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5 1 STRAW tubes

6 The option for STRAW tubes is similar as in NA62 experiment with one main
7 difference – the length is twice longer(5m versus 2.1m).

8 The next table. 1 describe STRAW tube options.

Parameter name	Value
wire	$30\mu m$ gold-plated Tungsten
straw length	5m
Voltage	1750V
inner tube radius	9.8 mm
wire medium density	19.3 g/cm^3
Wire tension	$\sim 90 \text{ g}$
Working tube gas mixture	Ar70% CO ₂ 30%

Табл. 1: STRAW tube parameters

9 2 Signal

10 Computer program Garfield [1] is designed for detailed simulation of two- and
11 three-dimensional drift chambers. So we will perform STRAW tube studies using
12 this program.

13 Charged particle create elector-ion pairs wile traverse the drift tube. Electrons
14 under affecting the electric field drift to the wire anode 1. During the travel they
15 increase their energy and invoke avalanche. Therefore they produce a measurable
16 signal.

17 Initial electrons drift to the wire due to the electrical field between the wire
18 and the tube wall. Electrons ionize gas molecules due to the high electric field
19 around the wire, especially near the wire when the electric becomes very strong.
20 Subsequently readout electronics process the signal induced on the wire.

21 The event registers if signal reach some a threshold voltage (Fig. 2). So the
22 value of threshold is a key factor on the way of searching optimal setting for
23 signal processing procedure.

24 We have to set threshold as low as possible but enough far from noise to
25 achieve highest value of relation true/false detected track and tube efficiency.

26 A variation of the signal height introduces a variation in the time when the
27 signal passes the threshold and is considered to be the main contribution to the
28 STRAW tracker resolution.

29 In the track reconstruction software(GARFIELD [1] an effective TR-relation
30 is used. It only describes the relation between the drift time and the distance
31 from the track to the wire, which differs from the distance to the ionization
32 cluster. The shape of the TR-relation is defined by the drift velocity of the
33 ionization cluster inside the straw. The electric field increases towards the wire,

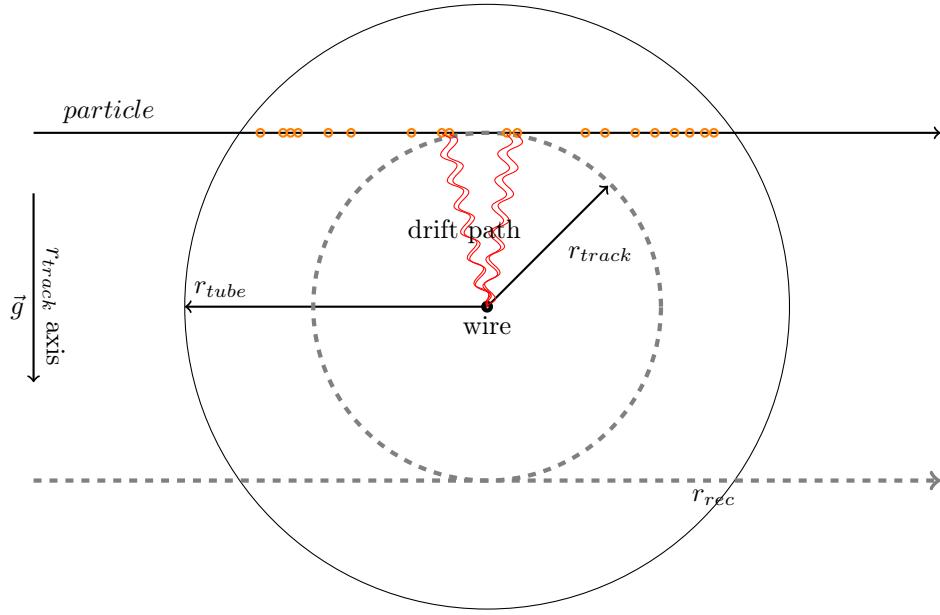


Рис. 1: Schematic view of a particle passing the straw and producing ionization clusters. The ionization cluster electrons drift to the wire and induce the signal. Only the earliest signal is detected. The closest distance from the track to the wire, r_{track} , and radius of the straw, $r_{tube} = 2.45\text{mm}$, are also indicated.

³⁴ leading to a non linear TR-relation. Currently almost parabolic dependence is
³⁵ used, and easily can be fitted by function (??).

³⁶ The drift time versus the unbiased distance distribution and the result of the
³⁷ fit are shown in Fig. 13a. Noise hits under the main distribution, i.e. at earlier
³⁸ times, are due to primary or secondary particles (δ -rays) passing the straw at a
³⁹ closer distance to the wire, consequently producing an earlier signal.

⁴⁰ Muon μ was chosen as test particle for simulation with energy 1GeV. You
⁴¹ can see some of typical tracks from the μ through the tube Fig.7a,7b. Initial
⁴² clusters along the track are marked by orange points on the figure.

⁴³ 2.1 Leakage noise

⁴⁴ Every time we deal with different kind of noise. Basically it is noise from leakage
⁴⁵ current through readout electronics.

