

# Simulating a One-Dimensional Heisenberg Model

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## 1 Introduction

### 1.1 Description of Model

The Heisenberg model is a way of simulating a ferromagnetic system, and is a simplified model of magnetism. It is closely related to the Ising model, with the major difference being that the spin states are represented by a three-dimensional vector rather than a simple up-down state as seen in the Ising model. It is seen as an extension of the Ising model, and has been used to solve various problems in physics, including physical problems dealing with quarks. In this paper I follow a paper written by C.G Windsor et al (Referenced in section "References") and try to recreate the same algorithm. I also do further testing on various data as outlined in the "Computational details" section of that paper.

## 1.2 Mathematical Formulation of Model

In this model, I simulate a one dimensional Heisenberg model, also sometimes called a Heisenberg Chain, or Heisenberg Spin Chain. The Hamiltonian of this one dimensional Heisenberg model is:

$$\hat{H} = -J \sum_{i=1}^N \mu_i \cdot \mu_{i+1} - h \sum_{i=1}^N \mu_i \quad (1)$$

where  $J$  is the coupling constant, and the sum is over the nearest neighbours of  $i$ . I also use the Metropolis algorithm, as described in the paper "Equation of State Calculations by Fast Computing Machines", by Metropolis et al. Specifically, we will be modelling the Heisenberg XXX model, with no magnetic field. Thus, we can simplify and expand our Hamiltonian to:

$$\hat{H} = -1/2 \sum_{i=1}^N (J_x \mu_i^x \mu_{i+1}^x + J_y \mu_i^y \mu_{i+1}^y + J_z \mu_i^z \mu_{i+1}^z) \quad (2)$$

And, since we are using the Heisenberg XXX model, this can be further simplified to:

$$\hat{H} = -1/2J \sum_{i=1}^N (\mu_i^x \mu_{i+1}^x + \mu_i^y \mu_{i+1}^y + \mu_i^z \mu_{i+1}^z) \quad (3)$$

This is because in the XXX model,  $J_x = J_y = J_z$ .

## 1.3 Aims of this Paper

The aim of this paper is to describe how to implement the model, and explore what changing various constants in the equations will do to our final result.

## 2 Implementation and Simulation

### 2.1 Formulation of Simulation

Consider a chain with  $n$  spin sites, with  $n$  being some finite number, and each spin site occupied by a spin  $\mu^n$ . These spins are represented by a three-dimensional unit vector. To simplify things, we have no external magnetic field acting on this system, so our Hamiltonian can also be simplified to:

$$\hat{H} = -1/2J \sum_{i=1}^N \mu_i \cdot \mu_{i+1} \quad (4)$$

Initially, we choose a random spin configuration, and then begin making trial spins. The energy of our system is equal to our Hamiltonian, dictated in (2). If  $\Delta E < 0$  (the change in energy of the system after the trial spin is performed) then the trial spin is accepted. If  $\Delta E > 0$ , then the trial spin is accepted if and only if some transition probability is satisfied.

Our choice of transition probability will follow  $T(S_i \rightarrow S_j)$  given by

$$T(S_i \rightarrow s_j) = \frac{e^{-\Delta E/kT}}{1 + e^{-\Delta E/kT}} \quad (5)$$

We also need to change the spin direction carefully. We would like our model to be less volatile and be able to solve problems effectively. We want our spin change to be symmetrically distributed around the current spin direction to avoid this volatility.

## 2.2 The Algorithm

The algorithm for computing the total energy of our system can be executed in as little as 5 steps.

1. Create a random spin configuration for our chain. I did this by creating random  $x, y, z$  coordinates and then normalising the vector to a unit vector. This is done by:

$$\vec{V} = \frac{\vec{V}}{|\vec{V}|} \quad (6)$$

2. The energy of the system is calculated using our Hamiltonian in (2)
3. We compute a Monte Carlo move by generating a vector

$$(p_1, p_2, p_3) \quad (7)$$

where  $p_n \in (-1, 1)$ . In this case, I randomised these values, using a `random()` function. We then multiply this by a factor  $\Delta S_{max}$ . We do this to limit the maximum change of any spin, and keep it symmetrically distributed around the current spin direction. If

$$|\Delta S| > \Delta S_{max} \quad (8)$$

we compute a new  $\Delta S$  in the same way. I do this using an iterative process, calling the function again if (8) is not satisfied.

4. We do a trial spin on every node in the chain. This trial spin is only accepted if  $\Delta E < 0$  or if  $\Delta E > 0$  and the transition probability in (5) is satisfied.
5. We store the current move number, as well as the total energy of the system, as we're primarily interested in convergence rates in this paper.

Step 3 utilises the Metropolis Algorithm (Equation of State Calculations by Fast Computing Machines, Metropolis et al, 1953).

