

3D MAGNETIC MODELLING FOR ELIPSOIDS

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Abstract. In this work we will present results from a numerical modeling of the magnetic field and total field anomaly, represented by triaxial and prolate ellipsoidal sources. Such approach provide analytical results for anisotropy of magnetic susceptibility as well as for self-demagnetization effects, which can be easily adapted for distinctive geologic structures - hence being an useful tool for educational (e.g., potential methods and rock magnetism) and applied geophysics (e.g., characterization of high magnetic susceptibility, mineralized bodies) purposes. Numerical tests by means of a Python code (currently under development) allowed us to compared the effects of different geometries (ellipsoidal sources, spheres, dipole lines and elliptic cylinders), which were used to validate our computational implementation. This code will be freely available to the scientific community by the end of the year.

1 Introduction

In the 19th century, Dirichlet (1839) manage to solve the problem of the potential of a homogeneous ellipsoid with semi-axes $a \geq b \geq c$ by using the first and second Legendre's normal elliptic integrals. Later, Maxwell (1881) showed that only bodies of second degree can be uniformly magnetized when placed in a uniform field. He also presented the formulae of the demagnetizing field in a uniformly magnetized ellipsoid of revolution.

The formulas for the demagnetizing factors of the general ellipsoid alongside a table of values for different proportions of semi-axis was published by Osborn (1945). In the same year, Stoner (1945) presented a more simple and friendly formulas of the demagnetizing factors for an ellipsoid, with step-by-step deduction. This work includes formulas for different parametrizations for the ellipsoid (triaxial, prolate, oblate).

A decade later, Peake and Davy (1953) published a paper with the formulas for the external field of ellipsoid of revolution. Some years after that, Chang (1961) showed the formulas for the external field of triaxial, prolate and oblate ellipsoids. Both used the Dirichlet solution in terms of elliptical integrals for external fields. More recently, Tejedor et al. (1995), gave the calculations of the external field generated by ellipsoids using a variational method, in contrast to the usual method of Dirichlet.

In the decade of 1970 Farrar (1979) showed in his work the value of the ellipsoidal model in geophysics exploration, where it was used for analysis of magnetic anomalies in Tennant Creeks's gold mine in Australia. The ellipsoid was the most correct geometry to model bodies of lenticular pipe shape and the most appropriate to handle self-demagnetization of bodies of high susceptibility - magnetic susceptibility higher than 0.1 SI (Clark, 2014). Guo et al. (1998), showed three case studies of the importance of self-demagnetization of high susceptibilities bodies when modelling magnetic sources. A magnetic-iron

deposit, volcanic-hosted iron deposit and an ultramafic-hosted nickel and copper deposits. The same author years later, Guo et al. (2001), published another paper analysing the errors of what an uncorrected by demagnetization data can affect on anomalies interpretation.

The magnetic field resultant of an ellipsoid was presented the first time in full modeling in the space domain by Emerson et al. (1985), along a compilation of several others geometric shapes. Through a consistent notation this work tapped a hole in the literature that missed the ellipsoidal model. This model, however, consists in a division between prolate and oblate ellipsoids, with a few differences in the algorithm. A generalized model for triaxial ellipsoids was published by Clark et al. (1986). Both in Emerson et al. and Clark et al. works, they were the firsts to include the calculations of the magnetic field in a different coordinate system other than with the position $(0, 0, 0)$ in the center of the ellipsoid (through a matrix of rotation). However, the codes presented for HP-41C are outdated and doesn't have a description of the algorithm, diffculting the reproduction by the scientific community. This may be the reason why the magnetic field modelling of the ellipsoid, in general, is only used through commercial softwares.

