

AMATH 567: Problem Set 10

Jithin D. George, No. 1622555

December 7, 2016

1. (a)

$$\psi^+(x) = \lim_{\epsilon \rightarrow 0^+} \frac{1}{2\pi i} \int_{-\infty}^{\infty} \frac{f(t)}{t - (x + i\epsilon)} dt$$

The function inside the integral has a singularity

$$P^+ = \lim_{\epsilon \rightarrow 0^+} \frac{1}{2\pi i} \int_{-\infty}^{\infty} \frac{1}{t - (x + i\epsilon)} dt$$

$$P^+ \psi^+(x) = \lim_{\epsilon \rightarrow 0^+} \frac{1}{2\pi i} \int_{-\infty}^{\infty} \frac{\psi^+(t)}{t - (x + i\epsilon)} dt$$

We decide to do the integral in the complex plan with the upper semicircle as the contour because $\psi^+(z)$ is analytic in the upper half plane. Assuming the function on the arc decays as the arc goes to ∞ , we can ignore the integral along it.

$$P^+ \psi^+(x) = \lim_{\epsilon \rightarrow 0^+} \frac{1}{2\pi i} \text{Res}_{t=x+i\epsilon} \frac{\psi^+(t)}{t - (x + i\epsilon)} dt$$

$$P^+ \psi^+(x) = \lim_{\epsilon \rightarrow 0^+} \psi^+(x + i\epsilon) = \psi^+(x)$$

$$\psi^-(x) = \lim_{\epsilon \rightarrow 0^+} \frac{1}{2\pi i} \int_{-\infty}^{\infty} \frac{f(t)}{t - (x - i\epsilon)} dt$$

$$P^- = \lim_{\epsilon \rightarrow 0^+} \frac{1}{2\pi i} \int_{-\infty}^{\infty} \frac{1}{t - (x - i\epsilon)} dt$$

$$P^- \psi^-(x) = \lim_{\epsilon \rightarrow 0^+} \frac{1}{2\pi i} \int_{-\infty}^{\infty} \frac{\psi^-(t)}{t - (x - i\epsilon)} dt$$

We decide to do the integral in the complex plan with the lower semicircle as the contour because $\psi^-(z)$ is analytic in the lower half plane. Assuming the similar decay as before,

$$P^- \psi^-(x) = - \lim_{\epsilon \rightarrow 0^+} \frac{1}{2\pi i} \text{Res}_{t=x-i\epsilon} \frac{\psi^-(t)}{t - (x - i\epsilon)} dt$$

The minus sign comes because the integral in the lower semicircle is clockwise.

$$P^- \psi^-(x) = - \lim_{\epsilon \rightarrow 0^+} \psi^-(x - i\epsilon) = -\psi^-(x)$$

Also,

$$P^- \psi^+(x) = \lim_{\epsilon \rightarrow 0^+} \frac{1}{2\pi i} \int_{-\infty}^{\infty} \frac{\psi^+(t)}{t - (x - i\epsilon)} dt = 0$$

since there is no residue in the upper half plane.

Similarly,

$$P^+ \psi^-(x) = 0$$

(b)

$$\begin{aligned}
\psi^+(x) - \psi^-(x) &= \frac{1}{x^4 + 1} \\
P^+\psi^+(x) - P^+\psi^-(x) &= P^+\frac{1}{x^4 + 1} \\
P^+\psi^+(x) &= P^+\left(\frac{1}{x^4 + 1}\right) \\
\psi^+(x) &= P^+\left(\frac{1}{x^4 + 1}\right) \\
\psi^-(x) &= P^-\left(\frac{1}{x^4 + 1}\right) \\
\psi^+(x) &= \lim_{\epsilon \rightarrow 0^+} \frac{1}{2\pi i} \int_{-\infty}^{\infty} \frac{1}{(t^4 + 1)(t - (x + i\epsilon))} dt
\end{aligned}$$

