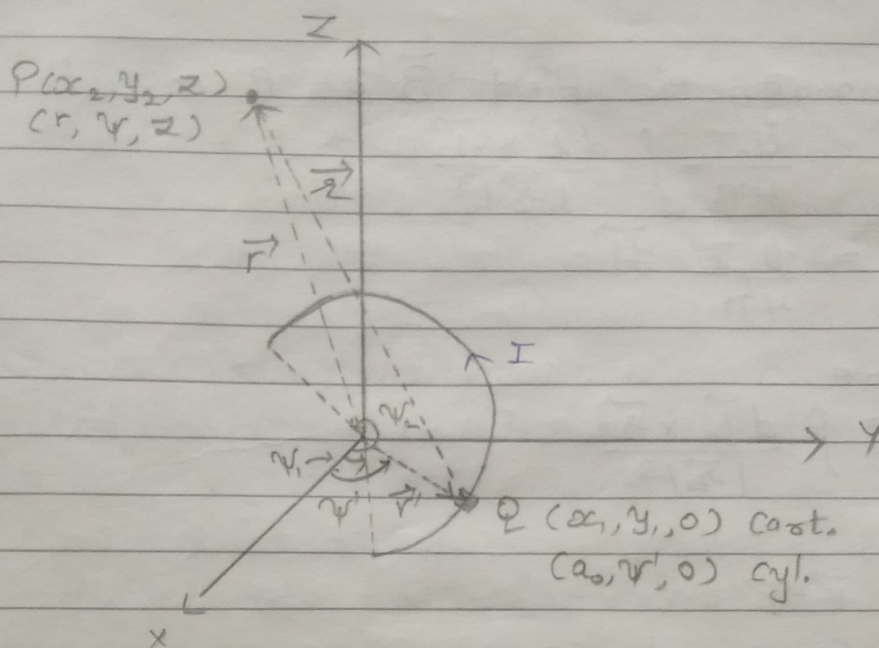


3

# Magnetic field due to a conducting

filament in the form of a circular arc

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→ A current carrying circular conducting filament of approximately zero thickness placed in the  $xy$ -Plane so that its center is at the origin and axis is  $z$ -axis of a right handed rectangular co-ordinate system. The magnetic field and vector  $\vec{p}$  will be calculated in Cartesian co-ordinate system.

→ Here,

(1) Source point,  $Q(a_0, \psi', 0)$

(2) Observation point,  $P(r, \psi, z)$

(3)  $a_0$  is the radius of the filament

(4)  $\psi_1$  is the angular distance of the nearest end from  $x$ -axis measured along the direction of current

(5)  $\psi_2$  is the angular distance of the farthest end from  $x$ -axis measured along the direction of current.

(6) Current  $I$  is flowing from  $\psi_1$  to  $\psi_2$



→ The magnetic field  $\vec{B}$  at P can be written as

$$\vec{B} = \frac{\mu_0 I}{4\pi} \int \frac{d\vec{l} \times \vec{r}}{|\vec{r}|^3}$$

$$= \frac{\mu_0 I}{4\pi} \vec{H}$$

where,

$$\vec{H} = \int \frac{d\vec{l} \times \vec{r}}{|\vec{r}|^3}$$

Here,

(i)  $\vec{r} = \vec{r} - \vec{r}'$

$$= (x_2, y_2, z) - (x_1, y_1, 0)$$

$$= (x_2 - x_1, y_2 - y_1, z) \quad \text{(in Cartesian Coordinates)}$$

take,  $x_2 = r \cos \psi$ ,  $x_1 = a_0 \cos \psi'$

$$y_2 = r \sin \psi, \quad y_1 = a_0 \sin \psi'$$

$$\therefore \vec{r} = (r \cos \psi - a_0 \cos \psi') \hat{i} + (r \sin \psi - a_0 \sin \psi') \hat{j} + z \hat{k}$$

(ii) Current element,  $d\vec{l}' = a_0 \psi' \hat{\psi}$

(in cylindrical coord. sys.)

