SECURE ASSISTED QUANTUM COMPUTATION

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Suppose Alice wants to perform some computation that could be done quickly on a quantum computer, but she cannot do universal quantum computation. Bob can do universal quantum computation and claims he is willing to help, but Alice wants to be sure that Bob cannot learn her input, the result of her calculation, or perhaps even the function she is trying to compute. We describe a simple, efficient protocol by which Bob can help Alice perform the computation, but there is no way for him to learn anything about it. We also discuss techniques for Alice to detect whether Bob is honestly helping her or if he is introducing errors.

1 Introduction

The idea of processing information stored in quantum states has spawned numerous cryptographic applications. A few examples include quantum key distribution [1], which allows remote parties to securely establish a shared list of random numbers; a fast quantum algorithm for factoring [2], which can be used to break certain classical cryptosystems; quantum secret sharing [3], by which a secret quantum state can be divided among several parties; quantum data hiding [4, 5, 6], which offers an information theoretically secure way of sharing a classical secret; quantum digital signatures [7], which can be used to authenticate documents; and secure quantum channels [8, 9, 10, 11], which allow secure transmission of quantum states.

But there are also a number of negative results about the possible cryptographic applications of quantum information, such as the impossibility of an unconditionally secure quantum protocol for bit commitment [12, 13]. A related result is the impossibility of "secure two-party computation," in which two parties collaborate to compute a function without revealing their inputs [14] (although in general, secure multi-party quantum computation is possible [15, 16]). However, this does not rule out all forms of collaborative computation by two parties: for example, what if one of the parties wishes to assist the other, with no possibility of learning the input or output of the computation? We will show that this kind of two-party computation can be done securely.

More precisely, the problem we will consider is the following: imagine that Alice would like to perform a quantum computation in secret, but although she can do some basic quantum gates, she does not have a full-fledged quantum computer. Bob, who runs a company that sells time on its quantum supercomputer, would like to supply Alice with the resources she needs to perform her computation. But Alice does not trust Bob; she needs to be absolutely sure that he cannot learn anything about her computation, i.e., that the state Bob sees at any time is the same, independent of her actual quantum state. Can they carry out an unconditionally secure protocol by which Bob can assist Alice? In this paper, we describe protocols that answer this question in the affirmative.

To fully specify our question, we must decide exactly what resources are allowed. Three kinds of resources must be specified: operations available to Alice, operations available to Bob, and ways in which they can communicate. We will always allow Bob to do universal quantum computation and make arbitrary quantum measurements, and we will allow bidirectional quantum communication. There are many possible restrictions on Alice's resources that might be of interest, but we will choose the most restrictive set under which Bob can help her do universal quantum computation. We will allow Alice to store quantum states and to route her qubits (i.e., perform the swap gate), but we will suppose that the only nontrivial gates she can perform are the *Pauli gates*,

$$X := \begin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix}, \quad Z := \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix} \tag{1}$$

(with which she can perform their product XZ, etc.). This gate set fails to be universal in two important ways. First, Alice cannot perform any interactions between qubits. Second, the single-qubit gates she can perform are restricted to the discrete set $\{I, X, Z, XZ\}$ (up to an overall phase^a), which forms a group (the *Pauli group*) under multiplication. In addition to the restrictions on her gates, we will suppose that Alice can only prepare the $|0\rangle$ state, and she cannot perform measurements. Alice also must be able to generate random classical bits (say, by flipping unbiased coins) and to perform Pauli gates conditioned on the values of these classical bits.

Note that a related question has been considered in the classical setting. As we argue in Section 5, there is no reasonable restriction on Alice's gate set under which she can carry out a classical protocol analogous to our quantum protocols. However, one can instead assume Alice has the ability to perform polynomial-time computation and ask whether she can securely receive help from an arbitrarily powerful Bob. For certain particular problems this is the case [17, 18], whereas for other problems (in particular, NP-hard ones), it is not [19, 18]. In contrast, our protocols make no computational assumptions, and will apply to *any* quantum circuit known to Alice, not just the computation of a particular function.

