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# MAT351 TUTORIAL 5

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PARTIAL DIFFERENTIAL EQUATIONS

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## Revisiting a question from Homework 2

Recall that in your second homework you were asked to do the following problem:

[Strauss 2.1.7] If both  $\phi$  and  $\psi$  are odd functions of  $x$ , show that the solution  $u(x, t)$  of the wave equation is also odd in  $x$  for all  $t$ .

Every student who did this problem correctly followed a similar method - use d'Alembert's formula. To that end, I want to provide an alternate solution which does not make use of such a solution formula.

We have that  $u(x, t)$  is solution of wave equation ( $u_{tt} = c^2 u_{xx}$ ,  $u(x, 0) = \phi(x)$ ,  $u_t(x, 0) = \psi(x)$ ). The IVP for  $u(-x, t)$  is

$$\begin{cases} u_{tt}(-x, t) = c^2 u_{xx}(-x, t) \\ u(-x, 0) = \phi(-x) = -\phi(x) \\ u_t(-x, 0) = \psi(-x) = -\psi(x) \end{cases}$$

We consider  $u(x, t) + u(-x, t)$ , and since the sum of solutions to wave equation is a solution, we have  $u(x, t) + u(-x, t)$  solves

$$\begin{cases} [u(x, t) + u(-x, t)]_{tt} = c^2 [u(x, t) + u(-x, t)]_{xx} \\ u(x, 0) + u(-x, 0) = \phi(x) - \phi(x) = 0 \\ u_t(x, 0) + u_t(-x, 0) = \psi(x) - \psi(x) = 0 \end{cases}$$

A solution satisfying these initial conditions is  $u(x, t) + u(-x, t) = 0$ . Then by uniqueness, it is the only solution and we have shown that  $u(-x, t) = -u(x, t)$  as desired.

*Remark.* One can use a similar argument and consider  $u(-x, t) - u(x, t)$  to prove that if both  $\phi$  and  $\psi$  are even functions of  $x$ , then  $u(x, t)$  is also even in  $x$  for all  $t$ .

## Diffusion on the half-line

Recall that the half-line Dirichlet problem

$$\begin{cases} v_t - kv_{xx} = 0 & \text{on } \{0 < x < \infty, 0 < t < \infty\} \\ v(x, 0) = \phi(x) \\ v(0, t) = 0 \end{cases}$$

has solution formula (via method of odd extension)

$$v(x, t) = \frac{1}{\sqrt{4\pi kt}} \int_0^\infty \left[ e^{-(x-y)^2/4kt} - e^{-(x+y)^2/4kt} \right] \phi(y) dy$$

Now, let's try and do similar with the half-line Neumann problem

$$\begin{cases} w_t - kw_{xx} = 0 & \text{on } \{0 < x < \infty, 0 < t < \infty\} \\ w(x, 0) = \phi(x) \\ w_x(0, t) = 0 \end{cases} \quad (2.1)$$

*Remark.* The Neumann boundary condition  $w_x(0, t) = 0$  can be viewed as insulating the side end for  $t > 0$  - there is no transfer of heat at  $x = 0$ .

Let  $\phi_{\text{even}}(x)$  be the even extension of  $\phi(x)$  to the whole real line, that is

$$\phi_{\text{even}}(x) = \begin{cases} \phi(x), & x \geq 0 \\ \phi(-x), & x \leq 0 \end{cases}$$

Then, the solution to

$$\begin{cases} u_t - ku_{xx} = 0 \\ u(x, 0) = \phi_{\text{even}}(x) \end{cases}$$

on the whole line is given (recall Tutorial 3) by  $u(x, t) = \frac{1}{\sqrt{4\pi kt}} \int_{-\infty}^\infty \exp(-(x-y)^2/4kt) \phi_{\text{even}}(y) dy$ . Then, our solution to (2.1) is the restriction  $w(x, t) = u(x, t)$  for  $x > 0$ . We now just need to

simplify the formula:

$$\begin{aligned}
 w(x, t) &= \frac{1}{\sqrt{4\pi kt}} \left[ \int_{-\infty}^0 e^{-(x-y)^2/4kt} \phi(-y) dy + \int_0^{\infty} e^{-(x-y)^2/4kt} \phi(y) dy \right] \\
 &= \frac{1}{\sqrt{4\pi kt}} \left[ - \int_{\infty}^0 e^{-(x+z)^2/4kt} \phi(z) dz + \int_0^{\infty} e^{-(x-y)^2/4kt} \phi(y) dy \right] \\
 &= \frac{1}{\sqrt{4\pi kt}} \left[ \int_0^{\infty} e^{-(x+z)^2/4kt} \phi(z) dz + \int_0^{\infty} e^{-(x-y)^2/4kt} \phi(y) dy \right] \\
 &= \frac{1}{\sqrt{4\pi kt}} \int_0^{\infty} \left[ e^{-(x-y)^2/4kt} + e^{-(x+y)^2/4kt} \right] \phi(y) dy, \quad x > 0
 \end{aligned}$$

Note that it only differs from the solution to the half-line Dirichlet problem by a minus sign. As an example, let's use this newly derived formula to solve (2.1) in the case where  $\phi(x) = 1$ . We have that

$$\begin{aligned}
 w(x, t) &= \frac{1}{\sqrt{4\pi kt}} \int_0^{\infty} e^{-(x-y)^2/4kt} dy + \frac{1}{\sqrt{4\pi kt}} \int_0^{\infty} e^{-(x+y)^2/4kt} dy \\
 &= \frac{-1}{\sqrt{\pi}} \int_{x/\sqrt{4kt}}^{-\infty} e^{-p^2} dp + \frac{1}{\sqrt{\pi}} \int_{x/\sqrt{4kt}}^{\infty} e^{-q^2} dq \\
 &= \frac{1}{\sqrt{\pi}} \int_{-\infty}^{x/\sqrt{4kt}} e^{-p^2} dp + \frac{1}{\sqrt{\pi}} \int_{x/\sqrt{4kt}}^{\infty} e^{-q^2} dq \\
 &= \frac{1}{\sqrt{\pi}} \left[ \frac{\sqrt{\pi}}{2} \left( \operatorname{erf} \left( \frac{x}{\sqrt{4kt}} \right) + 1 \right) \right] + \frac{1}{\sqrt{\pi}} \left[ \frac{\sqrt{\pi}}{2} \left( 1 - \operatorname{erf} \left( \frac{x}{\sqrt{4kt}} \right) \right) \right] \\
 &= 1
 \end{aligned}$$

where  $p = (x - y)/\sqrt{4kt}$  and  $q = (x + y)/\sqrt{4kt}$ .

### Strauss 3.1.2

Solve  $u_t = ku_{xx}$ ,  $u(x, 0) = 0$ ,  $u(0, t) = 1$  on the half-line  $0 < x < \infty$ .

Let  $v(x, t) := u(x, t) - 1$ , then  $v$  satisfies  $v_t = kv_{xx}$ ,  $v(x, 0) = -1$ ,  $v(0, t) = 0$  on the half line. By the solution formula, we thus have:

$$\begin{aligned}
 v(x, t) &= \frac{1}{\sqrt{4\pi kt}} \int_0^{\infty} \left[ e^{-(x-y)^2/4kt} - e^{-(x+y)^2/4kt} \right] (-1) dy \\
 &= \frac{-1}{\sqrt{\pi}} \int_{-x/\sqrt{4kt}}^{\infty} e^{-p^2} dp + \frac{1}{\sqrt{\pi}} \int_{x/\sqrt{4kt}}^{\infty} e^{-q^2} dq \\
 &= \frac{-1}{\sqrt{\pi}} \int_{-x/\sqrt{4kt}}^{x/\sqrt{4kt}} e^{-p^2} dp \\
 &= \frac{-2}{\sqrt{\pi}} \int_0^{x/\sqrt{4kt}} e^{-p^2} dp \\
 &= -\operatorname{erf} \left( \frac{x}{\sqrt{4kt}} \right)
 \end{aligned}$$

where  $p = (y - x)/\sqrt{4kt}$  and  $q = (x + y)/\sqrt{4kt}$ . Thus  $u(x, t) = 1 - \operatorname{erf}(x/\sqrt{4kt})$ .

## Bibliography

- [1] W. Strauss, *Partial Differential Equations: An Introduction*, 2nd edition, Wiley