

A model for nulling & mode changing in pulsars.

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Phenomena { mode changing : 不同时段, integrated prof 变
nulling : 有的时段 pulse 看不到

Same? ← single class of reasons

时标: a few pulse periods

properties → hours/days...

nulling { 轻微
 { 严重

reasons for nulling { "microphysics" in radiation mechanism
 { or? both?
 { whole magnetosphere's change

PSR B1931+24 & J1832+0029

(Kramer et al. 2006, Lyne 2009)

→ \dot{P} 不同, 在 nulling & 非 nulling 态

至少有一部分 nulling 伴随着磁层整体变化。

This paper: qualitative (semi-quantitative) model to explain nulling & mode change, and the \dot{P} change.

Recap: sph-down 能量分配

$$E_{\text{sph-down}} = \frac{1}{2} I \Omega^2, \quad \dot{E} = I \Omega \dot{\Omega} \quad \left\{ \begin{array}{l} < 10^{-3} \Rightarrow \text{radio emission} \\ \sim 10^{-1} \Rightarrow \text{HE emission/cold particle wind (non-radiative)} \end{array} \right.$$

⇒ $\dot{\Omega}$ 的变化不能仅因为有无 radio emission 导致

~ 变化 50%
B1931+24

less when radio invisible

days of nulling

many other psrs: ~ hours or less

\dot{P} 变化不易测 (timing)

Global changes: in magnetosphere structure

Recap: psr magnetosphere



accelerating E only in geometrically very small regions

$P = P_{\text{aJ}}, \quad j \leq j_{\text{aJ}}$
most magnetosphere: force-free

如何影响 \dot{P} : 改变 \dot{E} 改变 j & B

改磁层位形 或 改 current 分布

modeswitch & nulling:

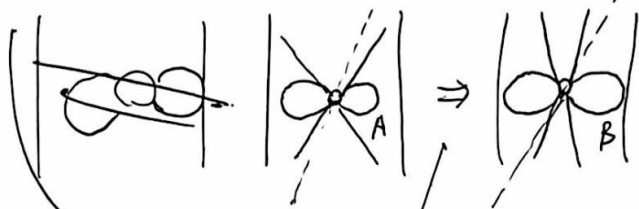
- switches between some quasi-stable states.
- ① 不同 closed field line zone
 - ② 不同 open field line 区 current distribution
 - ③ both

①

When will nullings happen?

- ① coherent condition broken
- ② coherent emission generated elsewhere.

② 另一种情形: radio emission zone shrink



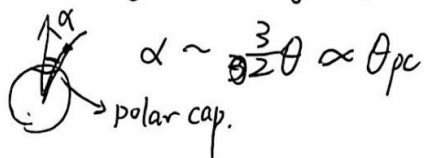
correlation between width & \dot{p}
in mode-switching pulsars?

另一种模型: simplified \dot{p} aligned rotator
• dipole \vec{B}

aim: to prove modest change in magnetosphere
 \Rightarrow large variations in \dot{p}

②.1 Shrinking-expanding
corotating zone

A: larger opening angle.



ω : \dot{p} 的一种写法

formulae:

① (a) \vec{B} 表示法

cylindrical: (ω, ϕ, Z)

$$\vec{B} = \frac{\nabla \psi \times \vec{e}_\phi}{\omega} + \frac{4\pi}{c} \cdot \frac{I}{\omega} \vec{e}_\phi$$

$$(B_\omega, B_\phi, B_z) = \left(-\frac{1}{\omega} \frac{\partial \psi}{\partial z}, \frac{4\pi}{c} \frac{I}{\omega}, \frac{1}{\omega} \frac{\partial \omega \psi}{\partial \omega} \right)$$

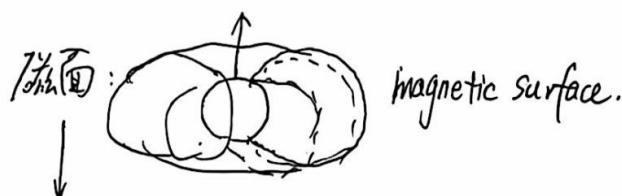
$$\nabla \times \vec{B} = \begin{vmatrix} \frac{1}{\omega} \vec{e}_\omega & \vec{e}_\phi & \frac{1}{\omega} \vec{e}_z \\ \frac{\partial}{\partial \omega} & \frac{\partial}{\partial \phi} & \frac{\partial}{\partial z} \\ -\frac{1}{\omega} \frac{\partial \psi}{\partial z} & \frac{4\pi}{c} \frac{I}{\omega} & \frac{1}{\omega} \frac{\partial \omega \psi}{\partial \omega} \end{vmatrix}$$