⁴⁶ As will be discussed further we analyse not the current invoked by particle
⁴⁷ but the output voltage from amplifier. In GARFIELD we able convolute input
⁴⁸ current $I(t)$ with electronic response function (1):

$$f_{resp} = A \cdot (e^{-t/0.005} - e^{-t/0.030}) \quad (1)$$

⁴⁹ Noise is very important for every calculations and it makes bit impact on

50 straw precision and straw efficiency. So we can't rely on results until we receive
 51 signal and noise from real STRAW tube prototypes.

52 Convolution smooth input current. Experiments that used to drift tubes
 53 (such as ??(advise of Iouri Guz)) say that the noise should have gauss distri-
 54 bution with RMS equal to a amplitude of signal from 2000 electron in the tube -
 55 electric noise charge (ENC). (**This part should be clarified more precisely. Would**
be good to include some results from noise measurements from STRAW tube
samples. In fig.2 you can see deposition from noise marked by blue line.

56 On the figure fig.2 The time stamp $Time = 0$ correspond to the time muon
 57 cross tube. The convolution function smooths and spreads input current. It
 58 mean that the output voltage in GARFIELD does not contain part of signal
 59 before hit event time stamp.

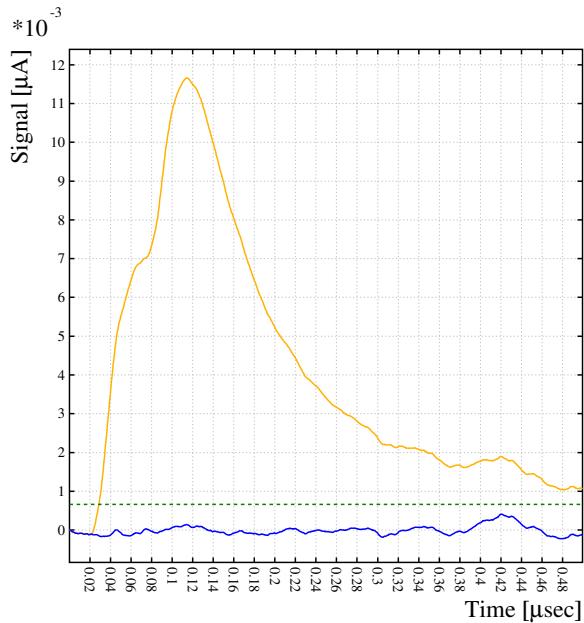


Рис. 2: Example of output signal $V(t)$ after convolution(front-end electronics) from central track(yellow line). The noise component of the same signal depicted by separate blue line. Grin dashed line is a threshold for trigering drift time and equal to 5σ of noise distribution.

62 2.2 STRAW efficiency

63 The interaction of charge particle with gas molecules nave probabilistic nature.
 64 For short distance tracks(somewhere at the tube periphery) the probability of
 65 tracks that do not produce any electron/ion pair becomes significantly high.

66 The number of produced ionization clusters directly affects the hit efficiency
 67 profile. [2] Smaller ionization length increase hit efficiency because of more
 68 ionization clusters per length unit are producing. In GARFIELD we can easily
 69 calculate amount of clusters per track. In fig. 4b you can see a distribution of
 70 number of clusters per central track for our STRAW tube. It mean that straw
 71 efficiency will be lower at the tube wall(see fig. ??).

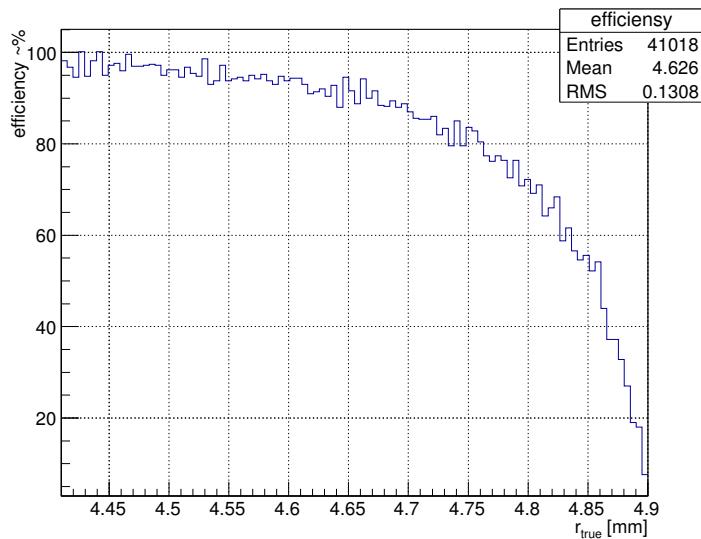


Рис. 3: Straw tube efficiency. Result of homogeneous penetrating periphery of tube by 50k events(scaled down by factor of 5. $\frac{50k \text{ events}}{100\text{bin}} = 500 \frac{\text{eventst}}{\text{bin}}$).

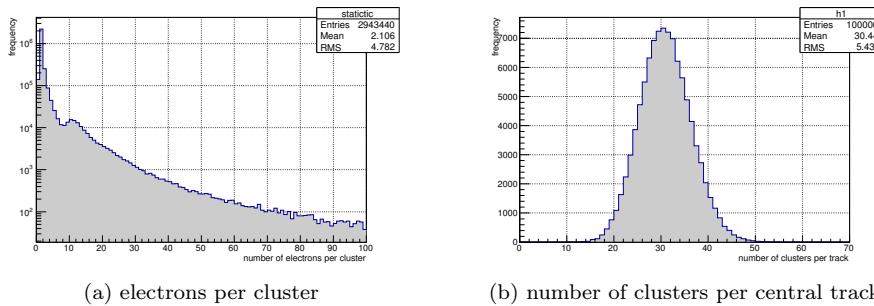


Рис. 4: There are big amount of graphs. So I'm trying to pair it. ere we can write something if needed. Some common description of (a) and (b) figures?

