

1 DG Error function

For a given set of N particles with positions \vec{r}_i , the distance geometry error function is given as:

$$\begin{aligned}
\text{errf}(\{\vec{r}_i\}) = & \underbrace{\sum_{i < j}^N \left[\max^2 \left(0, \frac{(\vec{r}_j - \vec{r}_i)^2}{U_{ij}^2} - 1 \right) + \max^2 \left(0, \frac{2L_{ij}^2}{L_{ij}^2 + (\vec{r}_j - \vec{r}_i)^2} - 1 \right) \right]}_{\text{Distance errors}} \\
& + \underbrace{\sum_{(S_\alpha, S_\beta, S_\gamma, S_\delta, U_V, L_V) \in C} \left[\max^2(0, V_{\alpha\beta\gamma\delta}(\{\vec{r}_i\}) - U_V) + \max^2(0, L_V - V_{\alpha\beta\gamma\delta}(\{\vec{r}_i\})) \right]}_{\text{Chiral errors}} \\
& + \underbrace{\sum_{(S_\alpha, S_\beta, S_\gamma, S_\delta, U_\phi, L_\phi) \in D} \max^2 \left(0, \left| \phi_{\alpha\beta\gamma\delta}(\{\vec{r}_i\}) + \begin{cases} 2\pi & \left| \phi < \frac{U_\phi + L_\phi - 2\pi}{2} \right| \\ 0 & \left| \phi < \frac{U_\phi + L_\phi - 2\pi}{2} \right| \\ -2\pi & \left| \phi > \frac{U_\phi + L_\phi + 2\pi}{2} \right| \end{cases} - \frac{U_\phi + L_\phi}{2} \right| - \frac{U_\phi - L_\phi}{2} \right)}_{\text{Dihedral errors}}.
\end{aligned} \tag{1}$$

Within the distance errors, the symbols U_{ij} and L_{ij} are the upper and lower distance bounds for the atoms i and j .

Within the chiral errors, C is a set of chiral constraint tuples consisting of the sets $S_\alpha, S_\beta, S_\gamma$ and S_δ . These contain mutually disjoint particle indices and contain at least one element. In further notation, \vec{s} denotes the average spatial position of all elements of a set, e.g. for S_α :

$$\vec{s}_\alpha = \frac{1}{|S_\alpha|} \sum_{i=1}^{|S_\alpha|} \vec{r}_{S_\alpha, i}, \tag{2}$$

where $|S_\alpha|$ denotes the number of elements in the set and $S_{\alpha, i}$ is the i -th element in the set.

The constraint tuple further consists of the scalars U_V and L_V , which are upper and lower bounds on the volume spanned by the average positions of the sets $S_\alpha, S_\beta, S_\gamma$ and S_δ . This volume is calculated in the symbol $V_{\alpha\beta\gamma\delta}$, which is the signed tetrahedron volume spanned by $\vec{s}_\alpha, \vec{s}_\beta, \vec{s}_\gamma$ and \vec{s}_δ :

$$V_{\alpha\beta\gamma\delta}(\{\vec{r}_i\}) = (\vec{s}_\alpha - \vec{s}_\delta)^T \cdot [(\vec{s}_\beta - \vec{s}_\delta) \times (\vec{s}_\gamma - \vec{s}_\delta)]. \tag{3}$$

It is important to note that tetrahedron volumes such as U_V are signed values. On odd permutations of indices, these quantities change sign.