2 Methodology

The total-field anomaly can be described as a difference between the total field vector and the induced magnetization of the crust of the Earth, that includes any anomalous distribution of magnetization that can occur by magnetized bodies in subsurface. This difference can be written as:

$$\Delta T_i^0 = ||T_i|| - ||F_i||, \quad (1)$$

where ΔT_i^0 is the observed vector of total-field anomaly in the i -th position (x_i, y_i, z_i) , $i = 1, \dots, N$, F_i is the geomagnetic field and $||\cdot||$ is the Euclidian norm. The total-field vector T_i is:

$$T_i = F_i + B_i, \quad (2)$$

and B_i is the total induced magnetization vector produced by all sources of anomalous susceptibility distribution. Considering F_i a constant vector, F_0 for local or regional scale, and that $||F_0|| \gg ||B_0||$, since, B_i is a small perturbation of the magnetic field F_i , we can approximate the euclidian norm of the vector T_i by a Taylor's expansion of first order:

$$||T_i|| \approx ||F_0 + B_i|| \approx ||F_0|| + F^T B_i \quad (3)$$

where T indicates transposition and

$$\hat{F} = \frac{F_o}{||F_0||} \quad (4)$$

is a unit vector that represents the gradient of the function $||T_i||$ in relation to the vector's components T_i . This way, we can approximate the Eq.(1) of total-field anomaly to:

$$\Delta T \approx \hat{F}^T B_i, \quad i = 1, \dots, N \quad (5)$$

2.1 The Forward model of a sphere

2.2 The Forward model of a triaxial ellipsoid

The implementation of the forward problem of a triaxial ellipsoid's magnetic field (three semi-axis $a > b > c$) is done in a new coordinate system, where its origin is the center of this ellipsoidal body.

This new coordinate system (x_1, x_2, x_3) is defined by the unit vectors $\hat{v}_h (h = 1, 2, 3)$ with the respect with the geographic axis x, y e z :

$$\hat{v}_1 = (l_1, m_1, n_1) = (-\cos \alpha \cos \delta, -\sin \alpha \cos \delta, -\sin \delta) \quad (6)$$

$$\begin{aligned} \hat{v}_2 = (l_2, m_2, n_2) = & (\cos \alpha \cos \gamma \sin \delta + \sin \alpha \sin \gamma, \\ & \sin \alpha \cos \gamma \sin \delta - \cos \alpha \sin \gamma, -\cos \gamma \cos \delta) \end{aligned} \quad (7)$$

$$\begin{aligned} \hat{v}_3 = (l_3, m_3, n_3) = & (\sin \alpha \cos \gamma - \cos \alpha \sin \gamma \sin \delta, \\ & -\cos \alpha \cos \gamma - \sin \alpha \sin \gamma \sin \delta, \sin \gamma \cos \delta) \end{aligned} \quad (8)$$

The angles referring to the unit vectors are determined by the orientations of the ellipsoid's semi-axis. The angle α is the azimuth of semi-major axis (a) plus 180° . While δ is the inclination of semi-major axis (a) in relation to the geographic plane. Lastly, γ is the angle between the semi-mid axis (b) and the vertical projection of the ellipsoid's center with the geographic plane.

Thus, the coordinates of the body's semi-axis are given by:

$$x_h = (x - x_c)l_h + (y - y_c)m_h + (z - z_c)n_h \quad (h = 1, 2, 3) \quad (9)$$

Where x_c, y_c e z_c are the coordinates of the ellipsoid's center in the geographic system x, y e z .

For an ellipsoid of semi-axis $a > b > c$, the equation that defines your surface is:

$$\frac{x_1^2}{(a^2 + s)} + \frac{x_2^2}{(b^2 + s)} + \frac{x_3^2}{(c^2 + s)} = 1 \quad (10)$$

The parameter s controls the ellipsoid form. When s gets close to ∞ the equation (10) tends to the sphere equation of radius $r = \sqrt{\lambda}$. When $s = -c^2$, the last term of the ellipsoid's equation is less than zero and it becomes the equation of a circle.