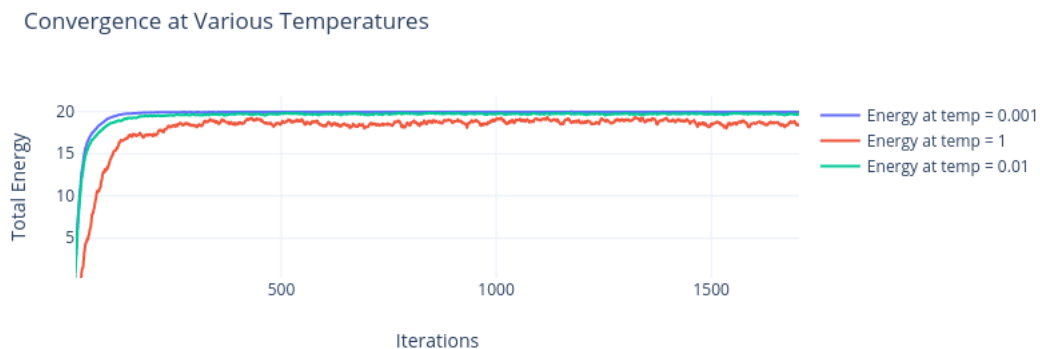
## 2.3 What our Model Solves

Our model solves a lot of practical and theoretical problems in statistical mathematics and physics. However, the objective of our model is instead to investigate what various properties the model has, and to play with some of the data. Particularly we are interested in seeing what a change in  $S_{max}$  does, as well as what changing the temperature will do to our model.  $S_{max}$  gives us mathematical insight, while changing temperatures gives us a physical insight. The mathematical insight is important to see why we have the stipulation on  $S_{max}$  and why it constitutes a Monte Carlo move.

## 2.4 Results

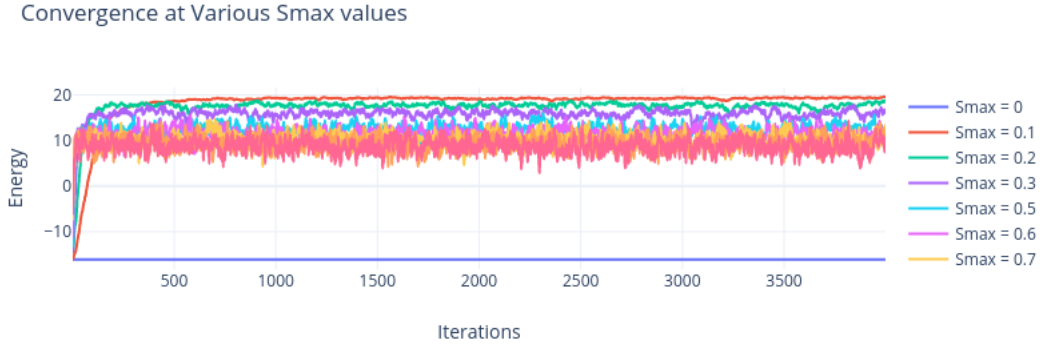
### 2.4.1 Temperature and Convergence

The total energy converges very quickly with a temperature between 0 and 1, and at high temperatures, converges more slowly.



### 2.4.2 Reducing $\Delta S_{max}$

A more interesting thing to explore is reducing the value of  $\Delta S_{max}$  used in equation (8), which will change the amount of spin our node is subject to during each iteration.



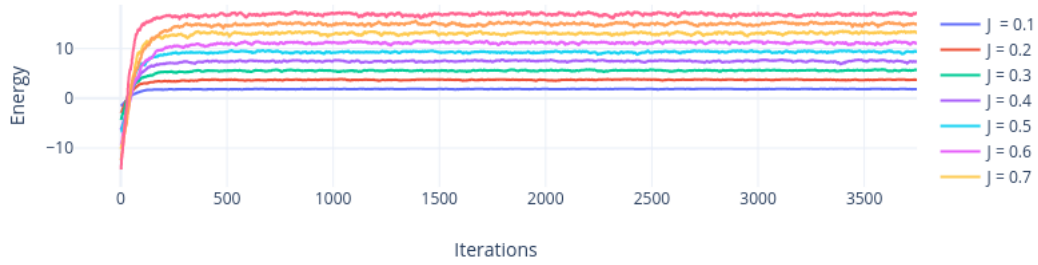
As we can see from the graph, with a low  $S_{max}$ , between  $(0, 0.1)$ , we have a less volatile model and it converges to a higher energy state.

This is an important part of this investigation, we have shown that setting an  $S_{max}$  as outlined in (8) is an important step in reducing the volatility of our model and having a better understanding of what we are modelling.

### 2.4.3 Changing our $J_x = J_y = J_z$

Changing our  $J$  value does nothing to change the convergence rate of our model. This only scales it to a higher/lower energy state.