Within the dihedral errors, D is a set of dihedral constraint tuples. Each tuple consists of particle index sets $S_\alpha, S_\beta, S_\gamma$ and S_δ . Exactly as for chiral errors, these sets do not intersect and each contain at least a single element. Furthermore, U_ϕ and L_ϕ are upper and lower bounds on the dihedral angle ϕ defined by the particle index sets:

$$\begin{aligned}
\phi_{\alpha\beta\gamma\delta}(\{\vec{r}_i\}) = & \text{atan2} \left(\frac{[(\vec{s}_\beta - \vec{s}_\alpha) \times (\vec{s}_\gamma - \vec{s}_\beta)] \cdot [(\vec{s}_\gamma - \vec{s}_\beta) \times (\vec{s}_\delta - \vec{s}_\gamma)]}{|(\vec{s}_\gamma - \vec{s}_\beta)|}, \right. \\
& \left. \frac{[(\vec{s}_\beta - \vec{s}_\alpha) \times (\vec{s}_\gamma - \vec{s}_\beta)] \cdot [(\vec{s}_\gamma - \vec{s}_\beta) \times (\vec{s}_\delta - \vec{s}_\gamma)]}{|(\vec{s}_\gamma - \vec{s}_\beta)|} \right)
\end{aligned} \tag{4}$$

Here we have used a three-vector dihedral angle definition and inserted the index set average-position differences that constitute the dihedral angle. Another definition of the dihedral angle with merely the inverse cosine is:

$$\phi_{\alpha\beta\gamma\delta}(\{\vec{r}_i\}) = \arccos \frac{[(\vec{s}_\alpha - \vec{s}_\beta) \times (\vec{s}_\beta - \vec{s}_\gamma)] \cdot [(\vec{s}_\beta - \vec{s}_\gamma) \times (\vec{s}_\delta - \vec{s}_\gamma)]}{|(\vec{s}_\alpha - \vec{s}_\beta) \times (\vec{s}_\beta - \vec{s}_\gamma)| |(\vec{s}_\beta - \vec{s}_\gamma) \times (\vec{s}_\delta - \vec{s}_\gamma)|}, \tag{5}$$

whose derivatives w.r.t. the constituting position vectors are considerably easier to evaluate.

2 Gradient of the error function

E is scalar-valued, so the columnar gradient of the error function is composed of partial derivatives to the individual position vectors \vec{r}_i :

$$\nabla E = \left(\frac{\partial E}{\partial \vec{r}_1}, \frac{\partial E}{\partial \vec{r}_2}, \dots, \frac{\partial E}{\partial \vec{r}_N} \right)$$

Each individual component ($\partial E / \partial \vec{r}_\xi$) is a vector whose components are the scalar derivatives:

$$\frac{\partial E}{\partial \vec{r}_\xi} = \begin{pmatrix} \partial E / \partial r_{\xi,x} \\ \partial E / \partial r_{\xi,y} \\ \partial E / \partial r_{\xi,z} \end{pmatrix}$$

We split the problem into the main terms:

$$\begin{aligned} \frac{\partial}{\partial \vec{r}_\xi} \text{erf}(\{\vec{r}_i\}) &= \underbrace{\frac{\partial}{\partial \vec{r}_\xi} \left(\sum_{i < j}^N \max^2 \left(0, \frac{(\vec{r}_j - \vec{r}_i)^2}{U_{ij}^2} - 1 \right) \right)}_{(1)} + \underbrace{\frac{\partial}{\partial \vec{r}_\xi} \left(\sum_{i < j}^N \max^2 \left(0, \frac{2L_{ij}^2}{L_{ij}^2 + (\vec{r}_j - \vec{r}_i)^2} - 1 \right) \right)}_{(2)} \\ &+ \underbrace{\frac{\partial}{\partial \vec{r}_\xi} \sum_{(S_\alpha, S_\beta, S_\gamma, S_\delta, U_V, L_V) \in C} \max^2(0, V_{\alpha\beta\gamma\delta}(\{\vec{r}_i\}) - U_V)}_{(3)} \\ &+ \underbrace{\frac{\partial}{\partial \vec{r}_\xi} \sum_{(S_\alpha, S_\beta, S_\gamma, S_\delta, U_V, L_V) \in C} \max^2(0, L_V - V_{\alpha\beta\gamma\delta}(\{\vec{r}_i\}))}_{(4)} \\ &+ \underbrace{\frac{\partial}{\partial \vec{r}_\xi} \sum_{(S_\alpha, S_\beta, S_\gamma, S_\delta, U_\phi, L_\phi) \in D} \max^2 \left(0, \left| \phi_{\alpha\beta\gamma\delta}(\{\vec{r}_i\}) + \begin{cases} 2\pi & \phi < \frac{U_\phi + L_\phi - 2\pi}{2} \\ 0 & \phi > \frac{U_\phi + L_\phi + 2\pi}{2} \\ -2\pi & \phi > \frac{U_\phi + L_\phi + 2\pi}{2} \end{cases} \right| - \frac{U_\phi + L_\phi}{2} \right| - \frac{U_\phi - L_\phi}{2} \right)}_{(5)}. \end{aligned}$$