$$= \frac{4\pi}{c} \frac{1}{\omega} \vec{e}_\omega \left(\frac{1}{\omega} \frac{\partial \omega \psi}{\partial \omega} - \frac{4\pi}{c} \frac{1}{\omega} \frac{\partial I}{\partial z} \right) - \vec{e}_\phi \left(-\frac{1}{\omega} \frac{\partial \omega \psi}{\partial z} + \frac{1}{\omega} \frac{\partial^2 \psi}{\partial \omega^2} + \frac{1}{\omega} \frac{\partial^2 \psi}{\partial z^2} \right) + \frac{1}{\omega} \vec{e}_z \left(\frac{4\pi}{c} \frac{1}{\omega} I + \frac{4\pi}{c} \frac{1}{\omega} \frac{\partial I}{\partial \omega} + \frac{1}{\omega} \frac{\partial^2 \psi}{\partial \phi^2} \right)$$

cylindrical (ρ, ϕ, z) $\frac{4\pi}{c} \frac{S}{R_c} \frac{S}{\rho} \vec{e}_\phi$

$$\vec{B} = \underbrace{\frac{\nabla \psi \times \vec{e}_\phi}{\rho}}_{\text{poloidal}} + \underbrace{\frac{4\pi}{c} \cdot \frac{I}{\rho} \cdot \vec{e}_\phi}_{\text{toroidal}}$$

$$(B_\rho, B_\phi, B_z) = \left(-\frac{1}{\rho} \frac{\partial \psi}{\partial z}, \frac{4\pi}{c} \frac{I}{\rho}, \frac{1}{\rho} \frac{\partial \rho \psi}{\partial \rho} \right)$$



空间中的位形: 用一个函数 ψ 表示.

$\psi(\rho, z) = \text{const}$ 表示一个磁面/磁线

flux $\Phi_{\text{mag}} = 2\pi \psi$ (轴对称)

依定义, 磁面的法向 \perp 磁切

$$\Rightarrow \vec{B} \cdot \nabla \psi = 0$$

同时, ψ 满足, $\nabla \Phi_0 = \Phi_0(\phi)$, 有

$$\nabla \psi \cdot \nabla \Phi_0 = 0$$

参考
BGI书



$$\mathbf{B} = \underbrace{\nabla\psi \times \nabla\phi_0}_{\text{Euler potential}} + \underbrace{g \cdot \nabla\phi_0}_{\text{general form}}$$

$g = g(\rho, z)$

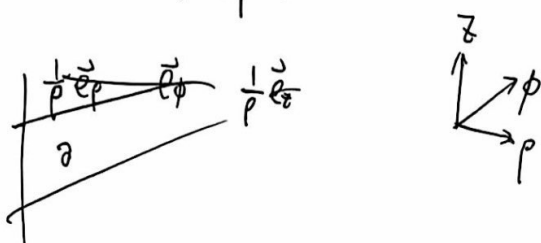
可以有环向分量

$$\frac{c}{4\pi} [\mathbf{E} \times \mathbf{B}]_{\text{pol}} = -\frac{\Omega}{c} I \vec{B}_{\text{pol}}$$

$$\nabla \cdot \mathbf{B} = \nabla \cdot (\nabla\psi \times \nabla\phi_0) + \nabla \cdot (g \nabla\phi_0)$$

$$\begin{aligned} \nabla\psi &= \frac{\partial\psi}{\partial\rho} \vec{e}_\rho + \frac{1}{\rho} \frac{\partial\psi}{\partial\phi} \vec{e}_\phi + \frac{\partial\psi}{\partial z} \vec{e}_z \\ &= \frac{\partial\psi}{\partial\rho} \vec{e}_\rho + 0 + \frac{\partial\psi}{\partial z} \vec{e}_z \end{aligned}$$

$$\nabla\phi_0 = \frac{1}{\rho} \frac{\partial\phi}{\partial\phi} \vec{e}_\phi$$



$$\nabla\psi \times \nabla\phi_0 = \frac{1}{\rho} \frac{\partial\psi}{\partial\rho} \frac{\partial\phi}{\partial\phi} \vec{e}_z - \frac{1}{\rho} \frac{\partial\psi}{\partial z} \frac{\partial\phi}{\partial\phi} \vec{e}_\rho$$