Taking the upper semicircle as the contour,

$$\begin{aligned}
\psi^+(x) &= \lim_{\epsilon \rightarrow 0^+} \left(\lim_{t \rightarrow \frac{1+i}{\sqrt{2}}} \frac{1}{(t + \frac{1+i}{\sqrt{2}})(t - \frac{1-i}{\sqrt{2}})(t - \frac{-1+i}{\sqrt{2}})(t - (x + i\epsilon))} \right) \\
&+ \lim_{\epsilon \rightarrow 0^+} \left(\lim_{t \rightarrow \frac{-1+i}{\sqrt{2}}} \frac{1}{(t + \frac{1+i}{\sqrt{2}})(t - \frac{1-i}{\sqrt{2}})(t - \frac{1+i}{\sqrt{2}})(t - (x + i\epsilon))} \right) \\
&+ \lim_{\epsilon \rightarrow 0^+} \left(\lim_{t \rightarrow x+i\epsilon} \frac{1}{(t + \frac{1+i}{\sqrt{2}})(t - \frac{1-i}{\sqrt{2}})(t - \frac{-1+i}{\sqrt{2}})(t - \frac{1-i}{\sqrt{2}})} \right) \\
\psi^+(x) &= \frac{1}{(2(1+i)i(1+i-\sqrt{2}x))} - \frac{1}{(2(-1+i)i(-1+i-\sqrt{2}x))} + \frac{1}{x^4 + 1} \\
\psi^-(x) &= \lim_{\epsilon \rightarrow 0^+} \frac{1}{2\pi i} \int_{-\infty}^{\infty} \frac{1}{(t^4 + 1)(t - (x - i\epsilon))} dt
\end{aligned}$$

Taking the lower semicircle as the contour,

$$\begin{aligned}
\psi^-(x) &= \lim_{\epsilon \rightarrow 0^+} \left(\lim_{t \rightarrow -\frac{1+i}{\sqrt{2}}} \frac{1}{(t - \frac{1+i}{\sqrt{2}})(t - \frac{1-i}{\sqrt{2}})(t - \frac{-1+i}{\sqrt{2}})(t - (x + i\epsilon))} \right) \\
&+ \lim_{\epsilon \rightarrow 0^+} \left(\lim_{t \rightarrow \frac{1-i}{\sqrt{2}}} \frac{1}{(t + \frac{1+i}{\sqrt{2}})(t - \frac{-1+i}{\sqrt{2}})(t - \frac{1+i}{\sqrt{2}})(t - (x + i\epsilon))} \right) \\
&+ \lim_{\epsilon \rightarrow 0^+} \left(\lim_{t \rightarrow x-i\epsilon} \frac{1}{(t + \frac{1+i}{\sqrt{2}})(t - \frac{1-i}{\sqrt{2}})(t - \frac{-1+i}{\sqrt{2}})(t - \frac{1-i}{\sqrt{2}})} \right) \\
\psi^-(x) &= \frac{1}{(2(1+i)i(1+i+\sqrt{2}x))} - \frac{1}{(2(1-i)i(1-i-\sqrt{2}x))} + \frac{1}{x^4 + 1} \\
\psi^+(x) - \psi^-(x) &= \frac{1}{(2(1+i)i(1+i-\sqrt{2}x))} - \frac{1}{(2(-1+i)i(-1+i-\sqrt{2}x))} \\
&- \frac{1}{(2(1+i)i(1+i+\sqrt{2}x))} + \frac{1}{(2(1-i)i(1-i-\sqrt{2}x))}
\end{aligned}$$

From Mathematica, we get that the right hand side is the partial fraction expansion of $\frac{1}{x^4+1}$. So,

$$\psi^+(x) - \psi^-(x) = \frac{1}{x^4 + 1}$$

(c)

$$\psi^+(x) = \lim_{\epsilon \rightarrow 0+} \frac{1}{2\pi i} \int_{-\infty}^{\infty} \frac{f(t)}{t - (x + i\epsilon)} dt$$

C_R is a tiny circle centered at x with radius δ . There is no singularity in the region A because the singularity is currently at $x + i\epsilon$. So, the integral along the dotted line is the same as the integral along C_R .

$$\begin{aligned} \psi^+(x) &= \lim_{\epsilon \rightarrow 0+} \frac{1}{2\pi i} \lim_{\delta \rightarrow 0} \left(\int_{-\infty}^{x-\delta} \frac{f(t)}{t - (x + i\epsilon)} dt + \int_{x+\delta}^{\infty} \frac{f(t)}{t - (x + i\epsilon)} dt + \int_{C_R} \frac{f(t)}{t - (x + i\epsilon)} dt \right) \\ &= \lim_{\epsilon \rightarrow 0+} \frac{1}{2\pi i} \int_{-\infty}^{\infty} \frac{f(t)}{t - (x + i\epsilon)} dt + \lim_{\epsilon \rightarrow 0+} \frac{1}{2\pi i} \lim_{\delta \rightarrow 0} \int_{C_R} \frac{f(t)}{t - (x + i\epsilon)} dt \end{aligned}$$

