

# Hyperbolcity-Preserving Well-Balanced Stochastic Galerkin Method for Shallow Water Equations

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## Abstract

We study the stochastic Galerkin (SG) method for stochastic parameterized shallow water equations. Our work comprises the following aspects:

- A hyperbolcity-preserving stochastic Galerkin formulation for the shallow water equations using only the conserved variables.
- A sufficient condition to preserve the hyperbolcity, which is a stochastic variant of the deterministic positivity condition.
- A computationally tractable condition to guarantee the hyperbolcity.
- A central-upwind scheme that preserves both the hyperbolcity and the well-balanced property at discrete time levels.

## Motivations

- Uncertainties can enter the shallow water system, for example, via the noisy measurement of the bottom.
- A SG formulation of shallow water equations is not necessarily hyperbolic.
- A non-well-balanced scheme may lead to spurious oscillations on relatively coarse grid.

## Stochastic Parameterized Shallow Water System

$$\begin{aligned} \frac{\partial h}{\partial t} + \frac{\partial q^x}{\partial x} + \frac{\partial q^y}{\partial y} &= 0, \\ \frac{\partial q^x}{\partial t} + \frac{\partial}{\partial x} \left( \frac{(q^x)^2}{h} + \frac{gh^2}{2} \right) + \frac{\partial}{\partial y} \left( \frac{q^x q^y}{h} \right) &= -gh \frac{\partial B}{\partial x}, \\ \frac{\partial q^y}{\partial t} + \frac{\partial}{\partial x} \left( \frac{q^x q^y}{h} \right) + \frac{\partial}{\partial y} \left( \frac{(q^y)^2}{h} + \frac{gh^2}{2} \right) &= -gh \frac{\partial B}{\partial y}. \end{aligned}$$

$h$  is the water height,  $q_x$  and  $q_y$  are the  $x$ - and  $y$ -discharges, and  $B$  is the time-independent surface. All the variables are  $\xi$ -dependent random fields, e.g.,  $h = h(x, y, t, \xi)$ .

## Polynomial Chaos Expansion (PCE)

- An unknown random field  $z(x, y, t, \xi)$  is represented in the  $L^2_\rho$  orthonormal basis  $\{\phi_k\}_{k \in \mathbb{R}}$ , where  $\rho := \rho(\xi)$  is the density of the random parameter.

$$z(x, y, t, \xi) = \sum_{k=1}^{\infty} \hat{z}_k(x, y, t) \phi_k(\xi),$$

- $K$ -term truncated PCE:

$$\Pi_\Lambda[z] := \sum_{k=1}^K \hat{z}_k(x, y, t) \phi_k(\xi),$$

where  $\Lambda$  is the index set for the (possibly multivariate) polynomials, the cardinality of  $\Lambda$  is  $K$ , and  $\phi_1(\xi) = 1$ .

- A  $K$ -term PCE approximation to product of two random fields  $a$  and  $b$ :

$$\Pi_\Lambda[a, b] := \Pi_\Lambda[\Pi_\Lambda[a] \Pi_\Lambda[b]].$$

- A  $K$ -term PCE approximation to the ratio of two random fields  $a$  and  $b$ :

$$\Pi_\Lambda^\dagger[b/a] : \text{the solution to } \Pi_\Lambda[a, b/a] = \Pi_\Lambda[b].$$

## Notations

- $\hat{z} = (\hat{z}_1, \dots, \hat{z}_K)^\top$ .
- $\mathcal{P}(\hat{z}) := \sum_{k=1}^K \hat{z}_k \mathcal{M}_k$ ,  $(\mathcal{M}_k)_{\ell m} = \langle \phi_k, \phi_\ell \phi_m \rangle_\rho$ .
- It can be shown that  $\widehat{\Pi_\Lambda[a, b]} = \mathcal{P}(\hat{a})\hat{b}$ ,  $\widehat{\Pi_\Lambda^\dagger[b/a]} = \mathcal{P}^{-1}(\hat{a})\hat{b}$ .