All of the protocols we present are applications of a quantum version of the Vernam cipher [20], also known as private key encryption or the one-time pad. The classical Vernam cipher works as follows: suppose Alice wishes to securely send a bit b to Charlie. She and Charlie share another bit k, the key, that is randomly chosen to be either 0 or 1, each with probability $\frac{1}{2}$. Alice computes the message bit $m = b \oplus k$, where \oplus denotes addition modulo two, and sends it to Charlie. Since he knows k, he can compute $b = m \oplus k$. However, an eavesdropper (whom we shall call Eve) who does not know k can learn nothing about b since b and m have

^aFor simplicity, we identify operators that differ by an overall phase, since such phases are irrelevant in all the situations we consider. Equivalently, we could multiply all operators by a phase so that they have unit determinant.

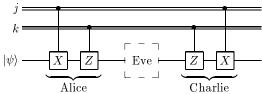


Fig. 1. Quantum circuit for the private quantum channel. The double lines represent randomly chosen classical bits j and k that are shared by Alice and Charlie. Whether the gates are performed or not depends on the values of j and k as indicated.

zero mutual information. To send multiple bits, Alice and Charlie can repeat this procedure, using a new random key bit for each message bit.

The private quantum channel is a quantum analogue of this protocol in which the key remains classical, but the channel is used to send quantum states [9, 10]. (Note that this differs from the quantum Vernam cipher, in which the key is also a quantum state [11].) In the private quantum channel, Alice and Charlie need to share two classical bits j and k for Alice to send her qubit. The circuit shown in Fig. 1 summarizes their protocol. Alice applies the unitary operator $Z^k X^j$ to her state $|\psi\rangle$ and sends the result to Charlie. In between, Eve may intercept the state, but since she doesn't know j or k, she sees the density matrix

$$\frac{1}{4} \sum_{j,k=0}^{1} Z^k X^j |\psi\rangle\langle\psi| X^j Z^k = \frac{I}{2}$$
 (2)

independent of $|\psi\rangle$. From Eve's perspective, Alice has applied the depolarizing channel, so Eve can learn nothing about the state. Although she can destroy the state or change it in some way, she cannot learn anything about it. Assuming she does nothing, Charlie will receive the state $Z^k X^j | \psi \rangle$. Since he knows the values of j and k, he can apply the inverse operation $X^{j}Z^{k}$ to recover the original state. If Alice wants to send Charlie n qubits, they can repeat the procedure independently for each qubit, using a total of 2n random classical bits as the key. From Eve's perspective, the density matrix of all n qubits is again maximally mixed, independent of Alice's state, so the procedure is secure.

We will show how the idea of a private quantum channel can be adapted to allow Bob to help Alice perform a quantum computation. The main idea behind these circuits is that Alice can use a private quantum channel to securely send qubits to herself. Since Alice is both the sender and the receiver, there is no need for her to distribute a key. Alice sends her qubits by way of Bob, who plays the part of the eavesdropper. However, instead of trying to learn Alice's state (which of course would be futile), he intentionally performs the gate Alice would like to be able to do. After he gives it back to Alice, she performs an appropriate decoding operation. Of course, this procedure is only useful if the decoding operation can be performed using only the restricted set of gates available to Alice, which is not necessarily the case for arbitrary gates.

We begin in Sec. 2 by showing how Bob can help Alice perform a measurement in the computational basis, and then show how he can help her complete her gate set to do universal quantum computation in Sec. 3. We give a unified description of k-round protocols for secure assisted gates in Sec. 4, and we discuss why they have no meaningful classical analogue in

4 Secure assisted quantum computation

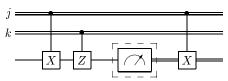


Fig. 2. Secure assisted computational basis measurement. The classical bits j and k are randomly chosen, and their values determine whether Alice applies certain gates. The meter inside a dashed box represents a computational basis measurement, the action performed by Bob if he is honest. Even if he does some other operation, he can learn nothing about Alice's data.

Sec. 5. In Sec. 6 we discuss the question of whether Alice can determine whether Bob is being honest. Finally, in Sec. 7 we conclude with a discussion and some open questions.

2 Secure assisted measurement

First, we describe how Bob can make a measurement for Alice. We begin by discussing a classical version of the protocol. Suppose Alice has a classical bit b, but although she can manipulate it, she cannot read its value. However, she can securely send her bit to Bob and ask him to read it for her. Alice chooses a key bit k at random and computes $b \oplus k$. She then gives the result to Bob, who reads the result and tells it to her. To determine the value of her original bit b, she simply flips the result if k = 1 and does nothing otherwise.