$$\begin{aligned} \nabla \cdot (\dots) &= \frac{1}{\rho} \frac{\partial}{\partial\rho} \left(\frac{1}{\rho} \frac{\partial\psi}{\partial z} \frac{\partial\phi}{\partial\phi} \right) + \frac{\partial}{\partial z} \left(\frac{1}{\rho} \frac{\partial\psi}{\partial\rho} \frac{\partial\phi}{\partial\phi} \right) \\ &= -\frac{1}{\rho} \frac{\partial\phi}{\partial\phi} \frac{\partial^2\psi}{\partial\rho\partial z} + \frac{1}{\rho} \frac{\partial\psi}{\partial\rho\partial z} \frac{\partial\phi}{\partial\phi} = 0 \end{aligned}$$

$$\begin{aligned} \nabla \cdot (g \nabla\phi_0) &= \nabla \cdot \left(\frac{g}{\rho} \frac{\partial\phi}{\partial\phi} \vec{e}_\phi \right) \\ &= \frac{1}{\rho} \frac{\partial}{\partial\rho} \left(\frac{g}{\rho} \frac{\partial\phi}{\partial\phi} \right) = \frac{g}{\rho^2} \frac{\partial\phi}{\partial\phi} = 0 \end{aligned}$$

$$\nabla\phi_0 = \frac{1}{\rho} \vec{e}_\phi, \quad \nabla\phi_0 = \frac{1}{\rho} \vec{e}_\phi$$

$$\Rightarrow \mathbf{B} = \frac{1}{\rho} \nabla\psi \times \vec{e}_\phi + \frac{1}{\rho} g \vec{e}_\phi$$

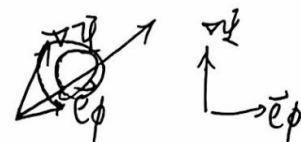
↓

$$\mathbf{B} = \frac{\nabla\psi \times \vec{e}_\phi}{\rho} + \frac{4\pi}{c} \frac{I}{\rho} \vec{e}_\phi$$

$$\vec{E} = -\frac{(\vec{\Omega} \times \vec{r})}{c} \times \mathbf{B} = -\frac{1}{c} (\vec{\Omega} \times (\rho \vec{e}_\phi + z \vec{e}_z)) \times \mathbf{B}$$

$$= -\frac{1}{c} \Omega \rho \vec{e}_\phi \times \left(\frac{1}{\rho} \nabla\psi \times \vec{e}_\phi + \frac{1}{\rho} g \vec{e}_\phi \right)$$

$$= -\frac{1}{c} \Omega \nabla\psi$$

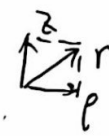


$$\Rightarrow P_r = -\frac{I\Omega}{c} \frac{1}{\rho r}$$

$$B_{\text{pol}} = -\frac{1}{\rho} \frac{\partial\psi}{\partial z} \vec{e}_\rho + \frac{1}{\rho} \frac{\partial\psi}{\partial\rho} \vec{e}_z$$

$$\vec{P}_{\text{pol}} = \frac{\Omega}{c} I \frac{1}{\rho} (\frac{\partial\psi}{\partial z} \vec{e}_\rho - \frac{\partial\psi}{\partial\rho} \vec{e}_z)$$

$$P_r = -\frac{I\Omega}{c} \frac{1}{\rho r} (z \frac{\partial\psi}{\partial\rho} - \rho \frac{\partial\psi}{\partial z})$$

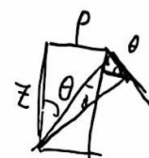


Energy loss: $dW = -r^2 P_r d\omega \rightarrow$ solid angle

$$\Rightarrow \frac{dW}{d\omega} = \frac{I\Omega}{c} \frac{r}{\rho} (z \frac{\partial\psi}{\partial\rho} - \rho \frac{\partial\psi}{\partial z})$$

$$W = \int_{4\pi} \frac{dW}{d\omega} d\omega$$

$$= \int_{4\pi} \frac{dW}{d\omega} \frac{\sin\theta d\theta d\phi}{\sqrt{z^2 + \rho^2}}$$