## Stochastic Galerkin (SG) Method

- Ansatz:

$$\begin{aligned} h &\simeq h_\Lambda := \sum_{k=1}^K \hat{h}_k(x, y, t) \phi_k(\xi), \\ q^x &\simeq q_\Lambda^x := \sum_{k=1}^K (\hat{q}^x)_k(x, y, t) \phi_k(\xi), \\ q^y &\simeq q_\Lambda^y := \sum_{k=1}^K (\hat{q}^y)_k(x, y, t) \phi_k(\xi), \end{aligned}$$

- Stochastic Galerkin method applies standard Galerkin procedure in the stochastic  $\xi$  space, which leads to a new system of partial differential equations with respect to the PCE coefficients.

## SG Projection of Nonlinear Terms

- $\frac{(q^x)^2}{h} = \frac{q^x}{h} q^x \longrightarrow \Pi_\Lambda \left[ \frac{(q_\Lambda^x)^2}{h_\Lambda} \right] = \Pi_\Lambda \left[ q_\Lambda^x \Pi_\Lambda^\dagger \left[ \frac{q_\Lambda^x}{h_\Lambda} \right] \right],$   
 $\frac{(q^y)^2}{h} = \frac{q^y}{h} q^y \longrightarrow \Pi_\Lambda \left[ \frac{(q_\Lambda^y)^2}{h_\Lambda} \right] = \Pi_\Lambda \left[ q_\Lambda^y \Pi_\Lambda^\dagger \left[ \frac{q_\Lambda^y}{h_\Lambda} \right] \right].$
- For  $q^x q^y / h$  in  $(q^x q^y / h)_x$ ,  
 $\frac{q^x q^y}{h} = (q^x) \frac{q^y}{h} \longrightarrow \Pi_\Lambda \left[ \frac{q_\Lambda^x q_\Lambda^y}{h_\Lambda} \right] = \Pi_\Lambda \left[ q_\Lambda^x \Pi_\Lambda^\dagger \left[ \frac{q_\Lambda^y}{h_\Lambda} \right] \right].$
- For  $q^x q^y / h$  in  $(q^x q^y / h)_y$ ,  
 $\frac{q^x q^y}{h} = (q^y) \frac{q^x}{h} \longrightarrow \Pi_\Lambda \left[ \frac{q_\Lambda^x q_\Lambda^y}{h_\Lambda} \right] = \Pi_\Lambda \left[ q_\Lambda^y \Pi_\Lambda^\dagger \left[ \frac{q_\Lambda^x}{h_\Lambda} \right] \right].$

## The Main Results

### Hyperbolcity-Preserving SG Formulation

$$\frac{\partial}{\partial t}(\widehat{U}) + \frac{\partial}{\partial x}(\widehat{F}(\widehat{U})) + \frac{\partial}{\partial y}(\widehat{G}(\widehat{U})) = \widehat{S}(\widehat{U}, \widehat{B}). \quad (1)$$

Here,  $\widehat{U} := (\widehat{h}^\top, \widehat{q}^x{}^\top, \widehat{q}^y{}^\top)^\top$ , and

$$\widehat{F}(\widehat{U}) = \begin{pmatrix} \mathcal{P}(\widehat{q}^x) \mathcal{P}^{-1}(\widehat{h}) \widehat{q}^x + \frac{1}{2} g \mathcal{P}(\widehat{h}) \widehat{h} \\ \mathcal{P}(\widehat{q}^x) \mathcal{P}^{-1}(\widehat{h}) \widehat{q}^y \end{pmatrix}, \quad \widehat{G}(\widehat{U}) = \begin{pmatrix} \mathcal{P}(\widehat{q}^y) \mathcal{P}^{-1}(\widehat{h}) \widehat{q}^x \\ \mathcal{P}(\widehat{q}^y) \mathcal{P}^{-1}(\widehat{h}) \widehat{q}^y + \frac{1}{2} g \mathcal{P}(\widehat{h}) \widehat{h} \end{pmatrix}, \quad \widehat{S}(\widehat{U}, \widehat{B}) = \begin{pmatrix} 0 \\ -g \mathcal{P}(\widehat{h}) \widehat{B}_x \\ -g \mathcal{P}(\widehat{h}) \widehat{B}_y \end{pmatrix}. \quad (2)$$

### Theorem (Hyperbolcity-preserving condition)

The system (1) is hyperbolic if the matrix  $\mathcal{P}(\widehat{h}) > 0$ .

