mathbf{p}} = \mathbf{0}$ (equation 14), where 237 $I_{P'}$ and $I_{D'}$ are identity matrices of order P' and D', respectively, and 0 is a vector of zeros. Due to the 238 presumed spatial configuration of the observed data and equivalent sources, a' and G' are the same for all 239 data windows. Hence, only the data vector \mathbf{d}^m is modified according to the position of the data window. 240 Note that equation 33 combines the potential-field transformation (equation 4) with the solution of the 241 undetermined problem (equation 23).

243 The method proposed by Leão and Silva (1989) can be outlined by the Algorithm 2. Note that Leão and Silva (1989) directly compute the transformed potential t_c^m at the central point of each data window without explicitly computing and storing an estimated for p^m (equation 32). It means that their method allows 245 computing a single potential-field transformation. A different transformation or the same one evaluated at 246 247 different points require running their moving-data window method again.

The total number of flops in Algorithm 2 depends on computing the $P' \times D'$ matrix B' (equation 34) in 248 line 6 and use it to define the $1 \times P'$ vector $(\mathbf{a}')^{\top} \mathbf{B}'$ (line 7) before starting the iterations and computing 249 an inner product (equation 33) per iteration. We consider that the total number of flops associated with B' 250 is obtained by the matrix-matrix product $\mathbf{G}'(\mathbf{G}')^{\top}$, its inverse and then the premultiplication by $(\mathbf{G}')^{\top}$. By 251 using table 1 and considering that inverse is computed via Cholesky factorization, we obtain that the total 252 number of flops for lines 6 and 7 is $2(D')^2P' + 7(D')^3/6 + 2(D')^2P'$. Then, the total number of flops for 253 Algorithm 2 is 254

$$f_{LS89} = 7/6(D')^3 + 4P'(D')^2 + M2P'.$$
 (35)

Soler and Uieda (2021) generalized the method proposed by Leão and Silva (1989) for irregularly spaced 255 data on an undulating surface. A direct consequence of this generalization is that a different submatrix 256 $\mathbf{G}^m \equiv \mathbf{G}[\mathbf{i}^m, \mathbf{j}^m]$ (equation 32) must be computed for each window. Differently from Leão and Silva 257

264

265

266

267

268

269

270271

272

273274

275

276

277278

279

280

281

282

283

284

285

286

287 288

290

258 (1989), Soler and Uieda (2021) store the computed $\tilde{\mathbf{p}}^m$ for all windows and subsequently use them to obtain 259 a desired potential-field transformation (equation 4) as the superposed effect of all windows. The estimated 260 $\tilde{\mathbf{p}}^m$ for all windows are combined to form a single $P \times 1$ vector $\tilde{\mathbf{p}}$, which is an estimate for original 261 parameter vector \mathbf{p} (equation 3). For each data window, Soler and Uieda (2021) solve an overdetermined 262 problem (equation 22) for $\tilde{\mathbf{p}}^m$ by using $\mathbf{H} = \mathbf{W}_q = \mathbf{I}_{P^m}$ (equations 9 and 13), \mathbf{W}_d^m (equation 12) equal to 263 a diagonal matrix of weights for the data inside the m-th window and $\bar{p} = \mathbf{0}$ (equation 14), so that

$$\left[(\mathbf{G}^m)^\top \mathbf{W}_d^m \mathbf{G}^m + \mu \mathbf{I}_{P'} \right] \tilde{\mathbf{p}}^m = (\mathbf{G}^m)^\top \mathbf{W}_d^m \mathbf{d}^m.$$
 (36)

Unlike Leão and Silva (1989), Soler and Uieda (2021) do not adopt a sequential order of the data windows; rather, they adopt a randomized order of windows in their iterations. The overall steps of the method proposed by Soler and Uieda (2021) are defined by the Algorithm 3. For convenience, we have omitted the details about the randomized window order and the normalization strategy employed by Soler and Uieda (2021). Note that this algorithm starts with a residuals vector r that is iteratively updated. The iterative algorithm in Soler and Uieda (2021) estimates a solution (\tilde{p}^m in equation 36) using the data and the equivalent sources that fall within a moving-data window; however, it calculates the predicted data and the residual data in the whole survey data. Next, the residual data that fall within a new position of the data window is used as input data to estimate a new solution within the data window which, in turn, is used to calculated a new predicted data and a new residual data in the whole survey data. Regarding the equivalent-source layout, Soler and Uieda (2021) proposed the block-averaged sources locations in which the survey area is divided into horizontal blocks and one single equivalent source is assigned to each block. Each single source per block is placed over the layer with its horizontal coordinates given by the average horizontal positions of observation points. According to Soler and Uieda (2021), the block-averaged sources layout may prevent aliasing of the interpolated values, specially when the observations are unevenly sampled. This strategy also reduces the number of equivalent sources without affecting the accuracy of the potential-field interpolation. Besides, it reduces the computational load for estimating the physical property on the equivalent layer. This block-averaged sources layout is not included in the Algorithm 3.

The computational cost of Algorithm 3 can be defined in terms of the linear system (equation 36) to be solved for each window (line 10) and the subsequent updates in lines 11 and 12. We consider that the linear system cost can be quantified by the matrix-matrix and matrix-vector products $(\mathbf{G}^m)^{\top} \mathbf{G}^m$ and $(\mathbf{G}^m)^{\top} \mathbf{d}^m$, respectively, and solution of the linear system (line 10) via Cholesky factorization (equations 25 and 26). The following updates represent a saxpy without scalar-vector product (line 11) and a matrix-vector product (line 12). In this case, according to table 1, the total number of flops associated with Algorithm 3 is given by:

$$f_{\text{SU21}} = M \left[\frac{1}{3} (P')^3 + 2(D'+1)(P')^2 + (4D'+1)P' \right] , \qquad (37)$$

289 where P' and D' represent, respectively, the average number of equivalent sources and data at each window.

6.2 Column-action update

We call the computational strategy *column-action update* because a single source is used to calculate the predicted data and the residual data in the whole survey data. Hence, a single column of the sensitivity matrix G (equation 3) is calculated iteratively.

Cordell (1992) proposed a computational strategy that was later used by Guspí and Novara (2009) and relies on first defining one equivalent source located right below each observed data d_i , $i \in \{1 : D\}$, at

a vertical coordinate $z_i + \Delta z_i$, where Δz_i is proportional to the distance from the *i*-th observation point 296 297 (x_i, y_i, z_i) to its closest neighbor. The second step consists in updating the physical property p_i of a single equivalent source, $j \in \{1 : D\}$ and remove its predicted potential field from the observed data vector d, 298 producing a residuals vector r. At each iteration, the single equivalent source is the one located vertically 299 beneath the observation station of the maximum data residual. Next, the predicted data produced by this 300 single source is calculated over all of the observation points and a new data residual r and the $D \times 1$ 301 parameter vector p containing the physical property of all equivalent sources are updated iteratively. During 302 each subsequent iteration, Cordell's method either incorporates a single equivalent source or adjusts an 303 existing equivalent source to match the maximum amplitude of the current residual field. The convergence 304 occurs when all of the residuals are bounded by an envelope of prespecified expected error. At the end, the 305 algorithm produces an estimate \tilde{p} for the parameter vector yielding a predicted potential field f (equation 306 3) satisfactorily fitting the observed data d according to a given criterion. Note that the method proposed 307 by Cordell (1992) iteratively solves the linear $G\tilde{p} \approx d$ with a $D \times D$ matrix G. At each iteration, only a 308 single column of G (equation 3) is used. An advantage of this column-action update approach is that the 309 full matrix G is never stored. 310

Algorithm 4 delineates the Cordell's method. Note that a single column $G[:, i_{max}]$ of the $D \times D$ matrix G (equation 3) is used per iteration, where i_{max} is the index of the maximum absolute value in \mathbf{r} . As pointed out by Cordell (1992), the method does not necessarily decrease monotonically along the iterations. Besides, the method may not converge depending on how the vertical distances Δz_i , $i \in \{1:D\}$, controlling the depths of the equivalent sources are set. According to Cordell (1992), the maximum absolute value r_{max} in \mathbf{r} decreases robustly at the beginning and oscillates within a narrowing envelope for the subsequent iterations.