72 From the figure ?? we can conclude that the efficiency of tube is 100% almost
 73 in whole region covered by tube except pre wall region which is quite small.

⁷⁴ Increasing the gas mixture density or increasing the tube radius for the same
⁷⁵ gas density can increase tube efficiency. Have to check this in feature studies.

⁷⁶ 3 Gain

⁷⁷ If multiplication occurs, the increase of the number of electrons per path ds is
⁷⁸ given by

$$dN = N\alpha ds \quad (2)$$

⁷⁹ The coefficient α is determined by the excitation and ionization cross sections
⁸⁰ of the electrons that have acquired sufficient energy in the field. It also depends
⁸¹ on the various transfer mechanism and electric field E and increases with the
⁸² field because the ionization cross-section goes up from threshold as the collision
⁸³ energy ε increases. As we can suppose the coefficient α is of big amount of
⁸⁴ parameters.

⁸⁵ The amplification factor G on a wire(that is more interesting for us) is given
⁸⁶ by integrating (2) between the point s_{min} where the field is just sufficient to
⁸⁷ start the avalanche and the wire radius a :

$$G = N/N_0 = \exp \int_{s_{min}}^a \alpha(s) ds \quad (3)$$

⁸⁸ GARFIELD can provide us by amplification factor G for any point of the
⁸⁹ tube(because G is coordinate dependent magnitude). The amplification factor
⁹⁰ is equal almost in whole tube space except neighbourhood near the wire because
⁹¹ electric field becomes significantly high only near the wire (see figs 5a, 5b). When
⁹² the wire is shifted from the center of the cube the electric field in area close to
⁹³ the wire is the same as in centered state. So the amplification factor G is quite
⁹⁴ similar in both cases.

⁹⁵ Implementation of gain value calculation is not so reliable in GARFIELD(especially
⁹⁶ fortran version). So we can reach better results using Garfield++ (which is newer
⁹⁷ and take into consideration more effects).

⁹⁸ On the Fig. ?? you can see that the gain $G(V)$ have precisely exponential
⁹⁹ dependence. This is frankly does not inspire confidence. The difference can be
¹⁰⁰ up to 100% (as Rob Veenhof - creator of GARFIELD [1] said).

¹⁰¹ 4 Wire sagging

¹⁰² Easy to predict that the displacement of the wire invokes distorting an electric
¹⁰³ field(see figs 5a,5b) and drift path for electrons/ions inside the tube(see fig.7a
¹⁰⁴ and fig.7b). The rt-relation for track reconstruction directly depend on the wire
¹⁰⁵ position in the tube. So rt-relation lose it's previous symmetry(see next sections).

¹⁰⁶ The direction of sagging is unpredictable when the wire is centered and the
¹⁰⁷ straw has vertical orientation. Impact of gravitation field into the wire does

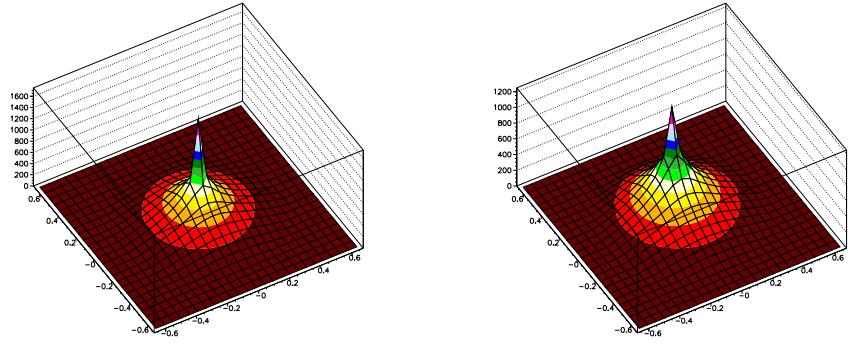


Рис. 5: Electric field intensity map for different wire position in the cube calculated in GARFIELD software. Conditions for those plots are described in table 1

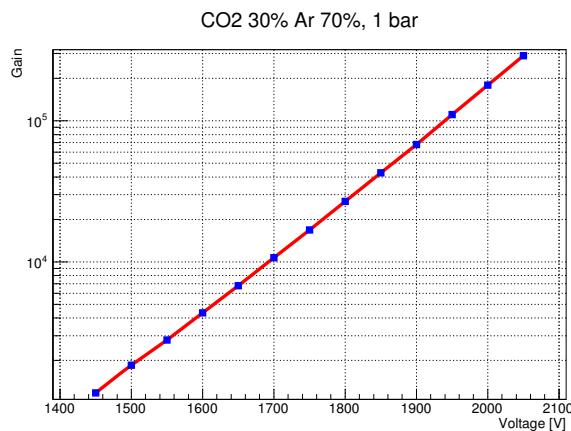


Рис. 6: Dependence of the gain of the voltage applied to the wire. The rest of STRAW tube settings you can find in table 1. need add gain(v) graph for shift wire STRAW tube for comparison

108 not make any effect in this state. But we can avoid this ambiguity by setting
109 straws horizontally. This condition is necessary to make track reconstruction
110 possible. Even when strung with a pulling force T close to the breaking limit,
111 wires in several metre long tubes will experience a gravitational sag that is large
112 in comparison with the achievable accuracy of drift tubes.