→ Must note that our current element is in cylindrical co-ordinate system and  $\vec{r}$  is in Cartesian co-ordinate system

→ Now, substitute these values in eq<sup>n</sup> of  $\vec{H}$

$$\therefore \vec{H} = \int_{\psi_1}^{\psi_2} \frac{a_0 \psi' \hat{\psi} \times [(r \cos \psi - a_0 \cos \psi') \hat{i} + (r \sin \psi - a_0 \sin \psi') \hat{j} + z \hat{k}]}{[(r \cos \psi - a_0 \cos \psi')^2 + (r \sin \psi - a_0 \sin \psi')^2 + z^2]^{3/2}}$$

→ Let's try to simplify  $\vec{H}$ ,

(iii)  $|\vec{r}|^3 = [(r \cos \psi - a_0 \cos \psi')^2 + (r \sin \psi - a_0 \sin \psi')^2 + z^2]^{3/2}$

$$= [r^2 \cos^2 \psi + a_0^2 \cos^2 \psi'^2 - 2r a_0 \cos \psi' \cos \psi$$

$$+ r^2 \sin^2 \psi + a_0^2 \sin^2 \psi'^2 - 2r a_0 \sin \psi' \sin \psi + z^2]^{3/2}$$



$$= [r^2 + a_0^2 + z^2 - 2ra_0(\cos\psi'\cos\psi + \sin\psi'\sin\psi)]^{3/2}$$

$$\therefore |\vec{r}|^3 = [r^2 + a_0^2 + z^2 - 2ra_0\cos(\psi - \psi')]^{3/2}$$

(iv) In this part we will convert the unit vector " $\hat{\psi}$ " in cartesian unit vector, so that we can do cross product.

→ transformation Equations (Cylindrical → Cartesian)

$$\hat{r} = \cos\theta \hat{i} + \sin\theta \hat{j}$$

$$\hat{\theta} = -\sin\theta \hat{i} + \cos\theta \hat{j}$$

→ Using second transformation Eq<sup>n</sup> we can write,

$$\hat{\psi} = -\sin\psi' \hat{i} + \cos\psi' \hat{j}$$

→ So, our current element becomes,

$$a_0 d\psi' \hat{\psi} = a_0 d\psi' (-\sin\psi' \hat{i} + \cos\psi' \hat{j})$$

(v) Now, we are ready to do cross product

$$\rightarrow d\vec{r}' \times \vec{r} = a_0 d\psi' (-\sin\psi' \hat{i} + \cos\psi' \hat{j})$$

$$\times [(r\cos\psi - a_0\cos\psi') \hat{i} + (r\sin\psi - a_0\sin\psi') \hat{j} + z \hat{k}]$$

$$= \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ -\sin\psi' & \cos\psi' & 0 \\ r\cos\psi - a_0\cos\psi' & r\sin\psi - a_0\sin\psi' & z \end{vmatrix} a_0 d\psi'$$

$$= (a_0 z \cos\psi' d\psi') \hat{i} + (a_0 z \sin\psi' d\psi') \hat{j}$$

$$+ \hat{k} a_0 d\psi' \begin{vmatrix} \sin\psi' & \cos\psi' \\ r\cos\psi - a_0\cos\psi' & r\sin\psi - a_0\sin\psi' \end{vmatrix}$$

→ Let's solve the determinant,

$$a_0 d\psi' \begin{vmatrix} -\sin\psi' & \cos\psi' \\ r\cos\psi - a_0\cos\psi' & r\sin\psi - a_0\sin\psi' \end{vmatrix}$$



$$\begin{aligned}
 &= a_0 d\psi' [(-r \sin \psi \sin \psi' + a_0 \sin^2 \psi') \\
 &\quad - (r \cos \psi \cos \psi' - a_0 \cos^2 \psi')] \\
 &= a_0 d\psi' [a_0 - r [\cos \psi \cos \psi' + \sin \psi \sin \psi']] \\
 &= a_0 (a_0 - r \cos(\psi - \psi')) d\psi'
 \end{aligned}$$

$$\therefore d\vec{r}' \times \vec{z} = (a_0 z \cos \psi' d\psi') \hat{i} + (a_0 z \sin \psi' d\psi') \hat{j} + a_0 (a_0 - r \cos(\psi - \psi')) d\psi' \hat{k}$$

→ Finally we can substitute all values of  $\vec{H}$

$$\therefore \vec{H} = \int_{\psi_1}^{\psi_2} \frac{(a_0 z \cos \psi' d\psi') \hat{i} + (a_0 z \sin \psi' d\psi') \hat{j} + a_0 (a_0 - r \cos(\psi - \psi')) d\psi' \hat{k}}{[r^2 + a_0^2 + z^2 - 2ra_0 \cos(\psi - \psi')]^{3/2}}$$