Switching the limits and using the baby limit theorem on the second term,

$$\begin{aligned} &= \lim_{\epsilon \rightarrow 0+} \frac{1}{2\pi i} \int_{-\infty}^{\infty} \frac{f(t)}{t - (x + i\epsilon)} dt + \frac{1}{2\pi i} \pi i f(x) \\ &= \frac{1}{2\pi i} \int_{-\infty}^{\infty} \frac{f(t)}{t - x} dt + \frac{1}{2} f(x) \end{aligned}$$

Similarly, for ψ^- , the contour C_R is in the upper plane making the integral along it negative.

$$\psi^- = \frac{1}{2\pi i} \int_{-\infty}^{\infty} \frac{f(t)}{t - x} dt - \frac{1}{2} f(x)$$

The first term is closely related to the Hilbert Transform

$$H(f(x)) = \frac{1}{\pi} \int_{-\infty}^{\infty} \frac{f(t)}{t - x} dt$$

$$\psi^+ = \frac{1}{2i} H(f(x)) + \frac{1}{2} f(x)$$

$$\psi^- = \frac{1}{2i} H(f(x)) - \frac{1}{2} f(x)$$

2. (a)

$$\begin{aligned} \hat{F}(k) &= \int_{-\infty}^{\infty} e^{-ikx} f(x) dx \\ &= \int_{-\infty}^0 e^{-ikx} 0 dx + \int_{0+}^{\infty} e^{-ikx} e^{-x} dx \\ &= 0 + \int_{0+}^{\infty} e^{(-ik-1)x} dx \\ &= \frac{1}{1 + ik} \end{aligned}$$

(b)

$$\hat{F}(k) = \frac{1}{1+ik}$$

This has a singularity at $k=i$.

$$f(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \frac{e^{ikx}}{1+ik} dk$$

Since this is tough to integrate on the real line, we use integrals along complex contours. The numerator is

$$e^{ikx} = e^{i(k_x + ik_y)x} = e^{ik_x x} e^{-k_y x}$$

When $x > 0$, this decays to 0 for positive k_y . So, we take the closed semicircle in the upper k plane as our contour. Since the integral over the arc decays to 0 for large k and there is only residue at i ,

$$\begin{aligned} \int_{-\infty}^{\infty} \frac{e^{ikx}}{1+ik} dk &= \text{Res}_{k=i} \frac{e^{ikx}}{1+ik} \\ &= 2\pi i \lim_{k \rightarrow i} (k-i) \frac{e^{ikx}}{1+ik} \\ &= 2\pi i \frac{e^{-x}}{i} \\ &= 2\pi e^{-x} \end{aligned}$$

When $x < 0$, we have to use negative k_y for the decay. So, we take the closed semicircle in the lower k plane as our contour. Since there is no residue in it,

$$\int_{-\infty}^{\infty} \frac{e^{ikx}}{1+ik} dk = 0$$

There is obviously a jump at $x=0$.

$$\begin{aligned} f(0) &= \frac{1}{2\pi} \int_{-\infty}^{\infty} \frac{1}{1+ik} dk = \frac{1}{2\pi} \int_{-\infty}^{\infty} \frac{1-ik}{1+k^2} dk = \frac{1}{2\pi} \int_{-\infty}^{\infty} \frac{1}{1+k^2} dk - \frac{i}{2\pi} \int_{-\infty}^{\infty} \frac{k}{1+k^2} dk \\ &= \frac{1}{2\pi} \int_{-\infty}^{\infty} \frac{1}{1+k^2} dk - \frac{i}{2\pi} \lim_{R \rightarrow \infty} \int_{-R}^R \frac{k}{1+k^2} dk \\ &= \frac{1}{2\pi} [\arctan(x)]_{-\infty}^{\infty} - 0 (\text{odd function}) \\ &= \frac{1}{2} \end{aligned}$$

$$f(x) = \begin{cases} e^{-x} & \text{if } x > 0 \\ \frac{1}{2} & \text{if } x = 0 \\ 0 & \text{if } x < 0 \end{cases}$$

3.

$$f(z) = -4z$$

$$g(z) = e^z - 1$$

On $|z| = 1$,

$$|f(z)| = 4$$

and

$$|g(z)| \leq |e^{e^{i\theta}}| + 1 = e + 1$$

Thus, by Roche's theorem, $f(z)+g(z)$ has same number of zeros inside the circle as $f(z)$ which has exactly one at 0. So, $e^z - 4z - 1$ has exactly one zero inside $|z| = 1$

4.

$$f(z) = z^4 + z^3 + 5z^2 + 2z + 4$$

We take the infinite quarter circle.