$$\rho = r \sin\theta, \quad r = \sqrt{\rho^2 + z^2}$$

$$\phi = \phi, \quad \theta = \arctan \frac{\rho}{z}$$

$$z = r \cos\theta$$

$$d\theta = \frac{1}{1 + (\frac{\rho}{z})^2} \left(\frac{d\rho}{z} - \frac{\rho dz}{z^2} \right)$$

$$= \frac{1}{z^2 + \rho^2} (z d\rho - \rho dz)$$

$$\psi_{\text{last}} = \psi|_{z=0}$$

$$= \mu \frac{1}{\rho}$$

(2)



$$W = \int_{4\pi} \frac{I\Omega}{c} \cdot \frac{r}{\rho} (z \partial_r \psi - \rho \partial_z \psi) \cdot \frac{\rho}{\sqrt{z^2 + \rho^2}} \cdot \frac{1}{\sqrt{z^2 + \rho^2}} (z d\rho - \rho dz)$$

$$d\rho = r \sin\theta \cdot d\theta$$

$$z = r \cos\theta$$

$$\rho = r \sin\theta$$

$$\partial_r \psi = (\partial_r \psi) \frac{\partial r}{\partial \rho} + (\partial_\theta \psi) \frac{\partial \theta}{\partial \rho} = (\partial_r \psi) \frac{1}{\sin\theta} + (\partial_\theta \psi) \cdot \frac{1}{r \cos\theta}$$

$$\partial_z \psi = (\partial_r \psi) \frac{\partial r}{\partial z} + (\partial_\theta \psi) \frac{\partial \theta}{\partial z} = (\partial_r \psi) \cdot \frac{1}{\cos\theta} - (\partial_\theta \psi) \cdot \frac{1}{r \sin\theta}$$

$$W = \int_{4\pi} \frac{I\Omega}{c} \cdot \frac{1}{\sin\theta} \cdot (r \cos\theta \cdot \partial_r \psi - r \sin\theta \partial_z \psi) \sin\theta d\theta d\phi$$

$$= \int_{4\pi} \frac{I\Omega}{c} \cdot \left(r \cdot \frac{\cos\theta}{\sin\theta} \cdot \partial_r \psi + \partial_\theta \psi \cdot \frac{r \sin\theta}{\cos\theta} \cdot \partial_r \psi + \partial_\theta \psi \right) d\theta d\phi$$

$$= 2 \int_{\psi_{\text{last}}}^{\psi} \frac{2\pi}{c} I \Omega_F d\psi$$

$I = I(\psi)$



$$I \propto -\psi \left(2 - \frac{\psi}{\psi_{\text{last}}} \right) \quad \text{Michel 1973b}$$

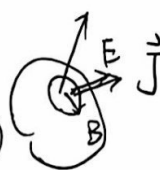
$$\Rightarrow \text{Diagram of a loop with a central dot} \quad J = 2\pi I$$

$$W \propto -\psi_{\text{last}}^2 \propto \rho_{\text{last}}^{-2}$$

$$\theta_{pc} = \sqrt{\frac{R_{\text{ns}}}{\rho_{\text{last}}}}$$

$$W \propto \theta_{pc}^4 \propto \alpha^4$$

$$P = I \cdot U$$



$$(2.2) \quad \rho_{\text{last}} \text{ 不变, } \vec{J} \text{ 变} \Rightarrow W \sim \int \vec{J} \cdot \vec{E} d\ell$$

$$S = \frac{4\pi}{c} \cdot \frac{R_{\text{ic}}}{\rho_b} \cdot I = S(\psi) \quad \sim \int \vec{J} \cdot \Omega \cdot \vec{B} \cdot r dr d\ell / c$$

$$\sim \vec{J} \cdot \vec{B} \cdot \pi R^2 \propto J \cdot \Omega$$

Pulsar Equation Michel 1973a,b

$$\psi_{zz} + \psi_{\theta\theta} - \frac{1}{\rho} \left(\frac{R_{\text{ic}}^2 + \rho^2}{R_{\text{ic}}^2 - \rho^2} \right) \cdot \psi_\rho + \frac{S' \cdot S}{R_{\text{ic}}^2 - \rho^2} = 0$$

$$W = |W_{\text{md}}| \cdot \int_0^{\psi_{\text{last}}} S d\psi$$

$$S = -\psi \left(2 - \frac{\psi}{\psi_{\text{last}}} \right)$$

$$- \int_0^{\psi_{\text{last}}} \psi \left(2 - \frac{\psi}{\psi_{\text{last}}} \right) d\psi \quad ?$$

$$- \frac{2}{3} \psi_{\text{last}}^2 + \frac{1}{3} \cdot \psi_{\text{last}}^2$$

$$S = - \left(\psi_{\text{last}} - \frac{\psi^2}{\psi_{\text{last}}} \right)$$



$J \propto \Omega$ (Blondford & Znajek 1977)
 $\Rightarrow W \propto J^2 \rightarrow \propto \Omega^4$

...