This classical procedure doesn't seem very useful since reading the value of a classical bit is usually an easy thing to do. But quantum measurement is a difficult task, so we can imagine a scenario in which Alice can coherently manipulate her qubits, but she cannot measure them. In such a situation, there is a quantum version of the above measurement protocol that allows Bob to make a measurement of Alice's state in the computational basis. Alice chooses two random bits j and k and applies the unitary operator $Z^k X^j$ to her state. She then gives the qubit to Bob. He can acquire no information from this state since from his point of view, by (2) the density matrix is maximally mixed, independent of Alice's actual state. However, if he measures the qubit in the computational basis and reports the result to Alice, she can determine the result of the corresponding measurement on her original state. The Z operator does not change the measurement result, and the X operator flips it, so Alice should flip the result if j = 1 and not flip it if j = 0. A quantum circuit for this protocol is shown in Fig. 2.

3 Secure assisted gates

We now describe how Bob can help Alice perform universal quantum computation. We do this by showing that she can perform a universal set of gates. In particular, we present circuits by which she can securely perform the Hadamard gate,

$$H := \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1\\ 1 & -1 \end{bmatrix}; \tag{3}$$

the $\pi/8$ gate,

$$T := \begin{bmatrix} 1 & 0 \\ 0 & \sqrt{i} \end{bmatrix}; \tag{4}$$

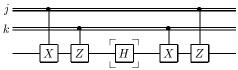


Fig. 3. Secure assisted Hadamard gate. The gate in the dashed box is what would be performed by an honest Bob, but if he performs some other operation, he will learn nothing about Alice's

and the controlled-NOT gate (a two-qubit interaction),

$$C := \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix} . \tag{5}$$

This gate set is universal for quantum computation in the sense that any unitary transformation can be approximated arbitrarily closely by some sequence of these gates [21]. Thus, the circuits for secure implementation of these gates can be used as subroutines to perform an arbitrary quantum computation.

The simplest of these constructions is the protocol for a secure Hadamard gate. Alice chooses two random classical bits j and k. To randomize her state so that Bob can learn nothing from it, Alice applies $Z^k X^j$ to her qubit. She then passes it to Bob. By (2), Bob's density matrix is maximally mixed, independent of Alice's actual state. If Bob is honest, he performs a Hadamard gate and hands the qubit back to Alice. Now she must correct her qubit so it is as if only the Hadamard were applied. Because XHZ = ZHX = H, Z can be undone by X and X can be undone by Z. Thus Alice can appropriately fix her state using only Pauli gates, regardless of the values of j and k. The resulting circuit, shown in Fig. 3, is equivalent to a Hadamard gate if Bob is honest. If he is dishonest, he can destroy Alice's qubit or give her the wrong result, but he can learn nothing from the state she gave him.

A similar procedure can be used to perform a controlled-NOT gate. Since this is a two qubit gate, Alice must choose four random classical bits j, k, l, m. She randomizes her state by applying $Z^k X^j$ to the first qubit and $Z^m X^l$ to the second. Then she gives the qubits to Bob, who is supposed to perform a controlled-NOT gate and return them to Alice. Note that by (2) applied to each of the two qubits, Bob's density matrix is maximally mixed, independent of Alice's state. Supposing that Bob performs the controlled-NOT gate as requested, Alice must correct the encoded qubits so that the overall interaction is a controlled-NOT. If j=1, then the target bit was inverted based on an inverted control bit, so she must apply X^{j} to the target. She then fixes the target bit by applying X^lZ^m and the control bit by applying X^jZ^k . However, if m=1, she has also performed a controlled-(-1) gate due to the anticommutation of X and Z. This can be fixed by applying Z^m to the control bit. The resulting circuit is shown in Fig. 4.

Although the $\pi/8$ gate is only a one-qubit gate, this operation is more complicated to implement: Alice and Bob must use a two-round protocol, as we prove in the next section. In total, Alice needs four random classical bits j, k, l, m. First, she randomizes her quantum state by applying $Z^k X^j$. She then gives it to Bob (whose density matrix, again by (2), is maximally mixed), and if he is honest, he applies a $\pi/8$ gate and gives it back to Alice. She

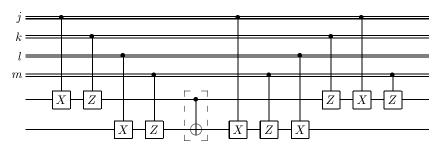


Fig. 4. Secure assisted controlled-NOT gate.