We estimate significant wire sagging (by comparison to the tube radius) because of wire attracts to the tube under affecting of gravitation and electric

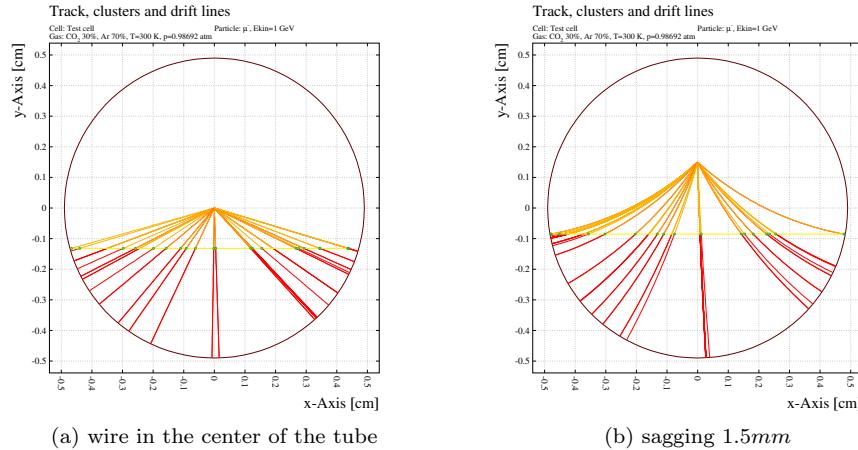


Рис. 7: An example of tracks from the on the tube for different position of the wire from GARFIELD simulations. Initial clusters marker by green. Drift lines for electrons marked by yellow, ions – red lines.

115 field force.

You can see a profile of wire sagging of 5m length wire in 1cm diameter straw tube and 1750V voltage on the fig.8 calculated in GARFIELD software [1].

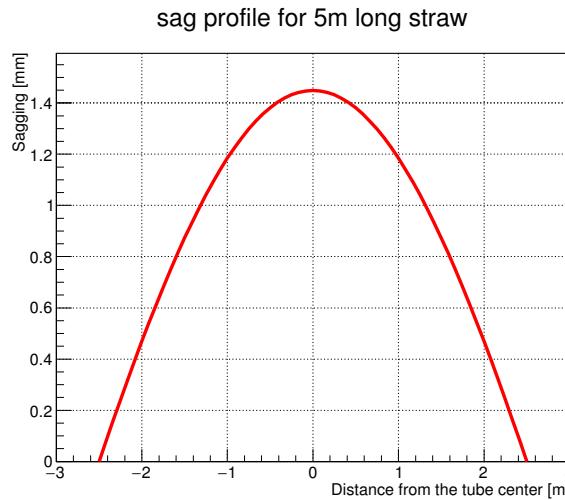


Рис. 8: Wire sag profile under electric and gravitation field calculated in GARFIELD-ELD. All options for this straw system are described in table 1.

The calibration of STRAW tube with sagged wire is more difficult by compari-

120 son to the mode without sagging.

121 Variation of wire tension, wire radius should be taken into account as high
122 affect factor for sag value.

123 5 Sag estimation

124 In this section we have to find out method for assessing sagging. This is key
125 step that makes track reconstruction procedure possible.

126 At first we have to think on data we can use for such kind of calculations.
127 Much attractable information we can extract from drift time distribution.

128 The wire sags under electric and gravitation force. Therefore the sag value
129 is differ along the tube(fig 8). But we can separate collected data for different
130 position along the tube. STRAW tube detector consist of several parallel layers
131 of tubes at some angle to each other. So we can easily fix longitudinal position for
132 tracks that cross several crossed tubes(at least two). Collimation is also possible
133 via scintillator triggering before and after STRAW tube.

134 Lets say we can install our STRAW tube into homogeneous particle flow and
135 save drift time distribution for some narrow section of the tube. These distri-
136 butions are different from each other(see example on Figure ??). The difference
137 between diagrams increasing with sag difference. So it is good possibility for sag
138 calibration.

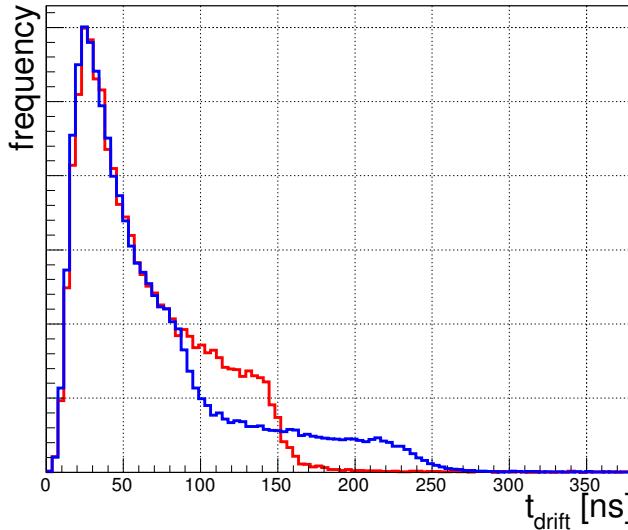


Рис. 9: Drift time distribution for a homogeneous irradiation with a centered wire (red) and for a wire offset of 0.9 mm (blue).