2.1 Distance error terms

We begin with (1). We use the chain rule:

$$\frac{\partial}{\partial \vec{r}_\xi} \sum_{i < j}^N \max^2 \left(0, \frac{(\vec{r}_j - \vec{r}_i)^2}{U_{ij}^2} - 1 \right) = 2 \sum_{i < j}^N \max \left(0, \frac{(\vec{r}_j - \vec{r}_i)^2}{U_{ij}^2} - 1 \right) \frac{\partial}{\partial \vec{r}_\xi} \max \left(0, \frac{(\vec{r}_j - \vec{r}_i)^2}{U_{ij}^2} - 1 \right). \quad (6)$$

If $\xi = i$, then:

$$\frac{\partial}{\partial \vec{r}_i} \frac{(\vec{r}_j - \vec{r}_i)^2}{U_{ij}^2} - 1 = -\frac{2}{U_{ij}^2} (\vec{r}_j - \vec{r}_i), \quad (7)$$

and likewise, but positive, for $\xi = j$. Consequently:

$$\frac{\partial}{\partial \vec{r}_\xi} \left(\max \left(0, \frac{(\vec{r}_j - \vec{r}_i)^2}{U_{ij}^2} - 1 \right) \right) = \frac{2}{U_{ij}^2} (\vec{r}_j - \vec{r}_i) \begin{cases} -1 & \xi = i \\ 1 & \xi = j \\ 0 & \text{else} \end{cases} = \frac{2}{U_{ij}^2} (\vec{r}_j - \vec{r}_i) (\delta_{\xi j} - \delta_{\xi i}), \quad (8)$$

where we have discarded the possibility of $\frac{(\vec{r}_j - \vec{r}_i)^2}{U_{ij}^2} - 1 < 0$ since this case is adequately covered by the first maximum function. δ is the Kronecker delta. So, in total:

$$\textcircled{1} = \sum_{i < j}^N (\delta_{\xi j} - \delta_{\xi i}) \underbrace{\frac{4}{U_{ij}^2} (\vec{r}_j - \vec{r}_i) \max \left(0, \frac{(\vec{r}_j - \vec{r}_i)^2}{U_{ij}^2} - 1 \right)}_{f(i,j)} \quad (9)$$

We can transform the summation further with Kronecker deltas:

$$\sum_{i < j}^N (\delta_{\xi j} - \delta_{\xi i}) f(i, j) = \sum_{i=1}^{\xi-1} f(i, \xi) - \sum_{j=\xi+1}^N f(\xi, j) \quad (10)$$

$$= \sum_{i=1}^{\xi-1} f(i, \xi) + \sum_{i=\xi+1}^N f(i, \xi) \quad (11)$$

$$= \sum_{i=1}^N (1 - \delta_{i\xi}) f(i, \xi), \quad (12)$$

where we have used that $f(i, j) = -f(j, i)$ (see Eq. 9). All in all:

$$\textcircled{1} = \sum_{i=1}^N (1 - \delta_{i\xi}) \frac{4}{U_{i\xi}^2} (\vec{r}_\xi - \vec{r}_i) \max \left(0, \frac{(\vec{r}_\xi - \vec{r}_i)^2}{U_{i\xi}^2} - 1 \right) \quad (13)$$