Guspí and Novara (2009) generalized Cordell's method to perform reduction to the pole and other transformations on scattered magnetic observations by using two steps. The first step involves computing the vertical component of the observed field using equivalent sources while preserving the magnetization direction. In the second step, the vertical observation direction is maintained, but the magnetization direction is shifted to the vertical. The main idea employed by both Cordell (1992) and Guspí and Novara (2009) is an iterative scheme that uses a single equivalent source positioned below a measurement station to compute both the predicted data and residual data for all stations. This approach entails a computational strategy where a single column of the sensitivity matrix G (equation 3) is calculated per iteration.

The total number of flops in Algorithm 4 consists in finding the maximum absolute value in vector $\bf r$ (line 6) before the while loop. Per iteration, there is a saxpy (line 11) and another search for the maximum absolute value in vector $\bf r$ (line 12). By considering that selecting the maximum absolute value in a $D \times 1$ vector is a $D \log_2(D)$ operation (e.g., Press et al., 2007, p. 420), the total number of flops in Algorithm 38 is given by:

$$f_{\text{C92}} = D\log(D) + \text{ITMAX} [2D + D\log_2(D)]$$
 (38)

6.3 Row-action update

318

319

320 321

322

323

324 325

331

332

333

334

335

336

337

We call the computational strategy *row-action update* because a single row of the sensitivity matrix G (equation 3) is calculated iteratively. Hence, the equivalent-layer solution is updated by processing a new datum (one matrix row) at each iteration. To reduce the total processing time and memory usage of equivalent-layer technique, Mendonça and Silva (1994) proposed a strategy called *equivalent data concept*. The equivalent data concept is grounded on the principle that there is a subset of redundant data that does not contribute to the final solution and thus can be dispensed. Conversely, there is a subset of observations,

called equivalent data, that contributes effectively to the final solution and fits the remaining observations (redundant data). Iteratively, Mendonça and Silva (1994) selected the subset of equivalent data that is substantially smaller than the original dataset. This selection is carried out by incorporating one data point at a time.

Mendonça and Silva (1994) proposes an algebraic reconstruction technique (ART) (e.g., van der Sluis 342 and van der Vorst, 1987, p. 58) to estimate a parameter vector $\tilde{\mathbf{p}}$ for a regular grid of P equivalent sources 343 344 on a horizontal plane z_0 . Such methods iterate on the linear system rows to estimate corrections for the parameter vector, which may substantially save computer time and memory required to compute and store 345 the full linear system matrix along the iterations. The convergence of such row-update methods depends 346 on the linear system condition. The main advantage of such methods is not computing and storing the 347 full linear system matrix, but iteratively using its rows. In contrast to ART-type algorithms, the rows in 348 Mendonça and Silva (1994) are not processed sequentially. Instead, in Mendonça and Silva (1994), the 349 rows are introduced according to their residual magnitudes (maximum absolute value in r), which are 350 computed based on the estimate over the equivalent layer from the previous iteration. The particular ART 351 method proposed by Mendonça and Silva (1994) considers that 352

$$\mathbf{d} = \begin{bmatrix} \mathbf{d}_e \\ \mathbf{d}_r \end{bmatrix} , \quad \mathbf{G} = \begin{bmatrix} \mathbf{G}_e \\ \mathbf{R}_r \end{bmatrix} , \tag{39}$$

where \mathbf{d}_e and \mathbf{d}_r are $D_e \times 1$ and $D_r \times 1$ vectors and \mathbf{G}_e and \mathbf{G}_r are $D_e \times P$ and $D_r \times P$ matrices, respectively. Mendonça and Silva (1994) designate \mathbf{d}_e and \mathbf{d}_r as, respectively, equivalent and redundant data. With the exception of a normalization strategy, Mendonça and Silva (1994) calculate a $P \times 1$ estimated parameter vector $\tilde{\mathbf{p}}$ by solving an underdetermined problem (equation 23) involving only the equivalent data \mathbf{d}_e (equation 39) for the particular case in which $\mathbf{H} = \mathbf{W}_p = \mathbf{I}_P$ (equations 9 and 13), $\mathbf{W}_d = \mathbf{I}_{D_e}$ (equation 12) and $\bar{p} = \mathbf{0}$ (equation 14), which results in

$$(\mathbf{F} + \mu \, \mathbf{I}_{D_e}) \, \mathbf{u} = \mathbf{d}_e$$

$$\tilde{\mathbf{p}} = \mathbf{G}_e^{\top} \mathbf{u} , \qquad (40)$$

where \mathbf{F} is a computationally-efficient $D_e \times D_e$ matrix that approximates $\mathbf{G}_e \mathbf{G}_e^{\top}$. Mendonça and Silva (1994) presume that the estimated parameter vector $\tilde{\mathbf{p}}$ obtained from equation 40 leads to a $D_r \times 1$ residuals vector

$$\mathbf{r} = \mathbf{d}_r - \mathbf{G}_r \tilde{\mathbf{p}} \tag{41}$$

362 having a maximum absolute value $r_{\text{max}} \leq \epsilon$, where ϵ is a predefined tolerance.

The overall method of Mendonça and Silva (1994) is defined by Algorithm 5. It is important noting that the number D_e of equivalent data in \mathbf{d}_e increases by one per iteration, which means that the order of the linear system in equation 40 also increases by one at each iteration. Those authors also propose a computational strategy based on Cholesky factorization (e.g., Golub and Van Loan, 2013, p. 163) for efficiently updating $(\mathbf{F} + \mu \mathbf{I}_{D_e})$ at a given iteration (line 16 in Algorithm 5) by computing only its new elements with respect to those computed in the previous iteration.

6.4 Reparameterization

369

Another approach for improving the computational performance of equivalent-layer technique consists in setting a $P \times Q$ reparameterization matrix **H** (equation 9) with Q << P. This strategy has been used

379

380

381

382

383 384

385

386

394

395 396

397

398

399

in applied geophysics for decades (e.g., Skilling and Bryan, 1984; Kennett et al., 1988; Oldenburg et al., 373 1993; Barbosa et al., 1997) and is known as *subspace method*. The main idea relies in reducing the linear system dimension from the original P-space to a lower-dimensional subspace (the Q-space). An estimate 374 $\tilde{\mathbf{q}}$ for the reparameterized parameter vector \mathbf{q} is obtained in the Q-space and subsequently used to obtain 375 an estimate \tilde{p} for the parameter vector p (equation 3) in the P-space by using equation 9. Hence, the key 376 aspect of this reparameterization approach is solving an appreciably smaller linear inverse problem for \tilde{q} 377 than that for the original parameter vector $\tilde{\mathbf{p}}$ (equation 3). 378

Oliveira Jr. et al. (2013) have used this approach to describe the physical property distribution on the equivalent layer in terms of piecewise bivariate polynomials. Specifically, their method consists in splitting a regular grid of equivalent sources into source windows inside which the physical-property distribution is described by bivariate polynomial functions. The key aspect of their method relies on the fact that the total number of coefficients required to define the bivariate polynomials is considerably smaller than the original number of equivalent sources. Hence, they formulate a linear inverse problem for estimating the polynomial coefficients and use them later to compute the physical property distribution on the equivalent layer.

The method proposed by Oliveira Jr. et al. (2013) consists in solving an overdetermined problem (equation 387 22) for estimating the polynomial coefficients $\tilde{\bf q}$ with ${\bf W}_d={\bf I}_D$ (equation 12) and $\bar q={\bf 0}$ (equation 14), so 388 that 389

$$\left(\mathbf{H}^{\top}\mathbf{G}^{\top}\mathbf{G}\,\mathbf{H} + \mu\,\mathbf{W}_{q}\right)\tilde{\mathbf{q}} = \mathbf{H}^{\top}\mathbf{G}^{\top}\,\mathbf{d}\,,\tag{42}$$

where $\mathbf{W}_q = \mathbf{H}^{ op} \mathbf{W}_p \mathbf{H}$ is defined by a matrix \mathbf{W}_p representing the zeroth- and first-order Tikhonov 390 regularization (e.g., Aster et al., 2019, p. 103). Note that, in this case, the prior information is defined in the 391 392 P-space for the original parameter vector p and then transformed to the Q-space. Another characteristic of their method is that it is valid for processing irregularly-spaced data on an undulating surface. 393

Mendonça (2020) also proposed a reparameterization approach for the equivalent-layer technique. Their approach, however, consists in setting H as a truncated singular value decomposition (SVD) (e.g., Aster et al., 2019, p. 55) of the observed potential field. Differently from Oliveira Jr. et al. (2013), however, the method of Mendonça (2020) requires a regular grid of potential-field data on horizontal plane. Another difference is that these authors uses $\mathbf{W}_q = \mathbf{I}_Q$ (equation 13), which means that the regularization is defined directly in the Q-space.