139 Then we have to bind each drift time distribution with appropriate sag value.

140 This is part of laboratory work when sag profile measurements can be performed
141 via optical method prior to the exposition.

142 Distributions on graph ?? contain GARFIELD simulations for some certain
143 wire(not for section of sagged wire)because of GARFIELD can handle only two-
144 dimensional tasks.

145 Lets say we have an equipment for scanning the tube to measure wire sagg-
146 ing profile. After profile measurements we divide our tube into sections. Wire
147 position within separate section should be within desired precision.

148 So we need divide our tube into 57 sections (see figure 10) if maximum of
149 wire offset(at the center of the tube) is equal to $1.45mm$ and desired precision
150 is $50\mu m$.

$$N_{halftube} = \frac{1.45mm}{50\mu m} = 29; \quad (4)$$

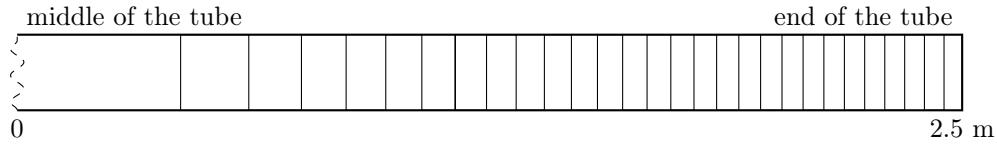


Рис. 10: Tube sectioning. Sag value at the tube center is $1.45mm$. Difference of wire sag value from section to section is $50\mu m$

151 Then we need an exposition of sufficient number of events for every of secti-
152 ons(at least 50k events). There can be troubles time of exposition time because
153 square of sections at the end of the tube is quite small. So the time of exposition
154 of distant sections will be inversely much longer.

155 The next step is to find dependence of dt-distribution shape with wire offset.
156 The point that we can evaluate matching between histograms via χ^2 criteria.
157 As we can see in the figure 11a the comparison of χ^2 has smooth dependence
158 across increasing of wire offset for high statistic histograms.

159 First steps for sag estimation are:

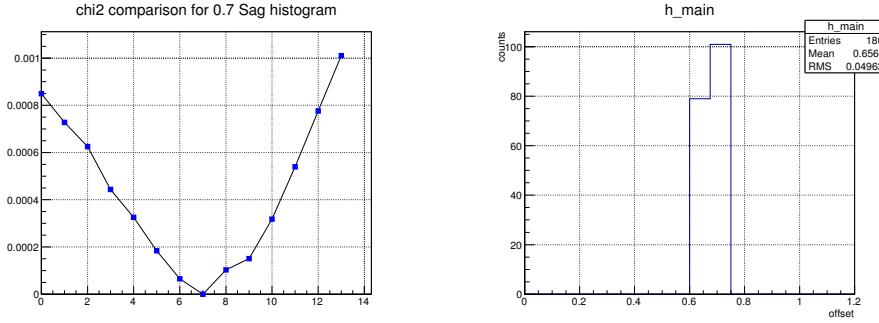
- 160 1. measure wire sag profile via optical method;
- 161 2. make a sectioning for wire sag profile;
- 162 3. collect enough amount of events for every of dt-distribution and save this
163 core distribution for further comparisons.
- 164 4. measure dt-distribution for new drift tube section that is subject of study.
- 165 5. calculate χ^2 criteria for this current dt-distribution with each of core di-
166 stribution.

167 **5.1 Finding most probable value of wire displacement for
168 certain point of the tube**

169 **5.2 Raw method**

170 The simplest method to find S is to equate it to the corresponding value of best
171 matched core DT-histogram.

172 On the figure 12b you can see distribution of such kind of reconstruction.
173 Even for 5k events td-distribution in this case the precision can be quite high(\sim
174 $50\mu m$).



(a) Series of χ^2 of comparison $0.7mm$ sag core td-distribution with each each of core histograms. 14 core histograms for sag dipason $0 \dots 1.3mm$ with step of $100\mu m$

(b) Distribution of wire offset reconstruction from 180 series 5k events each. 50k events for core template histograms. True bias is $0.063mm$. 1 bin = $0.1 mm$.

Рис. 11: Wire position(displacement) reconstruction

175 **5.3 "Minimum of χ^2 as linear approximation"**

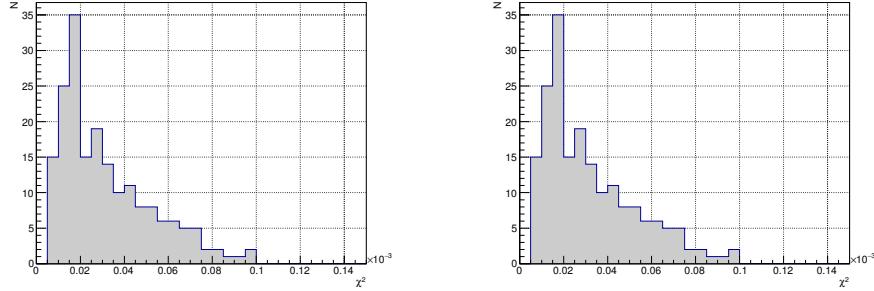
176 But what is most probable value of wire displacement S in this case? Probably
177 somewhere between them. So can we go in more clever way to reach better
178 result? Probably yes.