Let us continue with $\textcircled{2}$. The chain rule gives:

$$\frac{\partial}{\partial \vec{r}_\xi} \sum_{i < j}^N \max^2 \left(0, \frac{2L_{ij}^2}{L_{ij}^2 + (\vec{r}_j - \vec{r}_i)^2} - 1 \right) = 2 \sum_{i < j}^N \max \left(0, \frac{2L_{ij}^2}{L_{ij}^2 + (\vec{r}_j - \vec{r}_i)^2} - 1 \right) \frac{\partial}{\partial \vec{r}_\xi} \max \left(0, \frac{2L_{ij}^2}{L_{ij}^2 + (\vec{r}_j - \vec{r}_i)^2} - 1 \right) \quad (14)$$

For $\xi = i$,

$$\frac{\partial}{\partial \vec{r}_i} \left(\frac{2L_{ij}^2}{L_{ij}^2 + (\vec{r}_j - \vec{r}_i)^2} - 1 \right) = \frac{4L_{ij}^2 (\vec{r}_j - \vec{r}_i)}{(L_{ij}^2 + (\vec{r}_j - \vec{r}_i)^2)^2}, \quad (15)$$

and likewise, but negative, for $\xi = j$. Therefore,

$$\begin{aligned} \frac{\partial}{\partial \vec{r}_\xi} \left(\max \left(0, \frac{2L_{ij}^2}{L_{ij}^2 + (\vec{r}_j - \vec{r}_i)^2} - 1 \right) \right) &= \frac{4L_{ij}^2 (\vec{r}_j - \vec{r}_i)}{(L_{ij}^2 + (\vec{r}_j - \vec{r}_i)^2)^2} \begin{cases} 1 & \xi = i \\ -1 & \xi = j \\ 0 & \text{else} \end{cases} \\ &= \frac{4L_{ij}^2 (\vec{r}_j - \vec{r}_i)}{(L_{ij}^2 + (\vec{r}_j - \vec{r}_i)^2)^2} (\delta_{\xi i} - \delta_{\xi j}), \end{aligned}$$

where we have excluded the possibility of $\frac{2L_{ij}^2}{L_{ij}^2 + (\vec{r}_j - \vec{r}_i)^2} - 1 < 0$ since this case is covered by the first maximum function. Altogether:

$$\textcircled{2} = \sum_{i < j}^N (\delta_{\xi i} - \delta_{\xi j}) \underbrace{\frac{8L_{ij}^2 (\vec{r}_j - \vec{r}_i)}{(L_{ij}^2 + (\vec{r}_j - \vec{r}_i)^2)^2} \max \left(0, \frac{2L_{ij}^2}{L_{ij}^2 + (\vec{r}_j - \vec{r}_i)^2} - 1 \right)}_{g(i,j)}. \quad (16)$$

Transforming the summation:

$$\sum_{i < j}^N (\delta_{\xi i} - \delta_{\xi j}) g(i, j) = - \sum_{i=1}^{\xi-1} g(i, \xi) + \sum_{j=\xi+1}^N g(\xi, j) \quad (17)$$

$$= \sum_{i=1}^{\xi-1} g(\xi, i) + \sum_{i=\xi+1}^N g(\xi, i) \quad (18)$$

$$= \sum_{i=1}^N (1 - \delta_{i\xi}) g(\xi, i), \quad (19)$$

Where we have used $g(i, j) = -g(j, i)$ (see Eq. 16). All in all:

$$\textcircled{2} = \sum_{i=1}^N (1 - \delta_{i\xi}) \frac{8L_{\xi i}^2 (\vec{r}_i - \vec{r}_\xi)}{(L_{\xi i}^2 + (\vec{r}_i - \vec{r}_\xi)^2)^2} \max \left(0, \frac{2L_{\xi i}^2}{L_{\xi i}^2 + (\vec{r}_i - \vec{r}_\xi)^2} - 1 \right) \quad (20)$$