There is, however, a set of values for $s(\lambda, \mu, \nu)$, which are roots of the cubic equation:

$$s^3 + p_2 s^2 + p_1 s + p_0 = 0 \quad (11)$$

This set of roots, called elipsoidal coordinates, correspond to the parameters of a point (x_1, x_2, x_3) which are under the intersection of three ortogonal surfaces related to the body coordinates. Their expressions are:

$$\lambda = 2\sqrt{\left(\frac{-p}{3}\right)\cos\left(\frac{\theta}{3}\right) - \frac{p_2}{3}} \quad (12)$$

$$5 \quad \mu = -2\sqrt{\left(\frac{-p}{3}\right)\cos\left(\frac{\theta}{3} + \frac{\pi}{3}\right) - \frac{p_2}{3}} \quad (13)$$

$$\mu = -2\sqrt{\left(\frac{-p}{3}\right)\cos\left(\frac{\theta}{3} - \frac{\pi}{3}\right) - \frac{p_2}{3}} \quad (14)$$

Where:

$$p_0 = a^2b^2c^2 - b^2c^2x_1^2 - c^2a^2x_2^2 - a^2b^2x_3^2 \quad (15)$$

$$p_1 = a^2b^2 + b^2c^2 + c^2a^2 - (b^2 + c^2)x_1^2 - (c^2 + a^2)x_2^2 - (a^2 + b^2)x_3^2 \quad (16)$$

$$10 \quad p_2 = a^2 + b^2 + c^2 - x_1^2 - x_2^2 - x_3^2 \quad (17)$$

$$p = p_1 - \frac{p_2^2}{3} \quad (18)$$

$$q = p_0 - \frac{p_1p_2}{3} + 2\left(\frac{p_2}{3}\right)^3 \quad (19)$$

$$\theta = \cos^{-1} \left[\frac{-q}{2} \sqrt{\left(\frac{-p}{3}\right)^3} \right] \quad (20)$$

The calculation of the largest root λ of the equation (11) is essential, since the magnetic field depends on the spatial derivatives of the equation 10, where s admits the value of λ .

$$\frac{\partial \lambda}{\partial x_1} = \frac{2x_1/(a^2 + \lambda)}{\left(\frac{x_1}{a^2 + \lambda}\right)^2 + \left(\frac{x_2}{b^2 + \lambda}\right)^2 + \left(\frac{x_3}{c^2 + \lambda}\right)^2} \quad (21)$$

$$\frac{\partial \lambda}{\partial x_2} = \frac{2x_2/(b^2 + \lambda)}{\left(\frac{x_1}{a^2 + \lambda}\right)^2 + \left(\frac{x_2}{b^2 + \lambda}\right)^2 + \left(\frac{x_3}{c^2 + \lambda}\right)^2} \quad (22)$$

$$\frac{\partial \lambda}{\partial x_3} = \frac{2x_3/(c^2 + \lambda)}{\left(\frac{x_1}{a^2 + \lambda}\right)^2 + \left(\frac{x_2}{b^2 + \lambda}\right)^2 + \left(\frac{x_3}{c^2 + \lambda}\right)^2} \quad (23)$$

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The desmagnetization factors are given by:

$$N_1 = \frac{4\pi abc}{(a^2 - b^2)\sqrt{(a^2 - c^2)}} [F(\theta, k) - E(\theta, k)] \quad (24)$$

$$N_2 = \frac{4\pi abc\sqrt{(a^2 - c^2)}}{(a^2 - b^2)(b^2 - c^2)} \left[E(\theta, k) - \left(\frac{b^2 - c^2}{a^2 - c^2}\right) F(\theta, k) - \frac{c(a^2 - b^2)}{ab\sqrt{(a^2 - c^2)}} \right] \quad (25)$$

$$N_3 = \frac{4\pi abc}{(b^2 - c^2)\sqrt{(a^2 - c^2)}} \left[\frac{b\sqrt{(a^2 - c^2)}}{ac} - E(\theta, k) \right] \quad (26)$$