400 Before Oliveira Jr. et al. (2013) and Mendonça (2020), Barnes and Lumley (2011) also proposed a computationally efficient method for equivalent-layer technique based on reparameterization. A key 401 402 difference, however, is that Barnes and Lumley (2011) did not set a $P \times Q$ reparameterization matrix H (equation 9) with $Q \ll P$. Instead, they used a matrix **H** with $Q \approx 1.7 \, P$. Their central idea is setting a 403 reparameterization scheme that groups distant equivalent sources into blocks by using a bisection process. 404 405 This scheme leads to a quadtree representation of the physical-property distribution on the equivalent layer, so that matrix GH (equation 10) is notably sparse. Barnes and Lumley (2011) explore this sparsity in 406 solving the overdetermined problem for \tilde{q} (equation 42) via conjugate-gradient method (e.g., Golub and 407 408 Van Loan, 2013, sec. 11.3).

We consider an algorithm (not shown) that solves the overdetermined problem (equation 22) by combining 409 the reparameterization with CGLS method (Algorithm 1). It starts with a reparameterization step defined 410 by defining a matrix C = GH (equation 10). Then, the CGLS (Algorithm 1) is applied by replacing G 411 with C. In this case, the linear system is solved by the reparameterized parameter vector $\tilde{\mathbf{q}}$ instead of $\tilde{\mathbf{p}}$.

412

At the end, the estimated \tilde{q} is transformed into \tilde{p} (equation 15). Compared to the original CGLS show in

- 414 Algorithm 1, the proposed algorithm has the additional flops associated with the matrix-matrix product to
- 415 compute C and the matrix-vector product of equation 15 outside the while loop. Then, according to table 1,
- 416 the total number of flops given by:

$$f_{\text{reparam.}} = 2Q(DP + D + 1) + 2PQ + \text{ITMAX} [2Q(2D + 3) + 4D]$$
 (43)

417 6.5 Wavelet compression

- Previously to Barnes and Lumley (2011), the idea of transforming the dense matrix G (equation 3) into a
- 419 sparse one has already been used in the context of equivalent-layer technique. Li and Oldenburg (2010)
- 420 proposed a method that applies the discrete wavelet transform to introduce sparsity into the original dense
- 421 matrix G. Those authors approximate a planar grid of potential-field data by a regularly-spaced grid of
- 422 equivalent sources, so that the number of data D and sources P is the same, i.e., D = P. Specifically, Li
- 423 and Oldenburg (2010) proposed a method that applies the wavelet transform to the original dense matrix G
- 424 and sets to zero the small coefficients that are below a given threshold, which results in an approximating
- 425 sparse representation of G in the wavelet domain. They first consider the following approximation

$$\mathbf{d}_w \approx \mathbf{G}_s \, \mathbf{p}_w \,, \tag{44}$$

426 where

$$\mathbf{d}_w = \mathbf{W} \, \mathbf{d} \,, \quad \mathbf{p}_w = \mathbf{W} \, \mathbf{p} \,, \tag{45}$$

- 427 are the observed data and parameter vector in the wavelet domain; $\boldsymbol{\mathcal{W}}$ is a $D \times D$ orthogonal matrix
- 428 defining a discrete wavelet transform; and G_s is a sparse matrix obtained by setting to zero the elements of

$$G_w = \mathcal{W} G \mathcal{W}^{\top} \tag{46}$$

- 429 with absolute value smaller than a given threshold.
- Li and Oldenburg (2010) solve a normalized inverse problem in the wavelet domain. Specifically, they
- 431 first define a matrix

$$\mathbf{G}_L = \mathbf{G}_s \, \mathbf{L}^{-1} \tag{47}$$

432 and a normalized parameter vector

$$\mathbf{p}_L = \mathbf{L} \, \mathbf{p}_w \,, \tag{48}$$

- 433 where L is a diagonal and invertible matrix representing an approximation of the first-order Tikhonov
- 434 regularization in the wavelet domain. Then they solve an overdetermined problem (equation 22) to obtain
- 435 an estimate $\tilde{\mathbf{p}}_L$ for \mathbf{p}_L (equation 48), with \mathbf{G}_L (equation 47), $\mathbf{H} = \mathbf{I}_P$ (equations 9), $\mu = 0$ (equation 11),
- 436 $W_d = I_D$ (equation 12) and $\bar{p} = 0$ (equation 14) via conjugate-gradient method (e.g., Golub and Van
- 437 Loan, 2013, sec. 11.3). Finally, Li and Oldenburg (2010) compute an estimate \tilde{p} for the original parameter
- 438 vector given by

$$\tilde{\mathbf{p}} = \boldsymbol{\mathcal{W}}^{\top} \left(\mathbf{L}^{-1} \, \tilde{\mathbf{p}}_L \right) \,, \tag{49}$$

- 439 where the term within parenthesis is an estimate $\tilde{\mathbf{p}}_w$ of the parameter vector \mathbf{p}_w (equation 45) in the wavelet
- 440 domain and matrix \mathcal{W}^{\top} represents an inverse wavelet transform.

451

452

453

454

455

456 457

458

459 460

441 6.6 Iterative methods using the full matrix G

Xia and Sprowl (1991) introduced an iterative method for estimating the parameter vector $\tilde{\mathbf{p}}$ (equation 3), 442 which was subsequently adapted to the Fourier domain by Xia et al. (1993). Their method uses the full 443 444 and dense sensitivity matrix G (equation 3) (without applying any compression or reparameterization, for 445 example) to compute the predicted data at all observation points per iteration. More than two decades later, Siqueira et al. (2017) have proposed an iterative method similar to that presented by Xia and Sprowl (1991). 446 447 The difference is that Siqueira et al.'s algorithm was deduced from the Gauss' theorem (e.g., Kellogg, 1967, p. 43) and the total excess of mass (e.g., Blakely, 1996, p. 60). Besides, Siqueira et al. (2017) have included 448 a numerical analysis showing that their method produces very stable solutions, even for noise-corrupted 449 450 potential-field data.

The iterative method proposed by Siqueira et al. (2017) is outlined in Algorithm 6, presumes an equivalent layer formed by monopoles (point masses) and can be applied to irregularly-spaced data on an undulating surface. Note that the residuals $\bf r$ are used to compute a correction $\Delta \bf p$ for the parameter vector at each iteration (line 11), which requires a matrix-vector product involving the full matrix $\bf G$. Interestingly, this approach for estimating the physical property distribution on an equivalent layer is the same originally proposed by Bott (1960) for estimating the basement relief under sedimentary basins. The methods of Xia and Sprowl (1991) and Siqueira et al. (2017) were originally proposed for processing gravity data, but can be potentially applied to any harmonic function because they actually represent iterative solutions of the classical *Dirichlet's problem* or the *first boundary value problem of potential theory* (Kellogg, 1967, p. 236) on a plane.

Recently, Jirigalatu and Ebbing (2019) presented another iterative method for estimating a parameter 461 vector $\tilde{\mathbf{p}}$ (equation 3). With the purpose of combining different potential-field data, their method basically 462 modifies that shown in Algorithm 6 by changing the initial approximation and the iterative correction for 463 the parameter vector. Specifically, Jirigalatu and Ebbing (2019) replace line 4 by $\tilde{\mathbf{p}} = \mathbf{0}$, where $\mathbf{0}$ is a vector 464 of zeros, and line 9 by $\Delta \mathbf{p} = \omega \mathbf{G}^{\mathsf{T}} \mathbf{r}$, where ω is a positive scalar defined by trial and error. Note that 465 this modified approach requires two matrix-vector products involving the full matrix G per iteration. To 466 overcome the high computational cost of these two products, Jirigalatu and Ebbing (2019) set an equivalent 467 layer formed by prisms and compute their predicted potential field in the wavenumber domain by using the 468 Gauss-FFT technique Zhao et al. (2018). 469

Algorithm 6: entrywise product line 4(D), matrix-vector and subtraction line $6(2D^2 + D)$, entrywise product line 9(D), addition line 10(D), matrix-vector line $11(2D^2)$, subtraction line 12(D):

$$f_{\text{SOB17}} = 2D^2 + 2D + \text{ITMAX} \left(2D^2 + 3D\right)$$
 (50)