2.2 Chiral error terms

Next, we consider $\textcircled{3}$. Applying the chain rule gives:

$$\begin{aligned} & \frac{\partial}{\partial \vec{r}_\xi} \sum_{(S_\alpha, S_\beta, S_\gamma, S_\delta, U_V, L_V) \in C} \max^2(0, V_{\alpha\beta\gamma\delta}(\{\vec{r}_i\}) - U_V) \\ &= 2 \sum_{(S_\alpha, S_\beta, S_\gamma, S_\delta, U_V, L_V) \in C} \max(0, V_{\alpha\beta\gamma\delta}(\{\vec{r}_i\}) - U_V) \frac{\partial}{\partial \vec{r}_\xi} \max(0, V_{\alpha\beta\gamma\delta}(\{\vec{r}_i\}) - U_V) \end{aligned}$$

The partial derivatives of $V_{\alpha\beta\gamma\delta}(\{\vec{r}_i\})$ with respect to \vec{r}_ξ are split into five cases. The index ξ can be an element of one of the four sets $S_\alpha, S_\beta, S_\gamma, S_\delta$ (and only one, since they are mutually disjoint) or not. The derivative for the last case is zero. In the remaining cases, one average vector \vec{s} is a function of \vec{r}_ξ but the rest are not. The partial derivative of any average vector \vec{s} is:

$$\frac{\partial}{\partial \vec{r}_\xi} \vec{s}_\alpha = \frac{\partial}{\partial \vec{r}_\xi} \frac{1}{|S_\alpha|} \sum_{i=1}^{|S_\alpha|} \vec{r}_{S_\alpha, i} = \begin{cases} \frac{1}{|S_\alpha|} \mathbf{I}_3 & \xi \in S_\alpha \\ 0 & \text{else} \end{cases} \quad (21)$$

where \mathbf{I}_3 denotes the three dimensional identity matrix. The individual set membership cases are thus as follows:

$$\textcircled{S_\alpha} \quad \frac{\partial}{\partial \vec{r}_\xi} \left\{ (\vec{s}_\alpha - \vec{s}_\delta)^T \cdot [(\vec{s}_\beta - \vec{s}_\delta) \times (\vec{s}_\gamma - \vec{s}_\delta)] \right\} \quad (22)$$

$$= \frac{\partial}{\partial \vec{r}_\xi} \left\{ \vec{s}_\alpha^T \cdot [(\vec{s}_\beta - \vec{s}_\delta) \times (\vec{s}_\gamma - \vec{s}_\delta)] \right\} - \vec{0} \quad (23)$$

$$= \frac{1}{|S_\alpha|} [(\vec{s}_\beta - \vec{s}_\delta) \times (\vec{s}_\gamma - \vec{s}_\delta)] \quad (24)$$

In shorthand notation, in which all \vec{s} symbols are replaced by their subscripts:

$$\textcircled{S_\beta} \quad \frac{\partial}{\partial \vec{r}_\xi} (\vec{\alpha} - \vec{\delta})^T \cdot [(\vec{\beta} - \vec{\delta}) \times (\vec{\gamma} - \vec{\delta})] \quad (25)$$

$$= \frac{\partial}{\partial \vec{r}_\xi} (\vec{\alpha} - \vec{\delta})^T \cdot \left[\vec{\beta} \times \vec{\gamma} - \vec{\beta} \times \vec{\delta} - \vec{\delta} \times \vec{\gamma} + \underbrace{\vec{\delta} \times \vec{\delta}}_{=0} \right] \quad (26)$$

$$= \frac{\partial}{\partial \vec{r}_\xi} \left\{ \vec{\beta}^T \cdot [\vec{\gamma} \times (\vec{\alpha} - \vec{\delta})] - \vec{\beta}^T \cdot [\vec{\delta} \times (\vec{\alpha} - \vec{\delta})] \right\} - \vec{0} \quad (27)$$