Discussion:

① If milking/mode change pure geometrical?

\rightarrow Not really...



For this kind of psr, change of magnetospheric paras may cease radio emission...

② mode switching pulsars' \dot{P} change?

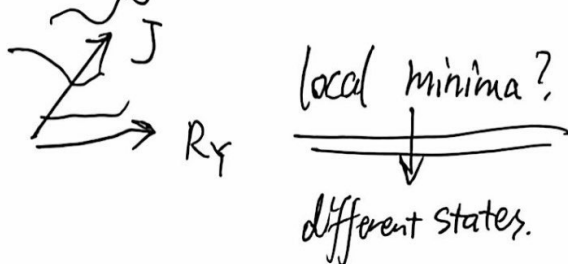
(measurements...)

\rightarrow with high timing noise.

③ Why states are quasi-stable?

$$(C \gg P_i)$$

total energy $E \Rightarrow E(R_p, J, \dots)$

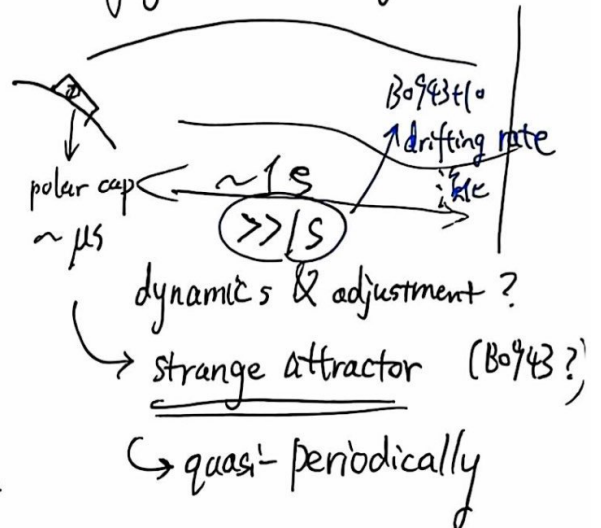


J distribution change $\Leftarrow V_{\text{polarcap}}$ change

With psr aging, $V_{\text{polarcap}} \downarrow$ (vocabulary)
 small ^{potential} variation leads to larger changes

④ psr magnetosphere:

highly non-linear system.



⑤ RRAT: pure geometrical?

\rightarrow 统计/分析-下 RRAT 分布?

<https://rratalog.github.io/rratalog>

~~没人干过...~~



③



$$\Psi = \Psi_0 \cdot \frac{z}{r} = \Psi_0 \cdot \frac{z}{\sqrt{\rho^2 + z^2}}$$

$$\Psi_{zz} = \Psi_0 \cdot \frac{\sqrt{\rho^2 + z^2} - z \cdot \frac{z}{\sqrt{\rho^2 + z^2}}}{\rho^2 + z^2} = \Psi_0 \cdot \frac{\rho^2}{(\rho^2 + z^2)^{3/2}}$$

~~$$\Psi_{zz} = \Psi_0 \cdot \frac{z}{(\rho^2 + z^2)^{3/2}}$$~~

$$\Psi_{zz} = \Psi_0 \cdot \rho^2 \cdot \frac{(-\frac{3}{2})}{(\rho^2 + z^2)^{5/2}} \cdot 2z$$

$$= -\Psi_0 \cdot \rho^2 \cdot \frac{3z}{(\rho^2 + z^2)^{5/2}}$$

$$\Psi_{\rho\rho} = \Psi_0 \cdot z \cdot \frac{-\frac{1}{2}}{(\rho^2 + z^2)^{3/2}} \cdot 2\rho$$

$$= -\Psi_0 \cdot \frac{\rho z}{(\rho^2 + z^2)^{3/2}}$$

$$\Psi_{\rho\rho} = -\Psi_0 \cdot \frac{z \cdot (\rho^2 + z^2)^{3/2} - \rho z \cdot \frac{3}{2} (\rho^2 + z^2)^{1/2} \cdot 2\rho}{(\rho^2 + z^2)^3}$$