472 6.7 Iterative deconvolution

473 Recently, Takahashi et al. (2020, 2022) proposed the convolutional equivalent-layer method, which explores the structure of the sensitivity matrix G (equation 3) for the particular case in which (i) there 474 is a single equivalent source right below each potential-field datum and (ii) both data and sources rely 475 on planar and regularly spaced grids. Specifically, they consider a regular grid of D potential-field data 476 at points (x_i, y_i, z_0) , $i \in \{1 : D\}$, on a horizontal plane z_0 . The data indices i may be ordered along the 477 x- or y-direction, which results in an x- or y-oriented grid, respectively. They also consider a single 478 equivalent source located right below each datum, at a constant vertical coordinate $z_0 + \Delta z$, $\Delta z > 0$. In 479 this case, the number of data and equivalent sources are equal to each other (i.e., D = P) and G (equation 480

- 481 3) assumes a doubly block Toeplitz (Jain, 1989, p. 28) or block-Toeplitz Toeplitz-block (BTTB) (Chan and
- 482 Jin, 2007, p. 67) structure formed by $N_B \times N_B$ blocks, where each block has $N_b \times N_b$ elements, with
- 483 $D = N_B N_b$. This particular structure allows formulating the product of G and an arbitrary vector as a fast
- 484 discrete convolution via Fast Fourier Transform (FFT) (Van Loan, 1992, section 4.2).
- Consider, for example, the particular case in which $N_B = 4$, $N_b = 3$ and D = 12. In this case, G
- 486 (equation 3) is a 12×12 block matrix given by

$$\mathbf{G} = \begin{bmatrix} \mathbf{G}^{0} & \mathbf{G}^{1} & \mathbf{G}^{2} & \mathbf{G}^{3} \\ \mathbf{G}^{-1} & \mathbf{G}^{0} & \mathbf{G}^{1} & \mathbf{G}^{2} \\ \mathbf{G}^{-2} & \mathbf{G}^{-1} & \mathbf{G}^{0} & \mathbf{G}^{1} \\ \mathbf{G}^{-3} & \mathbf{G}^{-2} & \mathbf{G}^{-1} & \mathbf{G}^{0} \end{bmatrix}_{D \times D},$$
(51)

- 487 where each block \mathbf{G}^n , $n \in \{(1 N_B) : (N_B 1)\}$, is a 3×3 Toeplitz matrix. Takahashi et al. (2020,
- 488 2022) have deduced the specific relationship between blocks G^n and G^{-n} and also between a given block
- 489 \mathbf{G}^n and its transposed $(\hat{\mathbf{G}}^n)^{\top}$ according to the harmonic function g_{ij} (equation 2) defining the element ij
- 490 of the sensitivity matrix G (equation 3) and the orientation of the data grid.
- 491 Consider the matrix-vector products

$$\mathbf{G}\,\mathbf{v} = \mathbf{w} \tag{52}$$

492 and

$$\mathbf{G}^{\top}\mathbf{v} = \mathbf{w} \,, \tag{53}$$

493 involving a $D \times D$ sensitivity matrix G (equation 3) defined in terms of a given harmonic function g_{ij}

494 (equation 2), where

$$\mathbf{v} = \begin{bmatrix} \mathbf{v}^0 \\ \vdots \\ \mathbf{v}^{N_B - 1} \end{bmatrix}_{D \times 1}, \quad \mathbf{w} = \begin{bmatrix} \mathbf{w}^0 \\ \vdots \\ \mathbf{w}^{N_B - 1} \end{bmatrix}_{D \times 1}, \tag{54}$$

are arbitrary partitioned vectors formed by N_B sub-vectors \mathbf{v}^n and \mathbf{w}^n , $n \in \{0 : (N_B - 1)\}$, all of them

96 having N_b elements. Equations 52 and 53 can be computed in terms of an auxiliary matrix-vector product

$$\mathbf{G}_c \, \mathbf{v}_c = \mathbf{w}_c \,, \tag{55}$$

497 where

$$\mathbf{v}_{c} = \begin{bmatrix} \mathbf{v}_{c}^{0} \\ \vdots \\ \mathbf{v}_{c}^{N_{B}-1} \\ \mathbf{0} \end{bmatrix}_{AD \times 1}, \quad \mathbf{w}_{c} = \begin{bmatrix} \mathbf{w}_{c}^{0} \\ \vdots \\ \mathbf{w}_{c}^{N_{B}-1} \\ \mathbf{0} \end{bmatrix}_{AD \times 1}, \quad (56)$$

498 are partitioned vectors formed by $2N_b \times 1$ sub-vectors

$$\mathbf{v}_c^n = \begin{bmatrix} \mathbf{v}^n \\ \mathbf{0} \end{bmatrix}_{2N_b \times 1}, \quad \mathbf{w}_c^n = \begin{bmatrix} \mathbf{w}^n \\ \mathbf{0} \end{bmatrix}_{2N_b \times 1}, \tag{57}$$

499 and G_c is a $4D \times 4D$ doubly block circulant (Jain, 1989, p. 28) or block-circulant circulant-block (BCCB)

- 500 (Chan and Jin, 2007, p. 76) matrix. What follows aims at explaining how the original matrix-vector products
- 600 (chair and 511, 2007, p. 70) matrix. What follows aims at explaining now the original matrix vector products 501 defined by equations 52 and 53, involving a $D \times D$ BTTB matrix **G** exemplified by equation 51, can be

- efficiently computed in terms of the auxiliary matrix-vector product given by equation 55, which has a $4D \times 4D$ BCCB matrix G_c .
- Matrix G_c (equation 55) is formed by $2N_B \times 2N_B$ blocks, where each block G_c^n , $n \in \{(1 N_B) : (N_B 1)\}$ is a $2N_b \times 2N_b$ circulant matrix. For the case in which the original matrix-vector product is that
- defined by equation 52, the first column of blocks forming the BCCB matrix G_c is given by

$$\mathbf{G}_{c}[:,:2N_{b}] = \begin{bmatrix} \mathbf{G}_{c}^{0} \\ \mathbf{G}_{c}^{-1} \\ \vdots \\ \mathbf{G}_{c}^{1-N_{B}} \\ \mathbf{0} \\ \mathbf{G}_{c}^{N_{B}-1} \\ \vdots \\ \mathbf{G}_{c}^{1} \end{bmatrix}_{AD \times 2N_{c}},$$

$$(58)$$

507 with blocks G_c^n having the first column given by

$$\mathbf{G}_{c}^{n}[:, 1] = \begin{bmatrix} \mathbf{G}^{n}[:, 1] \\ 0 \\ (\mathbf{G}^{n}[1, N_{b}: 2])^{\top} \end{bmatrix}_{2N_{b} \times 2N_{b}}, \quad n \in \{(1 - N_{B}): (N_{B} - 1)\},$$
 (59)

where G^n are the blocks forming the BTTB matrix G (equation 51). For the case in which the original matrix-vector product is that defined by equation 53, the first column of blocks forming the BCCB matrix G_c is given by

$$\mathbf{G}_{c}[:,:2N_{b}] = \begin{bmatrix} \mathbf{G}_{c}^{0} \\ \mathbf{G}_{c}^{1} \\ \vdots \\ \mathbf{G}_{c}^{N_{B}-1} \\ \mathbf{0} \\ \mathbf{G}_{c}^{1-N_{B}} \\ \vdots \\ \mathbf{G}_{c}^{-1} \end{bmatrix}_{4D \times 2N_{b}} , \qquad (60)$$

511 with blocks G_c^n having the first column given by

$$\mathbf{G}_{c}^{n}[:, 1] = \begin{bmatrix} (\mathbf{G}^{n}[1, :])^{\top} \\ 0 \\ \mathbf{G}^{n}[N_{b}: 2, 1] \end{bmatrix}_{2N_{b} \times 2N_{b}}, \quad n \in \{(1 - N_{B}): (N_{B} - 1)\}.$$
 (61)

- The complete matrix G_c (equation 55) is obtained by properly downshifting the block columns $G_c[:,:]$
- 513 $2N_b$] defined by equation 58 or 60. Similarly, the *n*-th block \mathbf{G}_c^n of \mathbf{G}_c is obtained by properly downshifting
- 514 the first columns $\mathbf{G}_c^{\ell}[:, 1]$ defined by equation 59 or 61.
- Note that G_c (equation 55) is a $4D \times 4D$ matrix and G (equation 51) is a $D \times D$ matrix. It seems weird
- 516 to say that computing $G_c v_c$ is more efficient than directly computing Gv. To understand this, we need first

to use the fact that BCCB matrices are diagonalized by the 2D unitary discrete Fourier transform (DFT) (e.g., Davis, 1979, p. 31). Because of that, G_c can be written as

$$\mathbf{G}_{c} = \left(\mathcal{F}_{2N_{B}} \otimes \mathcal{F}_{2N_{b}}\right)^{*} \Lambda \left(\mathcal{F}_{2N_{B}} \otimes \mathcal{F}_{2N_{b}}\right) , \tag{62}$$