The condition  $\mathcal{P}(\widehat{h}) > 0$  reduces to  $h > 0$  when the  $\xi$ -dependence is dropped from the system.

### Theorem (A computationally tractable condition)

Given  $\Lambda$ , let nodes  $\xi_m$  and weights  $\tau_m$  satisfying  $\{(\xi_m, \tau_m)\}_{m=1}^M$  represent any  $M$ -point positive quadrature rule that is exact on

$$P_\Lambda^3 := \text{span} \left\{ \prod_{n=1}^3 \phi_n \mid n \in [K] \right\}.$$

If

$$h_\Lambda(x, y, t, \xi_m) > 0 \quad \forall m = 1, \dots, M, \quad (3)$$

then the matrix  $\mathcal{P}(\widehat{h}) > 0$ .

In other words, we only need to ensure the positivity of the stochastic water heights at some quadrature points to preserve the hyperbolcity of (1).

## Second-Order Central-Upwind Scheme

Assuming uniform rectangular partition over a rectangular region,

$$\frac{d}{dt} \mathbf{U}_{i,j} = -\frac{\mathcal{F}_{i+\frac{1}{2},j} - \mathcal{F}_{i-\frac{1}{2},j}}{\Delta x} - \frac{\mathcal{G}_{i,j+\frac{1}{2}} - \mathcal{G}_{i,j-\frac{1}{2}}}{\Delta y} + \overline{\mathbf{S}}_{i,j},$$

where  $\mathbf{U}_{i,j}$  represent the cell averages of the vector  $\widehat{U}$  in rectangular cell  $\mathcal{C}_{i,j}$ .

- Source term:

$$\overline{\mathbf{S}}_{i,j} \approx \frac{1}{|\mathcal{C}_{i,j}|} \int_{\mathcal{C}_{i,j}} \widehat{S}(\widehat{U}, \widehat{B}) dx dy.$$

- Numerical fluxes:

$$\mathcal{F}_{i+\frac{1}{2},j} := \frac{a_{i+\frac{1}{2},j}^+ \widehat{F}(\mathbf{U}_{i,j}^E) - a_{i+\frac{1}{2},j}^- \widehat{F}(\mathbf{U}_{i+1,j}^W)}{a_{i+\frac{1}{2},j}^+ - a_{i+\frac{1}{2},j}^-} + \frac{a_{i+\frac{1}{2},j}^+ a_{i+\frac{1}{2},j}^-}{a_{i+\frac{1}{2},j}^+ - a_{i+\frac{1}{2},j}^-} [\mathbf{U}_{i+1,j}^W - \mathbf{U}_{i,j}^E],$$

$$\mathcal{G}_{i,j+\frac{1}{2}} := \frac{b_{i,j+\frac{1}{2}}^+ \widehat{G}(\mathbf{U}_{i,j}^N) - b_{i,j+\frac{1}{2}}^- \widehat{G}(\mathbf{U}_{i,j+1}^S)}{b_{i,j+\frac{1}{2}}^+ - b_{i,j+\frac{1}{2}}^-} + \frac{b_{i,j+\frac{1}{2}}^+ b_{i,j+\frac{1}{2}}^-}{b_{i,j+\frac{1}{2}}^+ - b_{i,j+\frac{1}{2}}^-} [\mathbf{U}_{i,j+1}^S - \mathbf{U}_{i,j}^N].$$