$$= -\Psi_0 \cdot \frac{z(\rho^2 + z^2) - 3\rho^2 z}{(\rho^2 + z^2)^{5/2}}$$

$$= -\Psi_0 \cdot \frac{z^3 - 2\rho^2 z}{(\rho^2 + z^2)^{5/2}}$$

$$\Psi_{zz} + \Psi_{\rho\rho} = -\Psi_0 \cdot \frac{z^3 + \rho^2 z}{(\rho^2 + z^2)^{5/2}}$$

$$= -\Psi_0 \cdot \frac{z}{(\rho^2 + z^2)^{3/2}}$$

$$-\Psi_0 \cdot \frac{z}{(\rho^2 + z^2)^{3/2}} + \Psi_0 \cdot \left(\frac{R_{LC}^2 + \rho^2}{R_{LC}^2 - \rho^2} \right) \cdot \frac{z}{(\rho^2 + z^2)^{3/2}} - \frac{S' \cdot S}{R_{LC}^2 - \rho^2} = 0$$

$$\frac{S' \cdot S}{R_{LC}^2 - \rho^2} = -\Psi_0 \cdot \frac{z}{(\rho^2 + z^2)^{3/2}} + \Psi_0 \cdot \frac{z}{(\rho^2 + z^2)^{3/2}} \cdot \frac{R_{LC}^2 + \rho^2}{R_{LC}^2 - \rho^2}$$

$$= \Psi_0 \cdot \frac{z}{(\rho^2 + z^2)^{3/2}} \cdot \frac{2\rho^2}{R_{LC}^2 - \rho^2}$$

$$\Rightarrow S' \cdot S = \Psi_0 \cdot \frac{2z\rho^2}{(\rho^2 + z^2)^{3/2}} = \Psi \cdot \frac{2\rho^2}{\rho^2 + z^2}$$

$$= \Psi \cdot 2 \cdot \left(1 - \frac{\Psi^2}{\Psi_0^2} \right)$$

$$\frac{dS^2}{d\Psi} = \Psi \cdot \left(1 - \frac{\Psi^2}{\Psi_0^2} \right)$$

$$S = A\Psi^2 + B\Psi + C$$

$$S' = 2A\Psi + B$$

$$S \cdot S' = 2A^2\Psi^3 + 3AB\Psi^2 + 2AB\Psi^2 + B^2\Psi$$

$$+ 2AC\Psi + BC$$

$$= 2A^2\Psi^3 + 3AB\Psi^2 + (B^2 + 2AC)\Psi + BC$$

$$2A^2 = \frac{2}{\Psi_0^2} \quad A = \frac{1}{\Psi_0}$$

$$B = 0$$

$$2AC = -2 \quad C = -\Psi_0$$

$$\Rightarrow S = \frac{1}{\Psi_0} \Psi^2 - \Psi_0$$

$$\vec{B} = \left(-\Psi_0 \cdot \frac{\rho}{(\rho^2 + z^2)^{3/2}}, \frac{1}{R_{LC}} (\Psi^2 - \Psi_0^2) \cdot \frac{1}{\rho}, -\Psi_0 \cdot \frac{z}{(\rho^2 + z^2)^{3/2}} \right)$$

= split monopole

赤道 $\Psi = 0$

$$\rho \vec{E} + \frac{1}{c} (\vec{j} \times \vec{B}) = 0 \quad \text{极区 } \Psi = \Psi_0$$

$$\Rightarrow (\nabla \cdot \vec{E}) \vec{E} + [(\nabla \times \vec{B}) \times \vec{B}] = 0$$

FFE E.M & Ohm's Law

$$\text{plug in } \vec{B} = \frac{\nabla \Psi \times \vec{e}_\phi}{\rho} + \frac{\Psi_0}{R_{LC}} \frac{S}{\rho} \vec{e}_\phi$$

$$\& \vec{E} + (\vec{\Omega} \times \vec{r}) \times \vec{B} = 0$$

$$\vec{j} = \sigma(\vec{E} + \vec{v} \times \vec{B}) \xrightarrow{\sigma \rightarrow \infty} \vec{E} + (\vec{\Omega} \times \vec{r}) \times \vec{B} = 0$$

\Rightarrow PSR E_ϕ $\sigma \rightarrow \infty$ 结果