519 where the symbol " \otimes " denotes the Kronecker product (e.g., Horn and Johnson, 1991, p. 243), \mathcal{F}_{2N_B} and

- 520 \mathcal{F}_{2N_b} are the $2N_B \times 2N_B$ and $2N_b \times 2N_b$ unitary DFT matrices (e.g., Davis, 1979, p. 31), respectively,
- 521 the superscritpt "*" denotes the complex conjugate and Λ is a $4D \times 4D$ diagonal matrix containing the
- 522 eigenvalues of G_c . Due to the diagonalization of the matrix G_c , equation 55 can be rewritten by using
- equation 62 and premultiplying both sides of the result by $(\mathcal{F}_{2N_B} \otimes \mathcal{F}_{2N_b})$, i.e.,

$$\Lambda \left(\mathcal{F}_{2N_B} \otimes \mathcal{F}_{2N_b} \right) \mathbf{v}_c = \left(\mathcal{F}_{2N_B} \otimes \mathcal{F}_{2N_b} \right) \mathbf{w}_c . \tag{63}$$

524 By following Takahashi et al. (2020), we rearrange equation 63 as follows

$$\mathcal{L} \circ (\mathcal{F}_{2N_B} \mathcal{V}_c \mathcal{F}_{2N_b}) = \mathcal{F}_{2N_B} \mathcal{W}_c \mathcal{F}_{2N_b}$$
(64)

- 525 where " \circ " denotes the Hadamard product (e.g., Horn and Johnson, 1991, p. 298) and \mathcal{L} , \mathcal{V}_c and \mathcal{W}_c are
- 526 $2N_B \times 2N_b$ matrices obtained by rearranging, along their rows, the elements forming the diagonal of Λ
- 527 (equation 62), vector \mathbf{v}_c and vector \mathbf{w}_c (equation 56), respectively. Then, by premultiplying both sides of
- equation 64 by $\mathcal{F}_{2N_B}^*$ and then postmultiplying both sides by $\mathcal{F}_{2N_b}^*$, we obtain

$$\mathcal{F}_{2N_B}^* \left[\mathcal{L} \circ \left(\mathcal{F}_{2N_B} \, \mathcal{V}_c \, \mathcal{F}_{2N_b} \right) \right] \, \mathcal{F}_{2N_b}^* = \mathcal{W}_c \,.$$
 (65)

- Finally, we get from equation 62 that matrix \mathcal{L} can be computed by using only the first column $G_c[:,1]$ of
- 530 the BCCB matrix G_c (equation 55) according to (Takahashi et al., 2020)

$$\mathcal{L} = \sqrt{4D} \, \mathcal{F}_{2N_B} \, \mathcal{C} \, \mathcal{F}_{2N_b} \,, \tag{66}$$

- where C is a $2N_B \times 2N_b$ matrix obtained by rearranging, along its rows, the elements of $G_c[:, 1]$ (equation
- 532 55). It is important noting that the matrices C and L (equation 66) associated with the BTTB matrix G
- 533 (equation 51) are different from those associated with \mathbf{G}^{\top} .
- The whole procedure to compute the original matrix-vector products Gv (equation 52) and $G^{\top}v$
- 635 (equation 53) consists in (i) rearranging the elements of the vector \mathbf{v} and the first column $\mathbf{G}[:,1]$ of matrix
- 536 G into the matrices \mathcal{V}_c and \mathcal{C} (equations 65 and 66), respectively; (ii) computing terms $\mathcal{F}_{2N_B} \mathcal{A} \mathcal{F}_{2N_b}$ and
- 537 $\mathcal{F}_{2N_B}^* \mathcal{A} \mathcal{F}_{2N_b}^*$, where \mathcal{A} is a given matrix, and a Hadamard product to obtain \mathcal{W}_c (equation 65); and (iii)
- retrieve the elements of vector w (equation 52) from \mathcal{W}_c (equation 65). It is important noting that the steps
- 539 (i) and (iii) do not have any computational cost because they involve only reorganizing elements of vectors
- and matrices. Besides, the terms $\mathcal{F}_{2N_B} \mathcal{A} \mathcal{F}_{2N_b}$ and $\mathcal{F}_{2N_B}^* \mathcal{A} \mathcal{F}_{2N_b}^*$ in step (ii) represent, respectively, the
- 2D Discrete Fourier Transform (2D-DFT) and the 2D Inverse Discrete Fourier Transform (2D-IDFT) of \mathcal{A} .
- 542 These transforms can be efficiently computed by using the 2D Fast Fourier Transform (2D-FFT). Hence,
- 543 the original matrix-vector products $\mathbf{G}\mathbf{v}$ (equation 52) and $\mathbf{G}^{\top}\mathbf{v}$ (equation 53) can be efficiently computed
- 544 by using the 2D-FFT.
- Algorithms 7 and 8 show pseudo-codes for the convolutional equivalent-layer method proposed by
- 546 Takahashi et al. (2020, 2022). Note that those authors formulate the overdetermined problem (equation
- 547 22) of obtaining an estimate \tilde{p} for the parameter vector p (equation 3) as an iterative deconvolution via

- 548 conjugate gradient normal equation residual (CGNR) Golub and Van Loan (2013, sec. 11.3) or conjugate
- 549 gradient least squares (CGLS) (Aster et al., 2019, p. 165) method. They consider $\mathbf{H} = \mathbf{I}_P$ (equation 9),
- 550 $\mu = 0$ (equation 11), $\mathbf{W}_d = \mathbf{W}_q = \mathbf{I}_P$ (equations 12 and 13) and $\bar{\mathbf{p}} = \mathbf{0}$ (equation 14). As shown by
- Takahashi et al. (2020, 2022), the CGLS produces stable estimates \tilde{p} for the parameter vector p (equation
- 552 3) in the presence of noisy potential-field data d. This is a well-known property of the CGLS method (e.g.,
- 553 Aster et al., 2019, p. 166).
- Algorithm 7: 2 FFT (FTT = kappa (4D)log2(4D)) lines 4 and 5, CONV (= 2FFT + 24D) + inner prod line
- 555 8, 3 saxpy (lines 12, 15 and 16), 2 inner prod (lines 14 and 17), 2 CONV lines 13 and 17:

$$f_{\text{CGLS}} = \kappa (16D) \log_2(4D) + 26D + \text{ITMAX} \left[\kappa (16D) \log_2(4D) + 58D \right].$$
 (67)

556 6.8 Direct deconvolution

The method proposed by Takahashi et al. (2020, 2022) can be reformulated to avoid the iterations of the conjugate gradient method. This alternative formulation consists in considering that $\mathbf{v} = \mathbf{p}$ and $\mathbf{w} = \mathbf{d}$ in equation 52, where \mathbf{p} is the parameter vector (equation 3) and \mathbf{d} the observed data vector. In this case, the equality "=" in equation 52 becomes an approximation" Then, equation 64 is manipulated to obtain

$$\mathcal{V}_c \approx \mathcal{F}_{2N_B}^* \left[\left(\mathcal{F}_{2N_B} \, \mathcal{W}_c \, \mathcal{F}_{2N_b} \right) \circ \check{\mathcal{L}} \right] \mathcal{F}_{2N_b}^* ,$$
 (68)

561 where

$$\overset{\sim}{\mathcal{L}} = \mathcal{L}^* \oslash (\mathcal{L} \circ \mathcal{L}^* + \zeta \mathbf{1}) ,$$
(69)