- Propagation speeds:

$$\begin{aligned} a_{i+\frac{1}{2},j}^- &= \min \left\{ \lambda_1 \left( \frac{\partial \widehat{F}}{\partial \widehat{U}}(\mathbf{U}_{i+1,j}^W) \right), \lambda_1 \left( \frac{\partial \widehat{F}}{\partial \widehat{U}}(\mathbf{U}_{i,j}^E) \right), 0 \right\}, \\ a_{i+\frac{1}{2},j}^+ &= \max \left\{ \lambda_{3K} \left( \frac{\partial \widehat{F}}{\partial \widehat{U}}(\mathbf{U}_{i+1,j}^W) \right), \lambda_{3K} \left( \frac{\partial \widehat{F}}{\partial \widehat{U}}(\mathbf{U}_{i,j}^E) \right), 0 \right\}, \\ b_{i,j+\frac{1}{2}}^- &= \min \left\{ \lambda_1 \left( \frac{\partial \widehat{G}}{\partial \widehat{U}}(\mathbf{U}_{i,j+1}^S) \right), \lambda_1 \left( \frac{\partial \widehat{G}}{\partial \widehat{U}}(\mathbf{U}_{i,j}^N) \right), 0 \right\}, \\ b_{i,j+\frac{1}{2}}^+ &= \max \left\{ \lambda_{3K} \left( \frac{\partial \widehat{G}}{\partial \widehat{U}}(\mathbf{U}_{i,j+1}^S) \right), \lambda_{3K} \left( \frac{\partial \widehat{G}}{\partial \widehat{U}}(\mathbf{U}_{i,j}^N) \right), 0 \right\}, \end{aligned}$$

$\mathbf{U}_{i,j}^{E,W,N,S}$  are the pointwise values of the second-order accurate, non-oscillatory piecewise linear reconstructions of  $\mathbf{U}_{i,j}$  at the midpoints of the boundaries, i.e.,

$$\begin{aligned} \mathbf{U}_{i,j}^E &= \mathbf{U}_{i,j} + \frac{\Delta x}{2} (\mathbf{U}_x)_{i,j}, & \mathbf{U}_{i,j}^W &= \mathbf{U}_{i,j} - \frac{\Delta x}{2} (\mathbf{U}_x)_{i,j}, \\ \mathbf{U}_{i,j}^N &= \mathbf{U}_{i,j} + \frac{\Delta y}{2} (\mathbf{U}_y)_{i,j}, & \mathbf{U}_{i,j}^S &= \mathbf{U}_{i,j} - \frac{\Delta y}{2} (\mathbf{U}_y)_{i,j}, \end{aligned}$$

## Hyperbolcity-Preserving Well-Balanced Central-Upwind Scheme

- Stochastic “lake-at-rest” state:  $\begin{cases} q_\Lambda^x = q_\Lambda^y \equiv 0, \\ h_\Lambda + \Pi_\Lambda[B] \equiv C(\xi), \end{cases} \Rightarrow \begin{cases} \widehat{q}^x = \widehat{q}^y \equiv \mathbf{0}, \\ \widehat{h} + \widehat{B} \equiv \widehat{C}. \end{cases}$
- The PCE vector  $\widehat{B}$  for the bottom function is replaced by its piecewise bilinear interpolant.
- The pointwise values of the reconstructions of the PCE of water surface  $\widehat{\eta}$  are reconstructed. The reconstructed water height are computed by  $\widehat{h} := \widehat{\eta} - \widehat{B}$ .
- The first moments  $\widehat{h}_1$  are “corrected” following a similar procedure to the central-upwind scheme for the deterministic shallow water equations.
- The PCE vectors  $\widehat{h}$  are filtered to satisfies the condition (3).

## Numerical simulations

Deterministic initial water surface:

$$\eta(x, y, 0, \xi) = \begin{cases} 1.01, & \text{if } 0.05 < x < 0.15, \\ 1, & \text{otherwise,} \end{cases}$$