→ The components of  $\vec{H}$ ,

$$(1) \vec{H}_x = a_0 z \int_{\psi_1}^{\psi_2} \frac{\cos \psi' d\psi'}{[r^2 + a_0^2 + z^2 - 2ra_0 \cos(\psi - \psi')]^{3/2}} \hat{i}$$

$$(2) \vec{H}_y = a_0 z \int_{\psi_1}^{\psi_2} \frac{\sin \psi' d\psi'}{[r^2 + a_0^2 + z^2 - 2ra_0 \cos(\psi - \psi')]^{3/2}} \hat{j}$$

$$(3) \vec{H}_z = a_0 \int_{\psi_1}^{\psi_2} \frac{(a_0 - r \cos(\psi - \psi'))}{[r^2 + a_0^2 + z^2 - 2ra_0 \cos(\psi - \psi')]^{3/2}} \hat{k}$$

→ To solve Eq<sup>n</sup> ①, ② and ③,

$$\text{Let } b^2 = a_0^2 + r^2 + z^2$$

$$\phi = \psi' - \psi$$

$$\text{also, } \phi_1 = \psi_1 - \psi$$

$$\Rightarrow d\phi = d\psi'$$

$$\phi_2 = \psi_2 - \psi$$

→ So our Eq<sup>n</sup> reduce to,

$$(4) H_x = a_0 z \int_{\phi_1}^{\phi_2} \frac{\cos(\phi + \psi) d\phi}{[b^2 - 2ra_0 \cos \phi]^{3/2}}$$



$$(5) \quad \mathcal{H}_y = a_0 z \int_{\phi_1}^{\phi_2} \frac{\sin(\phi + \psi) d\phi}{[b^2 - 2ra_0 \cos \phi]^{3/2}}$$

$$(7) \quad \mathcal{H}_z = a_0 \int_{\phi_1}^{\phi_2} \frac{(a_0 - r \cos \phi) d\phi}{[b^2 - 2ra_0 \cos \phi]^{3/2}}$$

→ Now,

$$\begin{aligned} (4) \quad \mathcal{H}_x &= a_0 z \int_{\phi_1}^{\phi_2} \frac{\cos(\phi + \psi) d\phi}{[b^2 - 2ra_0 \cos \phi]^{3/2}} \\ &= a_0 z \int_{\phi_1}^{\phi_2} \frac{(\cos \phi \cos \psi - \sin \phi \sin \psi) d\phi}{(b^2 - 2ra_0 \cos \phi)^{3/2}} \\ &= a_0 z \left( \cos \psi \int_{\phi_1}^{\phi_2} \frac{\cos \phi d\phi}{(b^2 - 2ra_0 \cos \phi)^{3/2}} - \sin \psi \int_{\phi_1}^{\phi_2} \frac{\sin \phi d\phi}{(b^2 - 2ra_0 \cos \phi)^{3/2}} \right) \end{aligned}$$

$$\begin{aligned} (5) \quad \mathcal{H}_y &= a_0 z \int_{\phi_1}^{\phi_2} \frac{(\sin \phi \cos \psi + \sin \psi \cos \phi) d\phi}{(b^2 - 2ra_0 \cos \phi)^{3/2}} \\ &= a_0 z \left( \cos \psi \int_{\phi_1}^{\phi_2} \frac{\sin \phi d\phi}{(b^2 - 2ra_0 \cos \phi)^{3/2}} + \sin \psi \int_{\phi_1}^{\phi_2} \frac{\cos \phi d\phi}{(b^2 - 2ra_0 \cos \phi)^{3/2}} \right) \end{aligned}$$

$$\begin{aligned} (6) \quad \mathcal{H}_z &= a_0 \int_{\phi_1}^{\phi_2} \frac{(a_0 - r \cos \phi) d\phi}{(b^2 - 2ra_0 \cos \phi)^{3/2}} \\ &= a_0^2 \int_{\phi_1}^{\phi_2} \frac{d\phi}{(b^2 - 2ra_0 \cos \phi)^{3/2}} - a_0 r \int_{\phi_1}^{\phi_2} \frac{\cos \phi d\phi}{(b^2 - 2ra_0 \cos \phi)^{3/2}} \end{aligned}$$