- 562 1 is a $4D \times 4D$ matrix of ones, " \oslash " denotes entrywise division and ζ is a positive scalar. Note that $\zeta = 0$
- leads to $1 \oslash \mathcal{L}$. In this case, the entrywise division may be problematic due to the elements of \mathcal{L} having
- absolute value equal or close to zero. So, a small ζ is set to avoid this problem in equation 69. Next, we use
- 565 \mathcal{L} to obtain a matrix \mathcal{V}_c from equation 68. Finally, the elements of the estimated parameter vector $\tilde{\mathbf{p}}$ are
- retrieved from the first quadrant of \mathcal{V}_c . This procedure represents a direct deconvolution (e.g., Aster et al.,
- 567 2019, p. 220) using a *Wiener filter* (e.g., Gonzalez and Woods, 2002, p. 263).
- Differently from the convolutional equivalent-layer method proposed by Takahashi et al. (2020, 2022),
- 569 the alternative direct deconvolution presented here produces an estimated parameter vector $\tilde{\mathbf{p}}$ directly from
- 570 the observed data d, in a single step, avoiding the conjugate gradient iterations. On the other hand, the
- alternative method presented here requires estimating a set of tentative parameter vectors $\tilde{\mathbf{p}}$ for different
- 572 predefined ζ . Besides, there must be criterion to chose the best $\tilde{\mathbf{p}}$ from this tentative set. This can be
- 573 made, for example, by using the well-known *L-curve* (Hansen, 1992). From a computational point of view,
- 574 the number of CGLS iterations in the method proposed by Takahashi et al. (2020, 2022) is equivalent to
- 575 the number of tentative estimated parameter vectors required to form the L-curve in the proposed direct
- 576 deconvolution.
- 577 Direct deconvolution: FFT to compute L, LL* eq 69 (24D), L*/(LL*) (24D), CONV (2FFT + 24D):

$$f_{\text{deconv.}} = \kappa (12D) \log_2(4D) + 72D$$
 (70)

CONFLICT OF INTEREST STATEMENT

- 578 The authors declare that the research was conducted in the absence of any commercial or financial
- 579 relationships that could be construed as a potential conflict of interest.

AUTHOR CONTRIBUTIONS

- 580 The Author Contributions section is mandatory for all articles, including articles by sole authors. If an
- 581 appropriate statement is not provided on submission, a standard one will be inserted during the production
- 582 process. The Author Contributions statement must describe the contributions of individual authors referred
- 583 to by their initials and, in doing so, all authors agree to be accountable for the content of the work. Please
- 584 see here for full authorship criteria.

FUNDING

- 585 Diego Takahashi was supported by a Post-doctoral scholarship from CNPq (grant 300809/2022-0) Valéria
- 586 C.F. Barbosa was supported by fellowships from CNPq (grant 309624/2021-5) and FAPERJ (grant
- 587 26/202.582/2019). Vanderlei C. Oliveira Jr. was supported by fellowships from CNPq (grant 315768/2020-
- 588 7) and FAPERJ (grant E-26/202.729/2018).

ACKNOWLEDGMENTS

- 589 We thank the brazillian federal agencies CAPES, CNPq, state agency FAPERJ and Observatório Nacional
- 590 research institute and Universidade do Estado do Rio de Janeiro.

DATA AVAILABILITY STATEMENT

- 591 The datasets generated for this study can be found in the frontiers-paper Github repository link:
- 592 https://github.com/DiegoTaka/frontiers-paper.

REFERENCES

- 593 Aster, R. C., Borchers, B., and Thurber, C. H. (2019). Parameter Estimation and Inverse Problems
- 594 (Elsevier), 3 edn.
- 595 Barbosa, V. C. F., Silva, J. B., and Medeiros, W. E. (1997). Gravity inversion of basement relief using
- approximate equality constraints on depths. *Geophysics* 62, 1745–1757
- 597 Barnes, G. and Lumley, J. (2011). Processing gravity gradient data. GEOPHYSICS 76, I33–I47. doi:10.
- 598 1190/1.3548548
- 599 Blakely, R. J. (1996). Potential Theory in Gravity and Magnetic Applications (Cambridge University
- 600 press)
- 601 Bott, M. H. P. (1960). The use of Rapid Digital Computing Methods for Direct Gravity Interpretation
- of Sedimentary Basins. Geophysical Journal International 3, 63–67. doi:10.1111/j.1365-246X.1960.
- 603 tb00065.x
- 604 Chan, R. H.-F. and Jin, X.-Q. (2007). An introduction to iterative Toeplitz solvers, vol. 5 (SIAM)
- 605 Cordell, L. (1992). A scattered equivalent-source method for interpolation and gridding of potential-field
- data in three dimensions. *Geophysics* 57, 629–636

- Dampney, C. N. G. (1969). The equivalent source technique. *GEOPHYSICS* 34, 39–53. doi:10.1190/1.
- 608 1439996
- 609 Davis, P. J. (1979). *Circulant matrices* (John Wiley & Sons, Inc.)
- 610 Emilia, D. A. (1973). Equivalent sources used as an analytic base for processing total magnetic field
- profiles. *GEOPHYSICS* 38, 339–348. doi:10.1190/1.1440344
- 612 Golub, G. H. and Van Loan, C. F. (2013). *Matrix Computations*. Johns Hopkins Studies in the Mathematical
- Sciences (Johns Hopkins University Press), 4 edn.
- 614 Gonzalez, R. C. and Woods, R. E. (2002). Digital Image Processing (Prentice Hall), 2 edn.
- 615 Gonzalez, S. P., Barbosa, V. C. F., and Oliveira Jr., V. C. (2022). Analyzing the ambiguity of the remanent-
- magnetization direction separated into induced and remanent magnetic sources. *Journal of Geophysical*
- 617 Research: Solid Earth 127, 1–24. doi:10.1029/2022JB024151
- 618 Guspí, F., Introcaso, A., and Introcaso, B. (2004). Gravity-enhanced representation of measured geoid
- 619 undulations using equivalent sources. *Geophysical Journal International* 159, 1–8. doi:10.1111/j.
- 620 1365-246X.2004.02364.x
- 621 Guspí, F. and Novara, I. (2009). Reduction to the pole and transformations of scattered magnetic data using
- Newtonian equivalent sources. *GEOPHYSICS* 74, L67–L73. doi:10.1190/1.3170690
- Hansen, P. C. (1992). Analysis of discrete ill-posed problems by means of the 1-curve. SIAM Review 34,
- 624 561–580. doi:10.1137/1034115
- 625 Hansen, R. O. and Miyazaki, Y. (1984). Continuation of potential fields between arbitrary surfaces.
- 626 GEOPHYSICS 49, 787–795. doi:10.1190/1.1441707
- 627 Horn, R. A. and Johnson, C. R. (1991). Topics in Matrix Analysis (Cambridge University Press), 1 edn.
- 628 Jain, A. K. (1989). Fundamentals of Digital Image Processing (Pearson), 1 edn.
- 629 Jirigalatu, J. and Ebbing (2019). A fast equivalent source method for airborne gravity gradient data.
- 630 Geophysics 84, G75–G82. doi:10.1190/GEO2018-0366.1
- 631 Kellogg, O. D. (1967). Foundations of Potential Theory (Springer-Verlag), reprint from the first edition of
- 632 1929 edn.
- 633 Kennett, B., Sambridge, M., and Williamson, P. (1988). Subspace methods for large inverse problems with
- 634 multiple parameter classes. *Geophysical Journal International* 94, 237–247
- 635 Leão, J. W. D. and Silva, J. B. C. (1989). Discrete linear transformations of potential field data. Geophysics
- 636 54, 497–507. doi:10.1190/1.1442676
- 637 Li, Y., Nabighian, M., and Oldenburg, D. W. (2014). Using an equivalent source with positivity for low-
- latitude reduction to the pole without striation. *GEOPHYSICS* 79, J81–J90. doi:10.1190/geo2014-0134.
- 639 1
- 640 Li, Y. and Oldenburg, D. W. (2010). Rapid construction of equivalent sources using wavelets.
- 641 GEOPHYSICS 75, L51–L59. doi:10.1190/1.3378764
- 642 Mendonça, C. A. (2020). Subspace method for solving large-scale equivalent layer and density mapping
- problems. *GEOPHYSICS* 85, G57–G68. doi:10.1190/geo2019-0302.1
- 644 Mendonça, C. A. and Silva, J. B. C. (1994). The equivalent data concept applied to the interpolation of
- potential field data. *Geophysics* 59, 722–732. doi:10.1190/1.1443630
- 646 Menke, W. (2018). Geophysical data analysis: Discrete inverse theory (Elsevier), 4 edn.
- 647 Oldenburg, D., McGillivray, P., and Ellis, R. (1993). Generalized subspace methods for large-scale inverse
- 648 problems. *Geophysical Journal International* 114, 12–20
- 649 Oliveira Jr., V. C., Barbosa, V. C. F., and Uieda, L. (2013). Polynomial equivalent layer. GEOPHYSICS 78,
- 650 G1-G13. doi:10.1190/geo2012-0196.1