Convergence at Various J values



### 3 Comparision to Other Simulations

In the paper "Computer simulation of the dynamics of onedimensional Heisenberg magnets, C G Windsor and JLocke-Wheaton 1976 J. Phys. C: Solid State Phys. 9 2749", comparing my simulation to the simulation used in the beginning stages of this paper, my program is exactly similar. They converge to similar rates. However, this paper uses a different transition probability as

$$T(S_i \rightarrow S_j) = e^{-\Delta E/kT} \quad (9)$$

I chose to use a different transition probability, as stated in (5) because it was the one used when we made an Ising model in class. I found that this transition probability would make my model less volatile and easier to manipulate than the transition probability given by the cited paper. The transition probability above is also more easily utilised by the Metropolis Algorithm defined in the paper "Equation of State Calculations by Fast Computing Machines, J. Chem. Phys. 21, 1087 (1953), Metropolis et al". The Metropolis Algorithm states that we should generate a candidate from an acceptance ratio

$$\mu = \frac{f'(x)}{f(x)} \quad (10)$$

which is not satisfied by the transition probability given in the paper, but is satisfied by the transition probability given in (5).

In the textbook, "An Introduction to Computer Simulation Methods Applications to Physical System Harvey Gould, Jan Tobochnik, and Wolfgang Christian" very little is given in the way of the Heisenberg model, but it is also what I used to find out that we needed to limit the amount of spin of each node per iteration. If I had not found this resource, my model would have been very volatile, or may have never converged to an equilibrium state.

## 4 Conclusion

### 4.1 Final Conclusions

A model was created which utilised Monte Carlo moves and the Metropolis Algorithm to simulate a Heisenberg chain XXX model. The major difference between this model and the Ising model is that the Heisenberg model has spin directions in three-dimensions.

It was found that there is much importance on the step of the algorithm which constitutes how much the node is spun during a Monte Carlo move. We found that if we increase the maximum amount that this node can be spun, we increase the volatility of the system. By reducing the value of  $S_{max}$ , we reduce the volatility and get a better representation of our system.

It was also found that the choice of coupling constant has very little effect on the structure



of our model, except for scaling of our model. With a higher coupling, our model has a higher total energy, and with a lower coupling it has a lower total energy. This is as expected, since  $J$  is a constant.

Lastly, we found that changing the temperature of our system affected the total energy convergence rate. As the temperature approaches zero, we see the system converging to its equilibrium much more quickly.

I would have liked to compare this to the Ising model and different rates of convergence, but in reality they solve two different things. The Ising model is usually good enough to solve problems of this type, and the Heisenberg model was developed to investigate phase transition properties.

## 4.2 Further Research

I would like to do more research on other Heisenberg models, including the XYY model, and the XYZ model. I am also interested in increasing the dimensions and moving from a Heisenberg chain to a lattice and seeing how this affects convergence rates, as well as looking at other problems that can be solved.

## 5 References

Computer simulation of the dynamics of onedimensional Heisenberg magnets, C G Windsor and J Locke-Wheaton 1976 J. Phys. C: Solid State Phys. 9 2749

Equation of State Calculations by Fast Computing Machines, J. Chem. Phys. 21, 1087 (1953),

Metropolis et al

An Introduction to Computer Simulation Methods Applications to Physical System Harvey Gould,  
Jan Tobochnik, and Wolfgang Christian