10 Where $F(\theta, k)$ e $E(\theta, k)$ are first and second Legendre's normal elliptical integrals, respectively. To calculate k and θ we have the following expressions:

$$k = \sqrt{\left(\frac{a^2 - b^2}{a^2 - c^2}\right)} \quad (27)$$

$$\theta = \cos^{-1}(c/a) \quad (0 \leq \theta \leq \pi/2) \quad (28)$$

The susceptibility tensor matrix is:

$$15 \quad k_{ij} = \sum_r k_r (L_r l_i + M_r m_i + N_r n_i)(L_r l_j + M_r m_j + N_r n_j) \quad (r = 1, 2, 3) \quad (29)$$

The Earth's field vector, F , and the remanent magnetization d , also must be converted to the body's coordinates:

$$F_h = F(l l_h + m m_h + n n_h) \quad (30)$$

$$(J_N)_h = J_N(l_N l_h + m_n m_h + n_N n_h) \quad (31)$$

Note that L_r, M_r, N_r ($r = 1, 2, 3$), l, m, n e l_N, m_N, n_N depends on their respective vectors of inclinations and declinations:

$$5 \quad L_R, l, l_N = \cos(dec) \cos(inc) \quad (32)$$

$$M_R, m, m_N = \sin(dec) \cos(inc) \quad (33)$$

$$N_R, n, n_N = \sin(inc) \quad (34)$$

In the case that the body has a very low susceptibility ($\chi < 0.1$ SI) the self-demagnetization is negligible and the resultant magnetic vector is given by:

$$10 \quad \tilde{J}_R = K\tilde{F} + \tilde{J}_{NRM} \quad (35)$$

For values bigger than 0.1 SI the resultant magnetic vector is:

$$\tilde{J}_{Rc} = A^{-1} \tilde{J}_R \quad (36)$$

Where:

$$A = I + KN = \begin{bmatrix} 1 + k_{11}N_1 & k_{12}N_2 & k_{13}N_3 \\ k_{21}N_1 & 1 + k_{22}N_2 & k_{23}N_3 \\ k_{31}N_1 & k_{23}N_2 & 1 + k_{33}N_3 \end{bmatrix} \quad (37)$$

15 This way, the components of the magnetic field produced by a triaxial ellipsoid in the body's coordinates outside of it is:

$$\Delta B_1 = f_1 \frac{\partial \lambda}{\partial x_1} - 2\pi abc J_1 A(\lambda) \quad (38)$$

$$\Delta B_2 = f_1 \frac{\partial \lambda}{\partial x_2} - 2\pi abc J_2 B(\lambda) \quad (39)$$

$$\Delta B_3 = f_1 \frac{\partial \lambda}{\partial x_3} - 2\pi abc J_3 C(\lambda) \quad (40)$$

where:

$$5 \quad f_1 = \frac{2\pi abc}{\sqrt{[(a^2 + \lambda)(b^2 + \lambda)(c^2 + \lambda)]}} \left[\frac{J_1 x_1}{a^2 + \lambda} + \frac{J_2 x_2}{b^2 + \lambda} + \frac{J_3 x_3}{c^2 + \lambda} \right] \quad (41)$$

$$A(\lambda) = \frac{2}{(a^2 - b^2)\sqrt{(a^2 - c^2)}} [F(\theta', k) - E(\theta', k)] \quad (42)$$

$$B(\lambda) = \frac{2\sqrt{(a^2 - c^2)}}{(a^2 - b^2)(b^2 - c^2)} \left[E(\theta', k) - \left(\frac{b^2 - c^2}{a^2 - c^2} \right) F(\theta', k) - \frac{k^2 \sin \theta' \cos \theta'}{\sqrt{(1 - k^2 \sin^2 \theta')}} \right] \quad (43)$$

$$C(\lambda) = \frac{2}{(b^2 - c^2)\sqrt{(a^2 - c^2)}} \left[\frac{\sin \theta' \sqrt{(1 - k^2 \sin^2 \theta')}}{\cos \theta'} - E(\theta', k) \right] \quad (44)$$