$$= \frac{1}{|S_\beta|} (\vec{\gamma} - \vec{\delta}) \times (\vec{\alpha} - \vec{\delta}) = -\frac{1}{|S_\beta|} (\vec{\alpha} - \vec{\delta}) \times (\vec{\gamma} - \vec{\delta}) \quad (28)$$

$$\textcircled{S_\gamma} \quad \frac{\partial}{\partial \vec{r}_\xi} (\vec{\alpha} - \vec{\delta})^T \cdot [\vec{\beta} \times \vec{\gamma} - \vec{\beta} \times \vec{\delta} - \vec{\delta} \times \vec{\gamma}] \quad (29)$$

$$= \frac{\partial}{\partial \vec{r}_\xi} \left\{ \vec{\gamma}^T \cdot [(\vec{\alpha} - \vec{\delta}) \times \vec{\beta}] - \vec{\gamma}^T \cdot [(\vec{\alpha} - \vec{\delta}) \times \vec{\delta}] \right\} - \vec{0} \quad (30)$$

$$= (\vec{\alpha} - \vec{\delta}) \times (\vec{\beta} - \vec{\delta}) \quad (31)$$

$$\textcircled{S_\delta} \quad \frac{\partial}{\partial \vec{r}_\xi} (\vec{\alpha} - \vec{\delta})^T \cdot [\vec{\beta} \times \vec{\gamma} - \vec{\beta} \times \vec{\delta} - \vec{\delta} \times \vec{\gamma}] \quad (32)$$

$$= \frac{\partial}{\partial \vec{r}_\xi} \vec{\alpha}^T \cdot [\vec{\beta} \times \vec{\gamma} - \vec{\beta} \times \vec{\delta} - \vec{\delta} \times \vec{\gamma}] - \frac{\partial}{\partial \vec{r}_\xi} \vec{\delta}^T \cdot [\vec{\beta} \times \vec{\gamma} - \vec{\beta} \times \vec{\delta} - \vec{\delta} \times \vec{\gamma}] \quad (33)$$

$$= \underbrace{\frac{\partial}{\partial \vec{r}_\xi} \vec{\alpha}^T \cdot (\vec{\beta} \times \vec{\gamma})}_{=0} - \frac{\partial}{\partial \vec{r}_\xi} \vec{\alpha}^T \cdot (\vec{\beta} \times \vec{\delta}) - \frac{\partial}{\partial \vec{r}_\xi} \vec{\alpha}^T \cdot (\vec{\delta} \times \vec{\gamma}) - \frac{\partial}{\partial \vec{r}_\xi} \vec{\delta}^T \cdot (\vec{\beta} \times \vec{\gamma}) + \underbrace{\frac{\partial}{\partial \vec{r}_\xi} \vec{\delta}^T \cdot (\vec{\delta} \times \vec{\gamma})}_{=0} + \underbrace{\frac{\partial}{\partial \vec{r}_\xi} \vec{\delta}^T \cdot (\vec{\beta} \times \vec{\delta})}_{=0} \quad (34)$$

$$= -\frac{\partial}{\partial \vec{r}_\xi} \vec{\delta}^T \cdot (\vec{\gamma} \times \vec{\alpha}) - \frac{\partial}{\partial \vec{r}_\xi} \vec{\delta}^T \cdot (\vec{\alpha} \times \vec{\beta}) - \frac{\partial}{\partial \vec{r}_\xi} \vec{\delta}^T \cdot (\vec{\beta} \times \vec{\gamma}) \quad (35)$$

$$= -\frac{1}{|S_\delta|} \vec{\gamma} \times \vec{\alpha} - \frac{1}{|S_\delta|} \vec{\alpha} \times \vec{\beta} - \frac{1}{|S_\delta|} \vec{\beta} \times \vec{\gamma} \quad (36)$$

$$= -\frac{1}{|S_\delta|} (\vec{\alpha} - \vec{\gamma}) \times (\vec{\beta} - \vec{\gamma}) \quad (37)$$