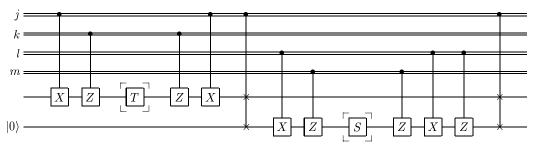


Fig. 5. Secure assisted $\pi/8$ gate. A vertical line with two \times symbols indicates a SWAP gate, classically controlled on the value of j as indicated.

can undo her randomization by applying X^jZ^k . The Z operation commutes with T, so it does not create any problems. However, $XTX = T^{\dagger}$ (up to an overall phase), which differs from T by $S := T^2$, a gate that Alice cannot perform. But she can have Bob do it for her, again encoding the qubit using (2), and since $S^2 = Z$, she will be able to undo the randomization herself. If she only asks Bob to help her perform S when j = 1, this tells him the value of j. But Alice can avoid revealing j by always asking Bob to participate in the second round, but operating on a dummy qubit when j = 0. The complete circuit is shown in Fig. 5. Assuming Bob is honest, this circuit is equivalent to the $\pi/8$ gate up to an irrelevant overall phase.

4 The Gottesman-Chuang hierarchy

In this section, we explain the existence of the gate constructions discussed in the previous section using a hierarchy of gates presented by Gottesman and Chuang [22]. We show that the gates that can be realized using a k-round protocol are exactly those in the (k+1)th level of this hierarchy.

A single-round protocol for a secure assisted gate U on n qubits requires that

$$D_j U E_j = U \quad \forall \ j \,, \tag{6}$$

where E_j, D_j are Alice's encoding and decoding operators. Let C_1 denote the Pauli group on n qubits; that is, all operators that are tensor products of n Pauli operators. Furthermore, let C_1^S denote those operators that can be written as a product of Pauli group elements and permutations of the qubits (i.e., products of SWAP gates). Gates in C_1^S are exactly the operations that Alice can perform without assistance, so $E_j, D_j \in C_1^S \,\forall j$. For simplicity, we will assume that Alice uses the encoding where E_j runs over all elements of C_1 , each

with probability $p_j = 2^{-n}$. (More generally, Alice could choose to include SWAP gates in her encoding, but this would not affect our conclusions.) The decoding operations are specified by

$$D_j = U E_j^{\dagger} U^{\dagger} \,. \tag{7}$$

Thus the requirement $D_j \in C_1^S$ allows us to classify the possible U's that can be realized in this way: U must come from the set

$$C_2^S := \{ U : \ U C_1 U^{\dagger} \subseteq C_1^S \} \,.$$
 (8)

It is not hard to show that in $fact^b$

$$C_2^S = C_2 := \{ U : UC_1 U^{\dagger} \subseteq C_1 \}. \tag{9}$$

The set C_2 forms a well known group, the Clifford group [23]. Any gates in this group, such as the Hadamard and CNOT gates, can be realized with a one-round protocol.

If Alice and Bob are willing to use a two-round protocol, they can do more gates. By the preceding discussion, Bob must do a gate from C_2 in the second round. However, in the first round, he can do any gate from

$$C_3 := \{ U : UC_1 U^{\dagger} \subseteq C_2 \}, \tag{10}$$

e.g., the $\pi/8$ gate (or the Toffoli or Fredkin gate). Different encoding/decoding pairs may require Bob to do different things in the second round, but if Alice uses the SWAP trick introduced in the secure $\pi/8$ gate construction, this need not provide him with any information. By induction, we see that in k rounds, Alice and Bob can securely perform any gate in the set

$$C_k := \{ U : UC_1 U^{\dagger} \subseteq C_{k-1} \}. \tag{11}$$

Note that gates in C_2 are not sufficient for universal computation. However, universal quantum computation can be done using C_3 gates.

At first glance, it may seem that the form of (6) is too restrictive. Why should Bob have to perform exactly the gate Alice wants? In other words, shouldn't we consider protocols of the form

$$D_j V E_j = U \quad \forall j \,, \tag{12}$$

where $V \neq U$? The answer is no, because if such a scheme exists, we can easily turn it into a scheme where V=U: we have $V=D_0^{\dagger}UE_0^{\dagger}$, which can be substituted into (12) to give

$$U = D_i D_0^{\dagger} U E_0^{\dagger} E_i \,. \tag{13}$$

By using the modified encoding operations $E'_j = E_0^{\dagger} E_j$ and decoding operations $D'_j = D_j D_0^{\dagger}$, Alice can ask Bob to perform U instead of V. Thus there is no loss of generality in the choice

^bI thank Wim van Dam for a discussion of this point.