Deterministic initial velocity field (0-discharge):

$$u(x, y, 0, \xi) = v(x, y, 0, \xi) = 0.$$

Stochastic bottom topography:

$$B(x, y, \xi) = 0.8e^{-5(x-0.9+0.1\xi^{(1)})^2 - 50(y-0.5+0.1\xi^{(2)})^2}.$$

Randomness:

$$\xi^{(1)} \sim \text{Beta}(4, 2), \quad \xi^{(2)} \sim \mathcal{U}(-1, 1).$$

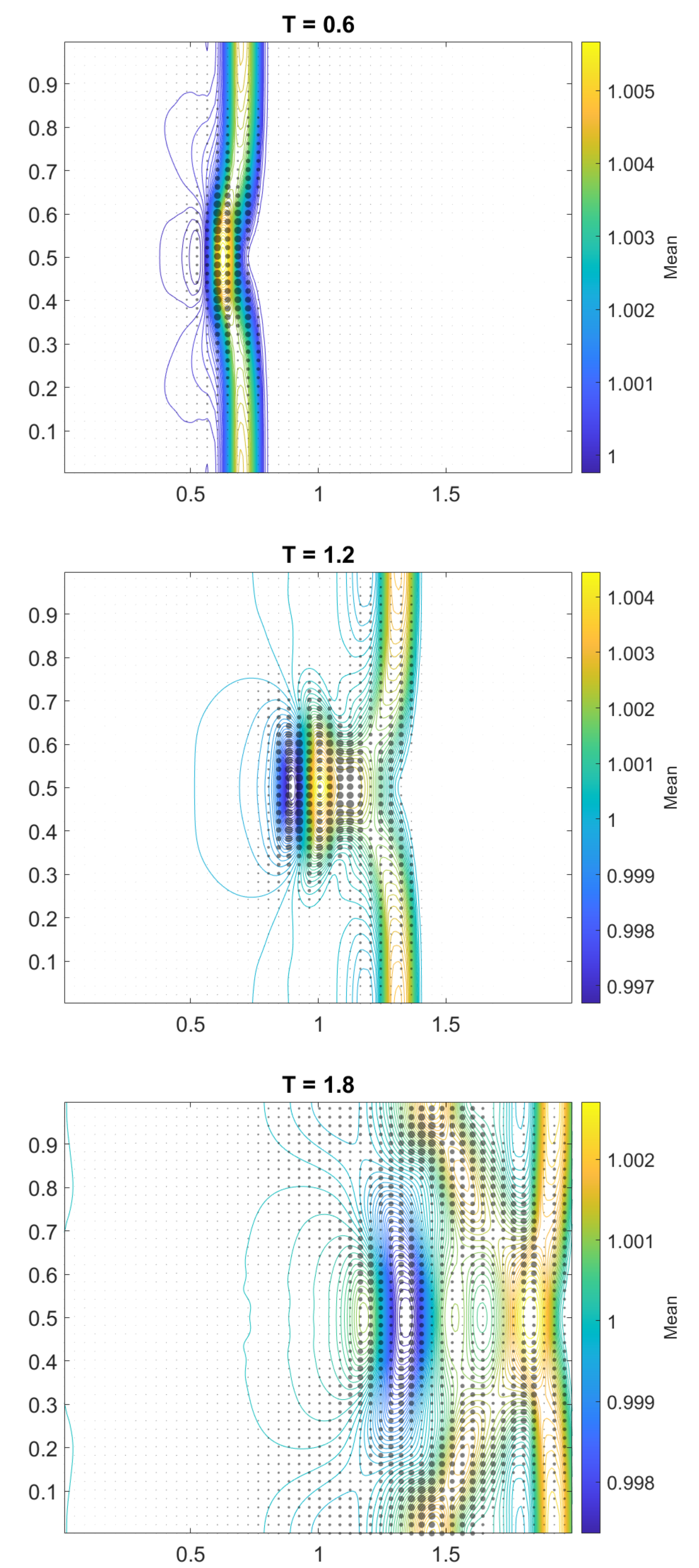


Figure 1: Numerical results at  $T = 0.6$  (top),  $T = 1.2$  (middle), and  $T = 1.8$  (bottom), respectively. The largest disks are corresponding to the standard deviation values 2.20e-3, 2.00e-3, and 1.20e-3, respectively. The index set  $\Lambda = \{(\nu^{(1)}, \nu^{(2)}) \in \mathbb{N}^2 \mid 0 \leq \nu^{(1)}, \nu^{(2)} \leq 3\}$ . The polynomial basis is chosen to be the tensor-product set.

## References

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