So, overall:

$$\textcircled{3} = \sum_{(S_\alpha, S_\beta, S_\gamma, S_\delta, U_V, L_V) \in C} 2 \max(0, V_{\alpha\beta\gamma\delta}(\{\vec{r}_i\}) - U_V) \left\{ \begin{array}{l|l} \frac{1}{|S_\alpha|} (\vec{s}_\beta - \vec{s}_\delta) \times (\vec{s}_\gamma - \vec{s}_\delta) & \xi \in S_\alpha \\ -\frac{1}{|S_\beta|} (\vec{s}_\alpha - \vec{s}_\delta) \times (\vec{s}_\gamma - \vec{s}_\delta) & \xi \in S_\beta \\ \frac{1}{|S_\gamma|} (\vec{s}_\alpha - \vec{s}_\delta) \times (\vec{s}_\beta - \vec{s}_\delta) & \xi \in S_\gamma \\ -\frac{1}{|S_\delta|} (\vec{s}_\alpha - \vec{s}_\gamma) \times (\vec{s}_\beta - \vec{s}_\gamma) & \xi \in S_\delta \\ 0 & \text{else} \end{array} \right\}. \quad (38)$$

For $\textcircled{4}$, the derivation is analog save for the sign of the great amount of cases, which we extrude from the sum:

$$\textcircled{4} = - \sum_{(S_\alpha, S_\beta, S_\gamma, S_\delta, U_V, L_V) \in C} 2 \max(0, L_V - V_{\alpha\beta\gamma\delta}(\{\vec{r}_i\})) \left\{ \begin{array}{l|l} \frac{1}{|S_\alpha|} (\vec{s}_\beta - \vec{s}_\delta) \times (\vec{s}_\gamma - \vec{s}_\delta) & \xi \in S_\alpha \\ -\frac{1}{|S_\beta|} (\vec{s}_\alpha - \vec{s}_\delta) \times (\vec{s}_\gamma - \vec{s}_\delta) & \xi \in S_\beta \\ \frac{1}{|S_\gamma|} (\vec{s}_\alpha - \vec{s}_\delta) \times (\vec{s}_\beta - \vec{s}_\delta) & \xi \in S_\gamma \\ -\frac{1}{|S_\delta|} (\vec{s}_\alpha - \vec{s}_\gamma) \times (\vec{s}_\beta - \vec{s}_\gamma) & \xi \in S_\delta \\ 0 & \text{else} \end{array} \right\}. \quad (39)$$

Both terms concerning the chiral error can be summarized as follows:

$$\begin{aligned} \textcircled{3} + \textcircled{4} &= \sum_{(S_\alpha, S_\beta, S_\gamma, S_\delta, U_V, L_V) \in C} 2 \max(0, V_{\alpha\beta\gamma\delta}(\{\vec{r}_i\}) - U_V) \{ \dots \} \\ &\quad - \sum_{(S_\alpha, S_\beta, S_\gamma, S_\delta, U_V, L_V) \in C} 2 \max(0, L_V - V_{\alpha\beta\gamma\delta}(\{\vec{r}_i\})) \{ \dots \} \\ &= \sum_{(S_\alpha, S_\beta, S_\gamma, S_\delta, U_V, L_V) \in C} 2 \{ \dots \} \left[\max(0, V_{\alpha\beta\gamma\delta}(\{\vec{r}_i\}) - U_V) - \max(0, L_V - V_{\alpha\beta\gamma\delta}(\{\vec{r}_i\})) \right], \end{aligned}$$

which ought to save some time in an implementation.