5 Secure assisted classical gates

It is interesting to contrast these protocols for secure assisted quantum gates with the corresponding problem for reversible classical computers. If Alice cannot do universal classical computation, is it possible for Bob to assist her in a secure way? The answer is no: secure assisted universality can only be achieved in a meaningful way in the quantum setting. This is essentially because the ability to do classical operations is assumed in the quantum case, so classically controlled quantum gates are no more difficult to implement than unconditional quantum gates; but a classically controlled classical gate can be a more powerful resource than the original classical gate. For example, the controlled-NOT gate is not universal for classical computation, but a controlled-controlled-NOT (more commonly known as a Toffoli gate) is universal.

To see why secure assisted classical universality is not possible, we introduce a classical analogue of the Gottesman-Chuang hierarchy. For n-bit gates, let \tilde{C}_1 denote gates that are tensor products of n identity and NOT gates. Also, let

$$\tilde{C}_k := \{ P : P\tilde{C}_1 P^{-1} \subseteq \tilde{C}_{k-1} \}$$
 (14)

denote the set of classical gates (permutations) that map \tilde{C}_1 gates into \tilde{C}_{k-1} gates under conjugation. Just as in the quantum case, gates in \tilde{C}_2 are not sufficient for universal computation, but there are gates from \tilde{C}_3 (e.g., the Toffoli and Fredkin gates) that are universal. Thus the natural restriction on Alice's gate set is to allow her to do only \tilde{C}_2 gates. In particular, since she must use gates that are controlled based on the values of her random key bits, she will be interested in using controlled- \tilde{C}_1 gates, all of which are in \tilde{C}_2 . However, with a single-round protocol, she can only use controlled- \tilde{C}_1 gates to build \tilde{C}_2 gates on her computational bits, which gives her no additional computational power. Since she cannot perform the controlled-SWAP (Fredkin) gate, she cannot securely perform a multi-round protocol. Thus there is no way for her to achieve secure assisted universality.

However, there are particular examples in which Alice and Bob can perform a computation securely even in the classical case. Some work along these lines was mentioned in Section 1, but a simple example that shows the advantage of having a particular problem in mind is the following. Suppose Alice would like to find a satisfying assignment for a Boolean formula containing n variables. Alice generates n random bits, one for each of the variables in the formula. If the bit corresponding to a particular variable is zero, she does nothing; if that bit is one, she inverts the variable wherever it appears in the formula. She then tells Bob the formula, and if he is honest, he gives her a satisfying assignment. To find a satisfying assignment for her original problem, she flips the bits corresponding to the variables that were inverted in the formula she gave to Bob. Although Bob can learn a lot about the structure of the problem, he cannot learn the particular satisfying assignment of her original problem.

6 Keeping Bob honest

The primary weakness of these protocols is that although he can learn nothing about the states Alice gives him, Bob can easily prevent her from performing her computation. He could simply not return her qubits, or worse yet, he could ruin the computation by performing the wrong

^cI thank Sam Gutmann for suggesting this example.

gate. This weakness is inherited from the private quantum channel, and there is no way to avoid it altogether. However, although she cannot force Bob to help her, there may be simple ways for Alice to detect whether Bob is cheating.

If Alice is asking Bob to help her solve a problem in the computational complexity class NP, there is a particularly simple way for her to check his honesty. Presumably, Alice is asking for Bob's help because the problem can be solved much faster on a quantum computer than on a classical computer—for example, she might ask him to help her perform Shor's algorithm to factor a number. But if Alice has access to a classical computer and the problem is in NP, she can easily check the solution to see if it is right. In the example of factoring, Alice simply multiplies the resulting factors to see if they give her input.

However, what if Alice cannot readily check the solution? Can she still efficiently detect whether Bob is cheating? Intuitively, it seems that Alice should be able to gain some confidence that Bob is behaving honestly by performing tests of his actions using a randomly chosen subset of her inputs. If these tests fully characterize Bob's operations, they can be used to bound the probability that Bob is cheating. Indeed, such a procedure provides an efficient way to check whether Bob is cheating in the restricted scenario in which he must act as a memoryless black box, as can be shown using ideas along the lines of [24]. However, in general, Bob could introduce errors adversarially rather than randomly, which presents a more difficult verification problem. We leave the general adversarial scenario as an open problem. A solution to this problem might also be relevant to fault-tolerant quantum computing, where the assumption that Bob acts as a memoryless black box corresponds to an assumption of independent errors, and the general adversarial scenario accommodates errors of the most general type (which are hopefully nevertheless small).