179 If dependence of χ^2 criteria of wire displacement S for near to the true
180 position region is linear(that certainly is not a true, but as first approximation)
181 than we can easily find this intermediate value of wire displacement.

182 There we proceed in two steps. The first is raw estimation of wire dis-
183 placement as in above mentioned method.

184 For the second step we need to know some additional estimations. The first
185 question is how small can be χ^2 it our case? Lets fix statistic on 50k events
186 for one DT-distribution. This value should a bit depend for different S . But
187 for now lets consider that it is a constant value. From the figure 12a you can
188 see distribution of χ^2 from comparison of 20 DT-distribution¹ for $S = 0.7mm$.
189 Mean value + RMS of distribution is $5.3 * 10^{-5}$. So if some of the χ^2 is higher
190 than this threshold than we go for second stage.

¹ pair comparison give us $C_{20}^2 = \frac{20!}{2^{10}18!} = 190$ different combinations



(a) χ^2 distribution of comparison of 20 DT-distributions diagrams each other ($C_{20}^2 = 190$ combinations)

(b) χ^2 distribution of comparison DT-distribution histograms 0.7mm S vs 0.8mm S here I estimate much narrow distribution. Image to be inserted

Рис. 12: χ^2 distributions.

191 On the figure 12b you can see distribution wire sag calculation for 180 histograms
 192 with 5k events statistic. Precision in this case $\sim 50\mu$. The algorithm
 193 of sag estimation is pretty simple: wire offset value equal to the best match
 194 between *test* and *core* histogram.

195 After we know sag value at some points of the tube or every where we can
 196 make one awesome collective analysis. The smoothing of wire offset value along
 197 the tube will give us much more precision results. Fitting of sag value at every
 198 point of $s(l)$ by some parabolic function should provide us the best results.

199 5.4 Wire sagging profile finding

200 to be completed ...

201 Here I would like to put total plot of wire sag profile and compare reconstructed
 202 profile with true profile.

203 6 Track reconstruction

204 The time between the track hit time stamp and the signal rising edge is a
 205 measure of *drift time* of these electrons. The relation between the *drift time*
 206 and the distance from the track to the center of the tube(wire while no sag for
 207 centered wire) is called *drift time - distance relation* or *tr-relation*.

208 The drift time t is a function of track position relative to the wire(so it's
 209 means the track position) and electric field along the drift trajectory.

210 Assumed that the working position for straws will be parallel to the particle
 211 bunch, and acceptance of particle spreading will not be significantly big. So
 212 tracks will be collinear each other within every separate STRAW tube unit.

213 Summing the above mentioned we have one dimension task – reconstruct

²¹⁴ tracks on vertical axis² (see examples of outcome tr-distribution $t = t(r, s = 0)$
²¹⁵ in Fig.13a and Fig.13b) even the wire sagging. Sagging will be always down
²¹⁶ thanks to gravitation force \vec{g} .

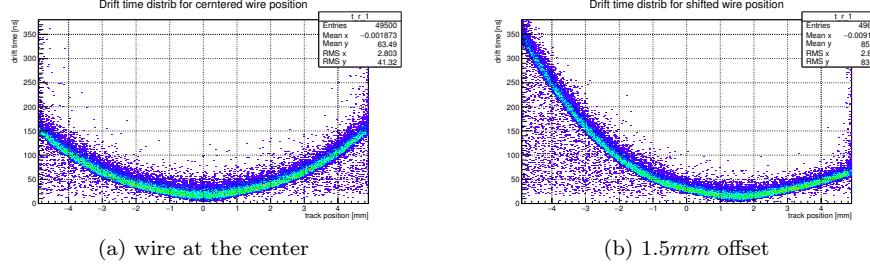


Рис. 13: Distribution of drift time t_{drift} as function of track position r_{track} relatively to the tube center

²¹⁷ The rt-relation is differ along the tube because different wire position s . Thus
²¹⁸ we have for the drift time

$$t_{drift} = t_{drift}(r_{track}, s) \quad (5)$$

²¹⁹ The idea to STRAW tube is to find the inverse dependence

$$r_{track} = r_{track}(t_{drift}, s) \quad (6)$$

²²⁰ From the section "Sag estimation" we can find sag profile for straw. Therefore
²²¹ the rt-calibration becomes 1 dimension less:

$$r = r(t, s = const) \quad (7)$$

²²² 6.1 How drift time resolution depend on wire offset?

²²³ Distorting of electric field inside the tube invoked by wire displacement from the
²²⁴ center position will make an effect on drift time. Here we are going to estimate
²²⁵ magnitude of drift time change.

²²⁶ As was noted above we make a binning for our data along the r_{track} (fig. 13a,
²²⁷ 13b). The resolution at every bin is RMS of every bit digram (fig. ??).