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# Heisenberg Chain

```
1  #ifndef HEISENBERGHEADERDEF
2  #define HEISENBERGHEADERDEF
3
4  #include <iostream>
5
6  class Heisenberg
7  {
8  private:
9      double** m_chain;
10     double* m_TrialSpin;
11     int m_nodes;
12     int max_iter;
13     double k_temp;
14     double m_j;
15     double m_smin;
16
17
18
19
20 public:
21     Heisenberg(int nodes=40, double temp=1, double j=1, double s_min=0.2, int max_iter=4000);
22     ~Heisenberg();
23     double Transition(double d_energy);
24     double CalculateEnergy();
25     double max(double x, double y, double z);
26     double deltaS();
27     void MonteCarloMove(double s_max);
28     double random(double m, double n);
29     void setInitialState();
30     void trialSpin(int element);
31     void printLines();
32
33 };
34
35
36 #endif
37
38
39 #include "Heisenberg.h"
40 #include <cmath>
41 #include <fstream>
42
43 Heisenberg::Heisenberg(int nodes, double temp, double j, double s_min, int max_iter)
44 {
45     m_nodes = nodes;
46     k_temp = temp;
47     m_j = j;
48     m_smin = s_min;
49
50     m_TrialSpin = new double[3];
51     m_chain = new double*[m_nodes];
```

```

52     for (int i = 0; i < m_nodes; i++)
53     {
54         m_chain[i] = new double[3];
55     }
56
57 }
58 Heisenberg::~Heisenberg()
59 {
60     for (int i = 0; i < m_nodes; i++)
61     {
62         delete[] m_chain[i];
63     }
64     delete[] m_chain;
65     delete[] m_TrialSpin;
66 }
67
68 double Heisenberg::Transition(double d_energy)
69 {
70     double value;
71     double result;
72     value = exp(-d_energy/k_temp);
73     result = value/(1+value);
74     return result;
75 }
76 double Heisenberg::CalculateEnergy()
77 {
78     double sum = 0;
79     for (int i = 0; i < m_nodes-1; i++)
80     {
81         for (int k = 0; k < 3; k++)
82         {
83             sum += m_chain[i][k]*m_chain[i+1][k];
84         }
85     }
86     for (int j = 0; j < 3; j++)
87     {
88         sum += m_chain[m_nodes-1][j]*m_chain[0][j];
89     }
90
91     return -0.5*m_j*sum;
92 }
93 void Heisenberg::setInitialState()
94 {
95
96
97     for (int i = 0; i < m_nodes; i++)
98     {
99         double a,b,c;
100         a = rand();
101         b = rand();
102         c = rand();
103         double magnitude = 1/(sqrt(pow(a,2)+pow(b,2)+pow(c,2)));
104         m_chain[i][0] = a*magnitude;
105         m_chain[i][1] = b*magnitude;
106         m_chain[i][2] = c*magnitude;
107     }

```

```

108 }
109
110 void Heisenberg::MonteCarloMove(double s_max)
111 {
112
113     double a = random(-1,1);
114     double b = random(-1,1);
115     double c = random(-1,1);
116
117     a = s_max*a;
118     b = s_max*b;
119     c = s_max*c;
120
121     if (sqrt(pow(a,2)+pow(b,2)+pow(c,2))>s_max)
122     {
123         MonteCarloMove(s_max);
124     }
125     else
126     {
127         m_TrialSpin[0] = a;
128         m_TrialSpin[1] = b;
129         m_TrialSpin[2] = c;
130     }
131 }
132
133
134 double Heisenberg::random(double m, double n)
135 {
136     return m + (rand() / ( RAND_MAX / (n-m)));
137 }
138 void Heisenberg::trialSpin(int element)
139 {
140     double EnergyOld = CalculateEnergy();
141
142
143
144     MonteCarloMove(m_smin);
145     double a = m_chain[element][0] + m_TrialSpin[0];
146     double b = m_chain[element][1] + m_TrialSpin[1];
147     double c = m_chain[element][2] + m_TrialSpin[2];
148
149     double magnitude = 1/sqrt(pow(a,2)+pow(b,2)+pow(c,2));
150
151     a = a * magnitude;
152     b = b * magnitude;
153     c = c * magnitude;
154
155     double holder_a = m_chain[element][0];
156     double holder_b = m_chain[element][1];
157     double holder_c = m_chain[element][2];
158
159     m_chain[element][0] = a;
160     m_chain[element][1] = b;
161     m_chain[element][2] = c;
162
163

```

```

164     double EnergyNew = CalculateEnergy();
165
166     double deltaEnerg = EnergyOld - EnergyNew;
167     double randomNum = random(0,1);
168
169     if (((deltaEnerg) > 0)&&(Transition(deltaEnerg)<randomNum))
170     {
171         m_chain[element][0] = holder_a;
172         m_chain[element][1] = holder_b;
173         m_chain[element][2] = holder_c;
174     }
175
176 }
177 void Heisenberg::printLines()
178 {
179
180     setInitialState();
181     std::ofstream myFile;
182     myFile.open("Project.dat");
183     for (int i = 0; i < max_iter; i++)
184     {
185         for (int j = 0; j < m_nodes; j++)
186         {
187             trialSpin(j);
188             myFile << i << "\t" << CalculateEnergy() << "\n";
189         }
190     }
191     myFile.close();
192 }
193 int main(int argc, const char** argv)
194 {
195     int i = 0;
196     double j = 0;
197     std::string s = "J";
198     while (i < 10)
199     {
200
201         char str[10];
202         sprintf(str,"%d.txt",i);
203         std::string filename = s;
204         filename.append(str);
205         std::ofstream myFile;
206
207         myFile.open(filename.c_str());
208
209         Heisenberg myHeis = Heisenberg(40,1,j,0.15,4000);
210         myHeis.setInitialState();
211
212
213         for (int i = 0; i < 4000; i++)
214         {
215             for (int j = 0; j < 40; j++)
216             {
217                 myHeis.trialSpin(j);
218             }
219             myFile << i << "\t" << myHeis.CalculateEnergy() << "\n";

```

```
220     }
221     myFile.close();
222     i += 1;
223     j += 0.1;
224 }
225 }
```