Discussion

We have shown that it is possible for a party who cannot do universal quantum computation (Alice) to have her computational power augmented by another party (Bob) without compromising the security of the computation. Furthermore, we have briefly discussed ways of detecting whether Bob is truly being helpful—a problem that deserves further study. The protocols we have described, and more general protocols for verifying the validity of Bob's actions, might prove useful for assuring the security of certain quantum information processing tasks.

We have focused on preventing Bob from obtaining information about the states Alice gives him, and we have not considered the information he might obtain from the particular gates Alice asks him to perform. We can imagine that she might want to prevent him from learning something this way. For example, in the context of programmable gates, Vidal and Cirac considered a different scenario in which Bob can learn Alice's input, but she does not want him to know the function she is trying to compute [25]. In the context of the present paper, it is not particularly difficult to prevent Bob from learning the function. The protocol can simply consist of Bob performing a fixed sequence of gates, cycling through Hadamard, CNOT, and $\pi/8$. If a particular gate is not needed, she can supply Bob with junk qubits. With this protocol, the number of gates is increased by at most a factor of three. Since Alice does not send any classical information to Bob to describe her circuit, and since we have already established that he can learn nothing from her quantum states, it follows that he cannot learn anything about which gates are being used. The only thing Bob can learn is the length of the protocol, i.e., the total number of gates Alice has him perform. Even this meager amount of information can be reduced (although not eliminated), since Alice is free to add additional unnecessary gate requests, so that Bob can only learn an upper bound on the number of gates in Alice's circuit.

Note that there is an analogy between programmable gates and secure assisted gates: whereas programmable gates generalize identity teleportation to gate teleportation, secure assisted gates generalize an identity private quantum channel to a private quantum channel that performs a gate. Furthermore, the Gottesman-Chuang hierarchy plays a similar role in the construction of gate teleportation circuits [26] and in showing that two-qubit measurements are universal for quantum computation [27].

There are many possible variants of this problem depending on the resources Alice and Bob are allowed to use. For example, if Alice is only allowed to perform single-qubit measurements, Bob can supply her with a cluster state [28]. Perhaps other resource limitations would lead to interesting forms of secure assisted quantum computation.

Finally, it might be useful to consider restricting the total amount of information transfer. We have assumed that Alice and Bob have an inexpensive quantum channel, so they can send quantum states back and forth as many times as they wish. But this may not be a realistic assumption. If Alice and Bob are connected only by a very slow or expensive channel—or perhaps only by a classical channel, with a small reserve of prior shared entanglement—can they still accomplish interesting computational tasks? In other words, can they perform secure remote quantum computation? We should not expect Bob to enable Alice to do secure universal computation on remote data, but she might nevertheless be able to perform certain tasks securely.

Acknowledgements

I thank Isaac Chuang, Wim van Dam, Edward Farhi, Jeffrey Goldstone, Sam Gutmann, Aram Harrow, Debbie Leung, Hoi-Kwong Lo, Ben Recht, Adam Smith, and Ronald de Wolf for helpful discussions. This work was supported by the Fannie and John Hertz Foundation, by the Cambridge–MIT Institute, by the National Security Agency and Advanced Research and Development Activity under Army Research Office contract number DAAD19-01-1-0656, and by the Department of Energy under cooperative research agreement DE-FC02-94ER40818.

References

- 1. C. H. Bennett and G. Brassard, Quantum cryptography: public key distribution and coin tossing, in Proc. IEEE International Conference on Computers, Systems, and Signal Processing, 175 (IEEE Press, Los Alamitos, CA, 1984).
- 2. P. W. Shor, Polynomial-time algorithms for prime factorization and discrete logarithms on a quantum computer, SIAM J. Comp. 26, 1484 (1997); see also quant-ph/9508027.
- R. Cleve, D. Gottesman, and H.-K. Lo, How to share a quantum secret, Phys. Rev. A 82, 648 (1999), quant-ph/9901025.
- 4. D. P. DiVincenzo, D. W. Leung, and B. M. Terhal, *Quantum data hiding*, IEEE Trans. Inf. Theory 48, 580 (2002), quant-ph/0103098.
- T. Eggeling and R.F. Werner, Hiding classical data in multi-partite quantum systems, Phys. Rev. Lett. 89, 097905 (2002), quant-ph/0203004.