- Press, W. H., Teukolsky, S. A., Vetterling, W. T., and Flannery, B. P. (2007). Numerical recipes: the art of 651 scientific computing (Cambridge University Press), 3 edn. 652
- Reis, A. L. A., Oliveira Jr., V. C., and Barbosa, V. C. F. (2020). Generalized positivity constraint on 653 magnetic equivalent layers. Geophysics 85, 1-45. doi:10.1190/geo2019-0706.1 654
- Roy, A. (1962). Ambiguity in geophysical interpretation. GEOPHYSICS 27, 90–99. doi:10.1190/1. 655 1438985 656
- Silva, J. B. C. (1986). Reduction to the pole as an inverse problem and its application to low-latitude 657 anomalies. GEOPHYSICS 51, 369–382. doi:10.1190/1.1442096 658
- Siqueira, F., Oliveira Jr., V. C., and Barbosa, V. C. F. (2017). Fast iterative equivalent-layer technique for 659
- gravity data processing: A method grounded on excess mass constraint. GEOPHYSICS 82, G57–G69. 660
- doi:10.1190/GEO2016-0332.1 661
- Skilling, J. and Bryan, R. (1984). Maximum entropy image reconstruction-general algorithm. Monthly 662 Notices of the Royal Astronomical Society, Vol. 211, NO. 1, P. 111, 1984 211, 111 663
- Soler, S. R. and Uieda, L. (2021). Gradient-boosted equivalent sources. Geophysical Journal International 664 227, 1768–1783. doi:10.1093/gji/ggab297 665
- Takahashi, D., Oliveira Jr., V. C., and Barbosa, V. C. (2022). Convolutional equivalent layer for magnetic 666 data processing. Geophysics 87, 1–59 667
- Takahashi, D., Oliveira Jr., V. C., and Barbosa, V. C. F. (2020). Convolutional equivalent layer for gravity 668 data processing. GEOPHYSICS 85, G129–G141. doi:10.1190/geo2019-0826.1 669
- 670 van der Sluis, A. and van der Vorst, H. A. (1987). Numerical solution of large, sparse linear algebraic
- systems arising from tomographic problems. In Seismic tomography with applications in global 671
- seismology and exploration geophysics, ed. G. Nolet (D. Reidel Publishing Company), chap. 3. 49–83 672
- Van Loan, C. F. (1992). Computational Frameworks for the fast Fourier transform. Frontiers in Applied 673 Mathematics (SIAM) 674
- Xia, J. and Sprowl, D. R. (1991). Correction of topographic distortion in gravity data. Geophysics 56, 675 676 537-541
- Xia, J., Sprowl, D. R., and Adkins-Heljeson, D. (1993). Correction of topographic distortions in potential-677 field data; a fast and accurate approach. Geophysics 58, 515-523. doi:10.1190/1.1443434 678
- Zhao, G., Chen, B., Chen, L., Liu, J., and Ren, Z. (2018). High-accuracy 3D Fourier forward modeling 679
- of gravity field based on the Gauss-FFT technique. Journal of Applied Geophysics 150, 294-303. 680
- 681 doi:10.1016/j.jappgeo.2018.01.002
- Zidarov, D. (1965). Solution of some inverse problems of applied geophysics. Geophysical Prospecting 682 683 13, 240–246. doi:10.1111/j.1365-2478.1965.tb01932.x

ALGORITHMS

Algorithm 1: Generic pseudo-code for the CGLS applied to the overdetermined problem (equation 22) for the particular case in which $\mathbf{H} = \mathbf{I}_P$ (equation 9 and subsection 2.2), $\mu = 0$ (equation 11), $\mathbf{W}_d = \mathbf{I}_D$ (equation 12) and $\bar{\mathbf{p}} = \mathbf{0}$ (equation 14), where \mathbf{I}_P and \mathbf{I}_D are the identities of order P and D, respectively.

```
Initialization:
 1 Compute G;
 2 Set \mathbf{r} = \mathbf{d} and compute \delta = \|\mathbf{r}\|/D;
 3 Compute \boldsymbol{\vartheta} = \mathbf{G}^{\top} \mathbf{r} and \rho_0 = \boldsymbol{\vartheta}^{\top} \boldsymbol{\vartheta};
 4 Set \tilde{\mathbf{p}} = \mathbf{0}, \tau = 0 and \boldsymbol{\eta} = \mathbf{0};
 5 \ m = 1;
     while (\delta > \epsilon) and (m < \text{ITMAX}) do
              Update \eta \leftarrow \vartheta + \tau \eta;
              Compute \nu = G \eta;
 8
              Compute v = \rho_0/(\boldsymbol{\nu}^{\top}\boldsymbol{\nu});
 9
              Update \tilde{\mathbf{p}} \leftarrow \tilde{\mathbf{p}} + v \, \boldsymbol{\eta};
10
              Update \mathbf{r} \leftarrow \mathbf{r} - \upsilon \, \boldsymbol{\nu} and \delta \leftarrow \|\upsilon \, \boldsymbol{\nu}\|/D;
11
              Compute \boldsymbol{\vartheta} = \mathbf{G}^{\top} \mathbf{r} and \rho = \boldsymbol{\vartheta}^{\top} \boldsymbol{\vartheta};
12
              Compute \tau = \rho/\rho_0;
13
              Update \rho_0 \leftarrow \rho;
14
              m \leftarrow m + 1;
15
16 end
```

Algorithm 2: Generic pseudo-code for the method proposed by Leão and Silva (1989).

```
1 Set the indices i^m for each data window, m \in \{1 : M\};
2 Set the indices \mathbf{j}^m for each source window, m \in \{1 : M\};
3 Set the constant depth z_0 + \Delta z_0 for all equivalent sources;
4 Compute the vector a' associated with the desired potential-field transformation;
5 Compute the matrix G':
6 Compute the matrix B' (equation 34);
7 Compute the vector (\mathbf{a}')^{\top} \mathbf{B}':
m = 1:
  while m < M do
       Compute t_c^m (equation 33); m \leftarrow m+1;
12 end
```

Algorithm 3: Generic pseudo-code for the method proposed by Soler and Uieda (2021).

```
Initialization:
 1 Set the indices i^m for each data window, m \in \{1 : M\};
 2 Set the indices j^m for each source window, m \in \{1 : M\};
 3 Set the depth of all equivalent sources;
 4 Set a D \times 1 residuals vector \mathbf{r} = \mathbf{d};
5 Set a P \times 1 vector \tilde{\mathbf{p}} = \mathbf{0};
 6 \ m = 1;
    while m < M 	ext{ do}
 7
           Set the matrix \mathbf{W}_d^m;
           Compute the matrix G^m;
 9
           Compute \tilde{\mathbf{p}}^m (equation 36);
10
           \tilde{\mathbf{p}}[\mathbf{j}^m] \leftarrow \tilde{\mathbf{p}}[\mathbf{j}^m] + \tilde{\mathbf{p}}^m;
11
           \mathbf{r} \leftarrow \mathbf{r} - \mathbf{G}[:, \mathbf{j}^m] \, \tilde{\mathbf{p}}^m;

m \leftarrow m + 1;
12
13
14 end
```

Algorithm 4: Generic pseudo-code for the method proposed by Cordell (1992).