On the real line, from 0 to ∞ , $f(z)$ is real. So, the change in argument is 0. On the arc,

$$\begin{aligned} f(z) &= R^4 e^{4i\theta} + R^3 e^{3i\theta} + 5R^2 e^{2i\theta} + 2R e^{i\theta} + 4 \\ &= R^4 \left(e^{4i\theta} + \frac{e^{3i\theta}}{R} + \frac{5e^{2i\theta}}{R^2} + \frac{2e^{i\theta}}{R^3} + \frac{4}{R^4} \right) \end{aligned}$$

As R goes to ∞ , $f(z)$ has an argument of 4θ .

From 0 to $\frac{\pi}{2}$, the change in argument is 2π .

On the y axis,

$$\begin{aligned} f(z) &= y^4 - iy^3 - 5y^2 + iy + 4 \\ \tan(\theta) &= \frac{-y^3 + y}{y^4 - 5y^2 + 4} \end{aligned}$$

For very large y , $\tan(\theta) = 0$ because of the term on the bottom. This matches with the value of $\theta = 2\pi$ which we got from the arc. Furthermore, the derivative is negative because of the dominance of y^3 in the numerator. So, the argument decreases.

At 0, $\tan(\theta) = 0$. So, the argument decreases to a value $k\pi$.

$y^4 - 5y^2 + 4$ has zeros at 2, -2, 1, -1. As y goes from ∞ to 0, $\tan(\theta)$ crosses over 2 of these at 2 and 1. Thus, $\tan(\theta)$ has two singularities on the y axis.

Looking at the graph of \tan , we see that after crossing two singularities, the argument comes to 0. Total change in argument is 0. Thus, by the argument principle, $f(z)$ has no roots in the first quadrant.

5. (a)

$$p(t) = \frac{tf'(t)}{f(t) - w}$$

This has a singularity when $f(t) = w$. There is only singularity for a particular w because $f(t)$ is bijective.

$$g(w) = \frac{1}{2\pi i} \oint_{C(z_0, R)} \frac{tf'(t)}{f(t) - w} dt$$

because $w = f(z)$, this evaluates to

$$\begin{aligned} &= \text{Res}_{t=z} \frac{tf'(t)}{f(t) - w} \\ &= \lim_{t \rightarrow z} (t - z) \frac{tf'(t)}{f(t) - w} \\ &= \lim_{t \rightarrow z} \frac{tf'(t)}{\frac{f(t) - w}{(t - z)}} \\ &= \frac{zf'(z)}{f'(z)} \\ &= z \end{aligned}$$

(b)

$$\begin{aligned} g(w) &= \frac{1}{2\pi i} \oint_{C(z_0, R)} \frac{tf'(t)}{f(t) - w} dt \\ &= \frac{1}{2\pi i} \oint_{C(z_0, R)} \frac{t(t+1)e^t}{te^t - w} dt \\ &= \frac{1}{2\pi i} \oint_{C(z_0, R)} \frac{(t+1)}{1 - \frac{w}{te^t}} dt \\ &= \frac{1}{2\pi i} \oint_{C(z_0, R)} (t+1) \left(1 + \frac{w}{te^t} + \left(\frac{w}{te^t}\right)^2 + \dots\right) dt \end{aligned}$$

Except the first term, all term have a singularity at $t=0$. So,

$$\begin{aligned} &= \frac{1}{2\pi i} \text{Res}_{t \rightarrow 0} (t+1) \left(\frac{w}{te^t} + \left(\frac{w}{te^t}\right)^2 + \dots\right) \\ &= w \frac{1}{2\pi i} \text{Res}_{t \rightarrow 0} \left(\frac{(t+1)}{te^t}\right) + w^2 \frac{1}{2\pi i} \text{Res}_{t \rightarrow 0} \frac{(t+1)}{e^{2t}} + \frac{w^3}{2!} \frac{1}{2\pi i} \text{Res}_{t \rightarrow 0} \frac{(t+1)}{e^{3t}} + \dots \\ g(w) &= \sum_0^{\infty} a_n w^n \\ a_0 &= 0 \\ a_n &= \frac{1}{2\pi i} \text{Res}_{t \rightarrow 0} \frac{(t+1)}{e^{nt}} \end{aligned}$$