→ Now take,

$$I_1 = \int_{\phi_1}^{\phi_2} \frac{\cos \phi d\phi}{(b^2 - 2ra_0 \cos \phi)^{3/2}}$$

$$I_2 = \int_{\phi_1}^{\phi_2} \frac{\sin \phi d\phi}{(b^2 - 2ra_0 \cos \phi)^{3/2}}$$

$$I_3 = \int_{\phi_1}^{\phi_2} \frac{d\phi}{(b^2 - 2ra_0 \cos \phi)^{3/2}}$$



→ So, Eq<sup>n</sup> (4), (5) and (6) can be written as

$$(4) \mathcal{H}_x = a_0 z (\cos \psi I_1 - \sin \psi I_2)$$

$$(5) \mathcal{H}_y = a_0 z (\cos \psi I_2 + \sin \psi I_1)$$

$$(6) \mathcal{H}_z = a_0 (a_0 I_3 - r I_1)$$

$$\rightarrow I_2 = \int_{\phi_1}^{\phi_2} \frac{\sin \phi \, d\phi}{[b^2 - 2ra_0 \cos \phi]^{3/2}} \quad \text{Case-1}$$

$$\text{Let } t = \cos \phi$$

$$\Rightarrow dt = -\sin \phi \, d\phi$$

$$\begin{aligned} \therefore I_2 &= \int \frac{-dt}{[b^2 - 2ra_0 t]^{3/2}} \\ &= -\int (b^2 - 2ra_0 t)^{-3/2} dt \end{aligned}$$

$$= -\frac{(b^2 - 2ra_0 t)^{-1/2}}{(-1/2)(-2ra_0)}$$

$$= -\frac{1}{ra_0} \left[ \frac{1}{\sqrt{b^2 - 2ra_0 \cos \phi}} \right]_{\phi_1}^{\phi_2}$$

$$\therefore I_2 = \frac{-1}{ra_0} \sum_{i=1}^2 (-1)^i \frac{1}{\sqrt{b^2 - 2ra_0 \cos \phi_i}} \quad \text{iff } r \neq 0$$

Case-2 if  $r=0$

$$\begin{aligned} \rightarrow I_2 &= \int_{\phi_1}^{\phi_2} \frac{\sin \phi \, d\phi}{b^3} \\ &= \frac{1}{(a_0^2 + z^2)^{3/2}} [-\cos \phi]_{\phi_1}^{\phi_2} \end{aligned}$$

$$\therefore I_2 = \frac{-1}{(a_0^2 + z^2)^{3/2}} \sum_{i=1}^2 (-1)^i \cos \phi_i$$



→ Sub. the values of  $\vec{H}$  components in  $\vec{B}$  we get,

$$(7) \quad B_x = \frac{\mu_0 I}{4\pi} a_0 z (\cos \psi I_1 - \sin \psi I_2)$$

$$(8) \quad B_y = \frac{\mu_0 I}{4\pi} a_0 z (\cos \psi I_2 + \sin \psi I_1)$$

$$(9) \quad B_z = \frac{\mu_0 I}{4\pi} a_0 (a_0 I_3 - r I_1)$$

$$\text{where } \rightarrow I_1 = \int_{\phi_1}^{\phi_2} \frac{\cos \phi \, d\phi}{(b^2 - 2ra_0 \cos \phi)^{3/2}}$$

$$\rightarrow \text{If } r=0, I_2 = \frac{-1}{(a_0^2 + z^2)^{3/2}} \sum_{i=1}^2 (-1)^i \cos \phi_i$$

$$\text{If } r \neq 0, I_2 = \frac{-1}{ra_0} \sum_{i=1}^2 (-1)^i \frac{1}{\sqrt{b^2 - 2ra_0 \cos \phi_i}}$$

$$\rightarrow I_3 = \int_{\phi_1}^{\phi_2} \frac{d\phi}{(b^2 - 2ra_0 \cos \phi)^{3/2}}$$

$$\text{and } \rightarrow \phi = \psi' - \psi \Rightarrow \phi_k = \psi_k - \psi, k=1,2$$

$$\rightarrow b^2 = a_0^2 + r^2 + z^2$$

Notes we can not solve  $I_1$  and  $I_3$  analytically.  
We can express  $I_1$  and  $I_3$  in term of Elliptic Integrals. Please see the book in references

### \* References &

- 1 I.S. Gradshteyn and I.M. Ryzhik, Table of Integrals, Series and Products. Page 180, 182, 179
- 2 Analytical derivation of magnetic field and vector potential due to different current carrying conducting geometries, IPR.
- 3 Special Thanks to Grattu Ramesh Sir.