Initialization:

```
1 Compute a D \times 1 vector \Delta z whose i-th element \Delta z_i is a vertical distance controlling the depth of
      the i-th equivalent source, i \in \{1 : D\};
 2 Set a tolerance \epsilon;
 3 Set a maximum number of iterations ITMAX;
 4 Set a D \times 1 residuals vector \mathbf{r} = \mathbf{d};
 5 Set a D \times 1 vector \tilde{\mathbf{p}} = \mathbf{0};
 6 Define the maximum absolute value r_{\text{max}} in \mathbf{r};
    while (r_{\text{max}} > \epsilon) and (m < \text{ITMAX}) do
          Define the coordinates (x_{\text{max}}, y_{\text{max}}, z_{\text{max}}) and index i_{\text{max}} of the observation point associated with
          \tilde{\mathbf{p}}[i_{\text{max}}] \leftarrow \tilde{\mathbf{p}}[i_{\text{max}}] + (r_{\text{max}} \Delta \mathbf{z}[i_{\text{max}}]);
10
          \mathbf{r} \leftarrow \mathbf{r} - (\mathbf{G}[:,i_{\mathtt{max}}]\,\tilde{\mathbf{p}}[i_{\mathtt{max}}]) ;
11
          Define the new r_{\text{max}} in \mathbf{r};
12
          m \leftarrow m + 1;
14 end
```

Algorithm 5: Generic pseudo-code for the method proposed by Mendonça and Silva (1994).

```
Initialization:
 1 Set a regular grid of P equivalent sources at a horizontal plane z_0;
 2 Set a tolerance \epsilon;
 3 Set a D \times 1 residuals vector \mathbf{r} = \mathbf{d};
 4 Define the maximum absolute value r_{max} in r;
 5 Define the index i_{max} of r_{max};
6 Define the list of indices i_r of the remaining data in r;
 7 Define \mathbf{d}_e = \mathbf{d}[i_{\text{max}}];
    Compute (\mathbf{F} + \mu \mathbf{I}_{D_e}) and \mathbf{G}_e;
     Compute \tilde{\mathbf{p}} (equation 40);
10 Compute \mathbf{r} = \mathbf{d}[\mathbf{i}_r] - \mathbf{G}[\mathbf{i}_r, :] \tilde{\mathbf{p}};
11 Define the maximum absolute value r_{\text{max}} in r;
    while (r_{\text{max}} > \epsilon) do
           Define the index i_{max} of r_{max};
Define the list of indices i_r of the remaining elements in r;
13
14
                      \left| egin{array}{c} \mathbf{d}_e \ \mathbf{d}[i_{	exttt{max}}] 
ight| \; ;
15
           Update (\mathbf{F} + \mu \mathbf{I}_{D_e}) and \mathbf{G}_e;
16
           Update \hat{\mathbf{p}} (equation 40);
17
           Update \mathbf{r} = \mathbf{d}[\mathbf{i}_r] - \mathbf{G}[\mathbf{i}_r,:] \tilde{\mathbf{p}};
18
           Define the maximum absolute value r_{max} in r;
19
20 end
```

Algorithm 6: Generic pseudo-code for the iterative method proposed by Siqueira et al. (2017). The symbol "o" denotes the entrywise or Hadamard product (e.g., Horn and Johnson, 1991, p. 298) and σ is a $P \times 1$ vector whose j-th element is the ratio of a predefined element of area centered at the j-th equivalent source and the term $2\pi\gamma$, where γ is the gravitational constant.

```
Initialization:
```

```
1 Set P equivalent sources on a horizontal plane z_0;
 2 Set a tolerance \epsilon;
 3 Set a maximum number of iterations ITMAX;
 4 Set an auxiliary vector \sigma;
 5 Compute \tilde{\mathbf{p}} = \boldsymbol{\sigma} \circ \mathbf{d};
 6 Compute G (equation 3);
 7 Compute \mathbf{r} = \mathbf{d} - \mathbf{G}\,\tilde{\mathbf{p}};
 8 Compute \delta = \|\mathbf{r}\|/D;
 9 m = 1:
10 while (\delta > \epsilon) and (m < ITMAX) do
          Compute \Delta \mathbf{p} = \boldsymbol{\sigma} \circ \mathbf{r};
11
          Update \tilde{\mathbf{p}} \leftarrow \tilde{\mathbf{p}} + \Delta \mathbf{p};
12
          Compute \nu = G \Delta p;
13
          Update \mathbf{r} \leftarrow \mathbf{r} - \boldsymbol{\nu};
14
          Compute \delta = \|\boldsymbol{\nu}\|/D;
15
          m \leftarrow m+1:
17 end
```

Algorithm 7: Generic pseudo-code for the convolutional equivalent-layer method proposed by Takahashi et al. (2020, 2022).

Initialization:

- 1 Set the regular grid of P equivalent sources on a horizontal plane z_0 ;
- 2 Set a tolerance ϵ and a maximum number of iterations ITMAX;
- 3 Compute the first column G[:, 1] and row G[1, :] of the sensitivity matrix G (equation 3) for the particular case in which it has a BTTB structure (equation 51);
- 4 Rearrange the elements of G[:, 1] into matrix C, compute its 2D-DFT via 2D-FFT and multiply by $\sqrt{4D}$ to obtain a matrix \mathcal{L}' (equation 66);
- 5 Rearrange the elements of G[1,:] into matrix C, compute its 2D-DFT via 2D-FFT and multiply by $\sqrt{4D}$ to obtain a matrix \mathcal{L}'' (equation 66);

```
6 Set \tilde{\mathbf{p}} = \mathbf{0};
 7 Set \mathbf{r} = \mathbf{d} and compute \delta = \|\mathbf{r}\|/D;
 8 Compute \boldsymbol{\vartheta} = \mathbf{G}^{\top} \mathbf{r} (Algorithm 8) and \rho_0 = \boldsymbol{\vartheta}^{\top} \boldsymbol{\vartheta};
     Set \tau = 0 and \eta = 0;
10 m = 1;
     while (\delta > \epsilon) and (m < \text{ITMAX}) do
11
             Update \eta \leftarrow \vartheta + \tau \eta;
12
             Compute \nu = G \eta (Algorithm 8);
13
             Compute v = \rho_0/(\boldsymbol{\nu}^{\top}\boldsymbol{\nu});
14
             Update \tilde{\mathbf{p}} \leftarrow \tilde{\mathbf{p}} + v \, \boldsymbol{\eta};
15
             Update \mathbf{r} \leftarrow \mathbf{r} - \upsilon \, \boldsymbol{\nu} and \delta \leftarrow \|\upsilon \, \boldsymbol{\nu}\|/D;
16
             Compute \boldsymbol{\vartheta} = \mathbf{G}^{\top} \mathbf{r} (Algorithm 8) and \rho = \boldsymbol{\vartheta}^{\top} \boldsymbol{\vartheta};
17
             Compute \tau = \rho/\rho_0;
18
             Update \rho_0 \leftarrow \rho;
19
             m \leftarrow m + 1;
20
21 end
```

Algorithm 8: Pseudo-code for computing the generic matrix-vector products given by equations 52 and 53 via fast 2D discrete convolution for a given vector \mathbf{v} (equation 54) and matrix \mathcal{L} (equation 66).

- Rearrange the elements of v (equations 52 and 54) into the matrix \mathcal{V}_c (equation 65);
- 2 Compute $\mathcal{F}_{2N_B} \mathcal{V}_c \mathcal{F}_{2N_b}$ via 2D-FFT;
- 3 Compute the Hadamard product with matrix \mathcal{L} (equation 66);
- 4 Compute 2D-IDFT via 2D-FFT to obtain matrix $\tilde{\mathcal{W}}_c$ (65);
- 5 Retrieve w (equations 52 and 54) from \mathbf{w}_c (equations 55–57);

8 TABLES

Reference	Term	flops
eq. 10	$_{ m GH}$	2DQP
eq. 15	$\mathbf{H} ilde{\mathbf{q}}$	2PQ
eq. 22	$\left(\mathbf{G}\mathbf{H}\right)^{\top}\underline{\left(\mathbf{G}\mathbf{H}\right)}$	$2Q^2D$
eq. 22	$\left(\mathbf{G}\mathbf{H}\right)^{\top}\boldsymbol{\delta}_{d}\;_$	2QD
eq. 23	$(\mathbf{G}\mathbf{H})\left(\mathbf{\underline{G}}\mathbf{H}\right)^{\perp}$	$2D^2Q$
eq. 23	$(\mathbf{G}\mathbf{H})^{\top}\mathbf{u}$	2QD
eq. 25	lower triangle of \mathcal{G}	$D^3/3 \text{ or } Q^3/3$
eq. 26	solve triangular systems	$2D^2$ or $2Q^2$
Alg. 1	$oldsymbol{\eta} \leftarrow oldsymbol{artheta} + au oldsymbol{\eta}$	2Q
Alg. 1	$\boldsymbol{\vartheta}^{\top}\boldsymbol{\vartheta}$	2Q

Table 1. Total number of flops associated with some useful terms according to Golub and Van Loan (2013, p. 12). The flops associated with equations 25 and 26 depends if the problem is over or underdetermined. Note that P=Q for the case in which $\mathbf{H}=\mathbf{I}_P$ (subsection 2.2). The term associated with Algorithm 1 is a vector update called saxpy (Golub and Van Loan, 2013, p. 4). The terms defined here are references to compute the total number of flops throughout the manuscript.