²²⁸ We are dealing with probabilistic nature of clustering that spread rt-relation
²²⁹ from thin line. The leakage noise is also present in calculation but the effect of
²³⁰ it is not very high (especially in this calculation).

²³¹ Every plot of output current (see fig. 2) consist of 1000 equidistant frames.
²³² The threshold is set to 5σ of noise. Leakage noise make effect on drift time
²³³ measurements in case its amplitude becomes higher than threshold value in
²³⁴ range from $t = 0$ to $t = t_{drift}$. At five-sigma there is only one chance in nearly
²³⁵ two million that a random fluctuation would yield the result. The drift time for

²An example of single track reconstruction which explains the approximate procedure of reconstruction you can see on Fig.1

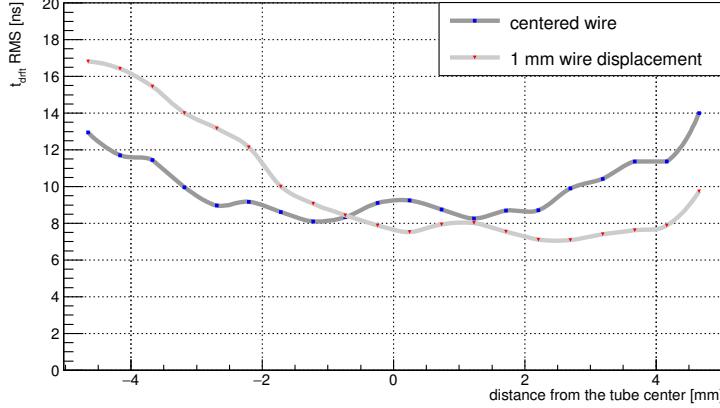


Рис. 14: Resolution of drift time as a function of distance from the wire.

tracks close to the tube edge can be up to 150ns and 300ns in case wire displaced. The probability to meet noise above threshold value is less than 0.02%.

Another source of noise points on tr-distribution comes from δ -electrons that cause secondary ionisation in tube volume. The impact do only those electrons which are emitted in the direction of the wire(see example on fig.15a).

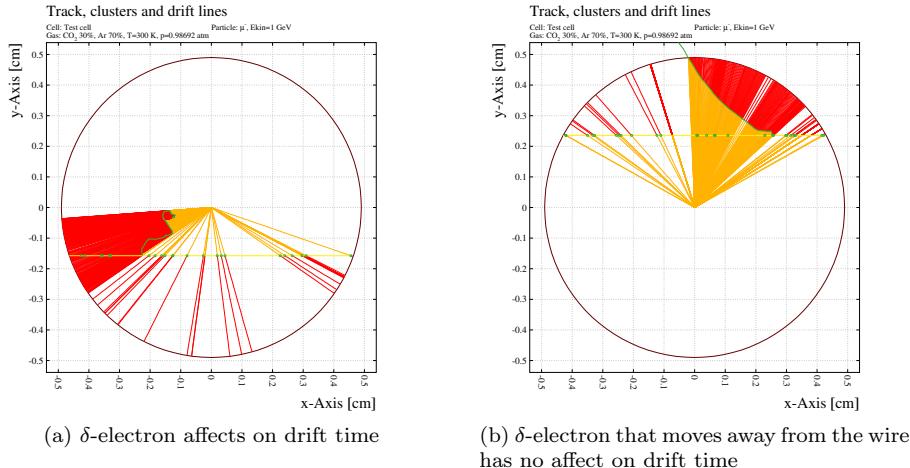


Рис. 15: Garfield simulation with δ -electron presence. Red lines - ion trajectory, yellow - electrons. Trajectory of δ -electron marked by green curve line.

The number of events out of TR-ralation because of δ -electrons is quite small. Especially percentage of events where δ -electrons make effect on drift time is less than 1% of total number of events in GARFIELD simulations.

244 Tube wall is very thin but particle still can cause δ -electrons when crossing
 245 it. GEANT4 studies show that such kind effect also presents in interaction of
 246 muon with tube volume, and percentage of events with δ -electron that affect
 247 drift time even less than 0.2%.

248 6.2 Finding of rt-relation

249 The rt-relation depict relation between drift time and track position. The idea is
 250 to find the best fit of give data to achieve higher resolution and avoid systematic
 251 errors.

252 The problem that we have to minimize influence of noise while fit. One
 253 suppose that the noise have approximately homogeneous distribution of points
 254 that locates below the main line of distribution. Consequently we can filter it
 255 by fitting only points from regions with local point density higher than some
 256 threshold value. Another way is to make a binning our distribution along the
 257 track position and fit every 1-D histogram by Gaussian. The fit points of Gaussian
 258 mean values by fit function.

259 Nevertheless our data contain very small amount of "non-track"points.

260 TR-relation is asymmetry relatively to the $r = 0$ almost in all cases except
 261 wire in the center of the tube. Therefore we have to calibrate for every of
 262 branches. It means we need to find two track positions for every of drift ti-
 263 me value and reject one of them in further data processing stages.

264 In previous section we found way to measure wire sag profile. So we can
 265 use this trick in present stage for separating data into "right" and "left" branch.
 266 Every of branches we will calibrate separately.