And:

$$10 \quad \theta' = \sin^{-1} \left(\frac{a^2 - c^2}{a^2 + \lambda} \right)^{0.5} \quad (0 \leq \theta' \leq \pi/2) \quad (45)$$

Both $F(\theta', k)$ as $E(\theta', k)$ are again the first and second normal Legendre's elliptical integrals. While $A(\lambda), B(\lambda), C(\lambda)$ are analytic solutions of the integrals of the potential equation of an ellipsoid. This problem was solved by Dirichlet in 1839 (Clark et al., 1986) for the gravitational potential given by:

$$U_i(x_1, x_2, x_3) = \pi abc G \rho \int_0^\infty \left[1 - \frac{x_1^2}{a^2 + u} - \frac{x_2^2}{b^2 + u} - \frac{x_3^2}{c^2 + u} \right] \frac{du}{R(u)} \quad (46)$$

15 That can be rewritten as:

$$U_i(x_1, x_2, x_3) = \pi abc G \rho \int_0^\infty [D(\lambda) - A(\lambda) - B(\lambda) - C(\lambda)] \frac{du}{R(u)} \quad (47)$$

Where G is the gravitational constant and ρ the body's density.

Using the equations from (38) to (44), we can rearrange them in a matrix format that will be more computationally efficient.

Thus, the magnetic field generated by the ellipsoid is:

$$5 \quad b_i^j = 2\pi a_j b_j c_j \times [M_i^j - D_i^j] \times J^j \quad j = 1, \dots, L \quad (48)$$

With:

$$M_i^j = \frac{1}{\sqrt{[(a^2 + \lambda)(b^2 + \lambda)(c^2 + \lambda)]}} \begin{bmatrix} \frac{\partial \lambda}{\partial x_1} \frac{x_1}{a^2 + \lambda} & \frac{\partial \lambda}{\partial x_1} \frac{x_2}{b^2 + \lambda} & \frac{\partial \lambda}{\partial x_1} \frac{x_3}{c^2 + \lambda} \\ \frac{\partial \lambda}{\partial x_2} \frac{x_1}{a^2 + \lambda} & \frac{\partial \lambda}{\partial x_2} \frac{x_2}{b^2 + \lambda} & \frac{\partial \lambda}{\partial x_2} \frac{x_3}{c^2 + \lambda} \\ \frac{\partial \lambda}{\partial x_3} \frac{x_1}{a^2 + \lambda} & \frac{\partial \lambda}{\partial x_3} \frac{x_2}{b^2 + \lambda} & \frac{\partial \lambda}{\partial x_3} \frac{x_3}{c^2 + \lambda} \end{bmatrix}_{3 \times 3} \quad (49)$$

$$D_i^j = \begin{bmatrix} A(\lambda) & 0 & 0 \\ 0 & B(\lambda) & 0 \\ 0 & 0 & C(\lambda) \end{bmatrix} \quad (50)$$

For $j = 1, \dots, L$, the number of ellipsoids been modeled and $i = 1, \dots, N$ for the i th element of the calculated field. The total
 10 magnetic field then is:

$$B_i = \sum_{j=1}^L b_i^j, \quad i = 1, \dots, N \quad (51)$$

We must remember that every calculation so far was done using the body's coordinates, however the notation in equation (48) the magnetic field is already in geographics coordinates. To calculate each component of the magnetic field's vector back:

$$\Delta B_x = \Delta B_1 l_1 + \Delta B_2 l_2 + \Delta B_3 l_3 \quad (52)$$

$$15 \quad \Delta B_y = \Delta B_1 m_1 + \Delta B_2 m_2 + \Delta B_3 m_3 \quad (53)$$

$$\Delta B_z = \Delta B_1 n_1 + \Delta B_2 n_2 + \Delta B_3 n_3 \quad (54)$$

2.2.1 HEADING

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3 Conclusions

5 TEXT

Appendix A

A1

Author contributions. TEXT

Acknowledgements. TEXT

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