- 7. D. Gottesman and I. L. Chuang, Quantum digital signatures, quant-ph/0105032.
- 8. C. H. Bennett, G. Brassard, C. Crépeau, R. Jozsa, A. Peres, and W. Wootters, *Teleporting an unknown quantum state via dual classical and EPR channels*, Phys. Rev. Lett. **70**, 1895 (1993).
- 9. A. Ambainis, M. Mosca, A. Tapp, and R. de Wolf, *Private quantum channels*, in Proc. 41st Annual Symposium on Foundations of Computer Science, 547 (IEEE Press, Los Alamitos, CA, 2000); see also quant-ph/0003101.
- P. O. Boykin and V. Roychowdhury, Optimal encryption of quantum bits, Phys. Rev. A 67, 042317 (2003), quant-ph/0003059.
- 11. D. W. Leung, Quantum Vernam cipher, Quant. Inf. Comp. 2, 14 (2001), quant-ph/0012077.
- 12. H.-K. Lo and H. F. Chau, Is quantum bit commitment really possible?, Phys. Rev. Lett. 78, 3410 (1997), quant-ph/9603004.
- 13. D. Mayers, Unconditionally secure quantum bit commitment is impossible, Phys. Rev. Lett 78, 3414 (1997), quant-ph/9605044.
- H.-K. Lo, Insecurity of quantum secure computations, Phys. Rev. A 56, 1154 (1997), quant-ph/9611031.
- H. F. Chau, Quantum-classical complexity-security tradeoff in secure multi-party computations, Phys. Rev. A 61, 032308 (2000), quant-ph/9901024.
- A. Smith, Multi-party Quantum Computation, M.S. thesis, Massachusetts Institute of Technology (2001), quant-ph/0111030.
- 17. J. Feigenbaum, Encrypting problem instances, or, Can you take advantage of someone without having to trust him?, Proc. CRYPTO '85, ed. H. C. Williams, Lecture Notes in Computer Science 218, 477 (1986).
- J. Feigenbaum and L. Fortnow, Random self-reducibility of complete sets, SIAM J. Comp. 22, 994 (1993).
- 19. M. Abadi, J. Feigenbaum, and J. Kilian, On hiding information from an oracle, J. Computer and System Sciences 39, 21 (1989).
- 20. G. S. Vernam, Cipher printing telegraph systems for secret wire and radio telegraphic communications, J. American Inst. Elec. Eng. 55, 109 (1926).
- 21. P. O. Boykin, T. Mor, M. Pulver, V. Roychowdhury, and F. Vatan, A new universal and fault-tolerant quantum basis, Inf. Proc. Lett. 75, 101 (2000), quant-ph/9906054.
- 22. D. Gottesman and and I. L. Chuang, Quantum teleportation is a universal computational primitive, Nature 402, 390 (1999), quant-ph/9908010.
- 23. D. Gottesman, *The Heisenberg representation of quantum computers*, in Group 22: Proceedings of the XXII International Colloquium on Group Theoretical Methods in Physics, eds. S. P. Corney, R. Delbourgo, and P. D. Jarvis, 32 (International Press, Cambridge, MA, 1999), quant-ph/9807006.
- 24. W. van Dam, F. Magniez, M. Mosca, and M. Santha, Self-testing of universal and fault-tolerant sets of quantum gates, in Proc. 32nd Annual ACM Symposium on the Theory of Computing, 688 (ACM Press, New York, 2000), quant-ph/9904108.
- G. Vidal and J. I. Cirac, Storage of quantum dynamics on quantum states: a quasi-perfect programmable quantum gate, quant-ph/0012067.
- X. Zhou, D. W. Leung, and I. L. Chuang, Methodology for quantum logic gate constructions, Phys. Rev. A 62, 052316 (2000), quant-ph/0002039.
- D. W. Leung, Two-qubit projective measurements are universal for quantum computation, quant-ph/0111122
- R. Raussendorf and H. J. Briegel, A one-way quantum computer, Phys. Rev. Lett. 86, 5188 (2001), quant-ph/0010033.