2.3 Dihedral error terms

Finally, we consider $\textcircled{5}$. The chain rule yields:

$$\frac{\partial}{\partial \vec{r}_\xi} \sum_{(\dots) \in D} \max^2(0, h(\phi)) = 2 \sum_{(\dots) \in D} \max(0, h(\phi)) \frac{\partial}{\partial \vec{r}_\xi} \max(0, h(\phi)), \quad (40)$$

where we have abbreviated the dihedral expression in the maximum function with $h(\phi)$. We can drop the wrapping maximum function in the derivative since the first maximum function adequately covers the case $h(\phi) < 0$. Employing the chain rule yet again, we note $\frac{\partial}{\partial \vec{r}_\xi} h(\phi) = \frac{\partial h}{\partial \phi} \frac{\partial \phi}{\partial \vec{r}_\xi}$. The first part is:

$$\frac{\partial h}{\partial \phi} = \frac{\partial}{\partial \phi} \underbrace{\left| \phi_{\alpha\beta\gamma\delta}(\{\vec{r}_i\}) + \begin{cases} 2\pi & \phi < \frac{U_\phi + L_\phi - 2\pi}{2} \\ 0 & \\ -2\pi & \phi > \frac{U_\phi + L_\phi + 2\pi}{2} \end{cases} - \frac{U_\phi + L_\phi}{2} \right|}_{w(\phi)} - \frac{U_\phi - L_\phi}{2} = \text{sgn}(w(\phi)). \quad (41)$$

The individual derivatives $\frac{\partial \phi}{\partial \vec{s}_\alpha}, \frac{\partial \phi}{\partial \vec{s}_\beta}, \frac{\partial \phi}{\partial \vec{s}_\gamma}, \frac{\partial \phi}{\partial \vec{s}_\delta}$ are given as:

$$\vec{f} = (\vec{s}_\alpha - \vec{s}_\beta), \quad \vec{g} = (\vec{s}_\beta - \vec{s}_\gamma), \quad \vec{h} = (\vec{s}_\delta - \vec{s}_\gamma), \quad \vec{a} = \vec{f} \times \vec{g}, \quad \vec{b} = \vec{h} \times \vec{g} \quad (42)$$

$$\begin{aligned} \frac{\partial \phi}{\partial \vec{s}_\alpha} &= -\frac{|\vec{g}|}{\vec{a}^2} \vec{a} \\ \frac{\partial \phi}{\partial \vec{s}_\beta} &= \frac{|\vec{g}|}{\vec{a}^2} \vec{a} + \frac{\vec{f}\vec{g}}{\vec{a}^2|\vec{g}|} \vec{a} - \frac{\vec{g}\vec{h}}{\vec{b}^2|\vec{g}|} \vec{b} \\ \frac{\partial \phi}{\partial \vec{s}_\gamma} &= -\frac{|\vec{g}|}{\vec{b}^2} \vec{b} + \frac{\vec{g}\vec{h}}{\vec{b}^2|\vec{g}|} \vec{b} - \frac{\vec{f}\vec{g}}{\vec{a}^2|\vec{g}|} \vec{a} \\ \frac{\partial \phi}{\partial \vec{s}_\delta} &= \frac{|\vec{g}|}{\vec{b}^2} \vec{b} \end{aligned} \quad (43)$$

The individual membership cases are thus:

$$\textcircled{S_i} \quad \frac{\partial \phi}{\partial \vec{r}_\xi} = \frac{\partial \phi}{\partial \vec{s}_i} \frac{\partial \vec{s}_i}{\partial \vec{r}_\xi} = \frac{1}{|S_i|} \frac{\partial \phi}{\partial \vec{s}_i} \quad (44)$$

And altogether:

$$\textcircled{5} = 2 \sum_{(\dots) \in D} \max(0, h(\phi)) \text{sgn}(w(\phi)) \left\{ \begin{array}{l|l} \frac{1}{|S_\alpha|} \frac{\partial \phi}{\partial \vec{s}_\alpha} & \xi \in S_\alpha \\ \frac{1}{|S_\beta|} \frac{\partial \phi}{\partial \vec{s}_\beta} & \xi \in S_\beta \\ \frac{1}{|S_\gamma|} \frac{\partial \phi}{\partial \vec{s}_\gamma} & \xi \in S_\gamma \\ \frac{1}{|S_\delta|} \frac{\partial \phi}{\partial \vec{s}_\delta} & \xi \in S_\delta \\ 0 & \text{else} \end{array} \right\} \quad (45)$$