267 Lets suppose we can fit every of tr-diagram by pair of analytic fit function
 268 (8):

$$t(r_{track}) = e^{a_0 + a_1 r_{track}} \quad (8)$$

269 If the figure ?? you can see tr-relation. Fitting is not perfect because of using
 270 simple fit function template (8). But we will use reverse to the (8) relation,
 271 because we have to find r_{track} from known t_{drift} . We can do it be because the
 272 aim of this studies is not a precision calibration but global evaluation affect of
 273 wire sagging into total result.

274 As you can see in the figure ?? red fit line does not cover whole drift time
 275 spectre. So events with drift time less than covered range(less than $\sim 20ns$)
 276 counts as track through the wire:

$$r_{track}(t_{drift} < t_{min}) = r_{wire\ pos} \quad (9)$$

277 where $t_{min} = \min(t_{drift}(r_{track}))$, $r_{track} = \overline{(-r_{tube}, r_{tube})}$. Respectively tracks
 278 with drift time higher than maximum of fit function range artificially counts as
 279 tracks with near tangents to the tube position $r_{track} = \pm r_{tube}$ (because efficiency
 280 decreases near the tube wall down to 20%).

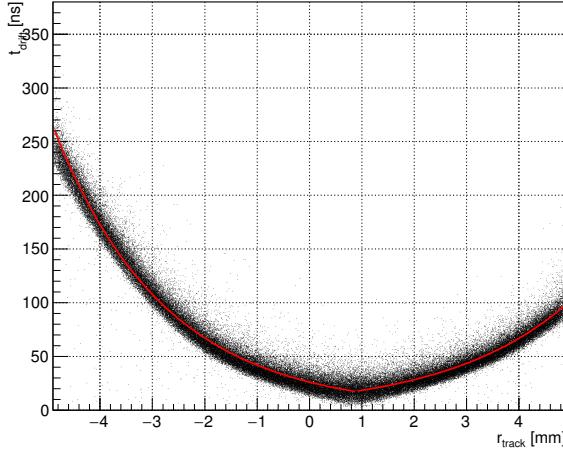


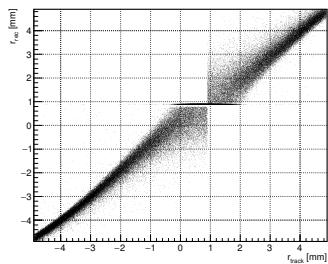
Рис. 16: TR-relation fitting for 0.9mm wire offset value

281 6.3 Track reconstruction precision

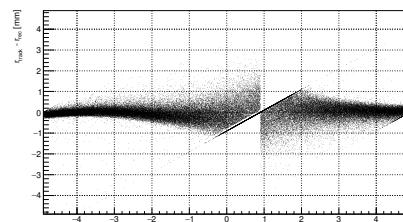
282 Obviously precision is head factor when during we decide design of detector.

283 The STRAW tube tracker should be as light as possible to avoid multiple
284 scattering on structural components of detector. But design should be changed
285 within reason if precision suffers from this³.

286 How precision of track reconstruction depends on wire position(wire di-
287 placement)?



(a) reconstructed track position r_{rec} as
function of true track position r_{track}



(b) $r_{track} - r_{reconstructed}$ as function of r_{track}

Рис. 17: Distributions of matching of track position to their reconstructed value.

³Especially design with no sagging works well for experiment NA62 [1]. But they have more than 2 times shorter straw when tube have insert in the middle of the tube. So sagging becomes negligible in this case.

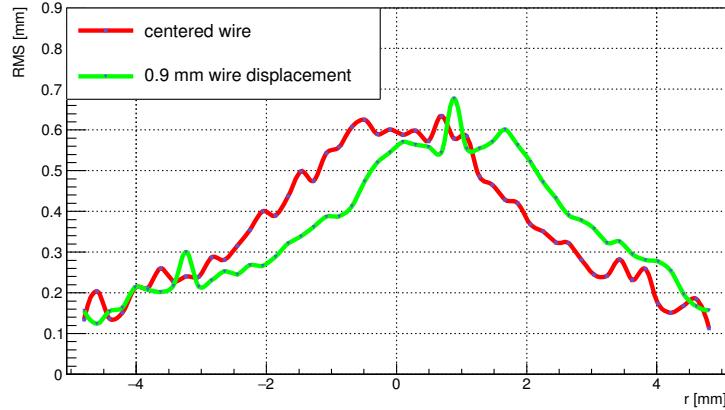


Рис. 18: Comparison of track reconstruction precision for two wire position. Value of precision at every point means RMS of data sample near corresponding track position r . Red line corresponds to the centered wire position, green line – to the 0.9mm sagged wire position.

288 As you can see on figure ?? there are no significant difference of track
 289 reconstruction precision between two mode of wire location despite of the increasing
 290 drift time for displaced wire position (with almost factor of two). The highest
 291 resolution ($\sim 0.1\text{mm}$) near the tube wall and worst value $\sim 0.6\text{mm}$ is near
 292 the wire because the clustering effect. Higher gas pressure should resolve this
 293 problem.

²⁹⁴ **Література**

- ²⁹⁵ [1] <http://garfield.web.cern.ch/garfield>
- ²⁹⁶ [2] thesis Kozlinskiy.pdf
- ²⁹⁷ [3] NA62 Technical Design Report from December 2010.